

Precision studies for Drell-Yan processes at NNLO

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RADCOR-LoopFest 2021, Florida State University, Tallahassee (via Zoom), May 19, 2021

Based on work done in collaboration with:

- *Precision studies for Drell-Yan processes at NNLO*
Sergey Alekhin, Adam Kardos, S. M. and Zoltan Trócsányi
[arXiv:2104.02400](https://arxiv.org/abs/2104.02400)

W - and Z-boson cross sections

Status (I)

Data

- High precision experimental data from LHC **ATLAS, CMS, LHCb** and Tevatron **D0** useful for determinations of parton distributions
 - statistically significant $N_{DP} = 172$ in **ABMP16**
- Differential distributions in decay lepton pseudo-rapidity extend kinematics to forward region
 - sensitivity to light quark flavors at $x \simeq 10^{-4}$
 - leading order kinematics with:
 $\sigma(W^+) \simeq u(x_2)\bar{d}(x_1)$ and $\sigma(W^-) \simeq d(x_2)\bar{u}(x_1)$;
 $\sigma(Z) \simeq Q_u^2 u(x_2)\bar{u}(x_1) + Q_d^2 d(x_2)\bar{d}(x_1)$
 - cf. DIS: $\sigma(\text{DIS}) \simeq q_u^2 u(x) + q_d^2 d(x)$

Status (II)

Theory

- Complete NNLO QCD corrections with fully differential kinematics to match experimental cuts \longrightarrow solved problem
 - combination of squared matrix elements with three different multiplicities of partons in final state
 - subtraction scheme to cancel of soft and collinear singularities upon integration over their phase space
- Public codes at NNLO including the leptonic decay
 - FEWZ (v3.1) (sector decomp.) Gavin, Li, Petriello, Quackenbush '12
 - DYNNLO (v1.5) (q_T slicing) Catani, Grazzini '07
 - MATRIX (v1.0.4) (q_T slicing) Grazzini, Kallweit, Wiesemann '17
 - MCFM at NNLO (N -jettiness slicing) Boughezal, Campbell, Ellis, Focke, Giele, Liu, Petriello, Williams '16
- Also:
 - (SHERPA-NNLO-FO (q_T slicing) Höche, Li, Prestel '14)
 - (DYTURBO (q_T slicing, based on DYNNLO (v1.5)) Camarda, Boonekamp, Bozzi et al. '19)

Regularization

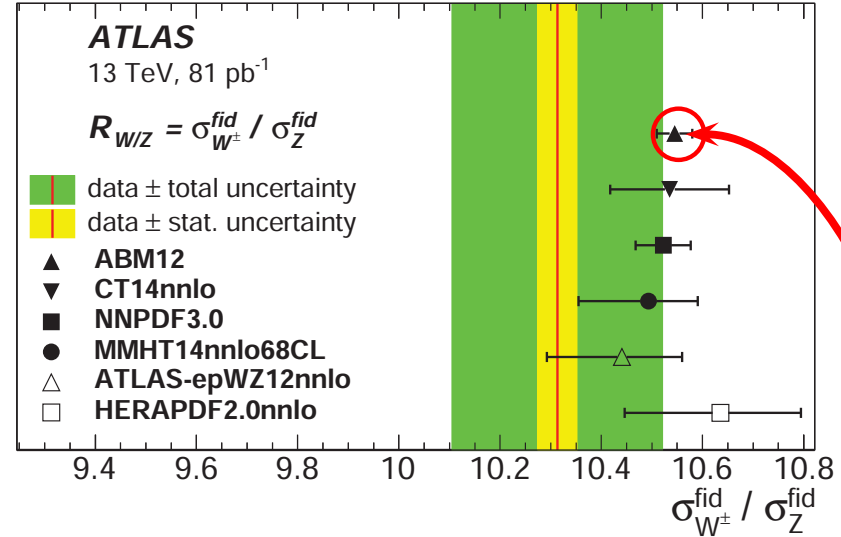
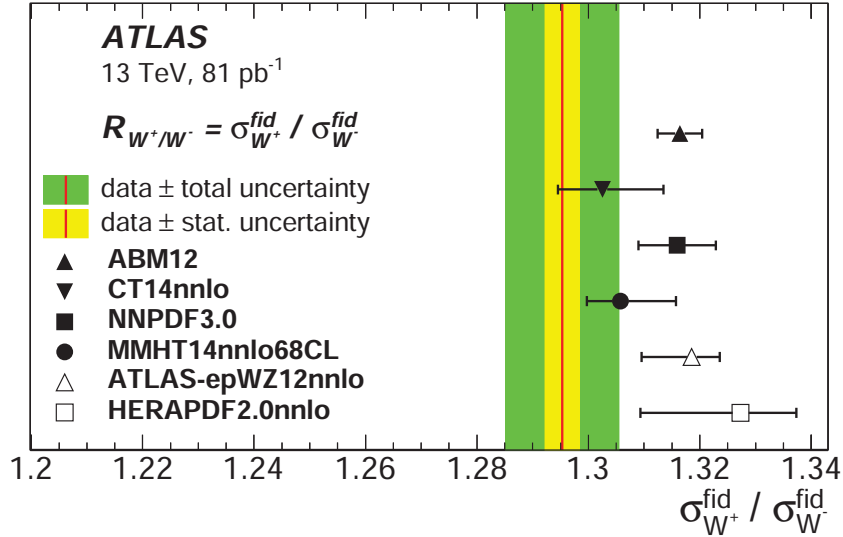
Subtraction

- removes all $\frac{1}{\epsilon^k}$ singularities in $d = 4 - 2\epsilon$ analytically; $k \leq 4$ at NNLO
 - Antenna subtraction Gehrmann-De Ridder, Gehrmann, Glover '05
 - Colourful subtraction Del Duca, Somogyi, Trocsanyi '05
 - Sector subtraction Czakon '10; Boughezal, Melnikov, Petriello '11; Czakon, Heymes '14
 - Projection to Born Cacciari, Dreyer, Karlberg, Salam, Zanderighi '15
 - Local analytic sector subtraction Magnea, Maina, Pelliccioli, Signorile-Signorile, Torrielli, Uccirati '18

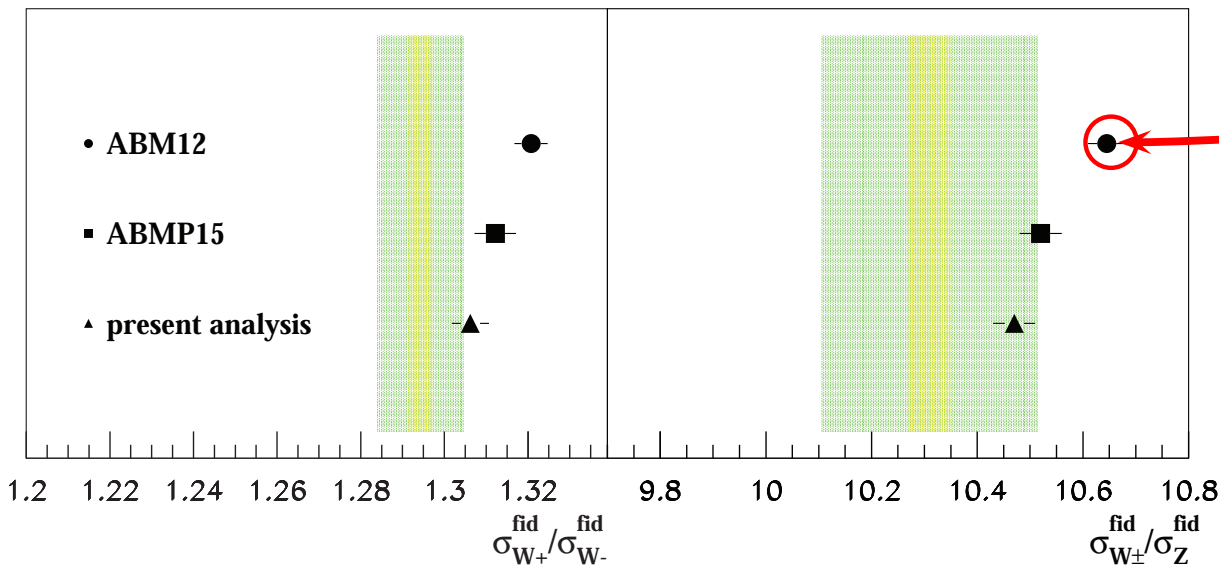
Slicing

- Cuts imposed in phase space; singularities \ln^k (some cut) need to cancel numerically; $k \leq 4$ at NNLO
 - q_T -subtraction Catani, Grazzini '07
 - N -jettiness subtraction Boughezal, Focke, Liu, Petriello '15; Gaunt, Stahlhofen, Tackmann, Walsh '15

Theory issues



ATLAS (13 TeV, 81 pb⁻¹) 1603.09222



DYNNLO (v1.5)

Catani, Grazzini '07

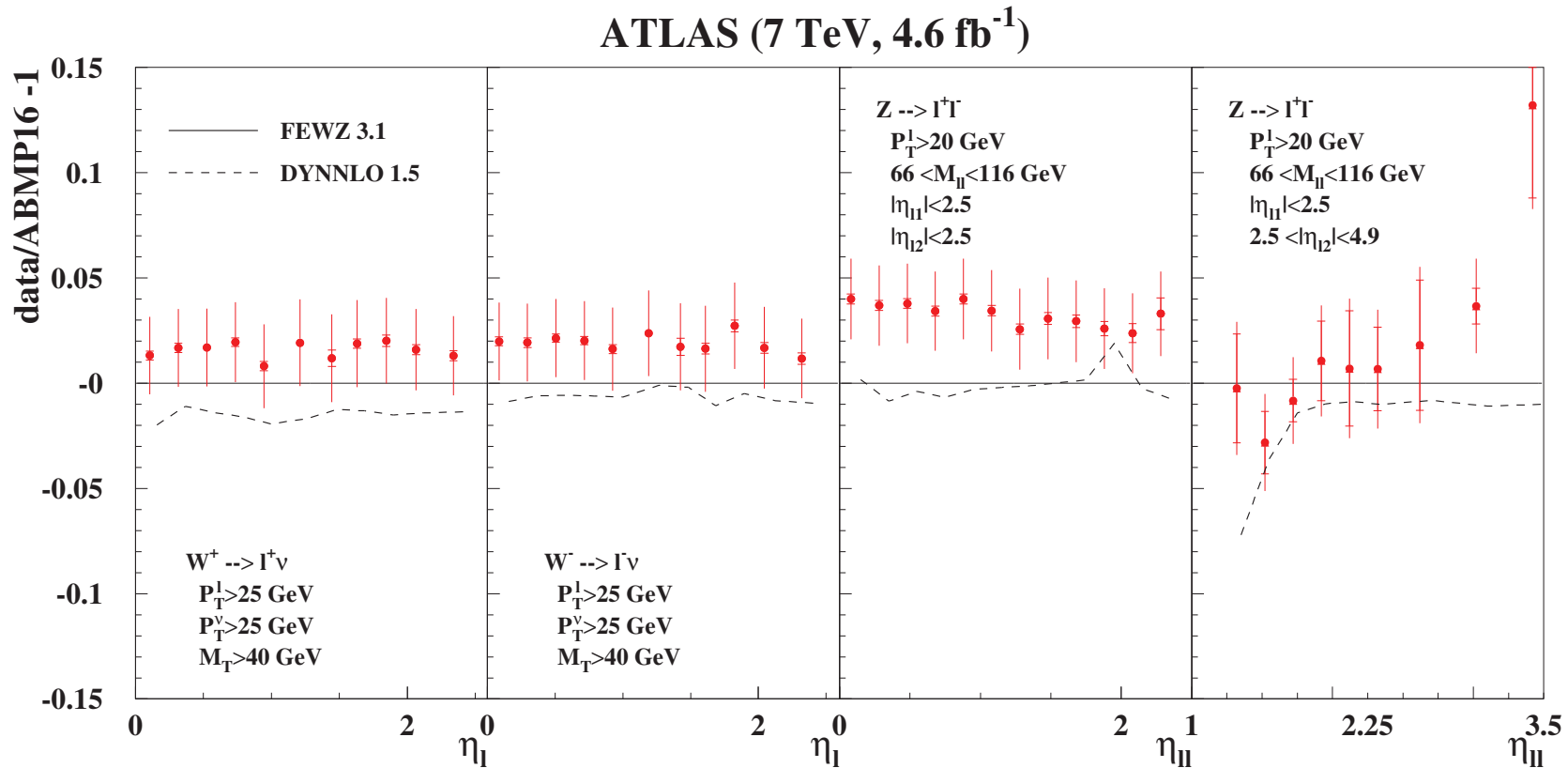
FEWZ (v3.1)

Gavin, Li, Petriello,
Quackenbush '12

- Differences at NNLO between **DYNNLO** and **FEWZ** up to $\mathcal{O}(1\%)$ or more

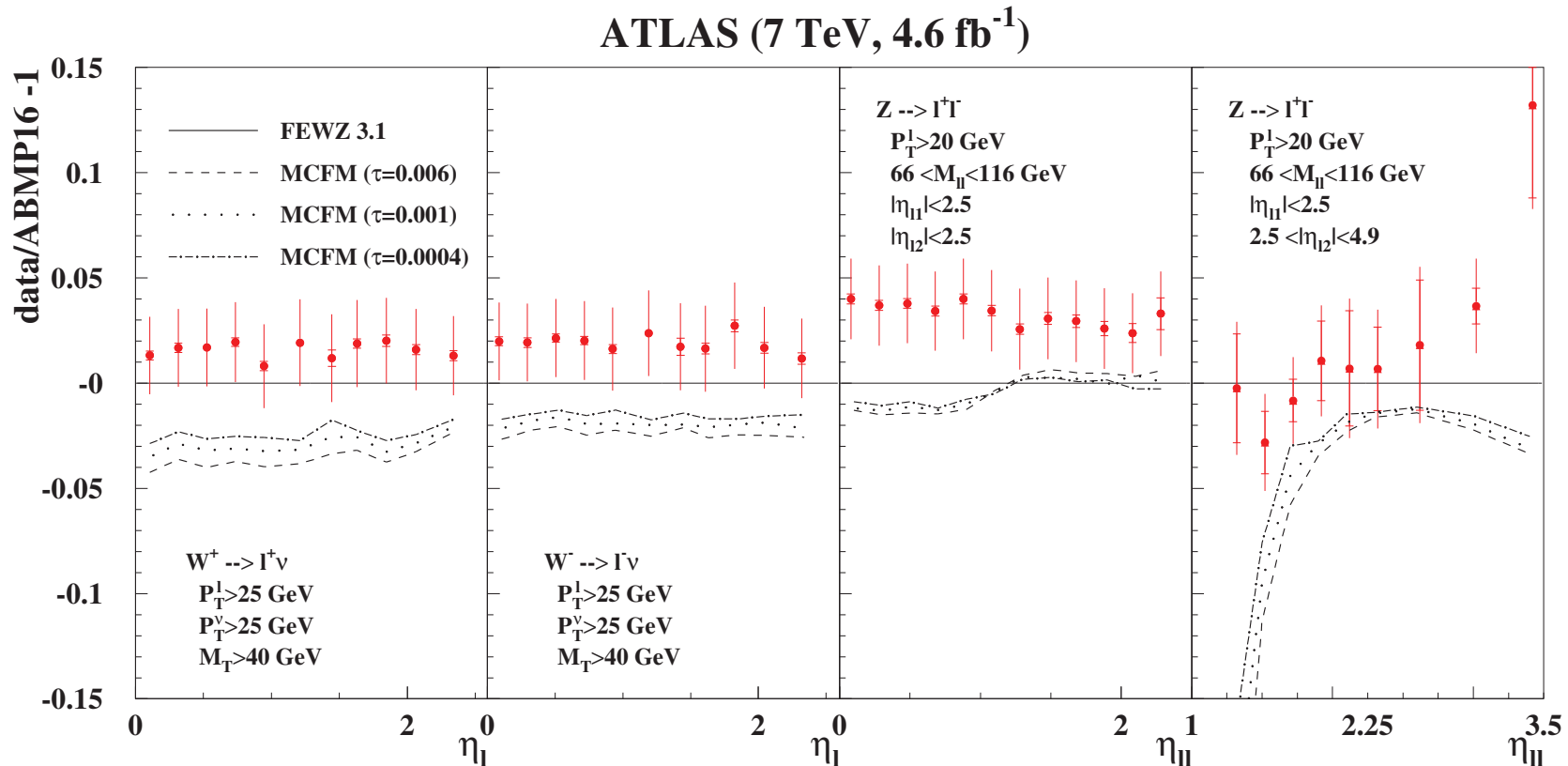
Benchmark computations

Benchmarking *DYNNLO*



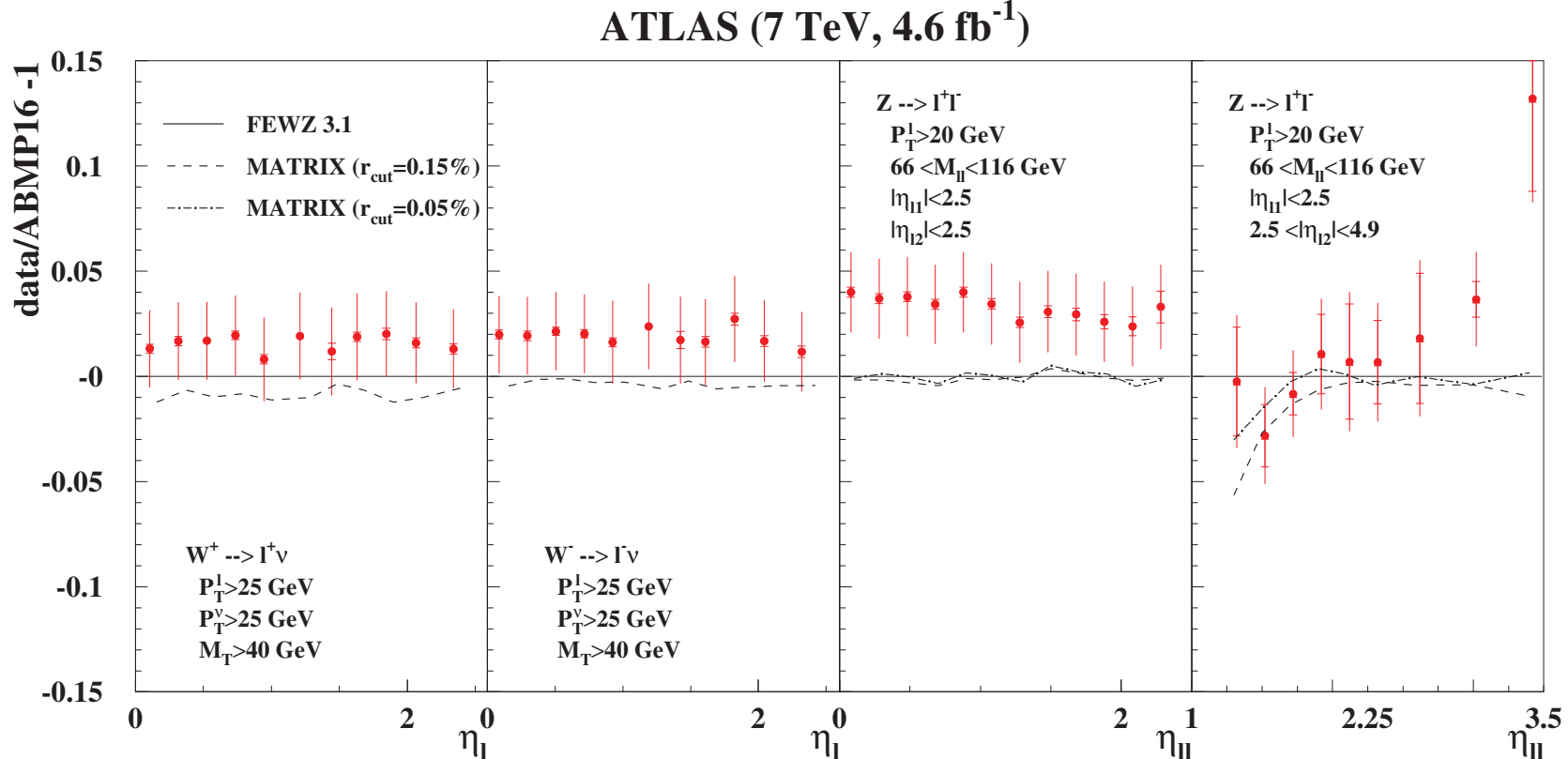
- Deviations of *DYNNLO* (q_T slicing) at NNLO
 - up to $\mathcal{O}(1\%)$ for W^+ -production
 - few per mill for W^- - and Z -production
 - up to $\mathcal{O}(2 - 3\%)$ for forward Z -production, $\mathcal{O}(10\%)$ in first bin

Benchmarking MCFM



- Substantial deviations of **MCFM** at NNLO (N -jettiness slicing)
- Variation of τ_{cut} (smallest value $\tau_{\text{cut}} = 4 \cdot 10^{-4}$ at computational limits)
 - up to $\mathcal{O}(3\%)$ for W^+ -production
 - up to $\mathcal{O}(2\%)$ for W^- - and Z -production
 - up to $\mathcal{O}(2 - 3\%)$ for forward Z -production, $\mathcal{O}(20\%)$ in first bin

Benchmarking MATRIX

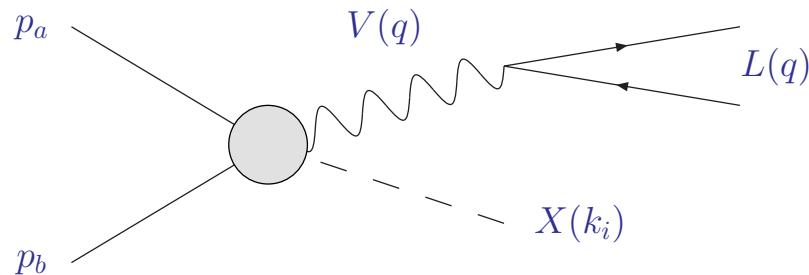


- Findings for **MATRIX** similar to **DYNNLO** (both based on q_T slicing)
 - up to $\mathcal{O}(1\%)$ for W^+ -production
 - few per mill for W^- - and Z -production
 - up to $\mathcal{O}(1 - 2\%)$ for forward Z -production, $\mathcal{O}(5\%)$ in first bin
- Variation of r_{cut} has similar effect as smaller τ_{cut} in **MC_{CFM}**

Slicing methods (I)

- W - and Z -boson production with decay to leptonic final state L
 - gauge boson mass $q^2 = Q^2$

$$a(p_a) + b(p_b) \rightarrow V(q) + X(k_i) \rightarrow L(q) + X(k_i)$$



- Slicing parameter τ for q_T -subtraction defined with $\vec{k}_{T,i}$ (transverse momenta of hadronic final states) Catani, Grazzini '07

$$\tau = q_T^2 / Q^2 = \left(\sum_i \vec{k}_{T,i} \right)^2 / Q^2$$

- Slicing parameter τ for leptonic 0-jettiness \mathcal{T}_0

Boughezal, Focke, Liu, Petriello '15; Gaunt, Stahlhofen, Tackmann, Walsh '15

$$\tau = \mathcal{T}_0 / Q = \sum_i \min \{ 2p_a \cdot k_i, 2p_b \cdot k_i \} / Q^2$$

Slicing methods (II)

- Definitions of τ vanish at Born level and resolve additional radiation in infrared-safe manner
- Phase space integration for cross section with cut for slicing τ_{cut}

$$\sigma = \int d\tau \frac{d\sigma}{d\tau} = \int^{\tau_{\text{cut}}} d\tau \frac{d\sigma}{d\tau} + \int_{\tau_{\text{cut}}} d\tau \frac{d\sigma}{d\tau} = \sigma(\tau_{\text{cut}}) + \int_{\tau_{\text{cut}}} d\tau \frac{d\sigma}{d\tau}$$

- Dependence of $d\sigma/d\tau$ on τ predicted from universal QCD factorization in soft and collinear limits
 - power corrections proportional to τ^{p-1} with $p > 0$ are integrable

$$\frac{d\sigma}{d\tau} \sim \delta(\tau) + \sum_i \left[\frac{\ln^i \tau}{\tau} \right]_+ + \sum_j \tau^{p-1} \ln^j \tau + \mathcal{O}(\tau^p)$$

- Analytical result for $\sigma(\tau_{\text{cut}})$

$$\sigma(\tau_{\text{cut}}) \sim 1 + \sum_i \ln^{i+1} \tau_{\text{cut}} + \sum_j \tau_{\text{cut}}^p \ln^j \tau_{\text{cut}} + \mathcal{O}(\tau_{\text{cut}}^{p+1})$$

- Subtraction scheme implemented via global subtraction term $\sigma^{\text{sub}}(\tau_{\text{cut}})$

$$\sigma = \sigma^{\text{sub}}(\tau_{\text{cut}}) + \int_{\tau_{\text{cut}}} d\tau \frac{d\sigma}{d\tau} + \Delta\sigma^{\text{sub}}(\tau_{\text{cut}})$$

Power corrections (I)

- Residual power corrections are numerically sizable

$$\Delta\sigma^{\text{sub}}(\tau_{\text{cut}}) = \sigma(\tau_{\text{cut}}) - \sigma^{\text{sub}}(\tau_{\text{cut}})$$

- difference between $\sigma(\tau_{\text{cut}})$ and global subtraction term $\sigma^{\text{sub}}(\tau_{\text{cut}})$

Scaling behavior of power corrections

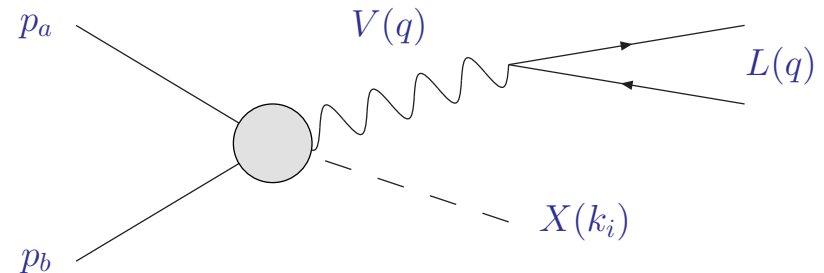
- Exponent p in power corrections τ_{cut}^p in $\sigma(\tau_{\text{cut}})$
- Stable gauge boson V
 - positive integer values for exponent $p = 1, 2, 3, \dots$
 - power corrections scale as $\tau = q_T^2/Q^2$ or \mathcal{T}_0/Q
- Decay of V to leptonic final state L
 - p_T -cuts on leptons in decay change power counting of power corrections
 - half-integers values $p = 1/2, 1, 3/2, \dots$ for exponent p
 - power corrections scale as $\sqrt{\tau} = q_T/Q$ or $\sqrt{\mathcal{T}_0/Q}$

Ebert, Tackmann '19

Cuts on leptons in decay (I)

- Experimental data has cuts on p_T of leptons
 - **ATLAS** data with $p_T^l \geq 20\text{GeV}$ for decay of Z -boson (or $p_T^{l/\nu} \geq 25\text{GeV}$ for W -boson)
- p_T -cuts break azimuthal symmetry in phase space integration
 - recall e.g. Drell-Yan process with $L(q) \rightarrow l^-(p_1) + l^+(p_2)$

$$a(p_a) + b(p_b) \rightarrow V(q) + X(k_i) \rightarrow l^-(p_1) + l^+(p_2) + X(k_i)$$



- Leptonic final state phase space Φ_L

$$\Phi_L(q_T) = \frac{1}{4\pi^2} \int_0^\pi d\phi \int_{-\infty}^\infty d\Delta y \frac{p_{T1}^2}{Q^2}$$
 - ϕ (azimuthal angle) and Δy (difference in rapidity between l^- and l^+)
- Lepton p_T -cuts $p_{T1}, p_{T2} \geq p_T^{\min}$ induce power corrections of order $\vec{k}_{T,i}$ for hadronic final state $X(k_i)$

Cuts on leptons in decay (II)

- p_T -cuts restrict phase space to $\min \{p_{T1}, p_{T2}\} \geq p_T^{\min}$
 - $(p_{T1})^2 = p_T^2$
 - $(p_{T2})^2 = (q_T + p_{T1})^2 = p_T^2 - 2p_T q_T \cos \phi + q_T^2$
- Azimuthal integration for small q_T under condition

$$\min \{p_T, p_T^2 - 2p_T q_T \cos \phi\} = \begin{cases} p_T, & \cos \phi \leq 0 \\ p_T^2 - 2p_T q_T \cos \phi, & \cos \phi > 0 \end{cases}$$

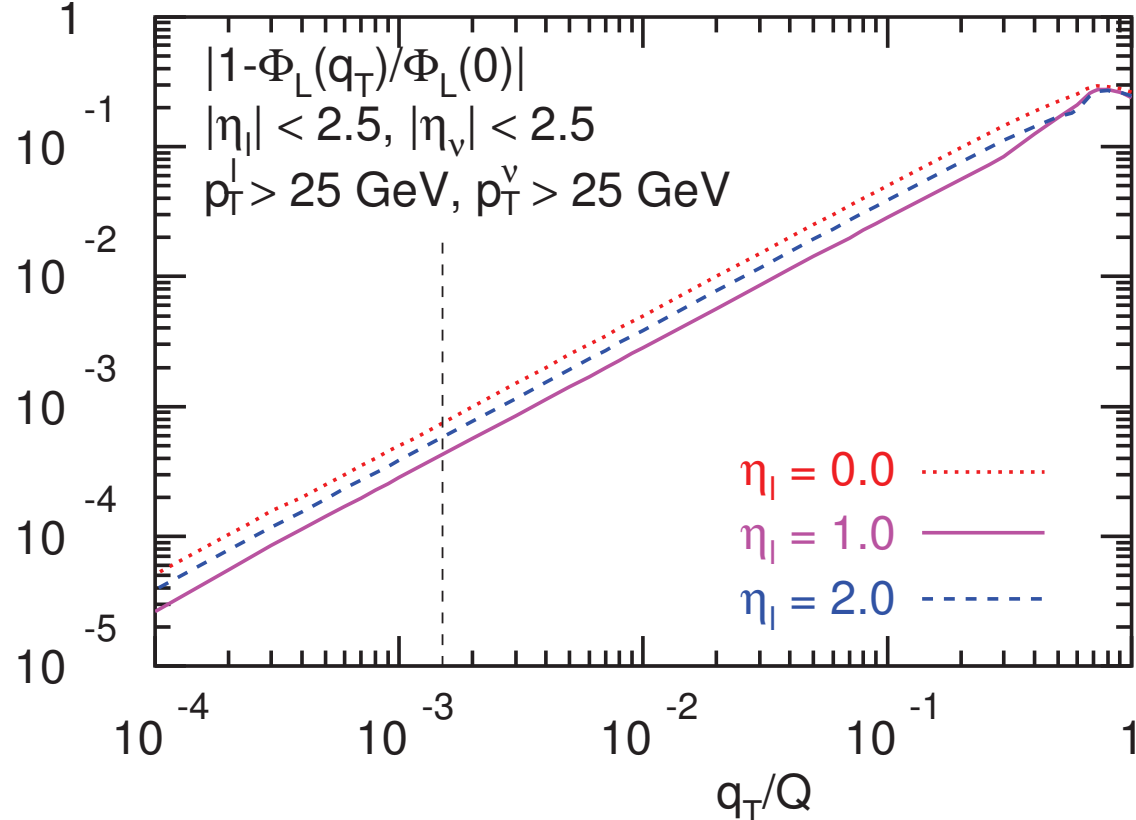
- Leptonic phase space $d\Phi_L$ after integration yields

$$d\Phi_L = \pi \sqrt{1 - (2p_T^{\min})^2/Q^2} - \frac{4}{\sqrt{1 - (2p_T^{\min})^2/Q^2}} \frac{p_T^{\min} q_T}{Q^2}$$

- η_l -cuts impose additional restriction on azimuthal angle ϕ
 - value ϕ^* for boundary of phase space depends on gauge-boson rapidity Y and on $\Delta y = \eta_1 - \eta_2$

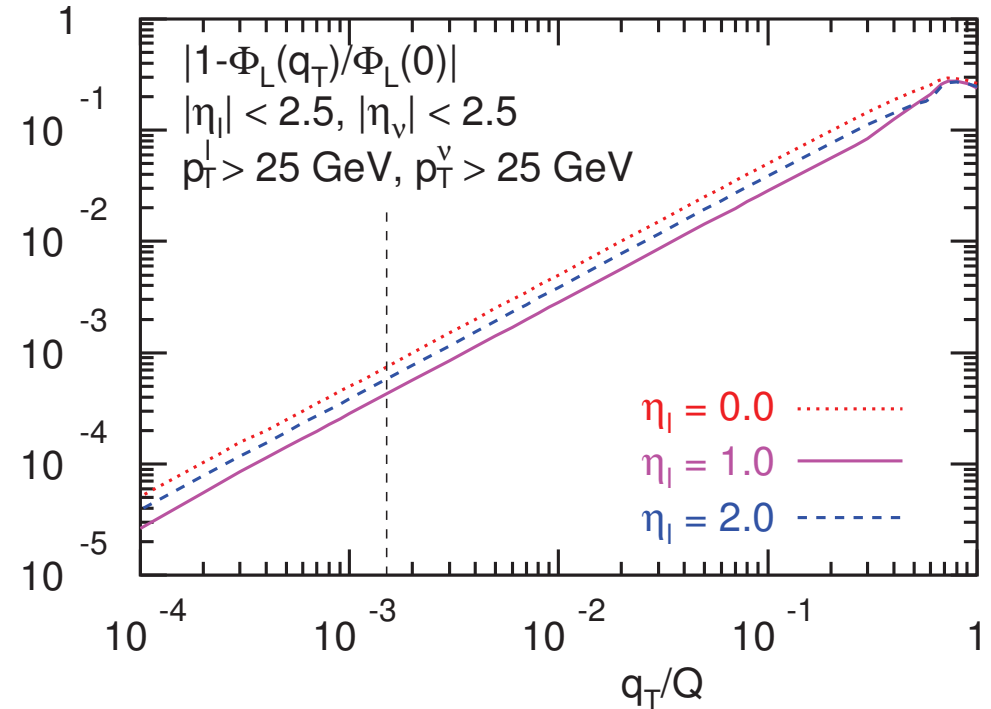
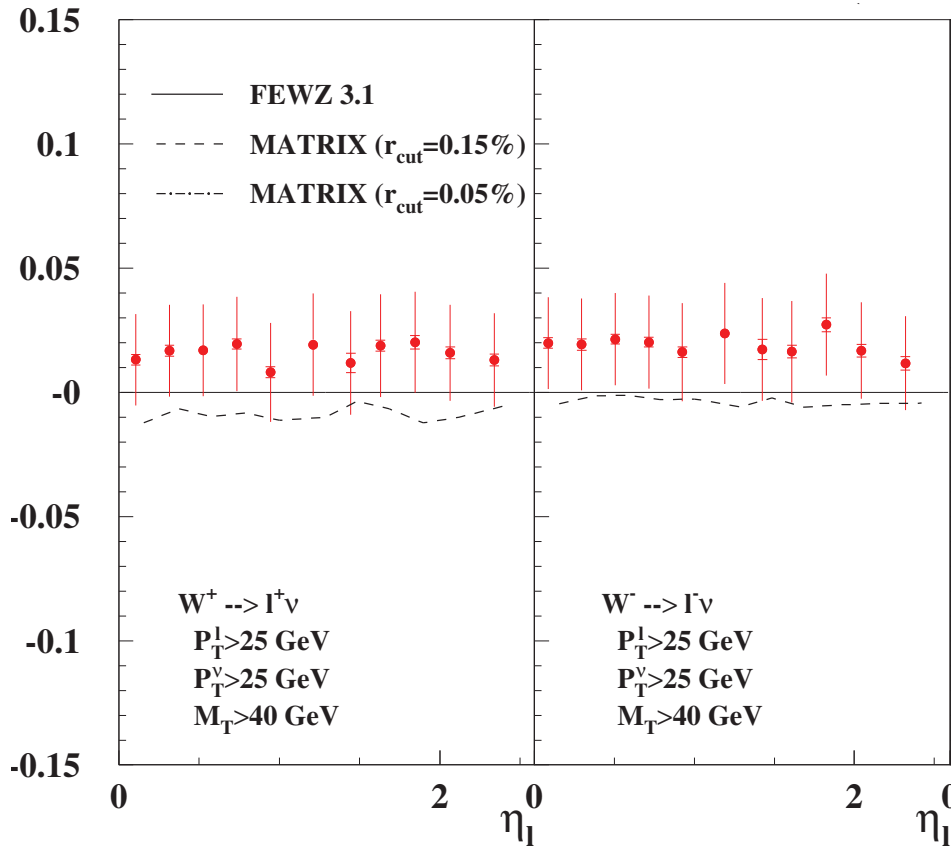
$$\cos \phi^* = \frac{Q}{2q_T} \frac{\sinh(2Y)}{\sinh(2Y + \Delta y)} + \mathcal{O}(q_T/Q)$$

Lepton decay phase space (I)



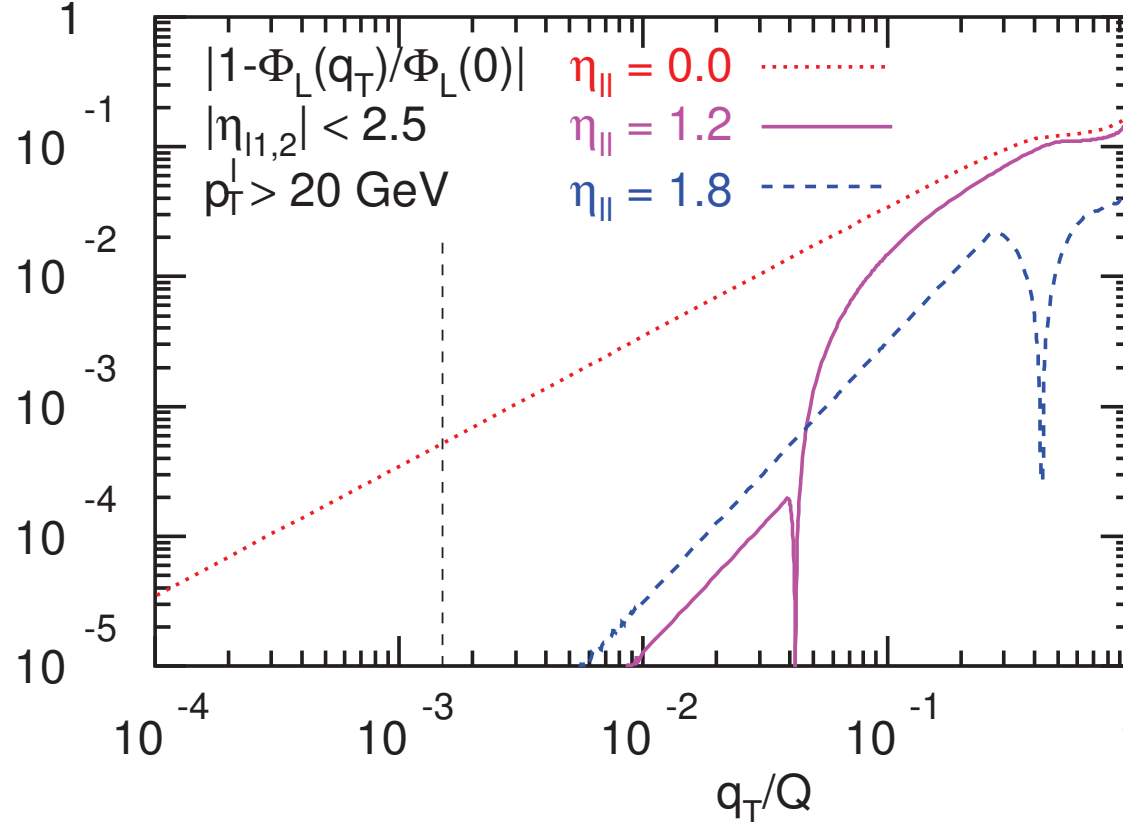
- Plot of $|1 - \Phi_L(q_T)/\Phi_L(0)|$ for W^\pm -production with cuts in ATLAS data
 - deviations of $\Phi_L(q_T)$ from Born level leading power results for $q_T = 0$
- Pseudo-rapidity distributions for W^\pm -production fix η_l
- p_T -cuts lead to linear power corrections in q_T
- MATRIX technical cutoff $r_{\text{cut}}^{\text{min}} = 0.15\%$ indicated by vertical dashed line

W -production with MATRIX



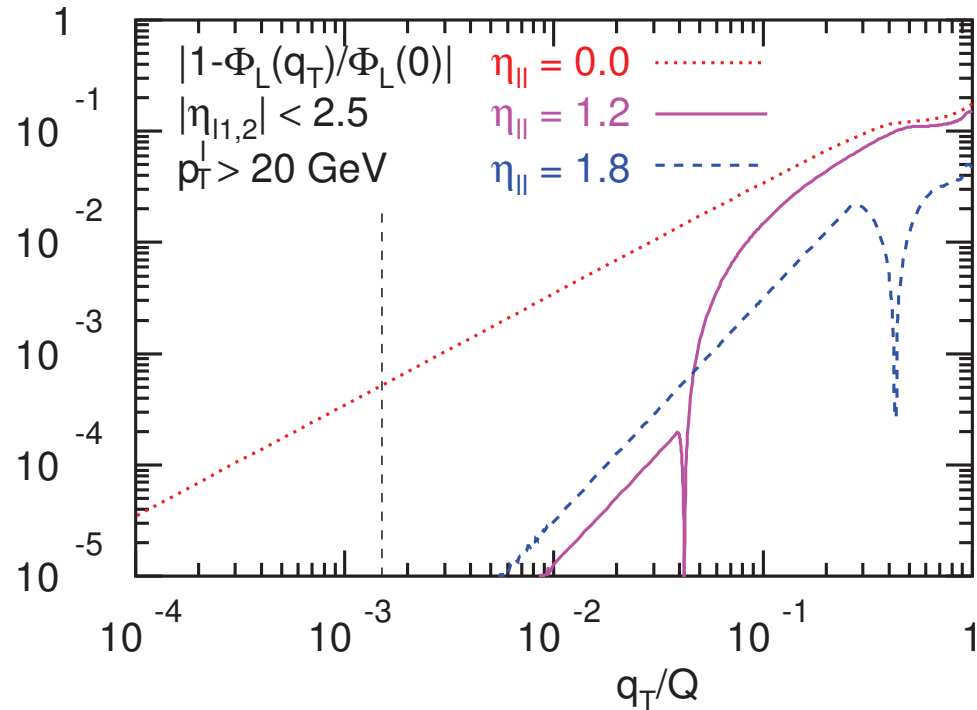
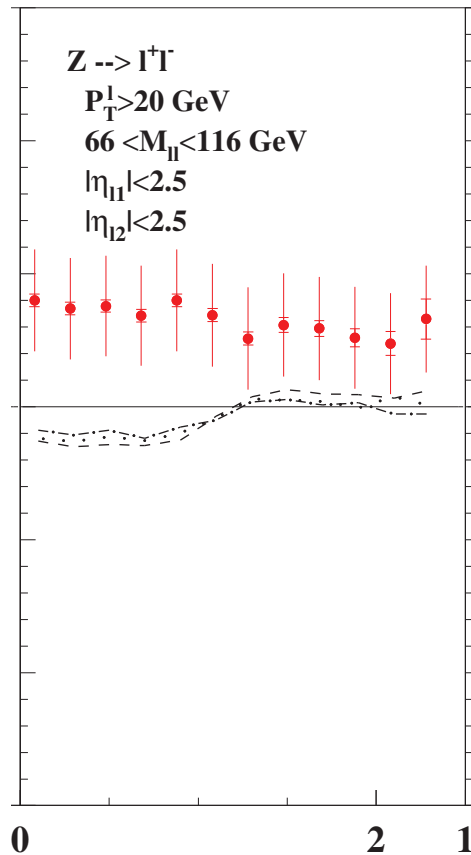
- Linear power corrections in q_T in leptonic phase space $d\Phi_L$ present in all η_l bins
- Deviations of **MATRIX** from **FEWZ** uniform across all η_l bins
 - up to $\mathcal{O}(1\%)$ for W^+ -production
 - few per mill for W^- -production

Lepton decay phase space (II)



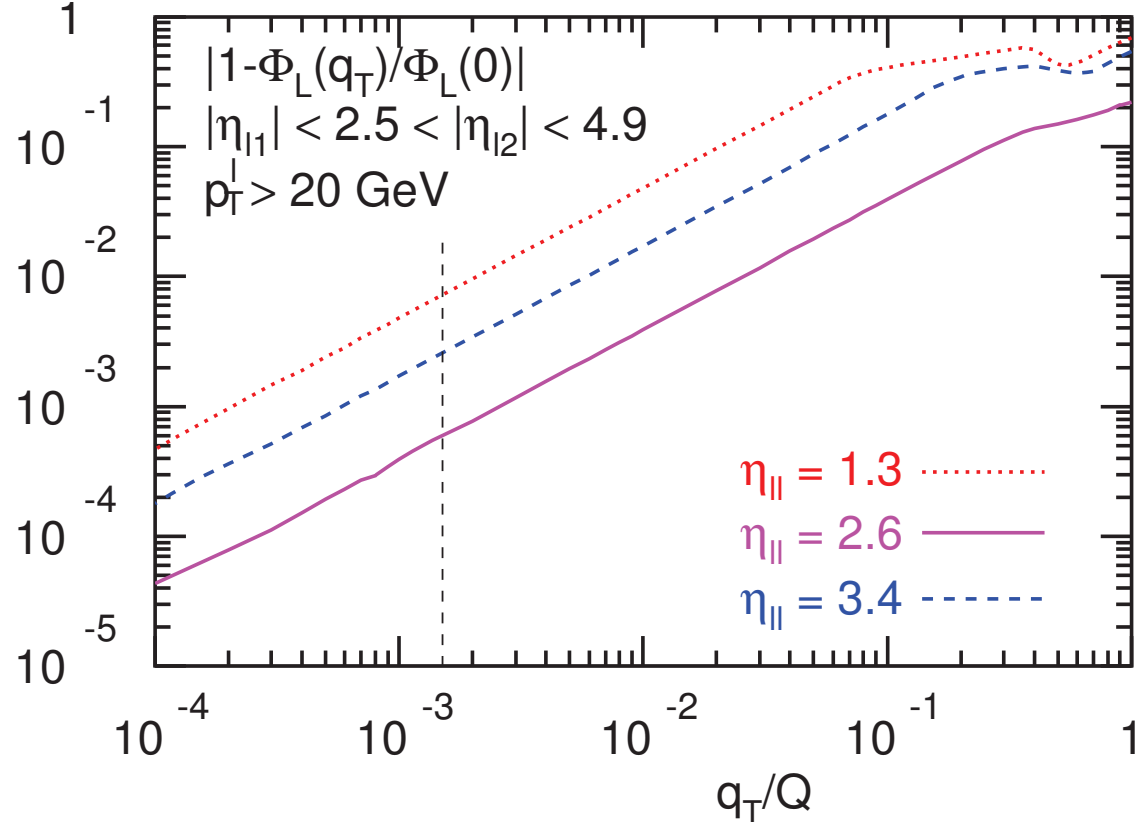
- Plot of $|1 - \Phi_L(q_T)/\Phi_L(0)|$ for central Z -production with **ATLAS** cuts
- Power corrections depend on gauge boson pseudo-rapidity $\eta_{||}$
- Small $\eta_{||}$: cuts on p_T dominate and lead to linear power corrections in q_T
- Large $\eta_{||}$: cuts on pseudo-rapidities η_{l_1}, η_{l_2} dominate
 - azimuthal symmetry is restored for small q_T , when $|\cos \phi^*| \geq 1$
 - transition between linear and quadratic behavior around $\eta_{||} \simeq 1.2$

Central Z-production with MCFM



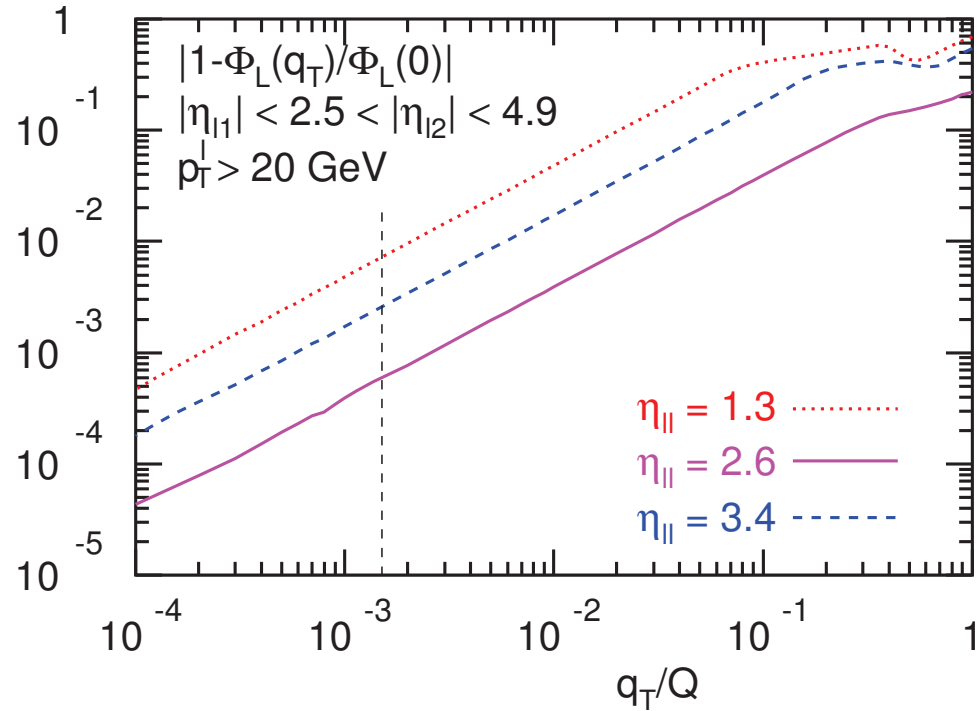
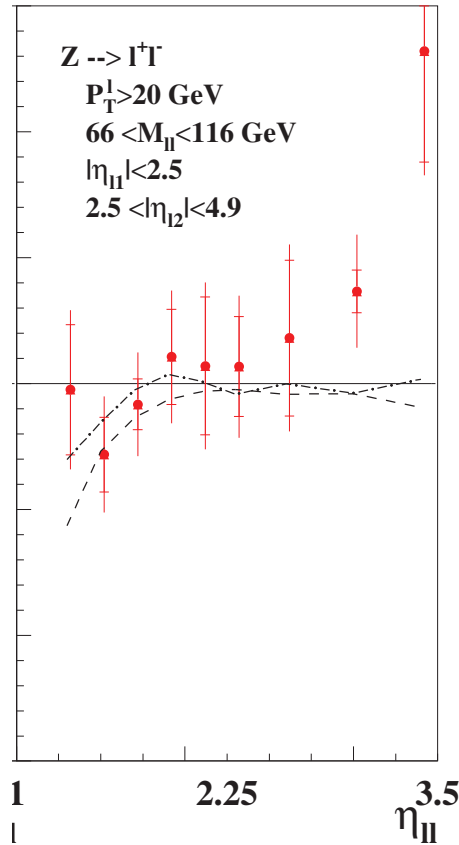
- Deviations of MCFM from FEWZ display particular pattern as function of pseudo-rapidity η_{ll}
 - consistent with appearance/disappearance of linear power corrections in q_T in leptonic phase space $d\Phi_L$

Lepton decay phase space (III)



- Plot of $|1 - \Phi_L(q_T)/\Phi_L(0)|$ for forward Z -production with ATLAS cuts
- Cuts on pseudo-rapidities $|\eta_{||1}| \leq 2.5 \leq |\eta_{||2}| \leq 4.9$ do no overlap
- p_T -cuts lead to linear power corrections in q_T
- Size of linear power corrections depend on pseudo-rapidity $\eta_{||}$
 - power corrections for small and large $\eta_{||}$ are order of magnitude larger than for W^\pm -production

Forward Z-production with MATRIX



- Linear power corrections in q_T in leptonic phase space $d\Phi_L$ present in all η_{ll} bins
 - $d\Phi_L(\eta_{ll} = 1.3) \gg d\Phi_L(\eta_{ll} = 3.6) \gg d\Phi_L(\eta_{ll} = 2.4)$
- Size of deviations of MATRIX from FEWZ display same pattern as function of pseudo-rapidity η_{ll}

Summary

W^{\pm} - and Z -boson production at the LHC

- Currently available public codes with short-comings
 - long run times
 - problems with precision (short-comings of slicing methods)
- Sizable differences between public codes
 - differences at least similar, sometimes larger than sizes of NNLO corrections
- Linear power corrections from cuts on final state momenta affect precision of NNLO results

Upshot

- Need for fast and precise public code for differential distribution at NNLO (per mill level accuracy)