

The tropical approach to numerical Feynman integration

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Talk based on [arXiv:2008.12310](https://arxiv.org/abs/2008.12310),
to appear in Annales de l'Institut Henri Poincaré D

Motivation

Quantum field theory

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perturbative expansions

$$\mathcal{O}(\hbar) = \sum_{n \geq 0} A_n \hbar^n = \sum_{\text{graphs } G} \frac{\phi(G)}{|\text{Aut } G|} \hbar^{L_G}$$

$$\text{where } \phi(G) = \prod_{\ell} \int d^D \mathbf{k}_{\ell} \prod_{e \in E} \frac{1}{D_e(\{\mathbf{k}\}, \{\mathbf{p}\}, m_e)}.$$

A computational question

$$\mathcal{O}(\hbar) = \sum_{n \geq 0} A_n \hbar^n$$

Lower orders A_0, A_1, A_2, \dots needed to interpret experimental data.

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What is the value of $A_0, A_1, A_2, A_3, \dots$?

How can we calculate them efficiently?

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Practical question

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How can we calculate them efficiently?

Associated 'meta' question

Is there an algorithm to compute A_n ?

What is its runtime?

$$\mathcal{O}(\hbar) = \sum_{n \geq 0} A_n \hbar^n = \sum_{\text{graphs } G} \frac{\phi(G)}{|\text{Aut } G|} \hbar^{L_G}$$

Runtime to compute A_n for n large:

$$\mathcal{O}\left(\underbrace{\alpha^n \Gamma(n + \beta)}_{\text{number of Feynman graphs}} \times \underbrace{F(n)}_{\text{time to evaluate a single Feynman integral}}\right)$$

→ number of Feynman graphs
with n loops for $n \rightarrow \infty$
(α and β depend on theory
and observable)

→ time it takes to
evaluate a single
Feynman integral
of order n

$F(n)$ = time to evaluate n -loop Feynman integral

'Analytic calculation':

- Unclear: What is an analytic answer for an integral?
- Can ask for an expression within a specific function space
- No function space is known that works for all F. integrals
⇒ complicated to formulate the question

⇒) What counts as an analytic answer?

analytic answer
to computation

(\Leftrightarrow)

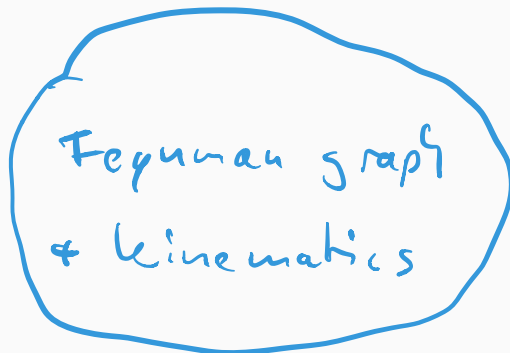
fast algorithm to
perform computation
numerically

In this talk

Cut the middle man!

→ direct numerical evaluation

INPUT



OUTPUT



Direct evaluation

$$\mathcal{O}(\hbar) = \sum_{n \geq 0} A_n \hbar^n = \sum_{\text{graphs } G} \frac{\phi(G)}{|\text{Aut } G|} \hbar^{L_G}$$

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Via the Schwinger trick we can rewrite the Feynman integral as

$$\phi(G) = \Gamma(\omega_G) \int_{\mathbb{P}_{>0}^{E-1}} \frac{\Omega}{\Psi_G(\mathbf{x})^{D/2}} \left(\frac{\Psi_G(\mathbf{x})}{\Phi_G(\mathbf{x})} \right)^{\omega_G}.$$

with $\omega_G = E - \frac{1}{2} Dh_1(G)$.

$$\int_{\mathbb{P}_{>0}^{E-1}} \frac{\Omega}{\Psi_G(\mathbf{x})^{D/2}} \left(\frac{\Psi_G(\mathbf{x})}{\Phi_G(\mathbf{x})} \right)^{\omega_G}$$

where

$$\int_{\mathbb{P}_{>0}^{E-1}} \frac{\Omega}{\Psi_G(\mathbf{x})^{D/2}} \left(\frac{\Psi_G(\mathbf{x})}{\Phi_G(\mathbf{x})} \right)^{\omega_G}$$

where

- Ω is the standard volume form on \mathbb{P}^{E-1} :

$$\Omega = \sum_{k=1}^E (-1)^k dx_1 \wedge \dots \wedge \widehat{dx_k} \wedge \dots \wedge dx_E.$$
- $\Psi_G = \sum_T \prod_{e \notin T} x_e$ (sum over spanning trees)
- $\Phi_G = \sum_F \|\mathbf{p}(F)\|^2 \prod_{e \notin F} x_e + \Psi_G \sum_e m_e^2 x_e$ (sum over 2-forests)
- Ψ_G and Φ_G are **homogeneous** polynomials in x_1, \dots, x_E .
- We assume that the integral exists.

$$\int_{\mathbb{P}_{>0}^{E-1}} \frac{\Omega}{\Psi_G^{D/2}(\mathbf{x})} \left(\frac{\Psi_G(\mathbf{x})}{\Phi_G(\mathbf{x})} \right)^{\omega_G}$$

$\Psi_G(\mathbf{x})$ and $\Phi_G(\mathbf{x})$ exhibit complicated geometric structures.

⇒ These integrals are hard to evaluate

⇒ These integrals are very interesting

$$\int_{\mathbb{P}_{>0}^{E-1}} \frac{\Omega}{\Psi_G^{D/2}(\mathbf{x})} \left(\frac{\Psi_G(\mathbf{x})}{\Phi_G(\mathbf{x})} \right)^{\omega_G}$$

Obstruction for direct numerical evaluation

Integrand has singularities on the boundary of $\mathbb{P}_{>0}^{E-1}$.

I.e. vanishing locus of Ψ_G and Φ_G meets the boundary of $\mathbb{P}_{>0}^{E-1}$.

⇒ Singularities have to be blown up first

$$\int_{\mathbb{P}_{>0}^{E-1}} \frac{\Omega}{\Psi_G^{D/2}(\mathbf{x})} \left(\frac{\Psi_G(\mathbf{x})}{\Phi_G(\mathbf{x})} \right)^{\omega_G} \rightarrow \int_{\mathbb{P}_{>0}^{n-1}} \frac{\prod_i p_i(\mathbf{x})^{\mu_i}}{\prod_j q_j(\mathbf{x})^{\nu_j}} \Omega$$

Sector decomposition approach

- Algorithms to perform blowups in the general case:
Binoth, Heinrich '03; Bogner, Weinzierl '07; (Hironaka 1964)
- Simple geometric interpretation:
Kaneko, Ueda '09

All algorithms are oblivious to the specific structures on the left!

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previous

Established solutions

$$\int_{\mathbb{P}_{>0}^{E-1}} \frac{\Omega}{\Psi_G^{D/2}(\mathbf{x})} \left(\frac{\Psi_G(\mathbf{x})}{\Phi_G(\mathbf{x})} \right)^{\omega_G} \rightarrow \int_{\mathbb{P}_{>0}^{n-1}} \frac{\prod_i p_i(\mathbf{x})^{\mu_i}}{\prod_j q_j(\mathbf{x})^{\nu_j}} \Omega$$

Numerical evaluation using sector decomposition for blowups:

- Runtime to evaluate the integral up to δ -accuracy

$$\approx \mathcal{O}(V^2 \cdot \delta^{-2})$$

where V is the number of monomials in $\frac{\prod_i p_i(\mathbf{x})^{\mu_i}}{\prod_j q_j(\mathbf{x})^{\nu_j}}$.

- For Feynman integrals V grows \approx exponentially with n .

Results MB 2020:

1. Numerical integration is an exercise in **tropical geometry**.
2. The general (oblivious) approach can be accelerated:

$$\mathcal{O}(V^2 \cdot \delta^{-2}) \quad \rightarrow \quad \mathcal{O}(V^2 + \delta^{-2})$$

⇒ achievable accuracy ‘decouples’ from integral complexity.

3. Euclidean Feynman integration can be accelerated extremely:

$$\mathcal{O}(V^2 \cdot \delta^{-2}) \approx \mathcal{O}(2^{cn} \cdot \delta^{-2}) \quad \rightarrow \quad \mathcal{O}(n2^n + \delta^{-2})$$

with $c \gg 1$ where n is the number of edges of the graph.

Theorem MB 2020

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- 3 loops is already a tough challenge for existing programs.
- New: ≥ 17 loops possible (with basic implementation).
- Caveat: Only Euclidean - no Minkowski regime (so far).

Interesting example

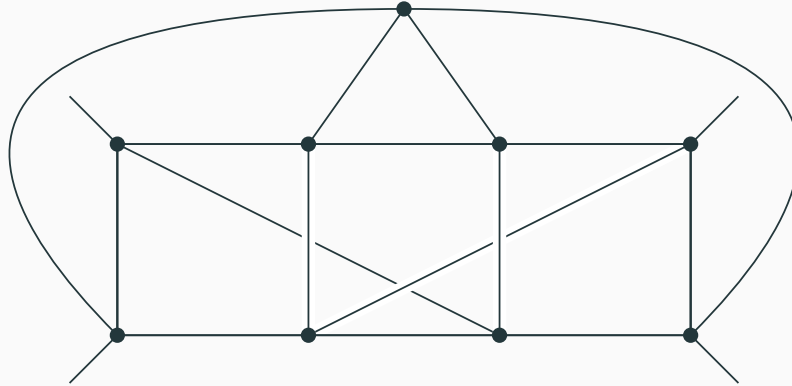


Figure 1: A non-generalized polylog/non-MZV 8-loop φ^4 -graph.

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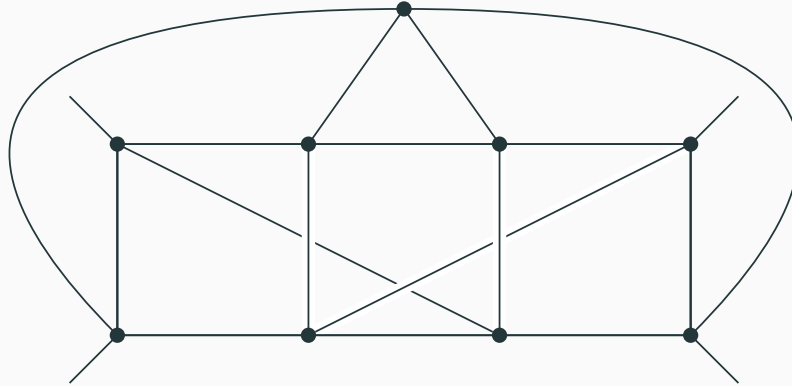


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$$\Gamma(\varepsilon) \int_{\mathbb{P}_{>0}^{E-1}} \frac{1}{\Psi_G(\mathbf{x})^{2-\varepsilon}} \left(\frac{\Psi_G}{\Phi_G} \right)^\varepsilon \Omega \approx \frac{1}{\varepsilon} 422.9610 \cdot (1 \pm 10^{-6}) + \dots$$

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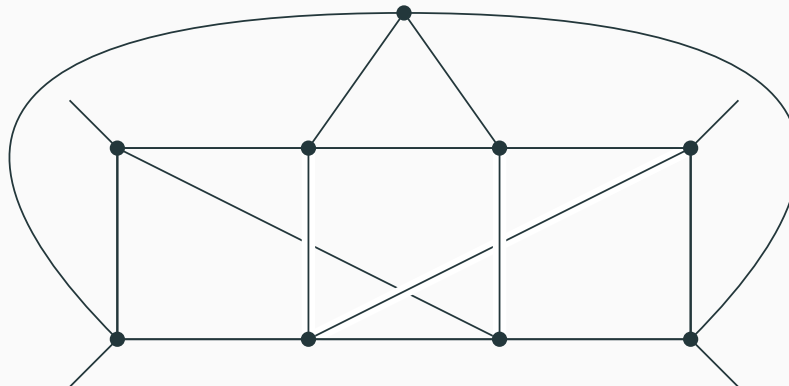


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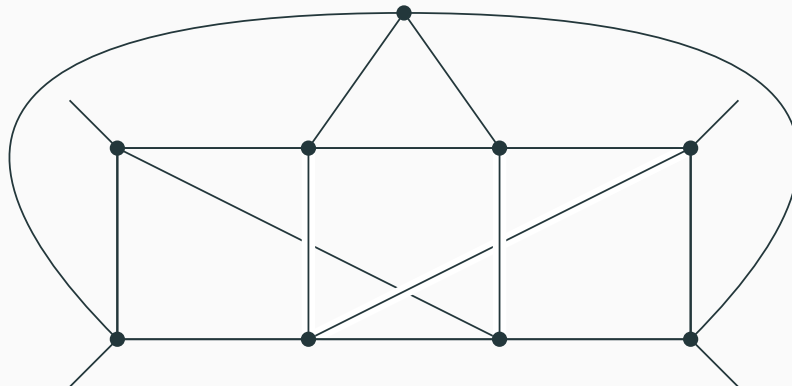


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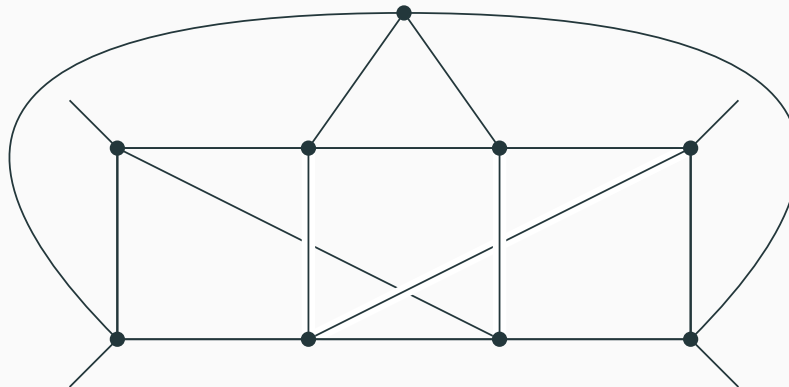


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- ~ 10 CPU secs to compute up to 10^{-3} -accuracy at 8 loops.
- ~ 30 CPU days to compute up to 10^{-6} -accuracy at 8 loops.
- Higher orders in ε can also be computed.

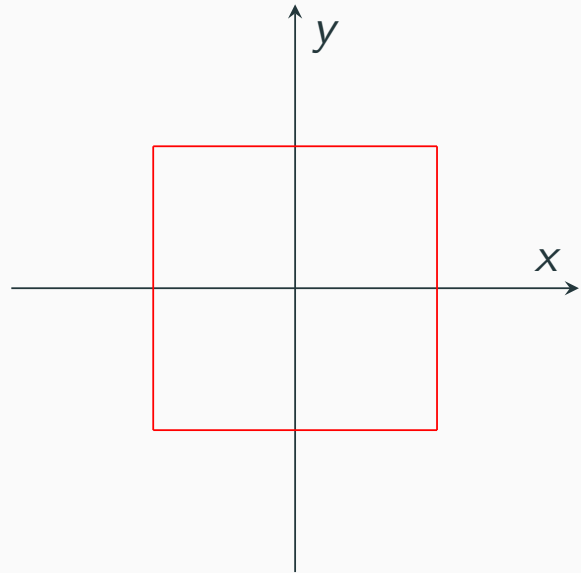
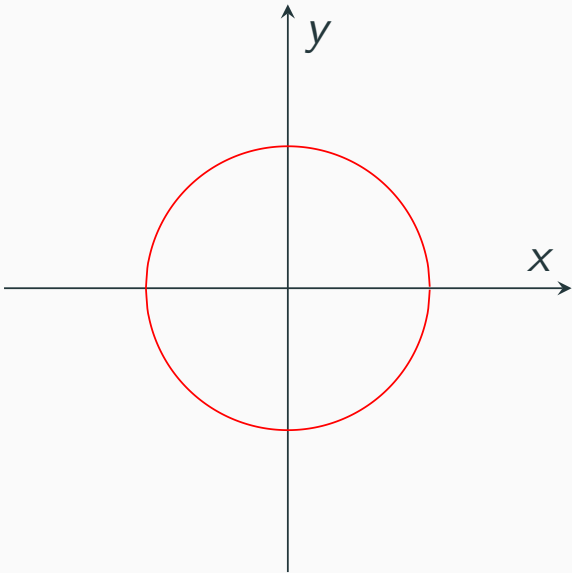
The Tropical Approach

Philosophy

Deform geometry to sacrifice smoothness for simplicity.

Various applications in algebraic geometry.

$$1 = x^2 + y^2 \quad \rightarrow \quad 1 = (x^2 + y^2)^{\text{tr}} = \max\{x^2, y^2\}$$



Application to Feynman graph polynomials

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$$\begin{aligned}\Psi_G &= \sum_T \prod_{e \notin T} x_e & \Rightarrow & \Psi_G^{\text{tr}} = \max_T \prod_{e \notin T} x_e \\ \Phi_G &= \sum_F \|p(F)\|^2 \prod_{e \notin F} x_e & \Rightarrow & \Phi_G^{\text{tr}} = \max_{\substack{F \\ \text{s.t. } \|p(F)\|^2 \neq 0}} \prod_{e \notin F} x_e\end{aligned}$$

Feynman integral:

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QFT tropicalization

Replace all instances of Ψ and Φ with their tropicalized versions.

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Panzer 2019; MB 2020

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 Hepp - Bound

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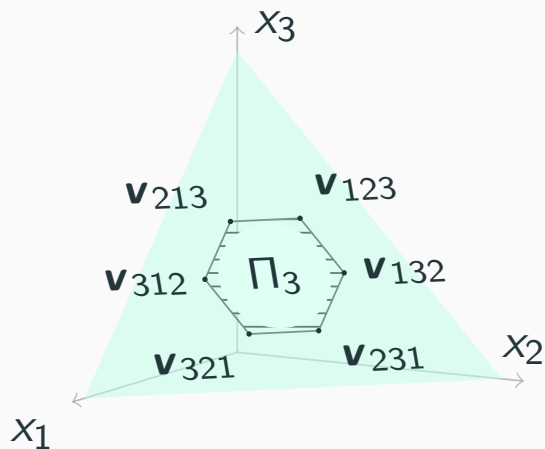
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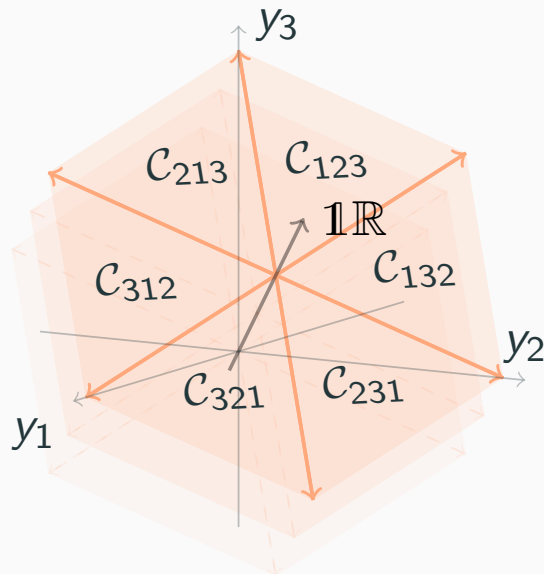
MB 2020

⇒ Better understanding of tropical geometry leads to faster numerical integration.

Relevant polytopes: Generalized permutahedra



(a) The permutahedron $\Pi_3 \subset \mathbb{R}^3$.



(b) Dual of Π_3 : The corresponding braid arrangement fan.

Gen. permutahedra are well understood thanks to **Postnikov 2008** and **Aguiar, Ardila 2017**.

Current limitations

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Problem: Non-Euclidean kinematic regions are not as fast, because

- The generalized permutahedron structure breaks down at singular momentum configurations (IR singularities).
- Φ_G can vanish in the integration domain (\Rightarrow analytic continuation is necessary).

Conclusions

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- Loop order ≈ 15 or ≈ 30 edges are easily possible.
- Applications:
 - Calculations in the Euclidean regime.
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 - (Massive) form factor calculations.
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CODE
AVAILABLE
AT

<https://github.com/michibo/tropical-feynman-quadrature>

