

Probing the connection between short gamma ray bursts and binary coalescing systems through gravitational waves

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21/04/2023

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Motivation for study

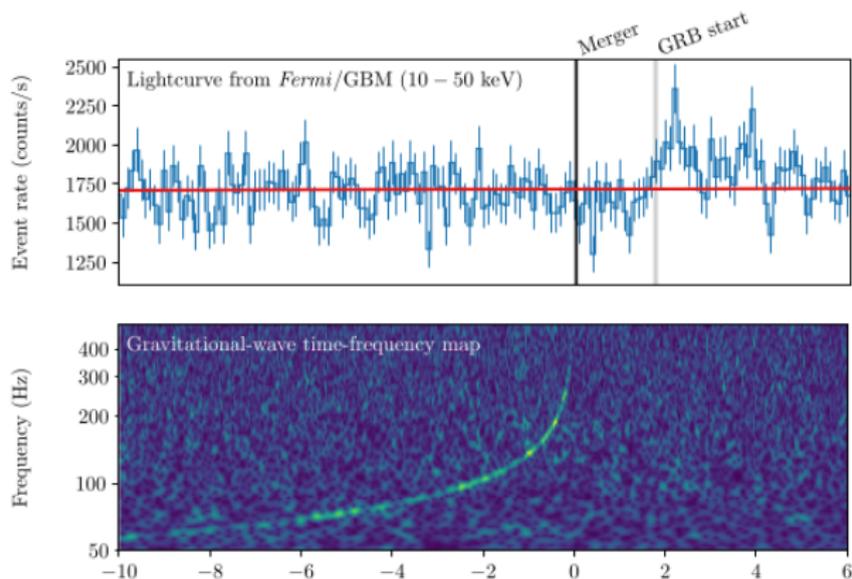
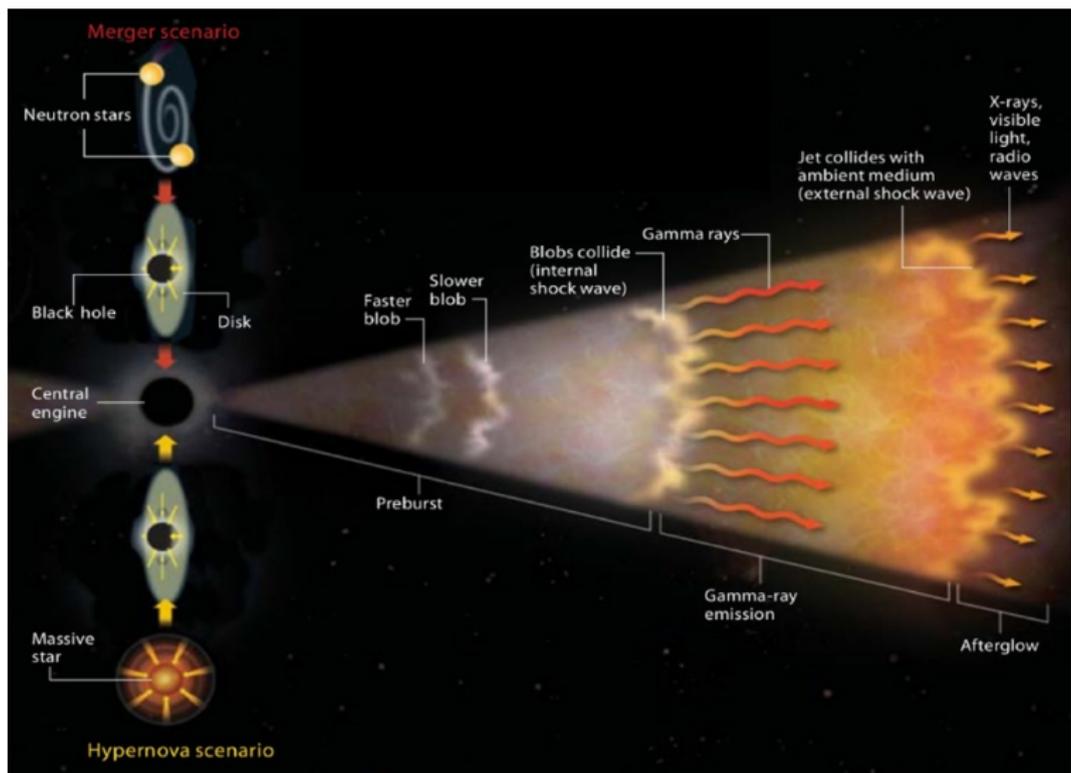


Figure: Fermi GBM light curve of GRB170817A above a time frequency map of GW170817 generated from LIGO Hanford and Livingston. A joint detection rate of $0.1 - 1.4 \text{ yr}^{-1}$ between LIGO and Fermi GBM was predicted. At LIGO's design sensitivity this climbed to $0.3 - 1.7 \text{ yr}^{-1}$ data.[1]

Motivation for study

- To date, the events of GW170817/GRB170817A has been the only joint detection of its kind so far.
- During LIGOS 2nd and 3rd observing runs O2 and O3, a second BNS merger GW190425 and Black Hole Neutron Star (BHNS) mergers GW200115_042309, GW200210_092254, GW190917_114636 were detected[2]. All of these events could be possible sources for a GRB. No EM counterpart for these events were detected.
- This study aims to find an explanation for the lack of joint detections through the O2 and O3 runs.

The Hypothesis



Gravitational waves

- Gravitational waves are travelling perturbations in spacetime caused by the acceleration of massive bodies. General relativity predicts the existence of 2 tensor polarisation modes:

$$h_+(t) = -\frac{1 + \cos^2 \iota}{2} \left(\frac{\mathcal{GM}}{c^2 D} \right) \left(\frac{t_c - t}{5\mathcal{GM}/c^3} \right)^{-1/4} \cos(2\phi_c - 2\phi(t - t_c; M, \mu))$$

$$h_\times = -\cos \iota \left(\frac{\mathcal{GM}}{c^2 D} \right) \left(\frac{t_c - t}{5\mathcal{GM}/c^3} \right)^{-1/4} \sin(2\phi_c + 2\phi(t - t_c; M, \mu))$$

- The GW strain as seen by a particular detector is given by

$$h(t) = h_+(t - t_c - t_0)F_+(\alpha, \delta, \Psi, t) + h_\times(t - t_c - t_0)F_\times(\alpha, \delta, \Psi, t) \quad (1)$$

¹reference [3]

Gravitational waves

- For short duration signal F_+ and F_\times are nearly constant. The GW strain seen by a particular detector can then be written as

$$h(t) = - \left(\frac{\mathcal{G}\mathcal{M}}{c^2 D_{eff}} \right) \left(\frac{t_0 - t}{5\mathcal{G}\mathcal{M}/c^3} \right)^{-1/4} \cos(2\phi_0 + 2\phi(t - t_0; M, \mu)) \quad (2)$$

- ϕ_0 is the termination phase which is given by the relation

$$2\phi_0 = 2\phi_c - \arctan \left(\frac{F_\times}{F_+} \frac{2 \cos \iota}{1 + \cos^2 \iota} \right) \quad (3)$$

- and D_{eff} is the effective distance given by

$$D_{eff} = D \left[F_+^2 \left(\frac{1 + \cos^2 \iota}{2} \right)^2 + F_\times \cos^2 \iota \right]^{-1/2} \quad (4)$$

²reference [3]

Analysing detector data

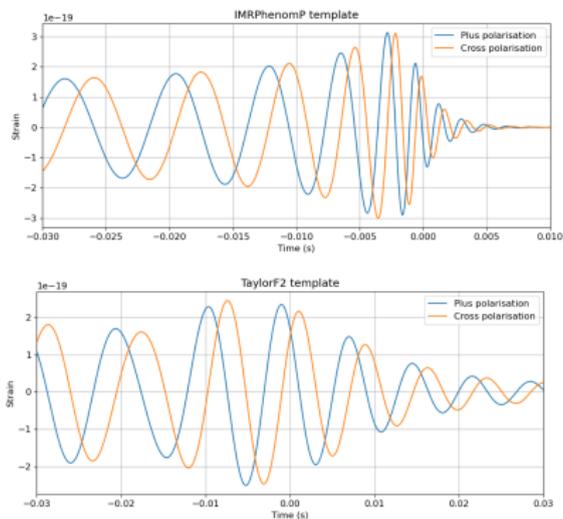


Figure: Gravitational waveform templates used in this study

Bayes theorem

- For a set of observations $\mathbf{d} = (d_1, \dots, d_n)$ and set of unknown parameters $\boldsymbol{\theta} = (\theta_1, \dots, \theta_n)$, the probability density of the values of $\boldsymbol{\theta}$ given the data \mathbf{d} is given by:

$$p(\boldsymbol{\theta}|\mathbf{d}) = \frac{L(\mathbf{d}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})}{\mathcal{Z}} = \frac{L(\mathbf{d}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})}{\int L(\mathbf{d}|\boldsymbol{\theta})\pi(\boldsymbol{\theta})d\boldsymbol{\theta}} \quad (5)$$

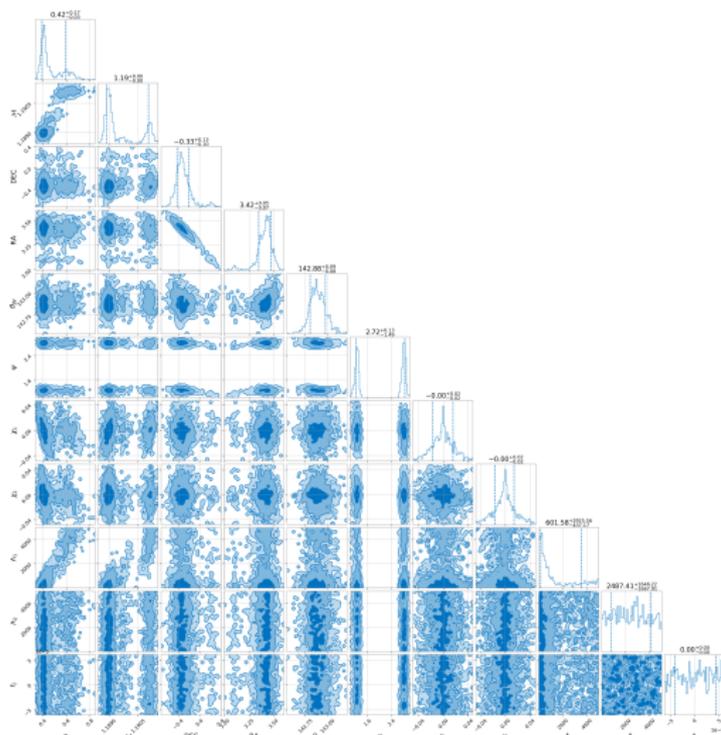
- where $L(\mathbf{d}|\boldsymbol{\theta})$ is the likelihood function of \mathbf{d} given $\boldsymbol{\theta}$. $\pi(\boldsymbol{\theta})$ is the prior probability density functions and \mathcal{Z} is the marginalised likelihood .
- By choosing a likelihood, a model for the GW is implicitly chosen. For example, a Gaussian likelihood for GW astronomy is given by

$$L(\mathbf{d}|\boldsymbol{\theta}) = \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left(-\frac{1}{2} \frac{(\mathbf{d} - h(\boldsymbol{\theta}))^2}{\sigma^2}\right) \quad (6)$$

Methodology

- For this study we performed Bayesian inference on the following GW events: GW170817, GW190425 (BNS events), GW190917_114636, GW200210_092254, GW2000115_042309 (BHNS events)
- We perform Bayesian inference using Bilby which is python based Bayesian inference library for GW astronomy [5]
- GW170817 has an observed EM counterpart GRB170817A. As a result, the inclination angle is well constrained. To test how effective pure GW analysis is using Bilby, we aim to obtain similar values for the inclination angle through pure GW analysis.
- In order to perform Bayesian analysis, we define a prior giving the distribution of the waveform parameters. Following convention, we set up two priors that represent a low spin and high spin case for the merger.

GW170817: Results (TaylorF2_Lowspin)



GW190425

- This is a BNS merger detected by a single detector (Livingston). The Hanford detector was offline during the event
- Component masses are $m_1 = 2.1 \pm_{0.4}^{0.5} M_{\odot}$ and $m_2 = 1.3 \pm_{0.2}^{0.3}$
- No trigger in the Virgo detector
- To analyse this event, the same set of waveform templates used on GW170817 were utilised (TaylorF2, IMRPhenomP, IMRPhenomD).
- 2 different low spin case priors were used with one having a uniform distribution between 0° and 90°
- For the high spin case a uniform distribution between 0° and 180° was chosen.

GW200115 and GW190917

- GW200115_042309 and GW190917 are BHNS events detected all through the LVC network with component masses.
- To analyse the signal a high or low spin prior case was not considered. Instead we considered a case of precessing spins with no consideration for tidal deformities in the neutron star.
- In this study we only made use of the gravitational waveform IMRPhenomP which is a waveform template that allows spin precession
- A new prior accommodating spin precession as well as the distance considerations was then set up

Summary of results

| Gravitational wave | Waveform | \mathcal{M} (M_{\odot}) | mass ratio |
|--------------------|-------------|-------------------------------|------------------------|
| GW170817 | IMRPhenomP | $1.20^{+0.0}_{-0.0}$ | $0.83^{+0.11}_{-0.11}$ |
| | TaylorF2 | $1.19^{+0.0}_{-0.0}$ | $0.42^{+0.17}_{-0.03}$ |
| | IMRPhenomD | $1.20^{+0.0}_{0.0}$ | $0.83^{+0.11}_{-0.11}$ |
| | LIGO result | 1.19 | (0.4, 0.8) |
| GW190425 | IMRPhenomP | $1.47^{+0.02}_{-0.0}$ | $0.43^{+0.40}_{-0.05}$ |
| | TaylorF2 | $1.47^{+0.00}_{-0.03}$ | $0.45^{+0.39}_{-0.05}$ |
| | IMRPhenomD | $1.47^{+0.02}_{-0.0}$ | $0.45^{+0.42}_{-0.06}$ |
| GW190917 | IMRPhenomP | $2.59^{+0.39}_{-0.17}$ | $0.33^{0.19}_{0.07}$ |
| | LIGO Result | $3.7^{+0.2}_{-0.2}$ | - |
| GW200115 | IMRPhenomP | $2.55^{+0.01}_{-0.00}$ | $0.34^{+0.15}_{-0.12}$ |
| | LIGO Result | $2.43^{+0.05}_{-0.07}$ | - |

Table: Chirp Mass and mass ratios estimated for events GW170817, GW190425, GW190917, GW200115

Summary of results

| Gravitational wave | Waveform | Low spin | High spin |
|--------------------|-------------|---------------------------------|---------------------------------|
| GW170817 | IMRPhenomP | $155.28^{\circ+15.99}_{-18.57}$ | $152.6^{\circ+18.65}_{-16.01}$ |
| | TaylorF2 | $142.88^{\circ+0.9}_{-0.8}$ | $152.41^{\circ+18.65}_{-15.82}$ |
| | IMRPhenomD | $155.21^{\circ+15.98}_{-18.56}$ | $155.57^{\circ+15.62}_{18.82}$ |
| | LIGO result | $146^{\circ+25}_{-27}$ | $152^{\circ+21}_{-27}$ |
| GW190425 | IMRPhenomP | $46.54^{\circ+28.96}_{-30.54}$ | $89.04^{\circ+63.11}_{-60.36}$ |
| | TaylorF2 | $44.64^{\circ+29.94}_{-28.45}$ | $98.03^{\circ+58}_{-63.39}$ |
| | IMRPhenomD | $46.09^{\circ+30.65}_{-32.38}$ | $87.95^{\circ+62.85}_{-63.11}$ |
| GW190917 | IMRPhenomP | $107.14^{\circ+34.37}_{67.18}$ | |
| GW200115 | IMRPhenomP | $64.74^{\circ+63.59}_{-40}$ | |

Table: Inclination angles estimated for events GW170817, GW190425, GW190917, GW200115

Discussions and conclusion

- The current set of results suggests that the binaries were orientated such that detection of the emitted GRB was not possible. The hypothesis still holds
- From existing joint detection predictions, 1 in 8 BNS mergers detected by the LVC network should have a GRB counterpart. Our current set of results is still in agreement with this prediction.
- Parameter degeneracies present the biggest challenges when it comes to parameter inference (e.g mass and spin degeneracy, luminosity distance and inclination angle degeneracy)
- to overcome these degeneracies, an independent observation of the parameter through a different messenger breaks the degeneracy
- In this study, without stricter constraints on either the luminosity distance or a smaller parameter space for the parameter of interest, reducing the uncertainty in the inferred value is not possible.

References I

- [1] B.P Abbott et al. “*Gravitational waves and gamma-rays from a binary neutron star merger: GW170817 and GRB 170817A*”. In: *The Astrophysical Journal Letters* 848.2 (2017), p. L13.
- [2] *GWTC event portal*. Last accessed 19 April 2023. 2023. URL: <https://gwosc.org/eventapi/html/GWTC/>.
- [3] B Allen et al. “*FINDCHIRP: An algorithm for detection of gravitational waves from inspiraling compact binaries*”. In: *Physical Review D* 85.12 (2012), p. 122006.
- [4] Benjamin P Abbott et al. “A guide to LIGO–Virgo detector noise and extraction of transient gravitational-wave signals”. In: *Classical and Quantum Gravity* 37.5 (2020), p. 055002.

References II

- [5] G Ashton et al. “BILBY: A user-friendly Bayesian inference library for gravitational-wave astronomy”. In: *The Astrophysical Journal Supplement Series* 241.2 (2019), p. 27.