

Search for dark photons with continuous wave methods

Andrew Miller

Pia Astone, Giacomo Bruno, Sebastien Clesse, Sabrina D'Antonio, Antoine Depasse, Federico De Lillo, Sergio Frasca, Iuri La Rosa, Paola Leaci, Cristiano Palomba, Ornella J. Piccinni, Lorenzo Pierini, and Luca Rei

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Outline

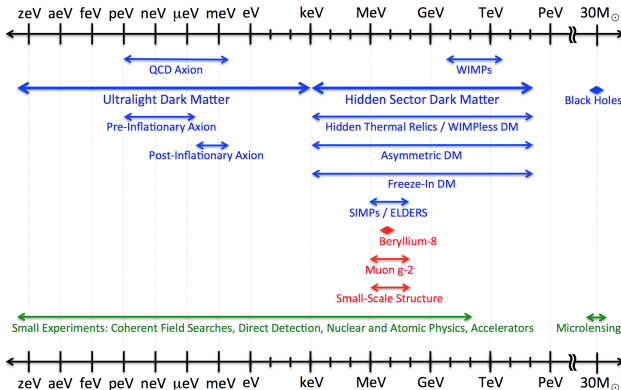
① Background

② Methods and limits

③ Conclusions

Dark matter

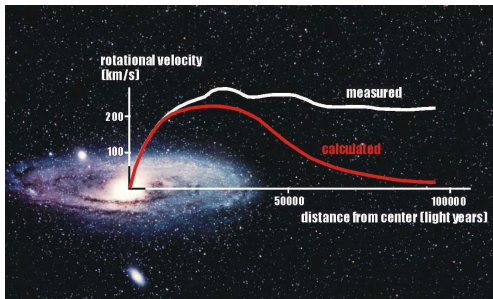
Dark Sector Candidates, Anomalies, and Search Techniques



- ▶ Any mass is possible, but GW detectors can probe ultra-heavy and ultralight regimes [3]

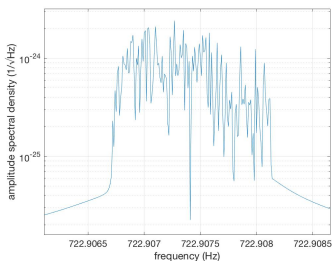
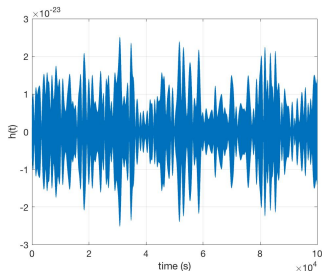
Signatures of ultralight dark matter

- ▶ Superradiantly formed axion clouds around black holes [1]
- ▶ Primordial black holes could be linked to dark matter [8, 13]
- ▶ Dark photons could couple to baryons in interferometers [11]
- ▶ Scalar dark matter could cause time-dependent changes in fundamental constants [6, 14, 2]



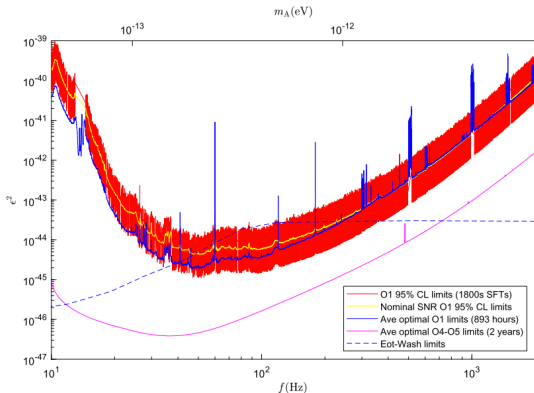
Dark photons

- ▶ Narrow-band, stochastic background: a superposition of plane waves with slightly different frequencies [11]
- ▶ Velocity of each dark photon follows a Maxwell-Boltzmann distribution $\Delta f \sim \frac{1}{2} \left(\frac{v_0}{c}\right)^2 f_0 \sim O(10^{-6}) f_0$ Hz
 - ▶ Continuous wave techniques, originally designed for isolated neutron stars, can be applied
 - ▶ Principle of detection: observe for a long time and take long Fast Fourier Transforms (FFTs)



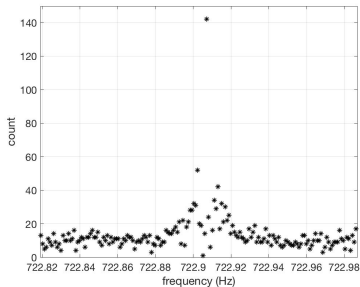
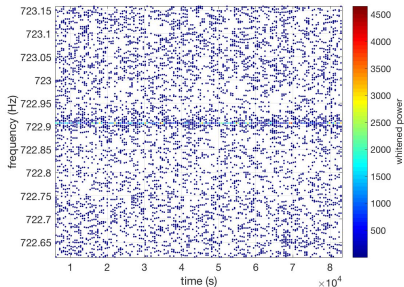
Cross-correlation method

- ▶ Separation between detectors \ll dark photon coherence length \rightarrow measurements should be almost identical
- ▶ Look for excess power above certain threshold, taking into account (1) misalignment and (2) spatial separation of detectors
- ▶ Upper limits in LIGO O1 data shown [7]



Continuous wave method to detect dark photons

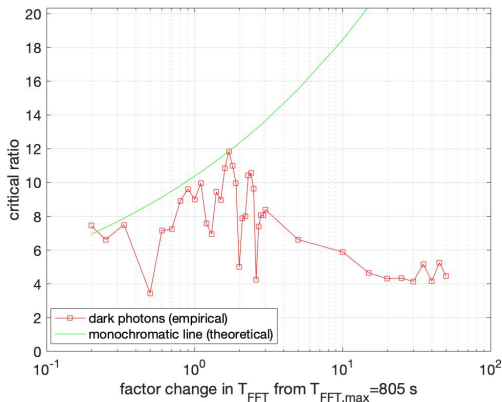
- ▶ Ideally, the FFT length should change based on the frequency/mass we consider
- ▶ A semicoherent method exists to search for boson clouds around black holes that varies FFT length [4]
 - ▶ The GW signal from a boson cloud was assumed to be monochromatic with a small random walk
- ▶ Follow up of dark photon dark matter signals [9]



Follow-up techniques

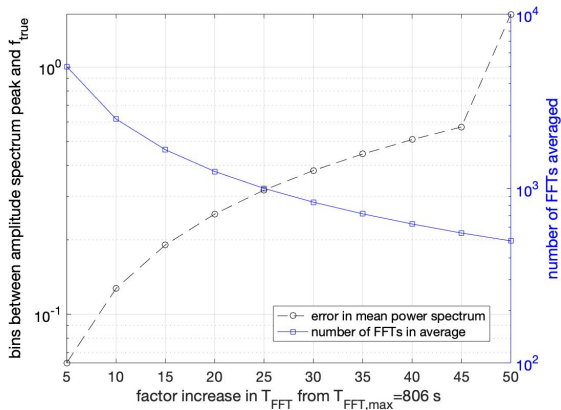
- ▶ Increase T_{FFT} and look for a *decrease* in signal-to-noise ratio
- ▶ Increase T_{FFT} and average power spectra, looking for a peak at $f_{true} = f_0(1 + \frac{1}{2}v_0^2/c^2)$
- ▶ Matched filter (in development)

Follow-up technique 1



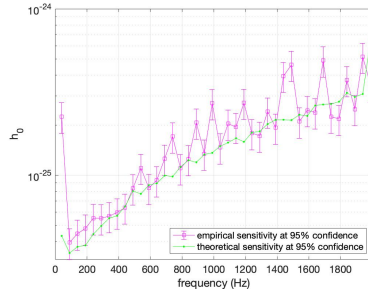
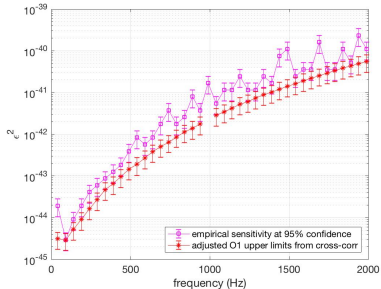
- ▶ $f_0 \simeq 744$ Hz; critical ratio falls off as T_{FFT} increases
- ▶ Injection done in LIGO O2 Livingston noise for whole O2 run; signal is recovered at 95% confidence

Follow-up technique 2



- ▶ Power spectra are averaged for different T_{FFT} . Error in bins from the simulated signal is plotted (no noise)

Sensitivity estimates



- ▶ Left: comparison of empirical sensitivity and adjusted upper limits in O1
 - ▶ We consider higher maximum ν for DPs \rightarrow lower T_{FFT}
 - ▶ Shorter T_{FFT} at higher frequencies
- ▶ Right: comparison of empirical and theoretical sensitivities

Conclusions

- ▶ Dark photons are a viable candidate of dark matter and can be searched with LIGO/Virgo
- ▶ Stochastic and continuous gravitational wave techniques are well-suited to detecting dark photon dark matter
- ▶ Combination of particle and gravitational wave physics
- ▶ End-to-end analysis scheme to detect dark photons
- ▶ Development of matched filter technique for follow-up step in progress
- ▶ Method to distinguish dark matter signals detected by interferometers in development

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Ultralight dark matter

- ▶ Axion-like particles
 - ▶ Pseudoscalar particle that solves the strong CP problem in QCD
 - ▶ Problem: no observed CP violations in QCD
 - ▶ Couples to gauge bosons (photons)
 - ▶ Also arise in compactifications of string theory [5]
- ▶ Dark photons: Gauge boson of the $U(1)_B$ or $U(1)_{B-L}$ group arising from [11]:
 - ① The misalignment mechanism
 - ▶ Field starts at nonzero, non-minimal vacuum value during inflation, then dissipates energy over time, eventually oscillating about the (very small) minimum in the potential
 - ② Tachyonic instability of a scalar field
 - ▶ The potential has a negative mass term, meaning perturbations to the field can cause energy emission in the form of particles
 - ③ Cosmic string network decays

Strain induced on detector

Start with a superposition of plane waves:

$$\vec{A} = \sum_{n=1}^N \vec{A}_{n,0} \sin \left(m \left(1 + \frac{1}{2} \left(\frac{v_n}{c} \right)^2 \right) t - \vec{k} \cdot \vec{x} + \phi_n \right) \quad (1)$$

$$\vec{E} = \partial_t \vec{A} \sim m \vec{A} \quad (2)$$

$$\vec{B} = \nabla \times \vec{A} \sim m v_j A_k \epsilon^{ijk} \quad (3)$$

$$\vec{a}_i = \frac{F}{M_i} = \epsilon e \frac{q_{D,i}}{M_i} \vec{A}_0 \cos(mt - \vec{k} \cdot \vec{x}_i) \quad (4)$$

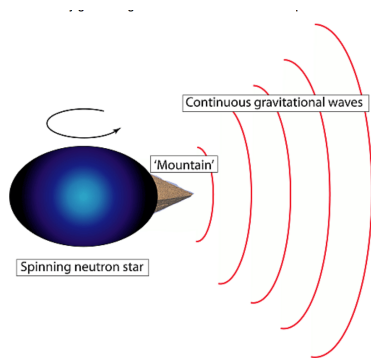
$|\vec{A}_0|$ normalized by current dark matter energy density, integrated over coherent time/space of signal

$$\Delta L(t) = [x_1(t) - x_2(t)] - [y_1(t) - y_2(t)] \quad (5)$$

ΔL contains amplitude and phase evolutions of the signal

Continuous gravitational waves

- ▶ Quasi-monochromatic, quasi-infinite signals
- ▶ Small frequency variations
- ▶ Principle of detection: observe for a long time and take long Fast Fourier Transforms



Stochastic Backgrounds

Normally the power spectrum of a stochastic gravitational wave background is:

$$S_{GW}(f) = \frac{3H_0^2}{2} f^{-3} \Omega_{GW}(f) \quad (6)$$

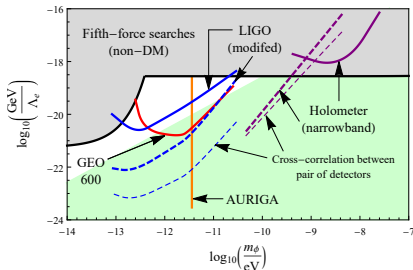
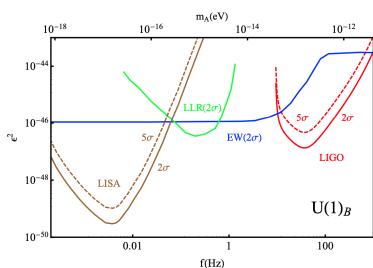
$$\Omega_{GW} = \frac{f}{\rho_c} \frac{d\rho_{GW}}{df} \quad (7)$$

For dark photons, $\rho_{GW} \propto f^2 h_0^2$:

$$S_{GW}(f) = \frac{h_0^2}{2\Delta f} \quad (8)$$

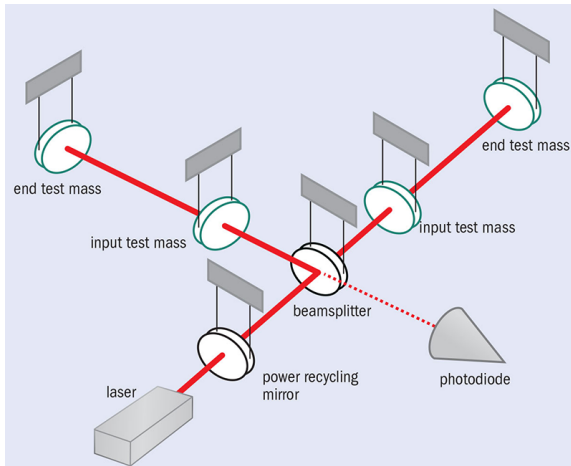
Almost constant background [11]

Projected constraints



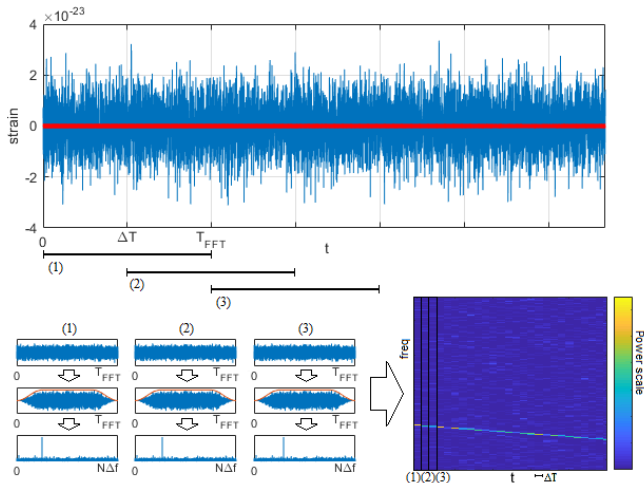
- ▶ Projected constraints on dark-photons (left) [12] and scalar dark matter (right) [6]
- ▶ Existing detectors can probe interesting portion of parameter space

LIGO/Virgo detectors



- Each incoming signal moves the mirrors uniquely, creating an interference pattern

Method overview



► Time series \rightarrow power spectrum, whitening \rightarrow spectrogram

Continuous wave method to detect dark photons

For each detector:

1. Cleaned, reduced analytic time series

2. Peakmap every 10 Hz with different T_{FFT}

3. Peakmap projection onto frequency axis

4. Candidate selection

5. Coincidences

6. Follow-up

- ▶ Ideally, the FFT length should change based on the frequency/mass we consider; based on [4]

Computational cost

- ▶ Search begins from 10 Hz BSD files [10]
- ▶ T_{FFT} chosen based on maximum frequency in each 10 Hz band
- ▶ Sensitivity loss $< 10\%$ compared to using 1 Hz bands
- ▶ Using 1 Xeon CPU E5-2695 v2 and analyzing a 1990-2000 Hz band takes ~ 350 s
- ▶ Extrapolating to 20-2000 Hz and a one year of observation \rightarrow 1.6 days
- ▶ Dividing the analysis on 2000 cores \rightarrow search takes $O(\text{mins})$