



Searching for Dark Matter Signals COHERENT Way



TEXAS A&M UNIVERSITY

Physics & Astronomy

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Magnificent CEvNS 2019

The PIT, Chapel Hill, NC, November 9st, 2019

In collaboration with B. Dutta, S. Liao, J.-C. Park, S. Shin and L. Strigari [arXiv:1906.10745, PRL submitted]

B. Dutta, S. Liao, J.-C. Park, S. Shin and L. Strigari, in progress

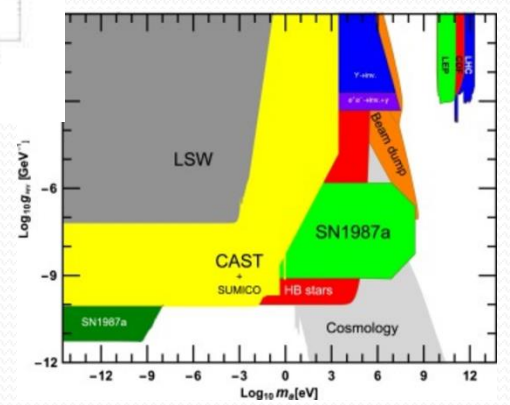
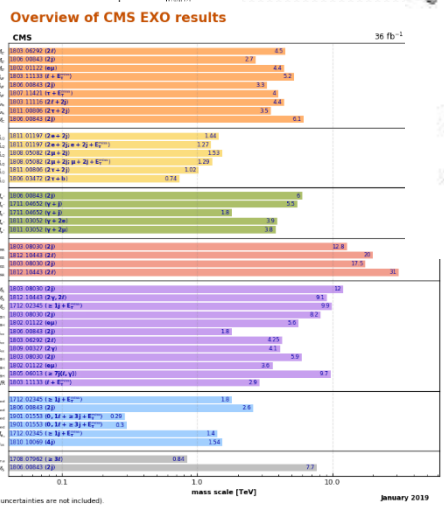
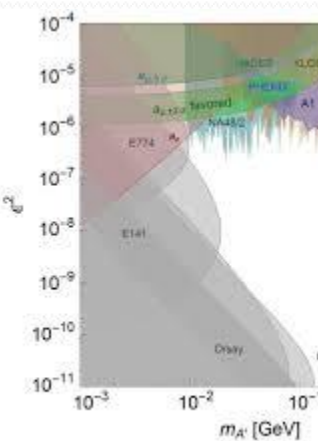
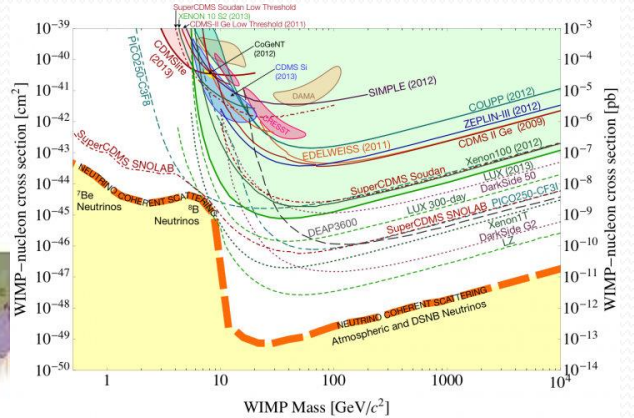
Hunt for New Physics: Current Status

ATLAS SUSY Searches - 95% CL Lower Limits

Model	Signature	$\sigma(\text{fb})^{-1}$	Mass limit	Reference
Inclusive Searches	\tilde{g}, \tilde{u}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{u}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\tau}_1$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\nu}_\tau$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
3-jet + missing	\tilde{g}, \tilde{u}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{u}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\tau}_1$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\nu}_\tau$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
2-jet + missing	\tilde{g}, \tilde{u}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{u}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\tau}_1$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\nu}_\tau$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
Monojet + missing	\tilde{g}, \tilde{u}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{u}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\tau}_1$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\nu}_\tau$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025

ATLAS Preliminary

Model	Signature	$\sigma(\text{fb})^{-1}$	Mass limit	Reference
Inclusive Searches	\tilde{g}, \tilde{u}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
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	\tilde{g}, \tilde{d}_L	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{d}_R	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{t}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_1	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	\tilde{g}, \tilde{b}_2	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\tau}_1$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025
	$\tilde{g}, \tilde{\nu}_\tau$	3.8×10^{-2}	1.8 TeV	ATLAS-CONF-2015-025

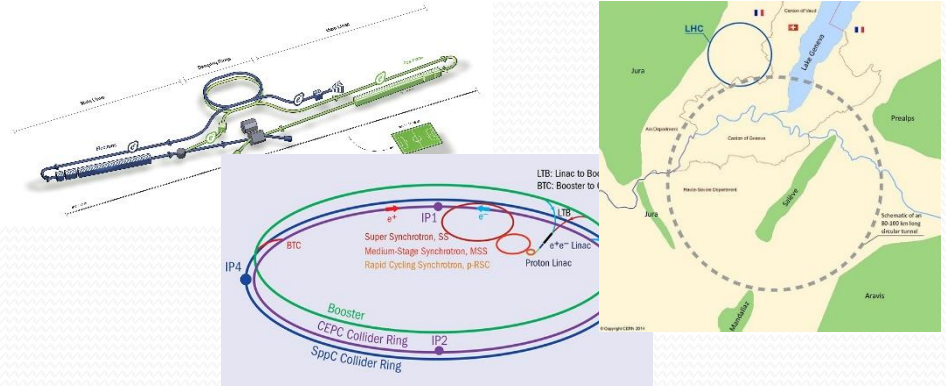


*Only a selection of the available mass lines on new states or phenomena is shown. Many of the limits are based on simplified models \tilde{g}, \tilde{u}_L for the astrophysical dark matter.

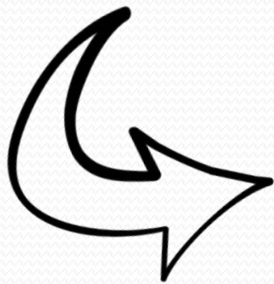
Hunt for New Physics: Future Directions



New physics searches at the LHC

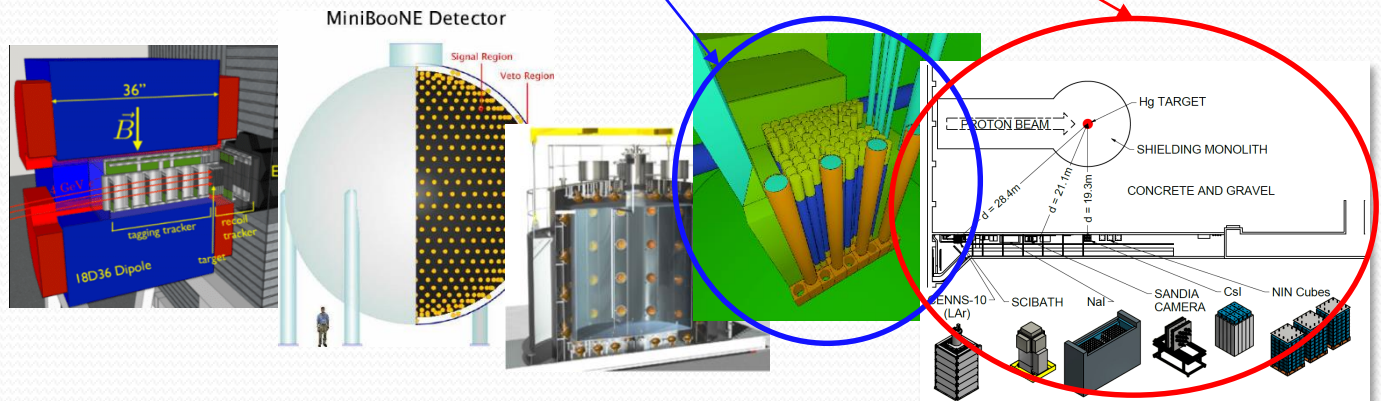


Higher-energy colliders, e.g., ILC, CEPC, CERN-FCC, etc



See S. Liao's talk for reactor exps.

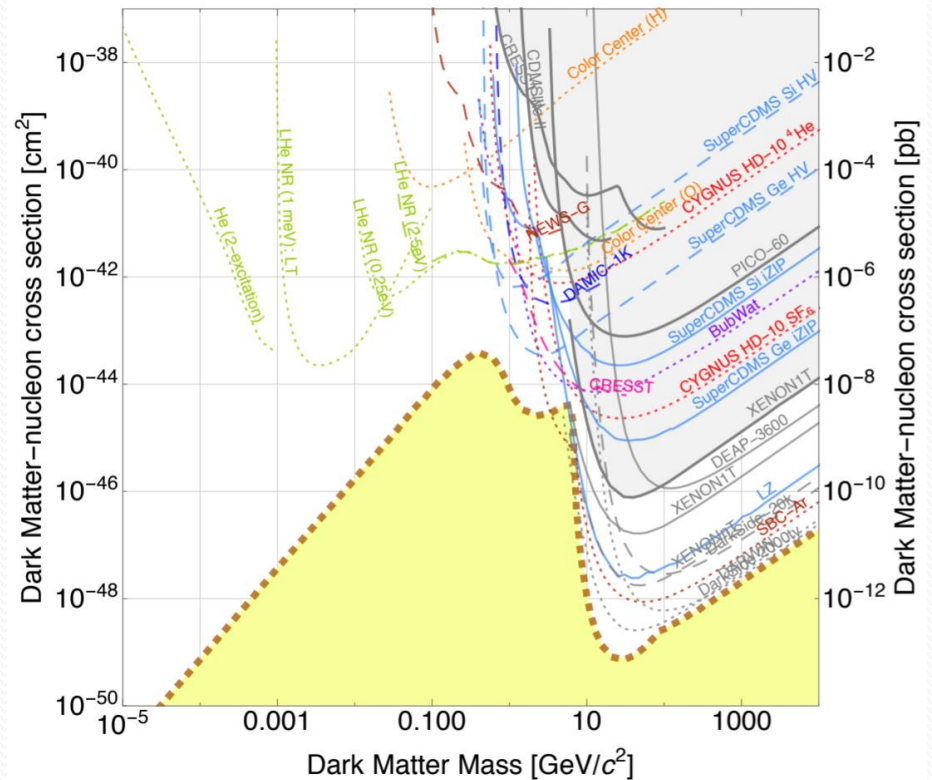
Today's focus



Various physics potentials in neutrino facilities, low-energy high-intensity experiments

Current Status of Dark Matter Searches

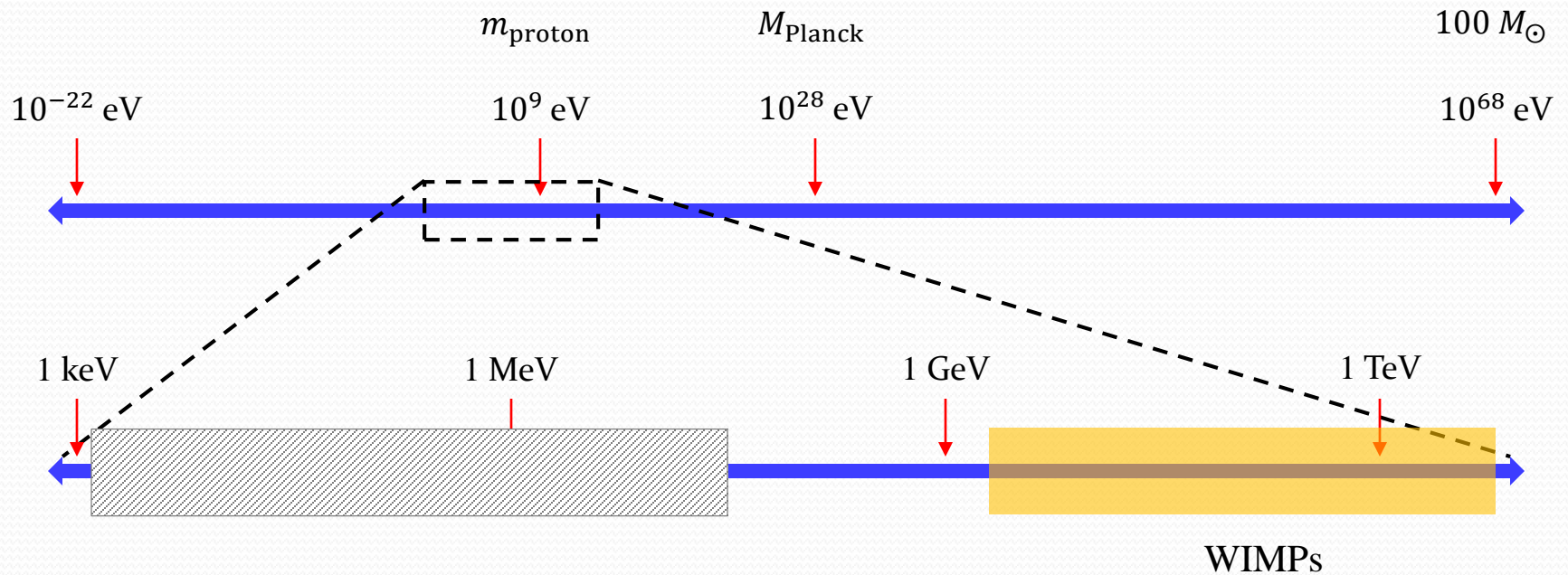
- **No observation** of DM signatures via non-gravitational interactions while many searches/interpretations designed/performed under nonrelativistic WIMP/WIMP-like scenarios
 - ⇒ merely excluding more parameter space in dark matter models



[US Cosmic Visions, Battaglieri et al (2017)]

*Time to pause, rethink and redesign our approach/search strategies, e.g., **COHERENT Way!***

The Dark Matter Landscape



- ✓ Probing dark sectors: **(Light) dark matter + new mediators**
- ✓ Less constrained by current searches

Light Dark-Sector Particle Models/Searches: Mediator

- ❑ Various light mediator scenarios have been proposed.
 - ✓ Dark matter scenarios based on hidden sectors: e.g., models of asymmetric dark matter, Sommerfeld enhancements motivated by SIMP, etc (see the review [Essig et al (2013)])
 - ✓ $g - 2$ of electron: 2.4σ discrepancy [Davoudiasl, Marciano (2018)]
 - ✓ Neutrino sector physics: new neutrino interactions to satisfy the MiniBooNE excess [Bertuzzo, Jana, Machado, Funchal (2018)]
 - ✓ Solutions of Yukawa coupling hierarchy problem [Dutta, Ghosh, Kumar (2019)]
 - ✓ See also US cosmic vision [Battaglieri et al (2017)]

- ❑ Light mediator searches at existing/future experiments, e.g., NA64, Belle I/II, Babar, SHiP, FASER, MATHSULA, SeaQuest

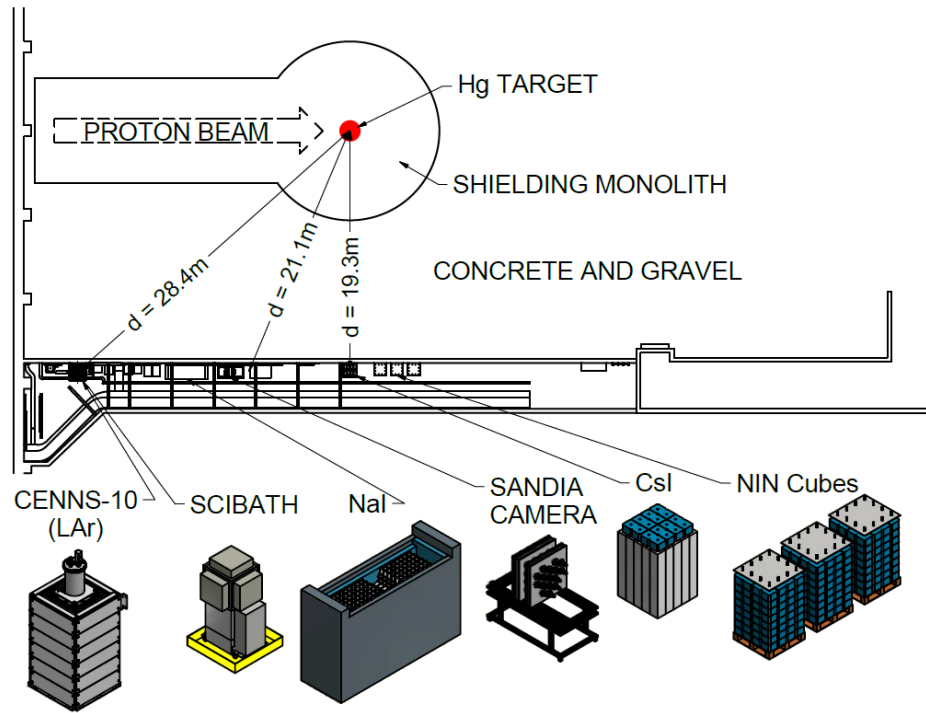
Light Dark-Sector Particle Models/Searches: Dark Matter

- ❑ Various light dark matter-involving phenomena has been studied.
 - ✓ Boosted dark matter scenarios [Agashe, Cui, Necib, Thaler (2014); Berger, Cui, Zhao (2014); Kong, Mohlabeng, Park (2014); DK, Park, Shin (2016)]
 - ✓ Fast-moving DM via induced nucleon decays [Huang, Zhao (2013)]
 - ✓ MeV-range DM indirect detection at gamma-ray telescopes [Boddy, Kumar (2015)]
 - ✓ Energetic cosmic-ray-induced (semi-)relativistic dark matter scenarios [Yin (2018); Bringmann, Pospelov (2018); Ema, Sala, Sato (2018); Dent, Dutta, Newstead, Shoemaker (2019)]
 - ✓ Ultra high energy cosmic ray phenomena [Bhattacharya, Gandhi, Gupta (2014); Heutier, DK, Park, Shin (2019)]
- ❑ Cosmogenic light dark matter searches at existing/future experiments, e.g., SK/HK, COSINE-100, ProtoDUNE, DUNE
- ❑ Beam-produced light dark matter searches at existing/future experiments, e.g., BDX, MicroBooNE, SeaQuest, LDMX, T2HKK, DUNE, SHiP, and proposals [Bjorken, Essig, Schuster, Toro (2009); Batell, Pospelov, Ritz (2009); deNiverville, Pospelov, Ritz (2011); Izaguirre, Krnjaic, Schuster, Toro (2014); Berlin, Gori, Schuster, Toro (2018), and many more]

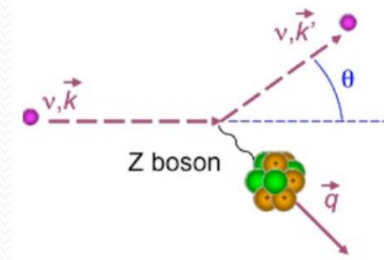
Goals

- ❑ How to **isolate (light) dark matter signal** events from the SM (neutrino) backgrounds with **timing spectra** at neutrino experiments, taking **COHERENT** as a benchmark experiment
- ❑ Application to the measurement data that COHERENT has released
- ❑ How to **interpret the result**

COHERENT Experiment: Primer

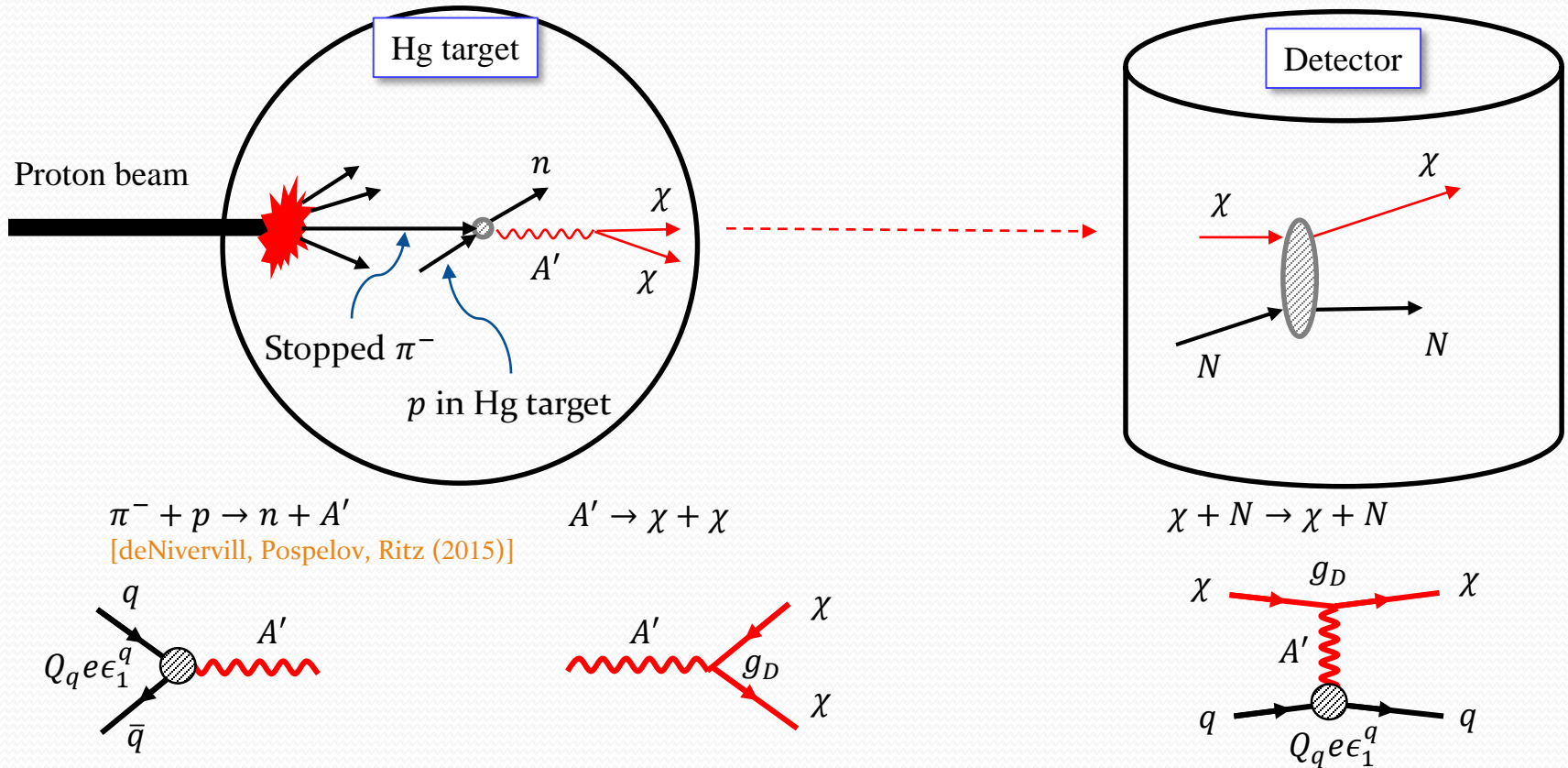


- ✓ Main mission (see M. Green's talk): first direct measurement of Coherent Elastic Neutrino-Nucleus Scattering (CEvNS).



- Prompt ν 's: $\pi \rightarrow \mu + \nu_\mu$
 - Delayed ν 's: $\mu \rightarrow e + \nu_\mu + \nu_e$
- ✓ ~ 1 GeV proton beam on Mercury target (pulse duration: 380 ns FWHM 60 Hz)
 - ✓ 5×10^{20} protons-on-target (POT) delivered per day

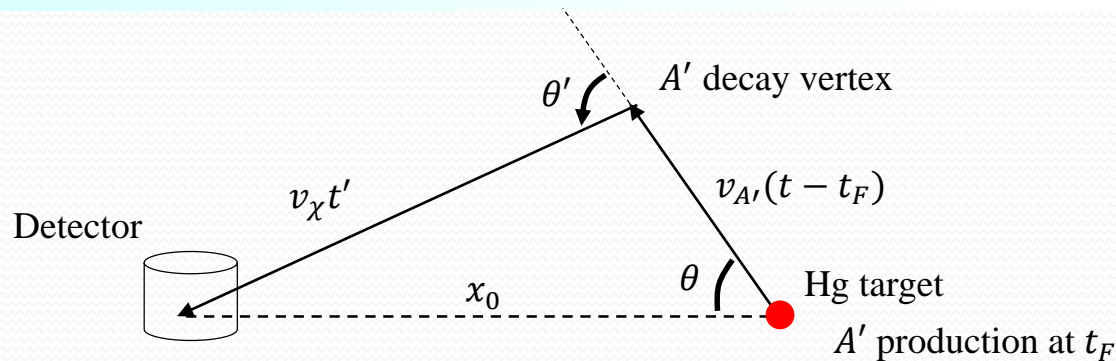
Dark Matter Scenarios in COHERENT



Cf.) Another (subdominant) process: charge exchange, $\pi^{-/+} + p/n \rightarrow \pi^0 + n/p$, $\pi^0 \rightarrow \gamma + A'$ [JSNS² TDR]

v-induced background vs. dark photon decay into dark matter utilizing timing measurements!

Timing Spectrum of Dark-Matter Events



Dark matter flux at the detector: $f(T) = dN_{\chi}/dT$

$$f(T) = \int_{-1}^1 d \cos \theta \int_0^{t_F^{\max}} dt_F \left| \frac{dT}{dt} \right|^{-1} \frac{d^2 N_{A'}}{dt d \cos \theta} \cdot w(\cos \theta') \cdot \mathcal{F}(t_F)$$

$$T = t + v_{\chi}^{-1} \sqrt{x_0^2 + v_{A'}^2 (t - t_F)^2 - 2x_0 v_{A'} (t - t_F) \cos \theta}$$

Model of π^- production timing (\propto POT)

$$\frac{d^2 N_{A'}}{dt d \cos \theta} = \frac{1}{2} \cdot \frac{1}{\tau_{A'}} e^{-\frac{t-t_F}{\tau_{A'}}} \Theta(t - t_F) \leftarrow \text{from the decay law}$$

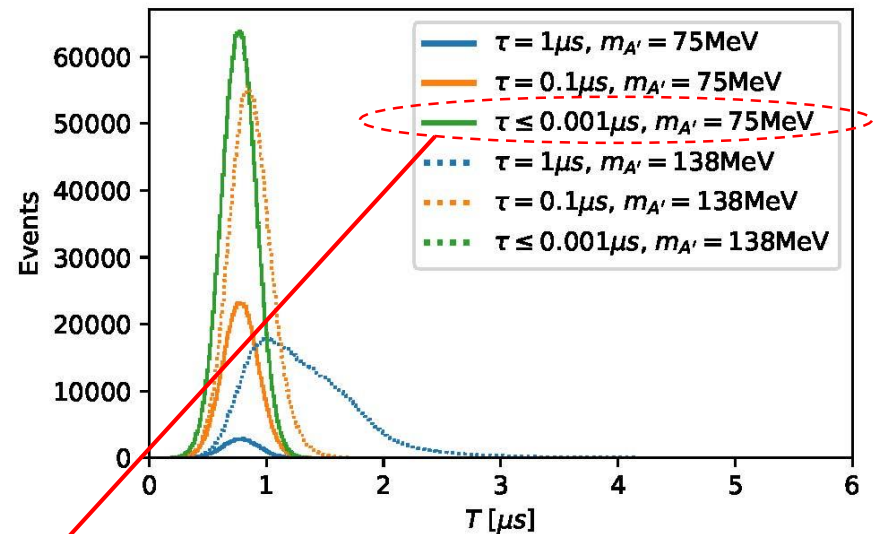
$$w(\cos \theta') = \frac{1}{2\pi (v_{\chi} t')^2} \left| \frac{d \cos \theta'}{d \cos \theta^*} \right|^{-1} \frac{dN_{A' \rightarrow \chi}}{d \cos \theta^*} \leftarrow \text{Probability that DM travels towards the detector}$$

Cf.) Search strategy with timing information at the LHC [Liu, Liu, Wang (2018)]

Parameter Space: Dark Photon

□ Various possibilities for dark photon A' (depending on ϵ_1^q and $m_{A'}$)

- Short-lived (large ϵ_1^q)
vs. Long-lived
- Relativistic
vs. Non-relativistic ($m_{A'} \sim 138$ MeV)



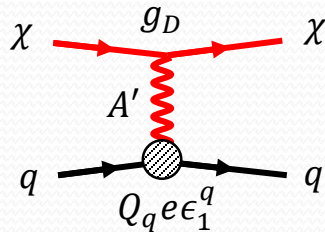
For $\tau < a$ few ns, we get maximum number of events

Dark matter events populate in prompt timing bins
(\Rightarrow Data in delayed timing bins as a control sample)

Parameter Space: Dark Matter

- Dark matter scatters off nucleus:

$$\frac{d\sigma}{dE_r} = \frac{e^2 \epsilon_q^2 g_D^2 Z^2 \cdot |F(2m_N E_r)|^2}{4\pi p_\chi^2 (2m_N E_r + M'^2)^2} \left\{ 2E_\chi^2 m_N \left(1 - \frac{E_r}{E_\chi} - \frac{m_N E_r}{2E_\chi^2} \right) + E_r^2 m_N \right\}$$



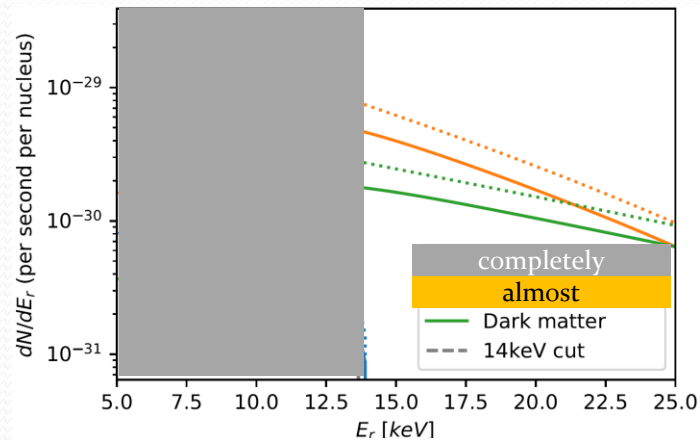
- ✓ In general, the scattering process could be mediated by a different particle (e.g., Baryon number gauged dark gauge boson [deNiverville, Pospelov, Ritz (2015)])
 $: A' \rightarrow V', m_{A'} \rightarrow M', Q_q e \epsilon_1^q \rightarrow Q_B e \epsilon_2^q, g_D = e \epsilon_1^D \rightarrow e \epsilon_2^D$

- Dark photon A' production to dark matter scattering can be described by **two variables**.

$$\epsilon \equiv \epsilon_1^q \epsilon_2^q \epsilon_2^D \sqrt{\text{BR}_{A' \rightarrow \chi\chi}} \text{ and } M'$$

Proposed Search Strategy

❑ A combination of energy and timing cuts (i.e., cut-and-count exp., less assumptions on the ν sector)

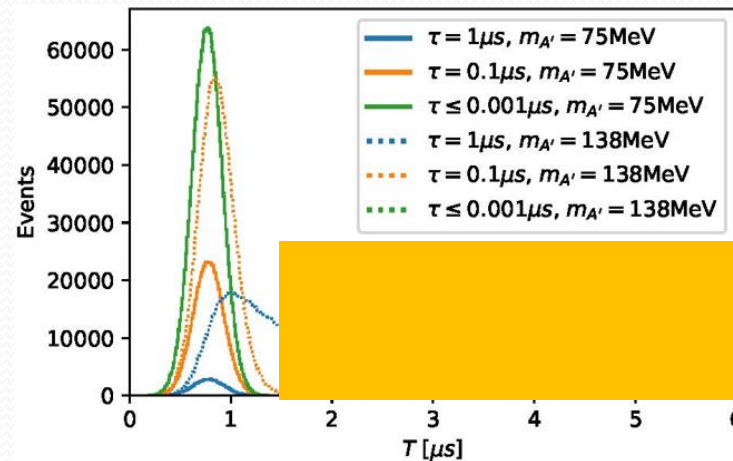
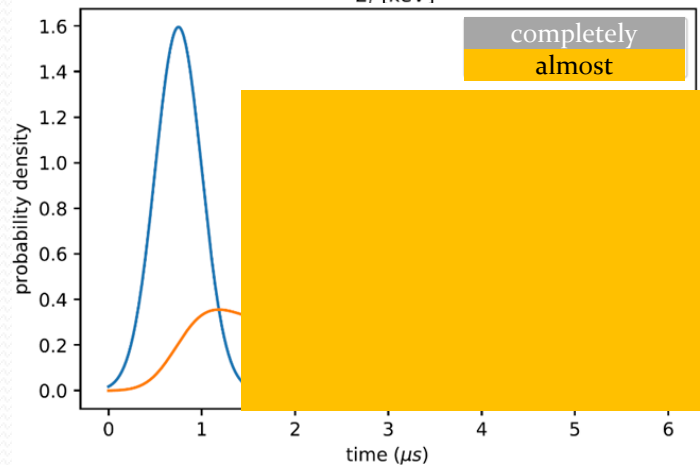


① $E_r > 14 \text{ keV}$

- Prompt neutrino: completely removed
- Delayed neutrino (and DM signal): still remains

② $T < 1.5 \mu\text{s}$

- Delayed neutrino: almost removed
- DM signal: still remains



Application to Existing CsI Data

❑ Data released by COHERENT: CsI 14.5 kg × 308 days = 4,466 kg·day [Akimov et al, 1804.09459]

❑ Analysis scheme (also following [Dutta, Liao, Sinha, Strigari (2019)] for background estimate)

- Fix the size of neutron distribution to $R_n = 4.7$ fm
- $14 \text{ keV} < E_r < 28 \text{ keV}$, $T < 1.5 \mu\text{s}$

$$F_N^{\text{Helm}}(q^2) = \frac{3j_1(qR_0)}{qR_0} \exp\left(-\frac{q^2 s^2}{2}\right)$$
$$R_n^2 = R_0^2 + 5s^2$$

97 : total events

- 49 : classified as the steady-state (SS) background
- 19 : identified as delayed neutrino events (SM)
- 0 : identified as prompt neutrino events (SM)
- 3 : beam-related neutron (BRN) backgrounds

26 : “Excess”

Significance ($R_n = 4.7$ fm): **2.4 σ**

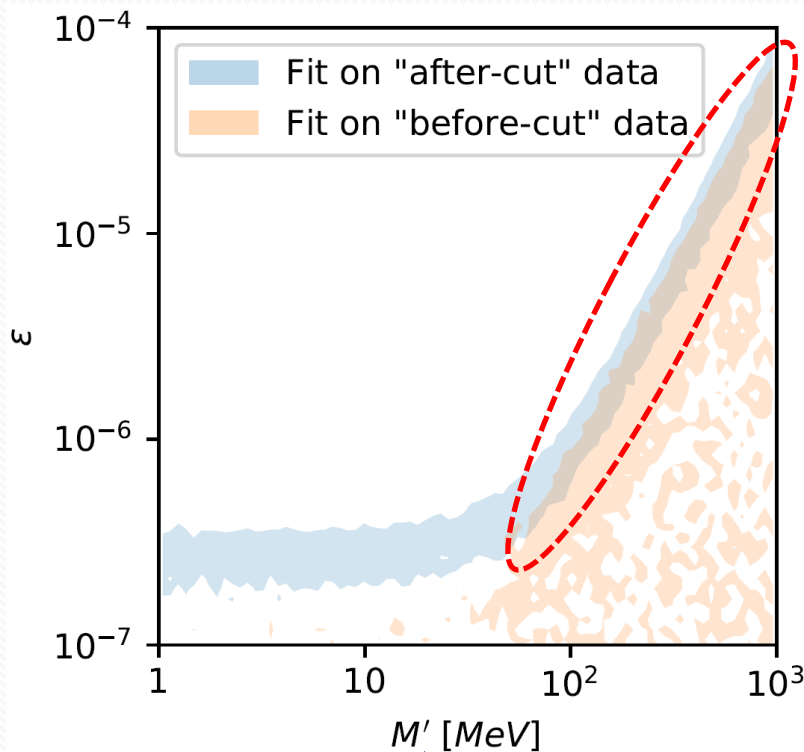
Significance ($R_n = 5.5$ fm): **3.0 σ**

$$\text{Significance} = \frac{\text{Excess}}{\sqrt{2\text{SS} + \text{BRN} + \text{SM}}} \text{ [COHERENT, 1708.01294]}$$

❑ Caveats: systematics on the SS background not considered, excess explained by other unidentified background (see also D. Pershey’s talk for the latest DM search)

Mild Excess? – Dark Matter Interpretation

- Fit to the excess after the cuts needs to fit the full data set (before the cuts).



- Baseline model point for the figure in the left:
 $\tau = 1 \text{ ns}, m_{A'} = 75 \text{ MeV}, m_\chi = 5 \text{ MeV}$

- Nevertheless, the figure holds for
 - ✓ $\tau \leq 4 \text{ ns}, m_{A'} < 138 \text{ MeV}$
 - ✓ $\tau \leq 30 \text{ ns}, m_{A'} \cong 138 \text{ MeV}$ (non-relativistic dark photon case)
 - ✓ Any $m_\chi < m_{A'}/2$

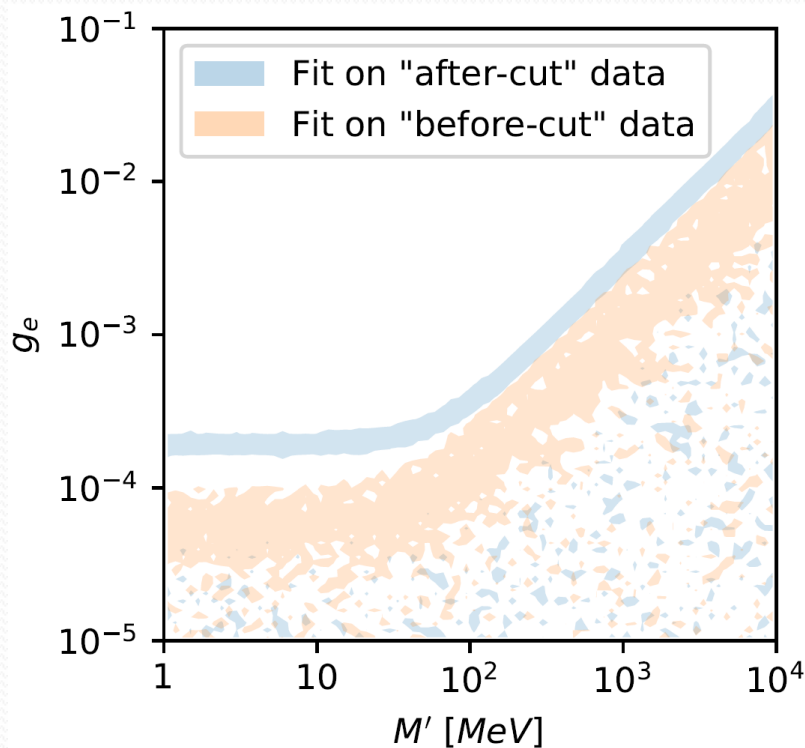
$$\epsilon = \epsilon_1^q \epsilon_2^q \epsilon_2^D \sqrt{\text{BR}_{A' \rightarrow \chi\chi}}$$

M' [MeV]

The mass of the DM-nucleus interaction mediator

Mild Excess? – Alternative Interpretation: NSI

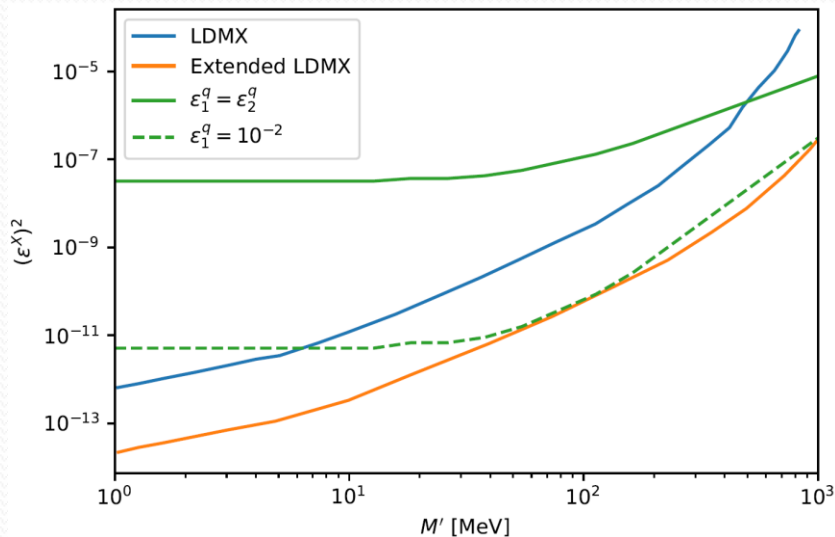
□ Example alternative new physics possibility, Non-Standard Interaction



- Benchmark case: non-zero coupling g_e , the NSI in the ν_e neutral-current interaction (along with a new mediator).
 - ⇒ No overlapping regions, especially the prompt timing bin (i.e., $T < 1.5 \mu\text{s}$) doesn't show a good fit. NSI affects the overall normalization of neutrino flux!
- The situation becomes even worse with $g_\mu \neq 0$, since it affects not only the delayed but the prompt spectrum.

No Excess? – Constraining Parameter Space

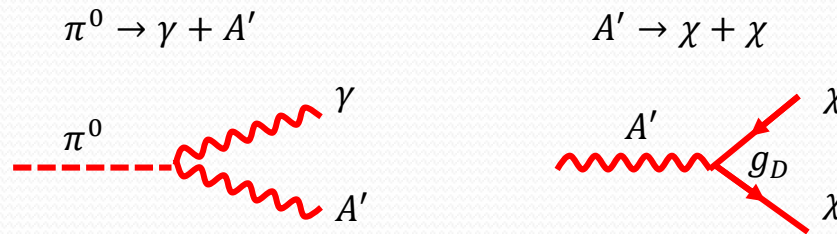
□ Assuming no excess is observed, we can constrain parameter space.



- $\alpha_D \equiv \frac{(e\epsilon_2^q)^2}{4\pi} = 0.5$
- M' : the mass of the DM-nucleus interaction mediator
- Solid green line: single mediator scenario, i.e., $\epsilon^X = \epsilon_1^q = \epsilon_2^q$
- Dashed green line: multi-mediator scenario. One of them is fixed to 10^{-2} (e.g., Baryon number gauged dark gauge boson [deNiverville, Pospelov, Ritz (2015)])
- For LDMX, ϵ^e in [LDMX, 1808.05219] identified as ϵ^X .
- Sensitivity reach is already better than DUNE, compared to the result in [De Romeri, Kelly, Machado (2019)].

Upcoming Work

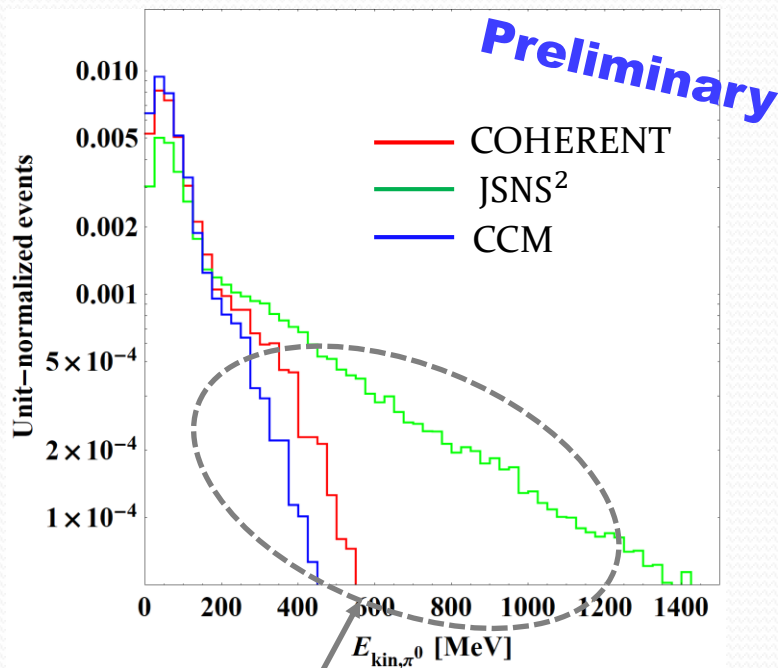
- We include the **contribution from neutral pions** and consider **two more similar experiments**, JSNS² (in J-PARC) and CCM (in Los Alamos).



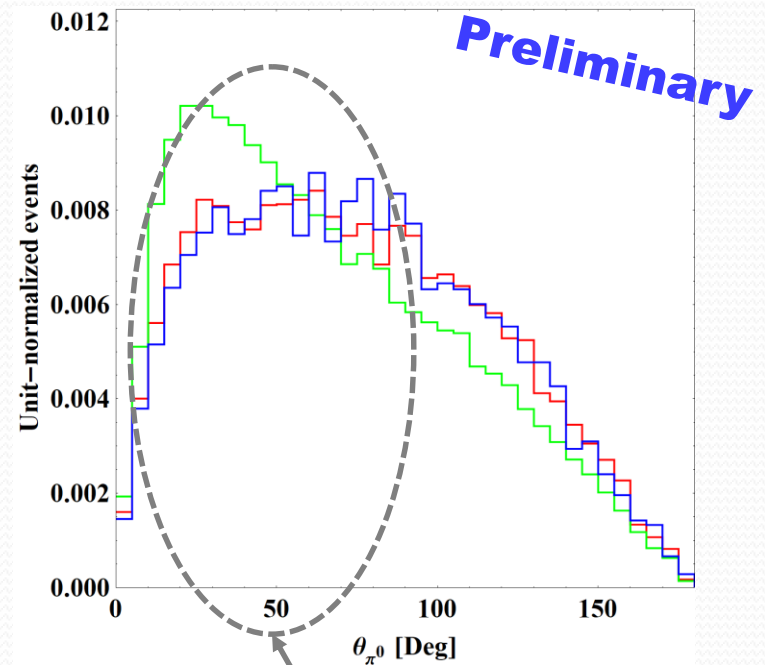
Experiment	E_{beam} [GeV]	POT [yr^{-1}]	Target	Detector (mass/distance/angle)
JSNS ² [11]	3	3.8×10^{22}	Hg	Gd-LS (17 ton/24 m/70°)
CCM [12]	0.8	1.8×10^{22}	W	LAr (7 ton/20 m)

Upcoming Work

- π^0 are not produced at rest. \Rightarrow Detector-level (e.g. GEANT) simulation is needed.



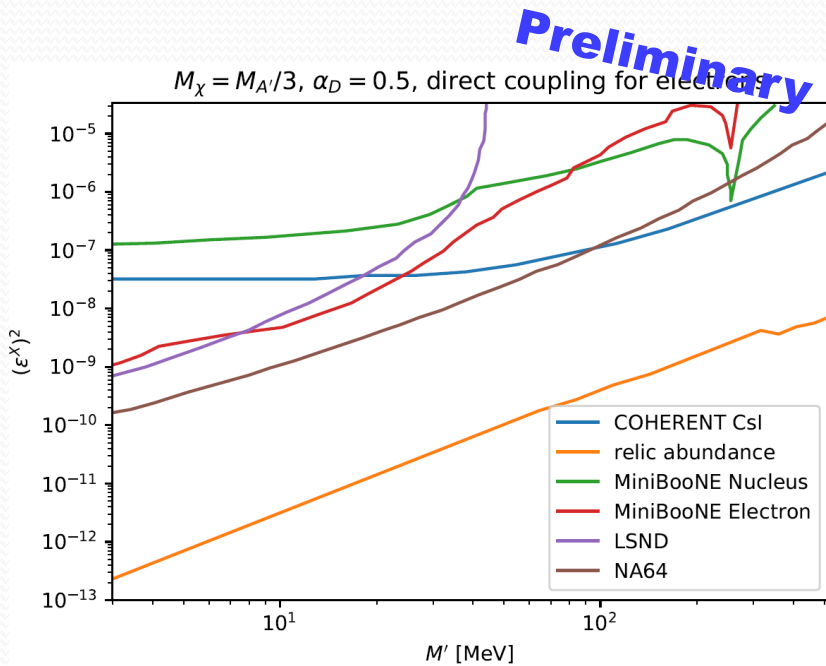
Dark matter can be more energetic.



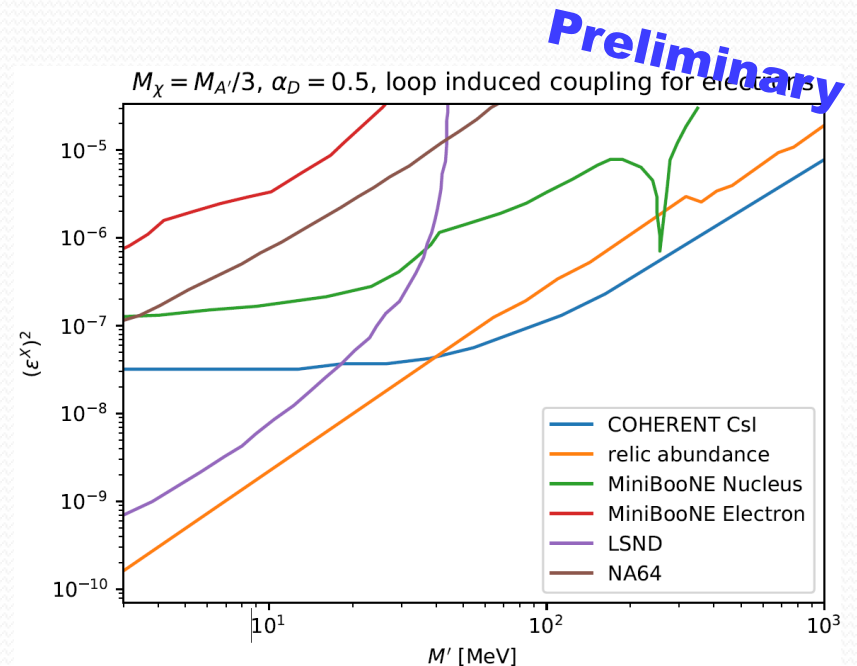
Dark matter flux is slightly forward-directed.

Improved Results

- Sensitivity reaches are improved.



[Single mediator]



[loop-induced coupling to e]

Conclusions

- ❑ Null signal observation at conventional dark matter detection experiments motivates us to look into mass scales other than that of WIMP.
- ❑ Models with **light mediators and dark matter** are interesting and receiving rising attention.
- ❑ Not only energy spectrum but also timing spectrum can be utilized in the search for light dark matter-induced signals: A **combination of timing and energy cuts can eliminate SM neutrino backgrounds** efficiently.
- ❑ The current Csl data shows a **$2.4 - 3\sigma$ excess** which can be explained by dark matter arising from dark photon decay.
- ❑ The experimental sensitivities for dark matter parameter space can be **better than the existing bounds**.
- ❑ More interesting results coming up. Stay tuned!

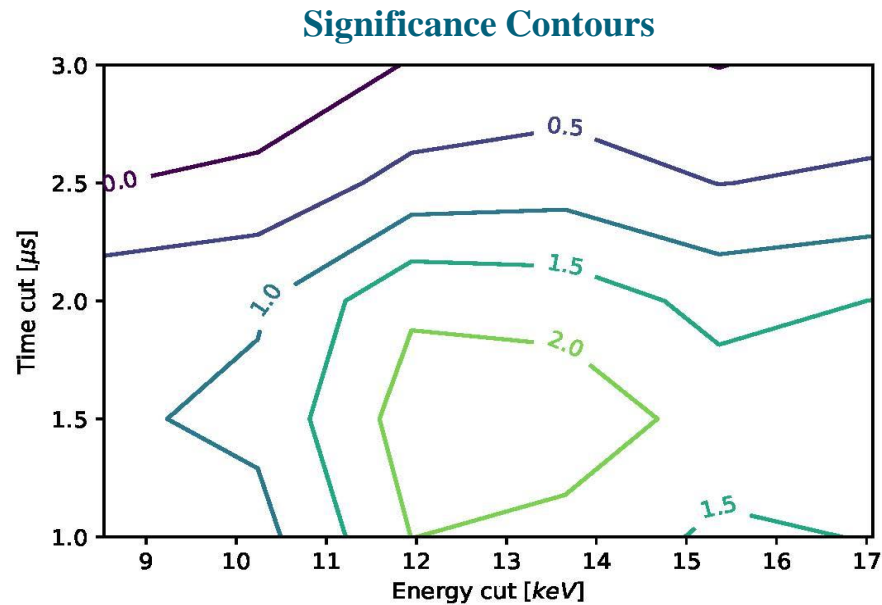


Thank you for your attention!



Bonus Slides

Cut Optimization



- $14 \text{ keV} < E_r < 28 \text{ keV}$
- $T < 1.5 \mu\text{s}$

Optimized set of cuts are independent of R_n

Preferred Parameter Values

- Single mediator scenario (i.e., the case where the dark photon A' mediates the scattering process of dark matter at the detector)

$$\epsilon = \epsilon_1^q \epsilon_2^q \epsilon_2^D \sqrt{\text{BR}_{A' \rightarrow \chi\chi}} \rightarrow (\epsilon^q)^2 \epsilon_2^D \sqrt{\text{BR}_{A' \rightarrow \chi\chi}} \rightarrow (\epsilon^q)^2 / e$$

$\epsilon^q \equiv \epsilon_1^q = \epsilon_2^q$ $\epsilon_2^D = 1/e, \text{BR} = 1$ for simplicity ($g_D \equiv e\epsilon_2^D$)

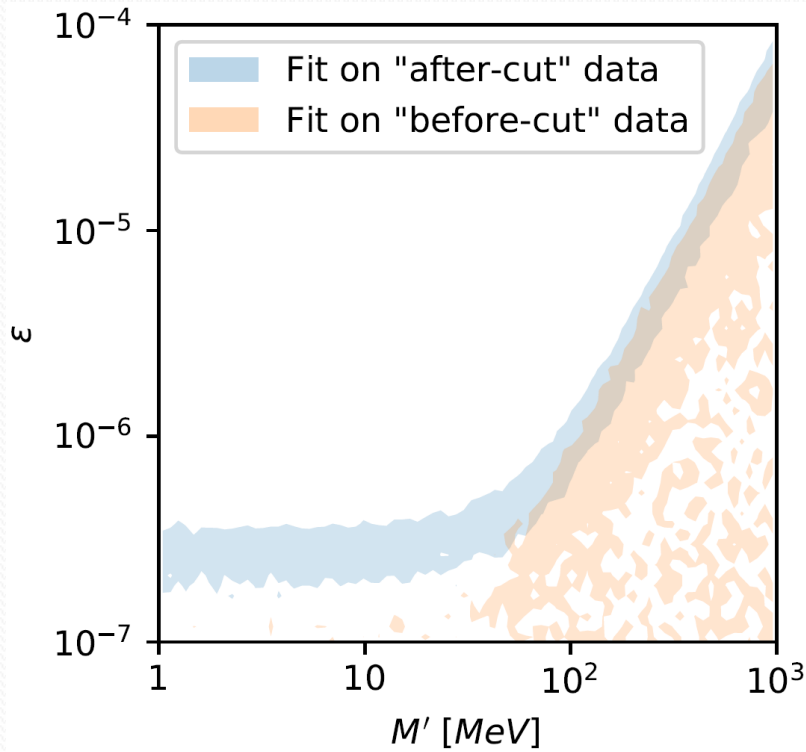
$M'(\text{MeV})$	50	75	100	1000
ϵ^q	3.5×10^{-4}	4.4×10^{-4}	5.5×10^{-4}	4.6×10^{-3}

- Multi-mediator scenario (i.e., the case where another mediator mediates the scattering process of dark matter at the detector)

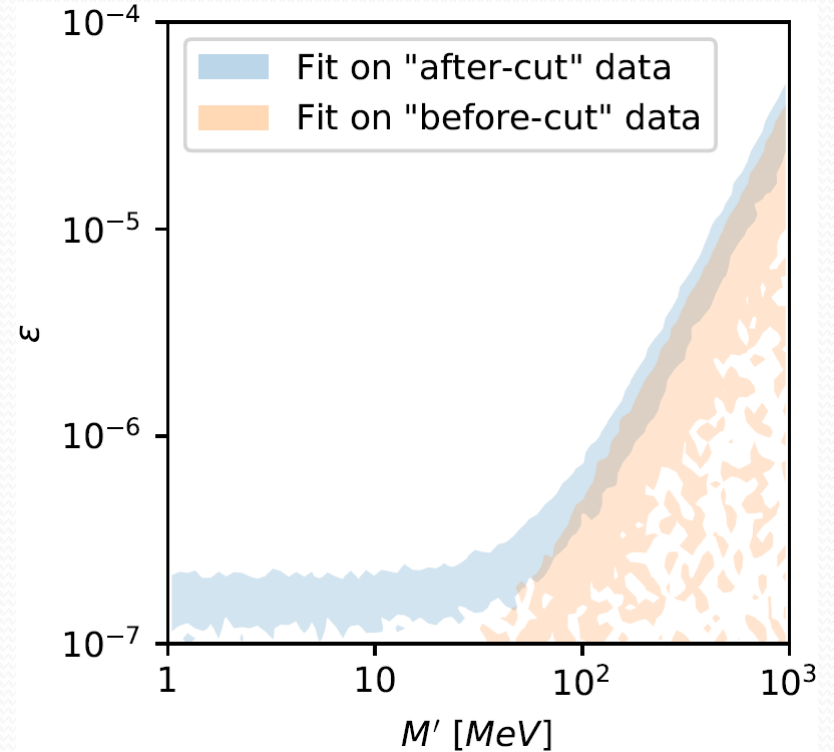
$\Rightarrow \epsilon^q$ above can be interpreted as $(\epsilon_1^q \epsilon_2^q \epsilon_2^D e)^{1/2}$

$\Rightarrow \epsilon_2^D$ is not the same as ϵ_1^D which controls τ of A' , so longer-lived dark photon could be possible.

Excess – Dark Matter Interpretation



[Without π^0 contribution]



[With π^0 contribution]