

Transitioning to Memory Burden: Detectable Small PBHs as Dark Matter

Sebastian Zell

Ludwig Maximilian University & Max Planck Institute for Physics, Munich

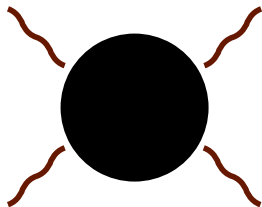
Work¹ with Gia Dvali, Marco Michel and Michael Zantedeschi

20th May 2025

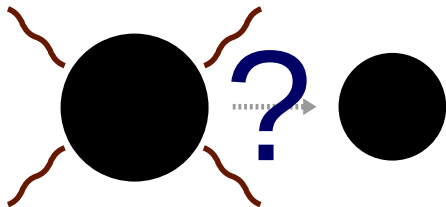
¹ *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

The Timescales of Quantum Breaking, Fortschr. Phys. **71** (2023) 2300163, arXiv:2306.09410. [News article “Where is the boundary to the quantum world?”].

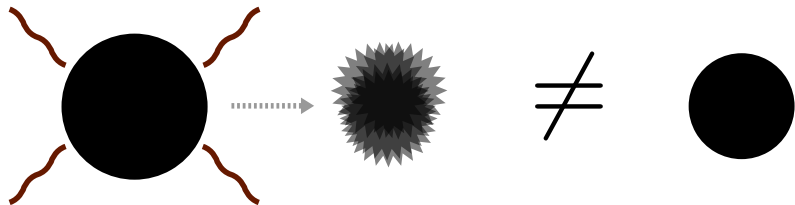
The fate of an old black hole



The fate of an old black hole

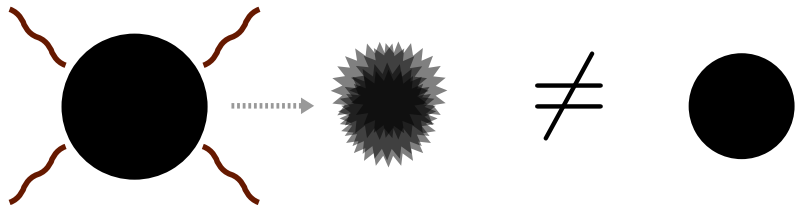


The fate of an old black hole



- ▶ Black hole evolution likely not self-similar

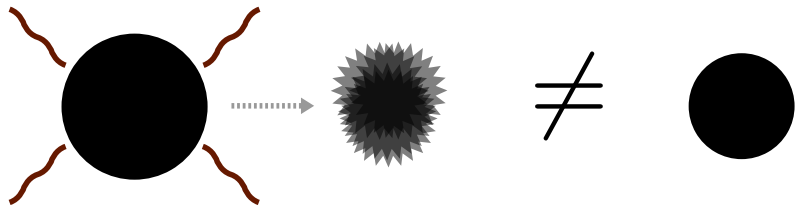
The fate of an old black hole



- ▶ Black hole evolution likely not self-similar
- ▶ Memory burden: indication for slowdown of evaporation²

²G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

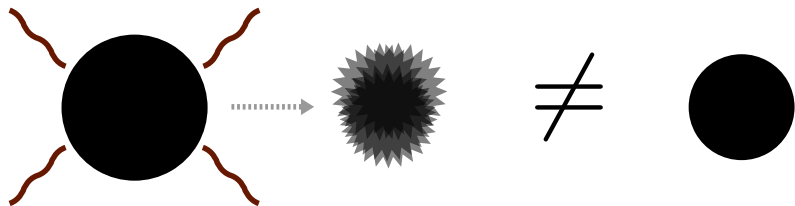
The fate of an old black hole



- ▶ Black hole evolution likely not self-similar
- ▶ Memory burden: indication for slowdown of evaporation²
- ▶ Small PBHs below 10^{15} g as dark matter

²G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

The fate of an old black hole



- ▶ Black hole evolution likely not self-similar
- ▶ Memory burden: indication for slowdown of evaporation²
- ▶ Small PBHs below 10^{15} g as dark matter
- ▶ Inflationary production compatible with CMB [talk B. Gladwyn]

²G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

PBH constraints

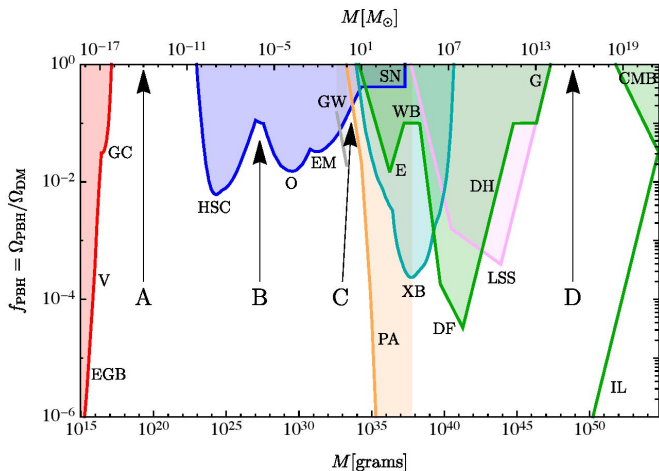


Figure from: B. Carr, F. Kühnel, *Primordial Black Holes as Dark Matter: Recent Developments*, arXiv:2006.02838.

PBH constraints

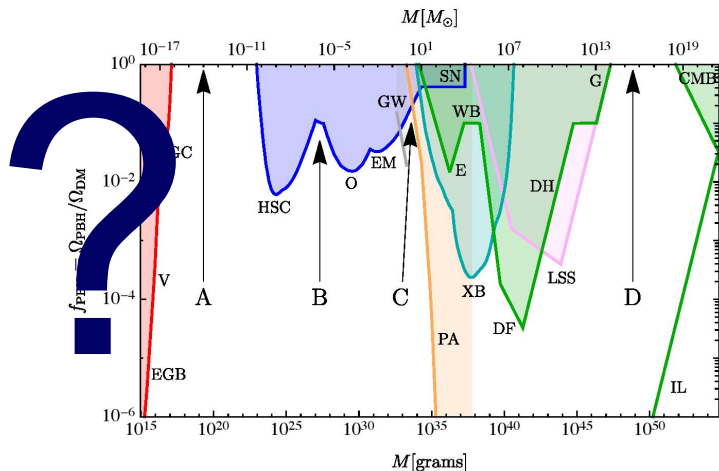


Figure from: B. Carr, F. Kühnel, *Primordial Black Holes as Dark Matter: Recent Developments*, arXiv:2006.02838.

A. Alexandre, G. Dvali, E. Koutsangelas, *New Mass Window for Primordial Black Holes as Dark Matter from Memory Burden Effect*, arXiv:2402.14069.

V. Thoss, A. Burkert, K. Kohri, *Breakdown of Hawking Evaporation opens new Mass Window for Primordial Black Holes as Dark Matter Candidate*, arXiv:2402.17823.

Detecting the transition to memory burden (MB)

Detecting the transition to memory burden (MB)

- 1 Results about MB
- 2 Transition to MB

Semi-classical evaporation

- ▶ Single scale

$$r_g \sim GM$$

Semi-classical evaporation

- ▶ Single scale

$$r_g \sim GM$$

- ▶ Dimensionless parameter³

$$S \sim \hbar^{-1} r_g M$$

³J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

Semi-classical evaporation

- ▶ Single scale

$$r_g \sim GM$$

- ▶ Dimensionless parameter³

$$S \sim \hbar^{-1} r_g M \sim 10^{22} \left(\frac{M}{10^6 \text{ g}} \right)^2$$

³J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

Semi-classical evaporation

- ▶ Single scale

$$r_g \sim GM$$

- ▶ Dimensionless parameter³

$$S \sim \hbar^{-1} r_g M \sim 10^{22} \left(\frac{M}{10^6 \text{ g}} \right)^2$$

- ▶ Hawking particle production:⁴

$$\Gamma_{\text{SC}} \sim \frac{1}{r_g}$$

³J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

⁴S. Hawking, *Particle Creation by Black Holes*, Commun. Math. Phys. **43** (1975).

Semi-classical evaporation

- ▶ Single scale

$$r_g \sim GM$$

- ▶ Dimensionless parameter³

$$S \sim \hbar^{-1} r_g M \sim 10^{22} \left(\frac{M}{10^6 \text{ g}} \right)^2$$

- ▶ Hawking particle production:⁴

$$\Gamma_{\text{SC}} \sim \frac{1}{r_g} \quad \left. \frac{dM}{dt} \right|_{\text{SC}} \sim \hbar r_g^{-1} \Gamma_{\text{SC}}$$

³J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

⁴S. Hawking, *Particle Creation by Black Holes*, Commun. Math. Phys. **43** (1975).

Semi-classical evaporation

- ▶ Single scale

$$r_g \sim GM$$

- ▶ Dimensionless parameter³

$$S \sim \hbar^{-1} r_g M \sim 10^{22} \left(\frac{M}{10^6 \text{ g}} \right)^2$$

- ▶ Hawking particle production:⁴

$$\Gamma_{\text{SC}} \sim \frac{1}{r_g} \quad \left. \frac{dM}{dt} \right|_{\text{SC}} \sim \hbar r_g^{-1} \Gamma_{\text{SC}}$$

- ▶ Semi-classical timescale of evaporation

$$\tau_{\text{SC}} \sim M \left(\left. \frac{dM}{dt} \right|_{\text{SC}} \right)^{-1}$$

³J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

⁴S. Hawking, *Particle Creation by Black Holes*, Commun. Math. Phys. **43** (1975).

Semi-classical evaporation

- ▶ Single scale

$$r_g \sim GM$$

- ▶ Dimensionless parameter³

$$S \sim \hbar^{-1} r_g M \sim 10^{22} \left(\frac{M}{10^6 \text{ g}} \right)^2$$

- ▶ Hawking particle production:⁴

$$\Gamma_{\text{SC}} \sim \frac{1}{r_g} \quad \left. \frac{dM}{dt} \right|_{\text{SC}} \sim \hbar r_g^{-1} \Gamma_{\text{SC}}$$

- ▶ Semi-classical timescale of evaporation

$$\tau_{\text{SC}} \sim M \left(\left. \frac{dM}{dt} \right|_{\text{SC}} \right)^{-1} \sim S r_g$$

³J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

⁴S. Hawking, *Particle Creation by Black Holes*, Commun. Math. Phys. **43** (1975).

MB (in one slide)⁵

- ▶ Microstate entropy:

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \sqrt{S} \sum_{k=1}^S \hat{n}_k$$

⁵G. Dvali, *A Microscopic Model of Holography: Survival by the Burden of Memory*, arXiv:1810.02336.

MB (in one slide)⁵

- ▶ Microstate entropy:

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

⁵G. Dvali, *A Microscopic Model of Holography: Survival by the Burden of Memory*, arXiv:1810.02336.

MB (in one slide)⁵

- ▶ Microstate entropy:

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^p$$

⁵G. Dvali, *A Microscopic Model of Holography: Survival by the Burden of Memory*, arXiv:1810.02336.

MB (in one slide)⁵

- ▶ Microstate entropy:

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^P \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^P \underset{=S}{\langle \hat{n}_0 \rangle} 0$$

⁵G. Dvali, *A Microscopic Model of Holography: Survival by the Burden of Memory*, arXiv:1810.02336.

MB (in one slide)⁵

- ▶ Microstate entropy:

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^p \underset{\langle \hat{n}_0 \rangle = S}{=} 0$$

- ▶ MB

$$\mu \sim \left| \frac{\partial \mathcal{E}_K}{\partial n_0} \right|$$

⁵G. Dvali, *A Microscopic Model of Holography: Survival by the Burden of Memory*, arXiv:1810.02336.

MB (in one slide)⁵

- ▶ Microstate entropy:

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^P \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^P \underset{=S}{\langle \hat{n}_0 \rangle} 0$$

- ▶ MB

$$\mu \sim \left| \frac{\partial \mathcal{E}_K}{\partial n_0} \right|$$

- ▶ Stored information ties system to initial state⁶

⁵G. Dvali, *A Microscopic Model of Holography: Survival by the Burden of Memory*, arXiv:1810.02336.

⁶MB also in inflation: G. Dvali, L. Eisemann, M. Michel, S. Z., *Universe's Primordial Quantum Memories*, arXiv:1812.08749.

Memory burden

MB: slowdown of evaporation⁷

additional material: MB

⁷G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

MB: slowdown of evaporation⁷

additional material: MB

- ▶ When does MB set in?

$$t_{\text{MB}} \sim qT_{\text{SC}}$$

⁷G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

MB: slowdown of evaporation⁷

additional material: MB

- ▶ When does MB set in?

$$t_{\text{MB}} \sim q\tau_{\text{SC}} \quad q \leq \frac{1}{2}$$

⁷G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

MB: slowdown of evaporation⁷

additional material: MB

- ▶ When does MB set in?

$$t_{\text{MB}} \sim q\tau_{\text{TSC}} \quad q \leq \frac{1}{2}$$

- ▶ Indications for early MB⁸

$$q \sim \frac{1}{\sqrt{S}}$$

⁷G. Dvali, L. Eiseemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

⁸M. Michel, S. Z., *The Timescales of Quantum Breaking*, arXiv:2306.09410.

Memory burden

- ▶ How drastic is the slowdown?

$$\Gamma \sim \frac{1}{S^k} \Gamma_{sc}$$

Memory burden

- ▶ How drastic is the slowdown?

$$\Gamma \sim \frac{1}{S^k} \Gamma_{\text{SC}}$$

- ▶ Quantum analogue system:⁹

$$k \sim 1, \dots, 3$$

⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

- ▶ How drastic is the slowdown?

$$\Gamma \sim \frac{1}{S^k} \Gamma_{\text{SC}}$$

- ▶ Quantum analogue system:⁹

$$k \sim 1, \dots, 3$$

- ▶ Full evolution

$$\frac{dM(t)}{dt} \sim r_g^{-1} \Gamma$$

⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

- ▶ How drastic is the slowdown?

$$\Gamma \sim \frac{1}{S^k} \Gamma_{\text{SC}}$$

- ▶ Quantum analogue system:⁹

$$k \sim 1, \dots, 3$$

- ▶ Full evolution

$$\frac{dM(t)}{dt} \sim r_g^{-1} \Gamma \sim \frac{dM(0)}{dt} \left\{ \begin{array}{l} 1 \\ M(t) \geq (1 - q)M_0 \end{array} \right.$$

⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

- ▶ How drastic is the slowdown?

$$\Gamma \sim \frac{1}{S^k} \Gamma_{\text{SC}}$$

- ▶ Quantum analogue system:⁹

$$k \sim 1, \dots, 3$$

- ▶ Full evolution

$$\frac{dM(t)}{dt} \sim r_g^{-1} \Gamma \sim \frac{dM(0)}{dt} \begin{cases} 1 & M(t) \geq (1 - q)M_0 \\ \frac{1}{S^k} & M(t) < M_{\text{MB}} \end{cases}$$

⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

- ▶ How drastic is the slowdown?

$$\Gamma \sim \frac{1}{S^k} \Gamma_{\text{SC}}$$

- ▶ Quantum analogue system:⁹

$$k \sim 1, \dots, 3$$

- ▶ Full evolution

$$\frac{dM(t)}{dt} \sim r_g^{-1} \Gamma \sim \frac{dM(0)}{dt} \begin{cases} 1 & M(t) \geq (1 - q)M_0 \\ ??? & \\ \frac{1}{S^k} & M(t) < M_{\text{MB}} \end{cases}$$

⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Memory burden

- ▶ How drastic is the slowdown?

$$\Gamma \sim \frac{1}{S^k} \Gamma_{\text{SC}}$$

- ▶ Quantum analogue system:⁹

$$k \sim 1, \dots, 3$$

- ▶ Full evolution

$$\frac{dM(t)}{dt} \sim r_g^{-1} \Gamma \sim \frac{dM(0)}{dt} \begin{cases} 1 & M(t) \geq (1-q)M_0 \\ ??? & \\ \frac{1}{S^k} & M(t) < M_{\text{MB}} \end{cases}$$

How does the black hole transition to MB?

⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Outline

Detecting the transition to memory burden

- 1 Results about MB
- 2 Transition to MB

The transition¹⁰

- ▶ Suppressed rate: additional material: transition to MB

$$\Gamma \sim \left(\frac{1}{S}\right)^{\mu r_g} \Gamma_{\text{sc}}$$

¹⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

The transition¹⁰

- Suppressed rate: additional material: transition to MB

$$\Gamma \sim \left(\frac{1}{S}\right)^{\mu r_g} \Gamma_{\text{sc}} \sim \exp\left(-\frac{(1-q)M_0 - M(t)}{\delta M_0}\right) \Gamma_{\text{sc}}$$

¹⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

The transition¹⁰

- ▶ Suppressed rate: additional material: transition to MB

$$\Gamma \sim \left(\frac{1}{S}\right)^{\mu r_g} \Gamma_{\text{sc}} \sim \exp\left(-\frac{(1-q)M_0 - M(t)}{\delta M_0}\right) \Gamma_{\text{sc}}$$

- ▶ “Steep” transition with width

$$\delta \sim O(10^{-1})q$$

¹⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

The transition¹⁰

- ▶ Suppressed rate: additional material: transition to MB

$$\Gamma \sim \left(\frac{1}{S}\right)^{\mu r_g} \Gamma_{\text{sc}} \sim \exp\left(-\frac{(1-q)M_0 - M(t)}{\delta M_0}\right) \Gamma_{\text{sc}}$$

- ▶ “Steep” transition with width

$$\delta \sim O(10^{-1})q$$

- ▶ Solve

$$\frac{dM(t)}{dt} \sim -r_g^{-1} \Gamma$$

¹⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

The transition¹⁰

- ▶ Suppressed rate: additional material: transition to MB

$$\Gamma \sim \left(\frac{1}{S}\right)^{\mu r_g} \Gamma_{\text{sc}} \sim \exp\left(-\frac{(1-q)M_0 - M(t)}{\delta M_0}\right) \Gamma_{\text{sc}}$$

- ▶ “Steep” transition with width

$$\delta \sim O(10^{-1})q$$

- ▶ Solve

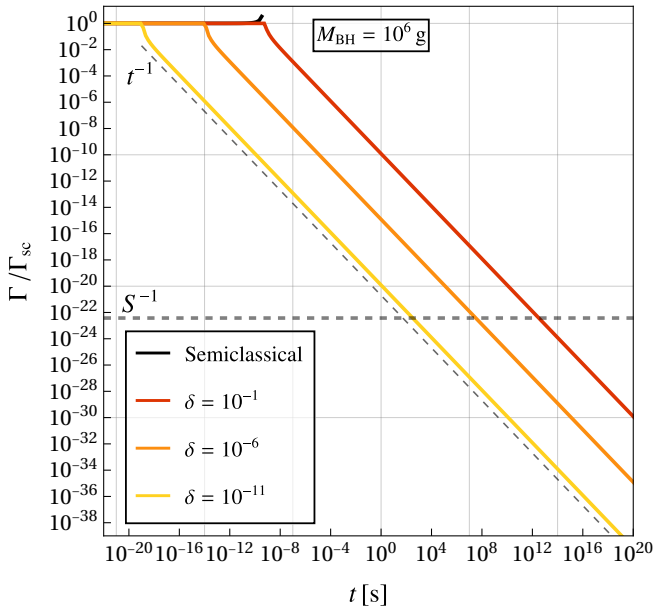
$$\frac{dM(t)}{dt} \sim -r_g^{-1} \Gamma$$

- ▶ Result: “slow” change of rate

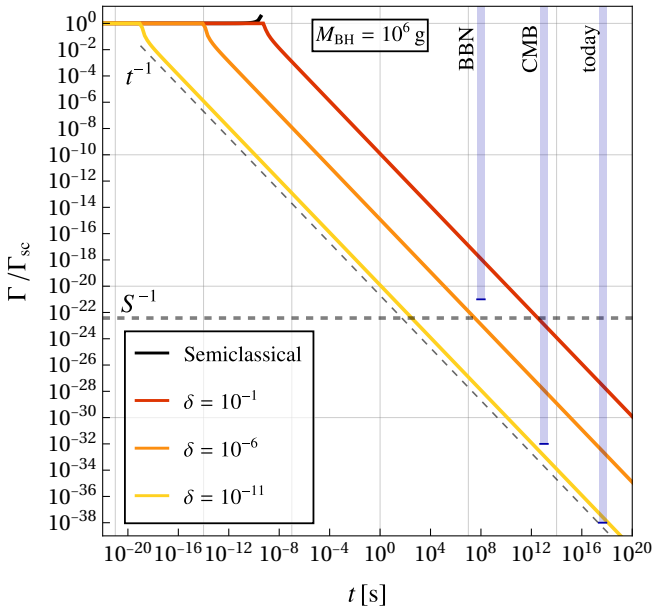
$$\frac{dM(t)}{dt} \sim -r_g^{-1} \Gamma_{\text{sc}} \frac{\delta \tau_{\text{SC}}}{t}$$

¹⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

Observational constraints



Observational constraints



Observational constraints

- ▶ Strongest constraint from today's Universe

Observational constraints

- ▶ Strongest constraint from today's Universe¹¹

$$\delta \lesssim 10^{-11}$$

¹¹See also G. Montefalcone et al., *Does Memory Burden Open a New Mass Window for Primordial Black Holes as Dark Matter?*, arXiv:2503.21005.

Observational constraints

- ▶ Strongest constraint from today's Universe¹¹

$$\delta \lesssim 10^{-11}$$

- ▶ Indication for early onset of MB

$$q \sim \frac{1}{\sqrt{S}} \quad \delta \sim O(10^{-1})q$$

¹¹See also G. Montefalcone et al., *Does Memory Burden Open a New Mass Window for Primordial Black Holes as Dark Matter?*, arXiv:2503.21005.

Observational constraints

- ▶ Strongest constraint from today's Universe¹¹

$$\delta \lesssim 10^{-11}$$

- ▶ Indication for early onset of MB

$$q \sim \frac{1}{\sqrt{S}} \quad \delta \sim O(10^{-1})q \quad M \gtrsim 10^6 \text{ g}$$

¹¹See also G. Montefalcone et al., *Does Memory Burden Open a New Mass Window for Primordial Black Holes as Dark Matter?*, arXiv:2503.21005.

Observational constraints

- ▶ Strongest constraint from today's Universe¹¹

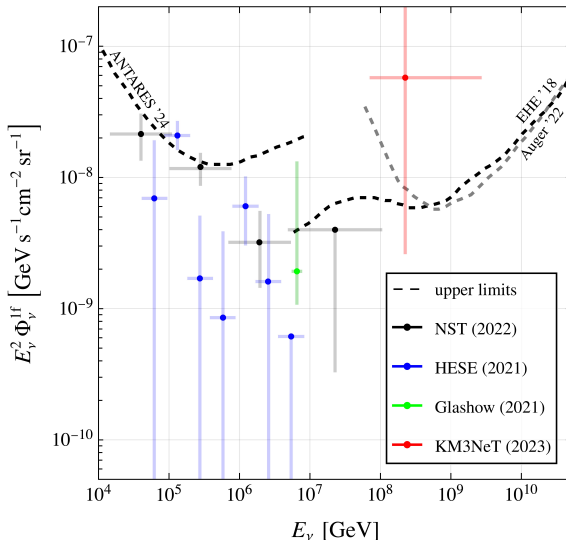
$$\delta \lesssim 10^{-11}$$

- ▶ Indication for early onset of MB

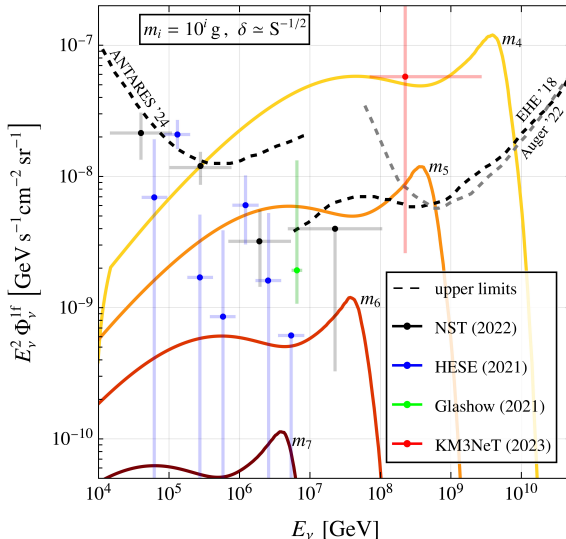
$$q \sim \frac{1}{\sqrt{S}} \quad \delta \sim O(10^{-1})q \quad M \gtrsim 10^6 \text{ g}$$

Small PBHs detectable today

¹¹See also G. Montefalcone et al., *Does Memory Burden Open a New Mass Window for Primordial Black Holes as Dark Matter?*, arXiv:2503.21005.

High-energetic neutrino signals¹²

¹²See also A. Boccia, F. Iocco, *A strike of luck: could the KM3-230213A event be caused by an evaporating primordial black hole?*, arXiv:2502.19245.

High-energetic neutrino signals¹²

¹²See also A. Boccia, F. Iocco, *A strike of luck: could the KM3-230213A event be caused by an evaporating primordial black hole?*, arXiv:2502.19245.

Conclusion

- ▶ MB opens new window for small PBHs as dark matter

Conclusion

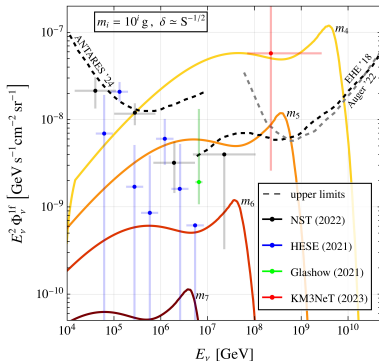
- ▶ MB opens new window for small PBHs as dark matter
- ▶ Transition crucial for constraints

Conclusion

- ▶ MB opens new window for small PBHs as dark matter
- ▶ Transition crucial for constraints
- ▶ Indication for early onset of MB

Conclusion

- ▶ MB opens new window for small PBHs as dark matter
- ▶ Transition crucial for constraints
- ▶ Indication for early onset of MB
- ▶ Small PBHs currently transitioning to MB detectable today



Analogue gravity

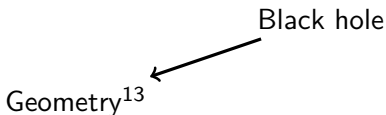
- ▶ Ideally: study evaporation without semi-classical limit

Analogue gravity

- ▶ Ideally: study evaporation without semi-classical limit
- ▶ Easier: analogue systems
 - ▷ Share important properties with gravity
 - ▷ Accessible for computations and experiments

Analogue gravity

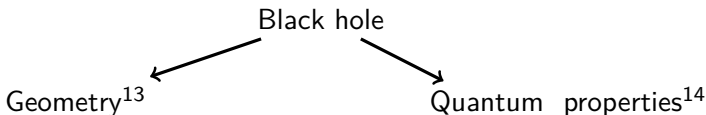
- ▶ Ideally: study evaporation without semi-classical limit
- ▶ Easier: analogue systems
 - ▷ Share important properties with gravity
 - ▷ Accessible for computations and experiments



¹³W. Unruh, *Experimental Black-Hole Evaporation?*, Phys. Rev. Lett. **46** (1981).
O. Lahav *et al.*, *Realization of a sonic black hole analogue in a Bose-Einstein condensate*, arXiv:0906.1337.

Analogue gravity

- ▶ Ideally: study evaporation without semi-classical limit
- ▶ Easier: analogue systems
 - ▷ Share important properties with gravity
 - ▷ Accessible for computations and experiments

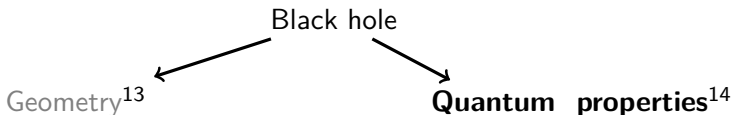


¹³W. Unruh, *Experimental Black-Hole Evaporation?*, Phys. Rev. Lett. **46** (1981).
O. Lahav *et al.*, *Realization of a sonic black hole analogue in a Bose-Einstein condensate*, arXiv:0906.1337.

¹⁴G. Dvali, C. Gomez, *Black Holes as Critical Point of Quantum Phase Transition*, arXiv:1207.4059.

Analogue gravity

- ▶ Ideally: study evaporation without semi-classical limit
- ▶ Easier: analogue systems
 - ▷ Share important properties with gravity
 - ▷ Accessible for computations and experiments



¹³W. Unruh, *Experimental Black-Hole Evaporation?*, Phys. Rev. Lett. **46** (1981).
O. Lahav *et al.*, *Realization of a sonic black hole analogue in a Bose-Einstein condensate*, arXiv:0906.1337.

¹⁴G. Dvali, C. Gomez, *Black Holes as Critical Point of Quantum Phase Transition*, arXiv:1207.4059.

Imitate entropy

- ▶ Entropy S : need e^S microstates

Imitate entropy

- ▶ Entropy S : need e^S microstates
- ▶ Natural explanation: S modes $\hat{a}_1^\dagger, \dots, \hat{a}_S^\dagger$
(with $[\hat{a}_i, \hat{a}_j^\dagger] = \delta_{ij}$)

Imitate entropy

- ▶ Entropy S : need e^S microstates
- ▶ Natural explanation: S modes $\hat{a}_1^\dagger, \dots, \hat{a}_S^\dagger$
(with $[\hat{a}_i, \hat{a}_j^\dagger] = \delta_{ij}$)
- ▶ Microstates

$$\left(\hat{a}_1^\dagger\right)^{\{0,1\}} \dots \left(\hat{a}_S^\dagger\right)^{\{0,1\}} |0\rangle$$

Imitate entropy

- ▶ Entropy S : need e^S microstates
- ▶ Natural explanation: S modes $\hat{a}_1^\dagger, \dots, \hat{a}_S^\dagger$
(with $[\hat{a}_i, \hat{a}_j^\dagger] = \delta_{ij}$)
- ▶ Microstates
$$\left(\hat{a}_1^\dagger\right)^{\{0,1\}} \dots \left(\hat{a}_S^\dagger\right)^{\{0,1\}} |0\rangle$$
- ▶ Crucial: all microstates must have similar energy

Enhanced memory storage¹⁵

► Hamiltonian

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \sqrt{S} \sum_{k=1}^S \underbrace{\hat{n}_k}_{\hat{a}_k^\dagger \hat{a}_k}$$

¹⁵ G. Dvali, *Critically excited states with enhanced memory and pattern recognition capacities in quantum brain networks: Lesson from black holes*, arXiv:1711.09079.
G. Dvali, M. Michel, S. Z., *Finding Critical States of Enhanced Memory Capacity in Attractive Cold Bosons*, arXiv:1805.10292.

Enhanced memory storage¹⁵

► Hamiltonian

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

¹⁵ G. Dvali, *Critically excited states with enhanced memory and pattern recognition capacities in quantum brain networks: Lesson from black holes*, arXiv:1711.09079.
G. Dvali, M. Michel, S. Z., *Finding Critical States of Enhanced Memory Capacity in Attractive Cold Bosons*, arXiv:1805.10292.

Enhanced memory storage¹⁵

- ▶ Hamiltonian

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^p$$

¹⁵G. Dvali, *Critically excited states with enhanced memory and pattern recognition capacities in quantum brain networks: Lesson from black holes*, arXiv:1711.09079.
G. Dvali, M. Michel, S. Z., *Finding Critical States of Enhanced Memory Capacity in Attractive Cold Bosons*, arXiv:1805.10292.

Enhanced memory storage¹⁵

► Hamiltonian

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^P \sum_{k=1}^S \hat{n}_k$$

► Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^P \underset{=S}{\langle \hat{n}_0 \rangle} 0$$

¹⁵ G. Dvali, *Critically excited states with enhanced memory and pattern recognition capacities in quantum brain networks: Lesson from black holes*, arXiv:1711.09079.
G. Dvali, M. Michel, S. Z., *Finding Critical States of Enhanced Memory Capacity in Attractive Cold Bosons*, arXiv:1805.10292.

Enhanced memory storage¹⁵

- ▶ Hamiltonian

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^P \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^P \underset{=S}{\langle \hat{n}_0 \rangle} 0$$

- ▶ Macrostate $\langle \hat{n}_0 \rangle = S$ has entropy S

¹⁵ G. Dvali, *Critically excited states with enhanced memory and pattern recognition capacities in quantum brain networks: Lesson from black holes*, arXiv:1711.09079.
G. Dvali, M. Michel, S. Z., *Finding Critical States of Enhanced Memory Capacity in Attractive Cold Bosons*, arXiv:1805.10292.

Enhanced memory storage¹⁵

- ▶ Hamiltonian

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^P \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} \left(\hat{a}_0^\dagger \hat{b} + \text{h.c.}\right)$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^P \underset{=S}{\langle \hat{n}_0 \rangle} 0$$

- ▶ Macrostate $\langle \hat{n}_0 \rangle = S$ has entropy S

¹⁵ G. Dvali, *Critically excited states with enhanced memory and pattern recognition capacities in quantum brain networks: Lesson from black holes*, arXiv:1711.09079.
G. Dvali, M. Michel, S. Z., *Finding Critical States of Enhanced Memory Capacity in Attractive Cold Bosons*, arXiv:1805.10292.

Enhanced memory storage¹⁵

- ▶ Hamiltonian

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^P \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} \left(\hat{a}_0^\dagger \hat{b} + \text{h.c.}\right)$$

- ▶ Effective energy gaps

$$\mathcal{E}_K \sim \sqrt{S} \left(1 - \frac{\langle \hat{n}_0 \rangle}{S}\right)^P \underset{=S}{\langle \hat{n}_0 \rangle} 0$$

- ▶ Macrostate $\langle \hat{n}_0 \rangle = S$ has entropy S

- ▶ Dictionary

\hat{n}_0 : carries mass

$\langle \hat{n}_0 \rangle = S$: black hole state

\hat{n}_k : carry entropy

\hat{n}_b : Hawking quanta

¹⁵ G. Dvali, *Critically excited states with enhanced memory and pattern recognition capacities in quantum brain networks: Lesson from black holes*, arXiv:1711.09079.
G. Dvali, M. Michel, S. Z., *Finding Critical States of Enhanced Memory Capacity in Attractive Cold Bosons*, arXiv:1805.10292.

Time evolution

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.})$$

Time evolution

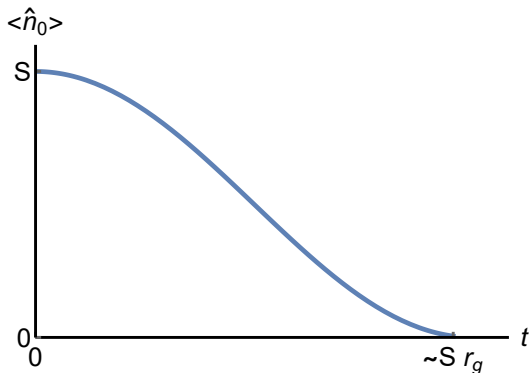
$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.})$$

$$\sum_{k=1}^S \langle \hat{n}_k \rangle = 0$$

Time evolution

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.})$$

$$\sum_{k=1}^S \langle \hat{n}_k \rangle = 0$$



Time evolution

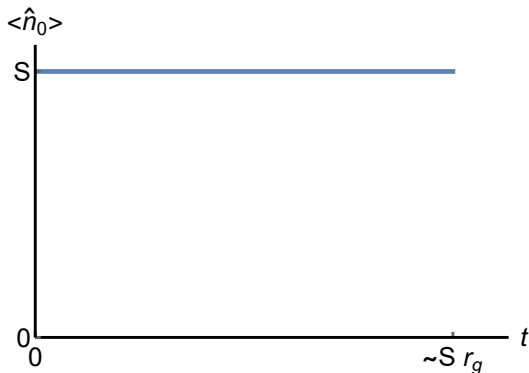
$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.})$$

$$\sum_{k=1}^S \langle \hat{n}_k \rangle \sim S$$

Time evolution

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.})$$

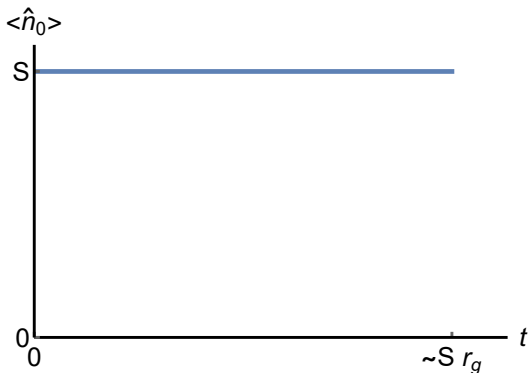
$$\sum_{k=1}^S \langle \hat{n}_k \rangle \sim S$$



Time evolution

$$\frac{\hat{\mathcal{H}}_S}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.})$$

$$\sum_{k=1}^S \langle \hat{n}_k \rangle \sim S$$



Memory burden:¹⁶ entropy prevents evaporation

¹⁶G. Dvali, *A Microscopic Model of Holography: Survival by the Burden of Memory*, arXiv:1810.02336.

Full model¹⁵

$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_{S_{>}} + \hat{n}_b + \frac{1}{S} \left(\hat{a}_0^\dagger \hat{b} + \text{h.c.} \right)$$

¹⁵G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

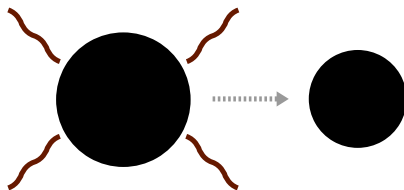
Full model¹⁵

$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_{S_{>}} + \hat{n}_b + \frac{1}{S} \left(\hat{a}_0^\dagger \hat{b} + \text{h.c.} \right) + \hat{\mathcal{H}}_{S_{<}} + \text{interactions}$$

¹⁵G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Full model¹⁵

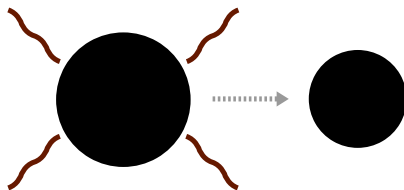
$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_{S_{>}} + \hat{n}_b + \frac{1}{S} \left(\hat{a}_0^\dagger \hat{b} + \text{h.c.} \right) + \hat{\mathcal{H}}_{S_{<}} + \text{interactions}$$



¹⁵G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Full model¹⁵

$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_{S_>} + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.}) + \hat{\mathcal{H}}_{S_<} + \text{interactions}$$

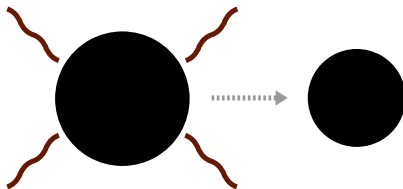


$$\langle \hat{n}_0 \rangle = S_> \longrightarrow \langle \hat{n}_0 \rangle = S_<$$

¹⁵ G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Full model¹⁵

$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_{S_{>}} + \hat{n}_b + \frac{1}{S} (\hat{a}_0^\dagger \hat{b} + \text{h.c.}) + \hat{\mathcal{H}}_{S_{<}} + \text{interactions}$$



$$\langle \hat{n}_0 \rangle = S_{>} \longrightarrow \langle \hat{n}_0 \rangle = S_{<}$$

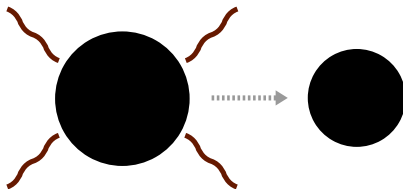
- Exact time evolution:¹⁶ transition suppressed dynamically

¹⁵ G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

¹⁶ M. Michel, S. Z., *TimeEvolver: A Program for Time Evolution With Improved Error Bound*, arXiv:2205.15346.

Full model¹⁵

$$\hat{\mathcal{H}} = \hat{\mathcal{H}}_{S_{>}} + \hat{n}_b + \frac{1}{S} \left(\hat{a}_0^\dagger \hat{b} + \text{h.c.} \right) + \hat{\mathcal{H}}_{S_{<}} + \text{interactions}$$



$$\langle \hat{n}_0 \rangle = S_{>} \quad \longrightarrow \quad \langle \hat{n}_0 \rangle = S_{<}$$

- ▶ Exact time evolution:¹⁶ transition suppressed dynamically
- ▶ Slowdown at the latest after half evaporation [back](#)

¹⁵ G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

¹⁶ M. Michel, S. Z., *TimeEvolver: A Program for Time Evolution With Improved Error Bound*, arXiv:2205.15346.

Black hole criticality

- ▶ Gravitational coupling

$$\alpha = \hbar G r_g^{-2}$$

Black hole criticality

- ▶ Gravitational coupling

$$\alpha = \hbar G r_g^{-2} = \frac{1}{S}$$

Black hole criticality

- ▶ Gravitational coupling

$$\alpha = \hbar G r_g^{-2} = \frac{1}{S}$$

- ▶ Black hole constituents

$$N = \frac{M}{\hbar r_g^{-1}}$$

Black hole criticality

- ▶ Gravitational coupling

$$\alpha = \hbar G r_g^{-2} = \frac{1}{S}$$

- ▶ Black hole constituents

$$N = \frac{M}{\hbar r_g^{-1}} = S$$

Black hole criticality

- ▶ Gravitational coupling

$$\alpha = \hbar G r_g^{-2} = \frac{1}{S}$$

- ▶ Black hole constituents

$$N = \frac{M}{\hbar r_g^{-1}} = S$$

- ▶ Critical collective coupling¹⁷

$$\alpha N = 1$$

¹⁷G. Dvali, C. Gomez, *Black Holes as Critical Point of Quantum Phase Transition*, [arXiv:1207.4059](https://arxiv.org/abs/1207.4059).

Imitate criticality

► Prototype model

$$\hat{\mathcal{H}} = \sum_{k=1}^Q \left(\hat{n}_k - \frac{\alpha}{4} \left(2\hat{n}_0\hat{n}_k + \hat{a}_0^{\dagger 2}\hat{a}_k^2 + \hat{a}_k^{\dagger 2}\hat{a}_0^2 \right) \right) + \frac{C_m}{2} \sum_{k=1}^Q \sum_{l=k+1}^Q f(k, l) \left(\hat{a}_k^{\dagger 2}\hat{a}_l^2 + \text{h.c.} \right) .$$

Imitate criticality

- ▶ Prototype model

$$\hat{\mathcal{H}} = \sum_{k=1}^Q \left(\hat{n}_k - \frac{\alpha}{4} \left(2\hat{n}_0\hat{n}_k + \hat{a}_0^{\dagger 2}\hat{a}_k^2 + \hat{a}_k^{\dagger 2}\hat{a}_0^2 \right) \right) + \frac{C_m}{2} \sum_{k=1}^Q \sum_{l=k+1}^Q f(k, l) \left(\hat{a}_k^{\dagger 2}\hat{a}_l^2 + \text{h.c.} \right) .$$

- ▶ Critical point

$$\alpha \langle \hat{n}_0 \rangle = 1$$

Imitate criticality

- ▶ Prototype model

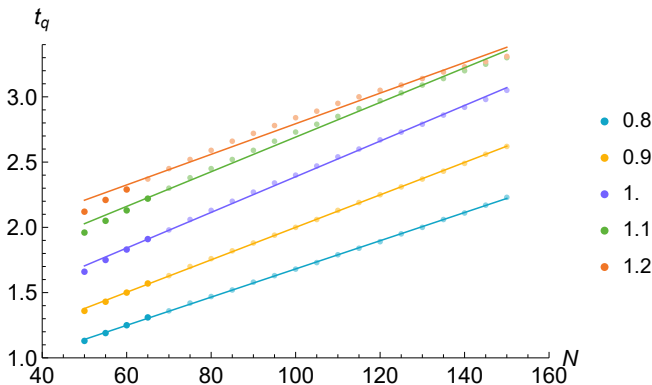
$$\hat{\mathcal{H}} = \sum_{k=1}^Q \left(\hat{n}_k - \frac{\alpha}{4} \left(2\hat{n}_0\hat{n}_k + \hat{a}_0^{\dagger 2}\hat{a}_k^2 + \hat{a}_k^{\dagger 2}\hat{a}_0^2 \right) \right) + \frac{C_m}{2} \sum_{k=1}^Q \sum_{l=k+1}^Q f(k, l) \left(\hat{a}_k^{\dagger 2}\hat{a}_l^2 + \text{h.c.} \right) .$$

- ▶ Critical point

$$\alpha \langle \hat{n}_0 \rangle = 1$$

- ▶ Study quantum break-time:¹⁸
timescale of breakdown of semi-classical approximation

¹⁸G. Dvali, C. Gomez, D. Flassig, A. Pritzel, *Scrambling in the Black Hole Portrait*, arXiv:1307.3458.

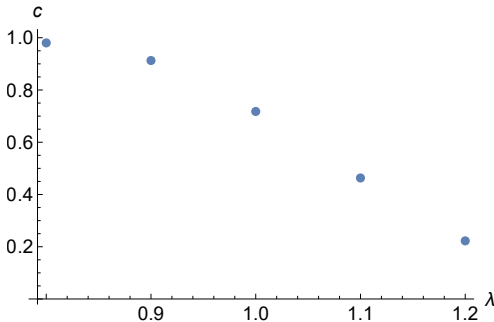
Result¹⁹

¹⁹ M. Michel, S. Z., *TimeEvolver: A Program for Time Evolution With Improved Error Bound*, arXiv:2205.15346.

M. Michel, S. Z., *The Timescales of Quantum Breaking*, arXiv:2306.09410.

Result¹⁹

$$t_q = a \cdot N^c + b ,$$

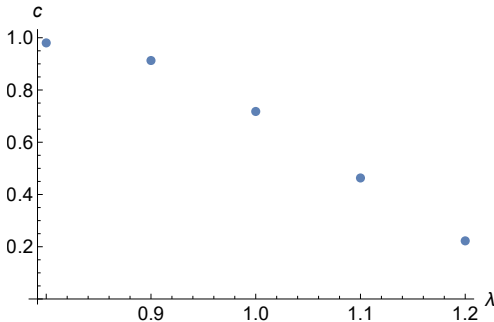


¹⁹ M. Michel, S. Z., *TimeEvolver: A Program for Time Evolution With Improved Error Bound*, arXiv:2205.15346.

M. Michel, S. Z., *The Timescales of Quantum Breaking*, arXiv:2306.09410.

Result¹⁹

$$t_q = a \cdot N^c + b,$$



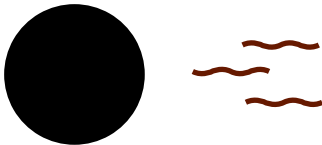
$$\lambda = \alpha N = 1 \quad \Rightarrow \quad c \approx 0.5$$

[back](#)

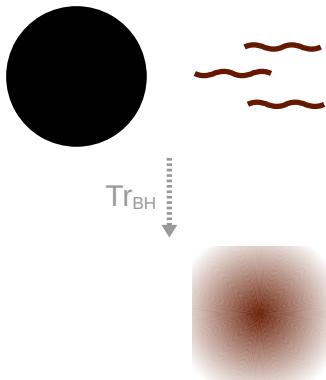
¹⁹ M. Michel, S. Z., *TimeEvolver: A Program for Time Evolution With Improved Error Bound*, arXiv:2205.15346.

M. Michel, S. Z., *The Timescales of Quantum Breaking*, arXiv:2306.09410.

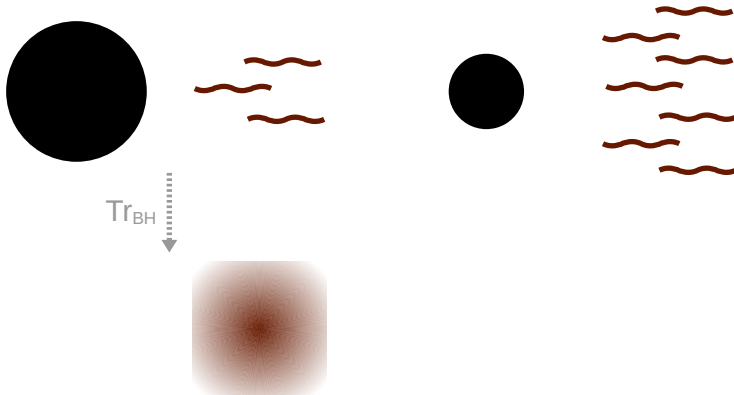
Page time²⁰



²⁰D. Page, *Information in black hole radiation*, arXiv:hep-th/9306083.

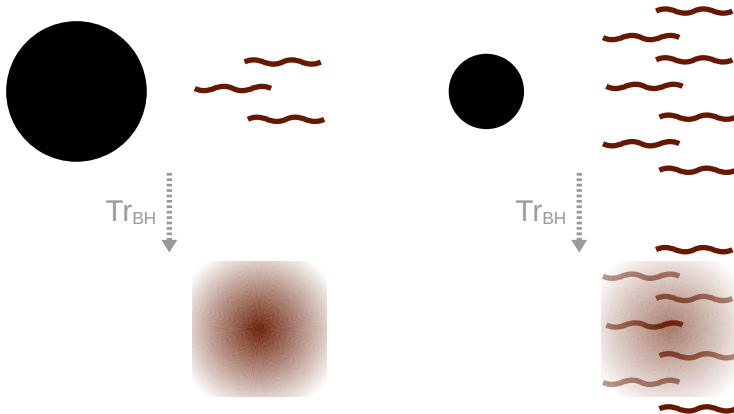
Page time²⁰

²⁰D. Page, *Information in black hole radiation*, arXiv:hep-th/9306083.

Page time²⁰

²⁰D. Page, *Information in black hole radiation*, arXiv:hep-th/9306083.

Page time²⁰



back

²⁰D. Page, *Information in black hole radiation*, arXiv:hep-th/9306083.

Classical black hole: no hair

- ▶ Geometry fully determined by mass

$$r_g \sim GM$$

Classical black hole: no hair

- ▶ Geometry fully determined by mass

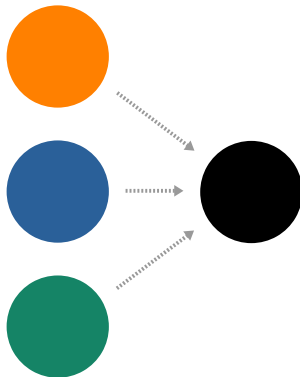
$$r_g \sim GM$$



Classical black hole: no hair

- ▶ Geometry fully determined by mass

$$r_g \sim GM$$

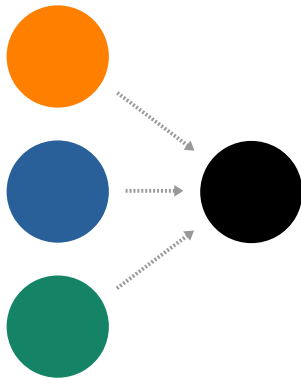


Classical black hole: no hair

- ▶ Geometry fully determined by mass

$$r_g \sim GM$$

- ▶ No hair outside²¹



²¹See e.g., P. Chrusciel, J. Costa, M. Heusler, *Stationary Black Holes: Uniqueness and Beyond*, arXiv:1205.6112.

Quantum black hole: entropy²²

► Entropy

$$S \sim \frac{r_g^2}{\hbar G}$$

²² J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

Quantum black hole: entropy²²

- ▶ Entropy

$$S \sim \frac{r_g^2}{\hbar G}$$

- ▶ Black holes quantum-mechanically distinct

²²J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

Quantum black hole: entropy²²

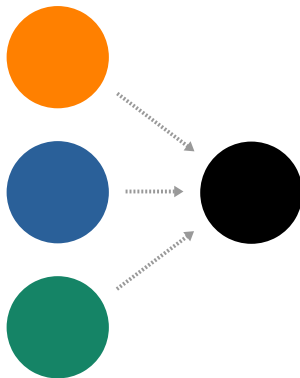
- ▶ Entropy

$$S \sim \frac{r_g^2}{\hbar G}$$

- ▶ Black holes quantum-mechanically distinct
- ▶ $\exp(S)$ different versions of a black hole of mass M

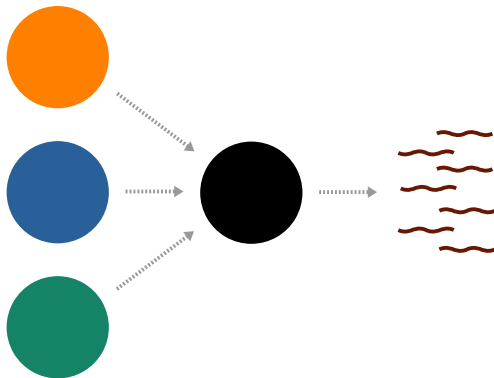
²² J. Bekenstein, *Black holes and entropy*, Phys. Rev. D **7** (1973).

Add evaporation¹⁷



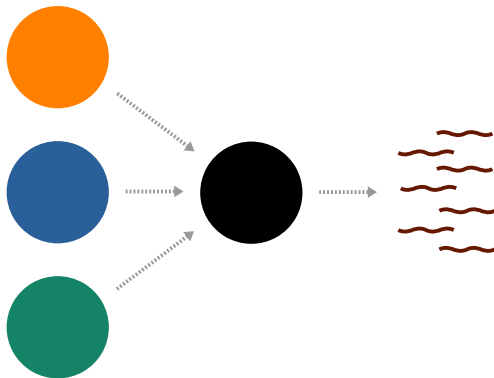
¹⁷S. Hawking, *Particle Creation by Black Holes*, *Commun. Math. Phys.* **43** (1975).

Add evaporation¹⁷



¹⁷S. Hawking, *Particle Creation by Black Holes*, *Commun. Math. Phys.* **43** (1975).

Add evaporation¹⁷

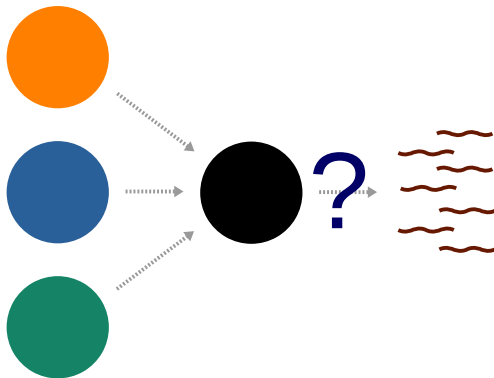


- ▶ Not unitary: information about initial state lost¹⁸

¹⁷ S. Hawking, *Particle Creation by Black Holes*, *Commun. Math. Phys.* **43** (1975).

¹⁸ S. Hawking, *Breakdown of predictability in gravitational collapse*, *Phys. Rev. D* **14** (1976).₂₅

Add evaporation¹⁷



- ▶ Not unitary: information about initial state lost¹⁸
- ▶ Question: how long is Hawking evaporation valid?

¹⁷S. Hawking, *Particle Creation by Black Holes*, *Commun. Math. Phys.* **43** (1975).

¹⁸S. Hawking, *Breakdown of predictability in gravitational collapse*, *Phys. Rev. D* **14** (1976).₂₅

Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass

Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission

$$\frac{\hbar r_g^{-1}}{M}$$

Breakdown of semi-classicality

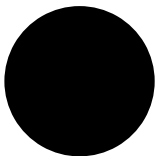
- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission

$$\frac{\hbar r_g^{-1}}{M} = \frac{1}{S}$$

Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission

$$\frac{\hbar r_g^{-1}}{M} = \frac{1}{S}$$

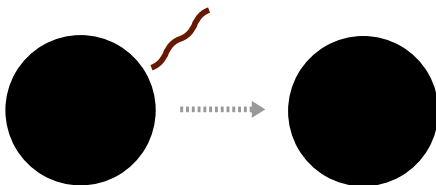


$|BH; M\rangle$

Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission

$$\frac{\hbar r_g^{-1}}{M} = \frac{1}{S}$$

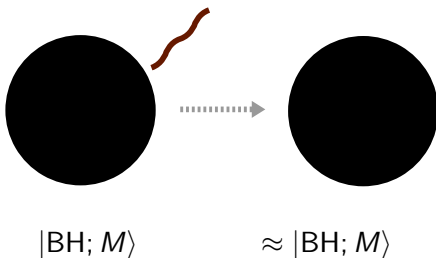


$|BH; M\rangle$

Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission

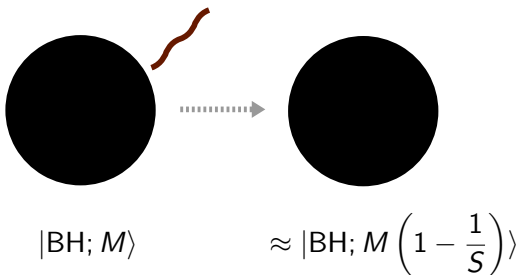
$$\frac{\hbar r_g^{-1}}{M} = \frac{1}{S}$$



Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission

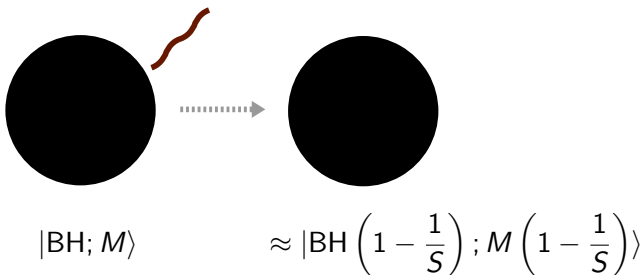
$$\frac{\hbar r_g^{-1}}{M} = \frac{1}{S}$$



Breakdown of semi-classicality

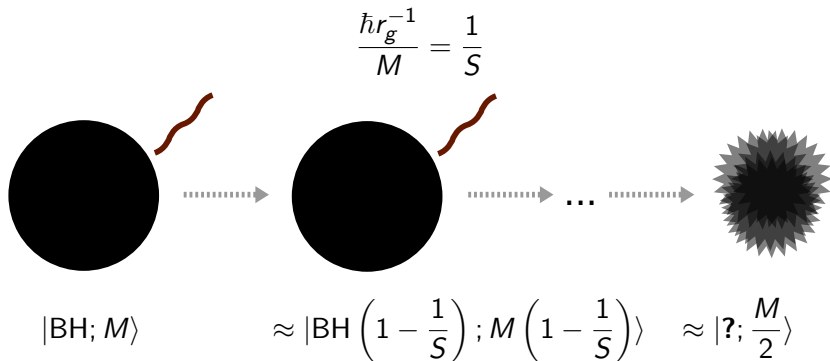
- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission

$$\frac{\hbar r_g^{-1}}{M} = \frac{1}{S}$$



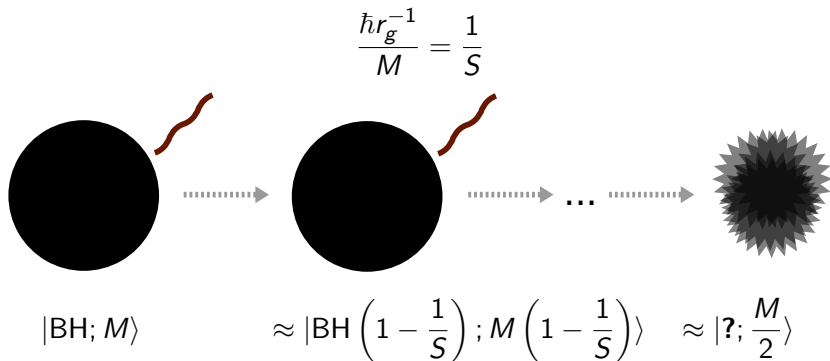
Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission



Breakdown of semi-classicality

- ▶ Semi-classical approximation: fixed mass
- ▶ Small correction after single emission



- ▶ Full breakdown of semi-classical description [back](#)

Microscopic MB¹⁹

► Microscopic model

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

¹⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Microscopic MB¹⁹

- ▶ Microscopic model

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gap

$$\mathcal{E}_K = \left(1 - \frac{n_0}{S}\right)^p \sqrt{S} r_g^{-1}$$

¹⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Microscopic MB¹⁹

- ▶ Microscopic model

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gap

$$\mathcal{E}_K = \left(1 - \frac{n_0}{S}\right)^p \sqrt{S} r_g^{-1}$$

- ▶ Memory burden

$$\mu \sim \left| \frac{\partial \mathcal{E}_K}{\partial n_0} \right| = p \left(1 - \frac{n_0}{S}\right)^{p-1} \sqrt{S} r_g^{-1} \sum_k n_k$$

¹⁹G. Dvali, L. Eise mann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Microscopic MB¹⁹

- ▶ Microscopic model

$$\frac{\hat{\mathcal{H}}}{r_g^{-1}} = \hat{n}_0 + \sqrt{S} \left(1 - \frac{\hat{n}_0}{S}\right)^p \sum_{k=1}^S \hat{n}_k$$

- ▶ Effective energy gap

$$\mathcal{E}_K = \left(1 - \frac{n_0}{S}\right)^p \sqrt{S} r_g^{-1}$$

- ▶ Memory burden

$$\mu \sim \left| \frac{\partial \mathcal{E}_K}{\partial n_0} \right| = p \left(1 - \frac{n_0}{S}\right)^{p-1} \sqrt{S} r_g^{-1} \sum_k^S n_k$$

- ▶ Critical value

$$q \equiv \frac{(S - n_0)_{\text{crit}}}{S} = \left(p\sqrt{S}\right)^{-1/(p-1)}$$

¹⁹G. Dvali, L. Eisemann, M. Michel, S. Z., *Black Hole Metamorphosis and Stabilization by Memory Burden*, arXiv:2006.00011.

Transition to MB²⁰

- ▶ Increased energy gap

$$\Delta N = \mu r_g = \frac{p\sqrt{S}}{2} \left(\frac{M_0 - M(t)}{M_0} \right)^{p-1}$$

²⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

Transition to MB²⁰

- ▶ Increased energy gap

$$\Delta N = \mu r_g = \frac{p\sqrt{S}}{2} \left(\frac{M_0 - M(t)}{M_0} \right)^{p-1}$$

- ▶ Suppressed emission

$$\Gamma = \left(\frac{1}{S} \right)^{\Delta N} \Gamma_{\text{sc}}$$

²⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

Transition to MB²⁰

- ▶ Increased energy gap

$$\Delta N = \mu r_g = \frac{p\sqrt{S}}{2} \left(\frac{M_0 - M(t)}{M_0} \right)^{p-1}$$

- ▶ Suppressed emission

$$\Gamma = \left(\frac{1}{S} \right)^{\Delta N} \Gamma_{\text{sc}}$$

- ▶ Approximate around onset of MB

$$\Gamma \simeq \exp \left(- \frac{(1-q)M_0 - M(t)}{\delta M_0} \right) \Gamma_{\text{sc}}$$

with

$$\delta = \frac{q}{(p-1) \ln S}$$

²⁰G. Dvali, M. Zantedeschi, S. Z., *Transitioning to Memory Burden: Detectable Small Primordial Black Holes as Dark Matter*, arXiv:2503.21740.

No dependence on q 