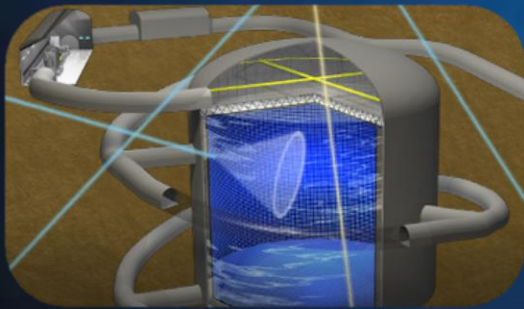
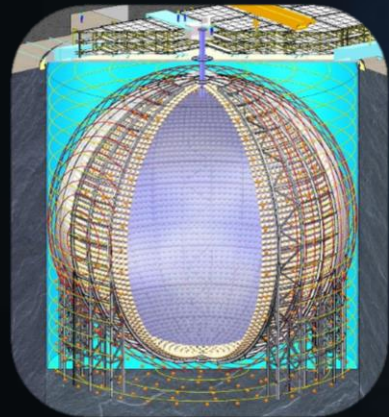


Exciting Prospects @ Neutrino Detectors

B. Dutta, W.C. Huang, D. Kim, J. Newstead, **JCP**, I. Shaukat Ali
[2401.09529]

Jong-Chul Park

CNU 충남대학교
CHUNGNAM NATIONAL UNIVERSITY



The XVIth Quark Confinement & Hadron Spectrum Conference
August 20 (2024)

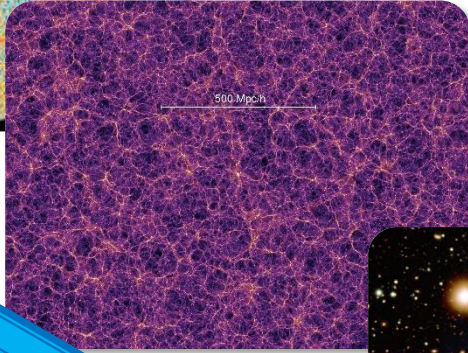
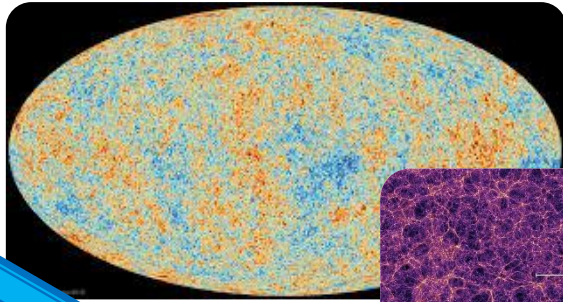
Outline

- ❖ **Dark Matter: WIMP?**
- ❖ **Nuclear Scattering**
- ❖ **Boosted Dark Matter (BDM)**
- ❖ **Exciting Prospects for Dark Matter**
- ❖ **Summary**

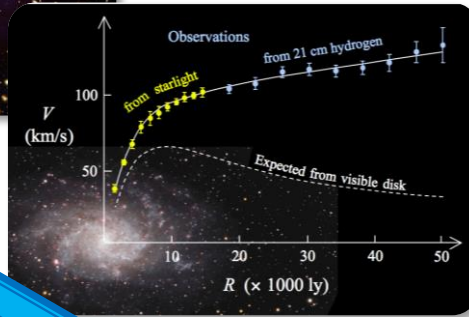
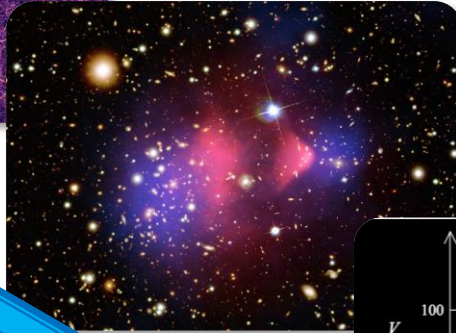


Dark Matter: WIMP?

Message from Cosmology: Dark Matter (DM)

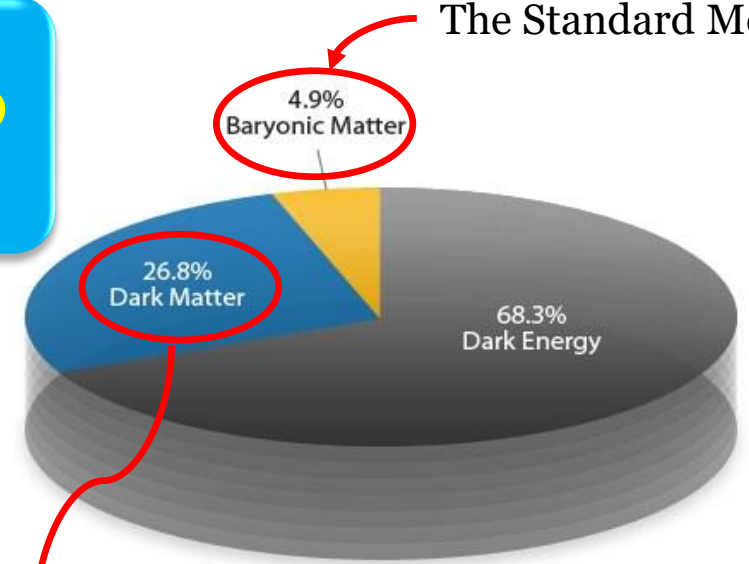


Dark Matter?



❖ **Modern cosmology:**

The Standard Model



❖ **Compelling paradigm:**

- ✓ Massive,
- ✓ Non-relativistic ($v \ll c$),
- ✓ Non-luminous (no/tiny EM interaction),
- ✓ Stable particles

Larger scale
Earlier

Many more other observations!

Smaller scale
Later

Classic Solution*: WIMP

Cosmological Lower Bound on Heavy-Neutrino Masses

Benjamin W. Lee^(a)

Fermi National Accelerator Laboratory, ^(b) Batavia, Illinois 60510

and

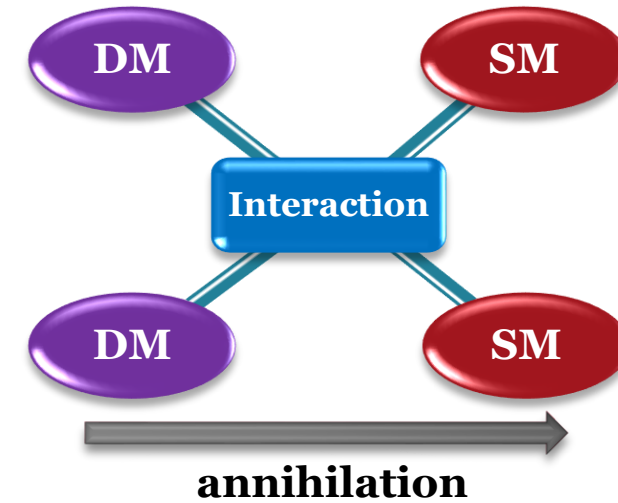
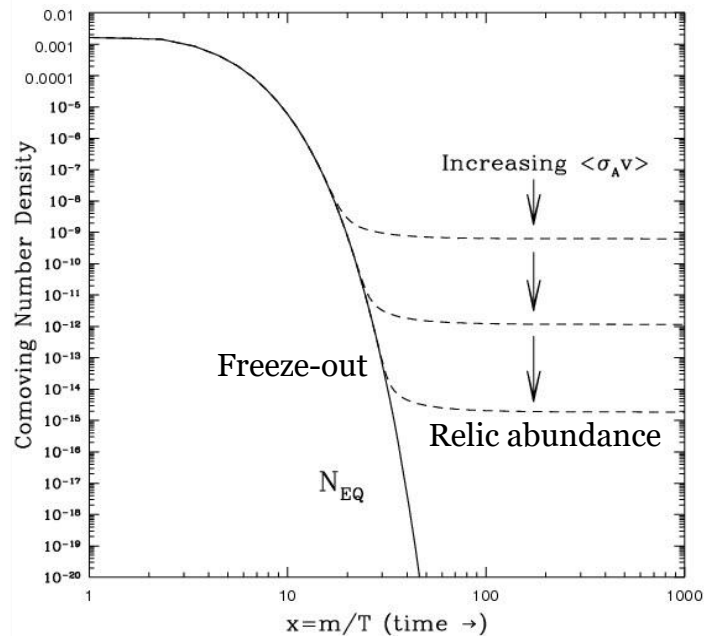
Steven Weinberg^(c)

Stanford University, Physics Department, Stanford, California 94305

(Received 13 May 1977)

The present cosmic mass density of possible stable neutral heavy leptons is calculated in a standard cosmological model. In order for this density not to exceed the upper limit of $2 \times 10^{-29} \text{ g/cm}^3$, the lepton mass would have to be *greater* than a lower bound of the order of 2 GeV.

PRL (1977)



- Correct thermal relic abundance:

$$\Omega h^2 \sim \frac{0.1 \text{ pb}}{\langle \sigma v \rangle} \text{ with } \langle \sigma v \rangle \sim \frac{\alpha_X^2 m_X^2}{M^4} \text{ (M: dark scale/mediator)}$$

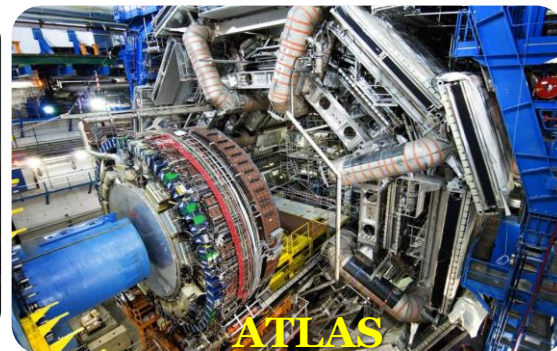
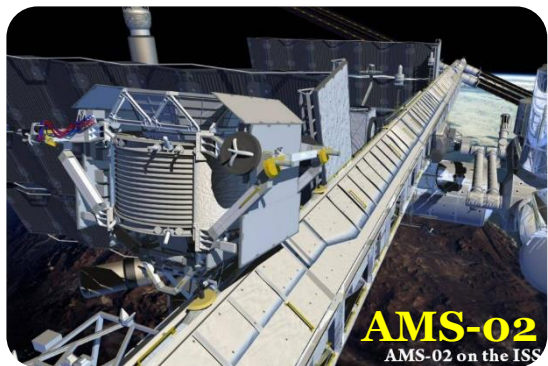
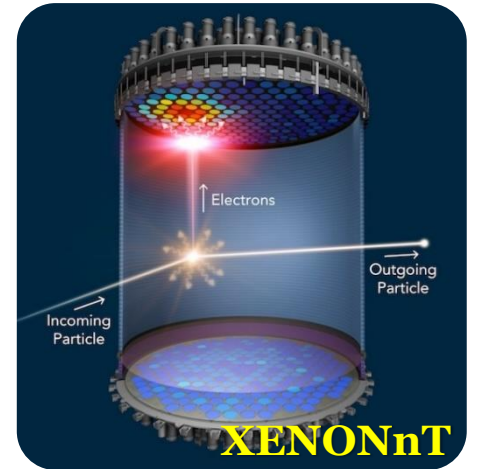
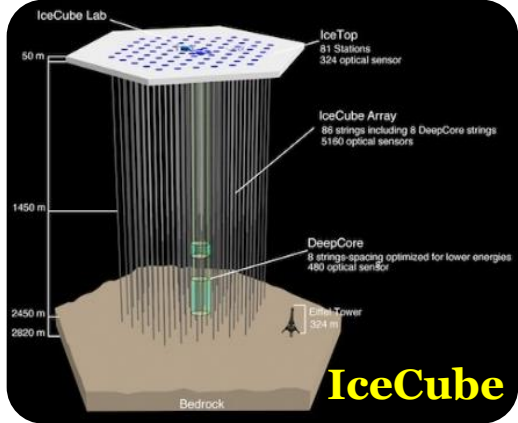
- Weak coupling → naturally weak scale mass:

~ 1 GeV – 100 TeV mass range favored

→ weak scale (new) physics

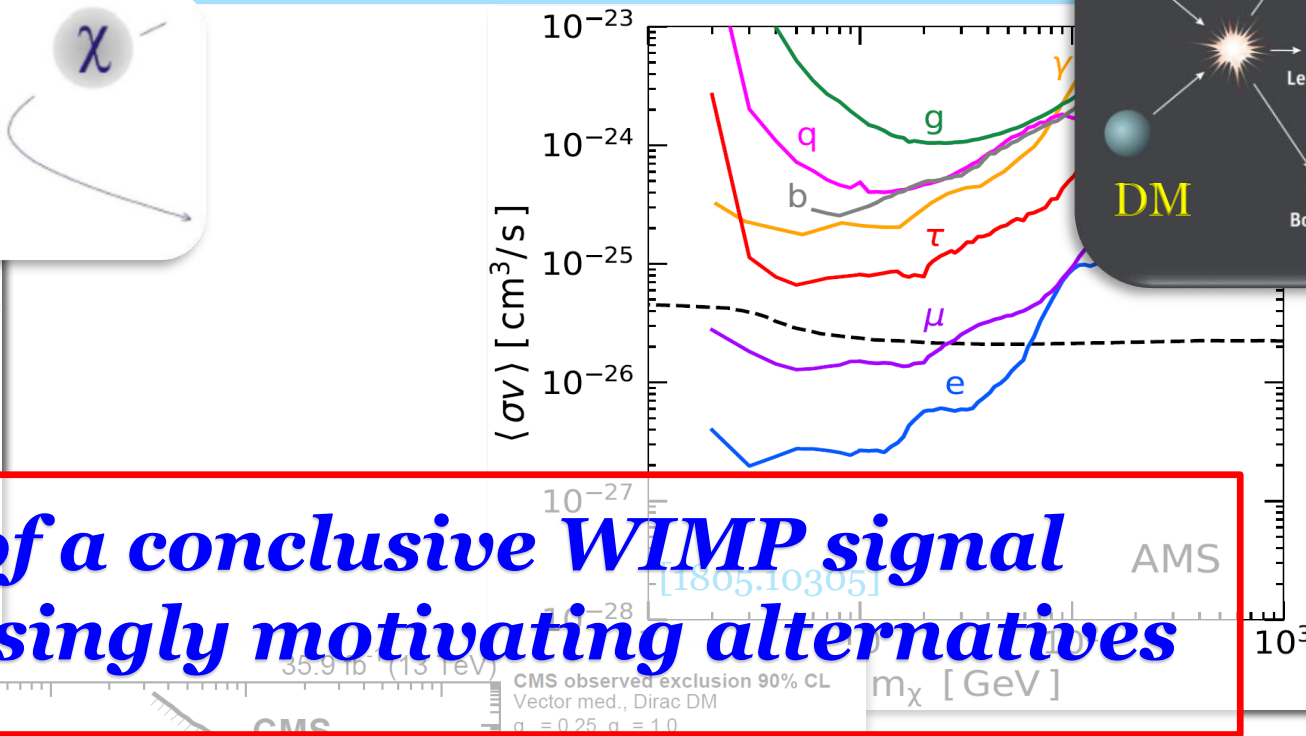
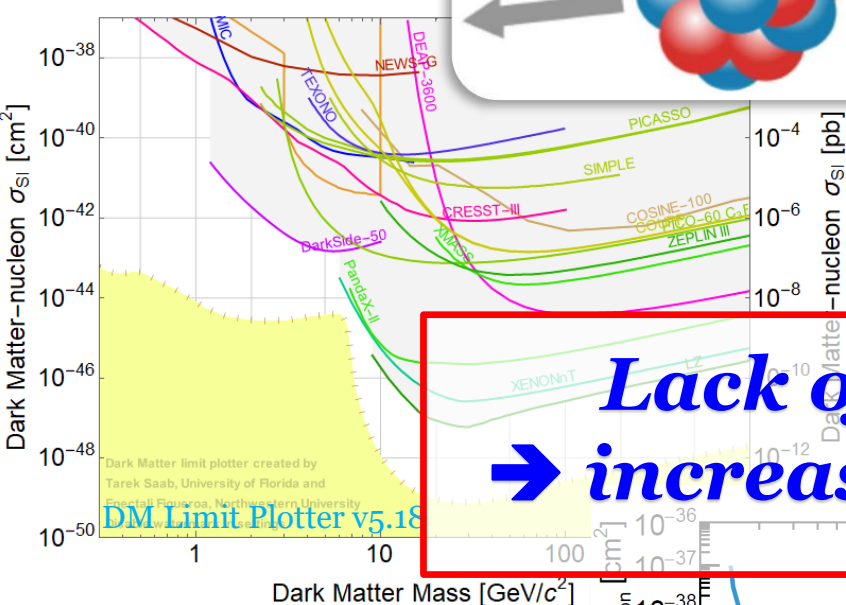
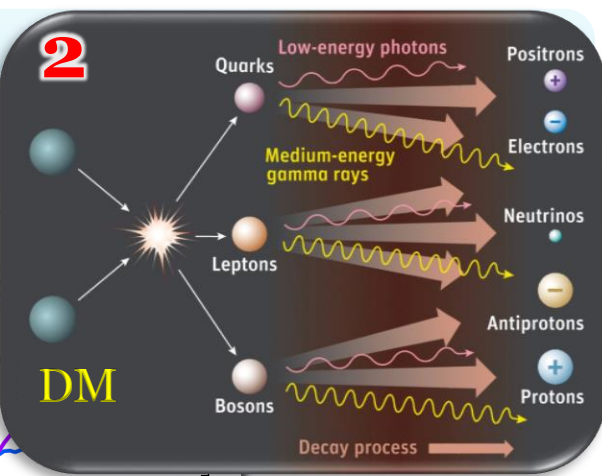
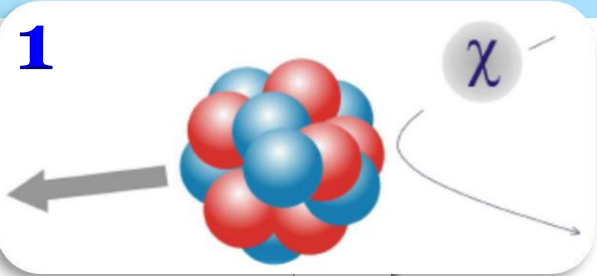
* Of course also **axion**

Diverging Efforts for WIMP Searches



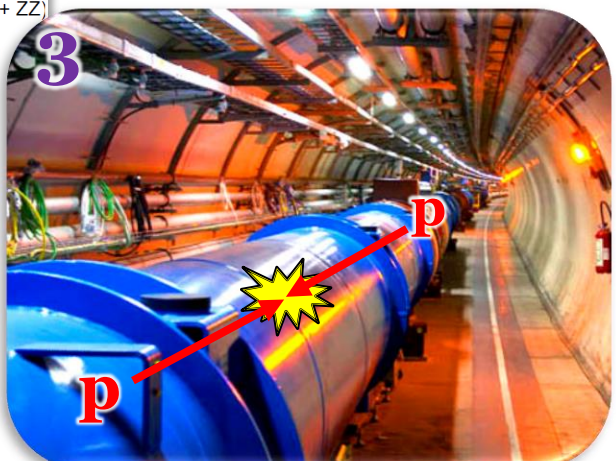
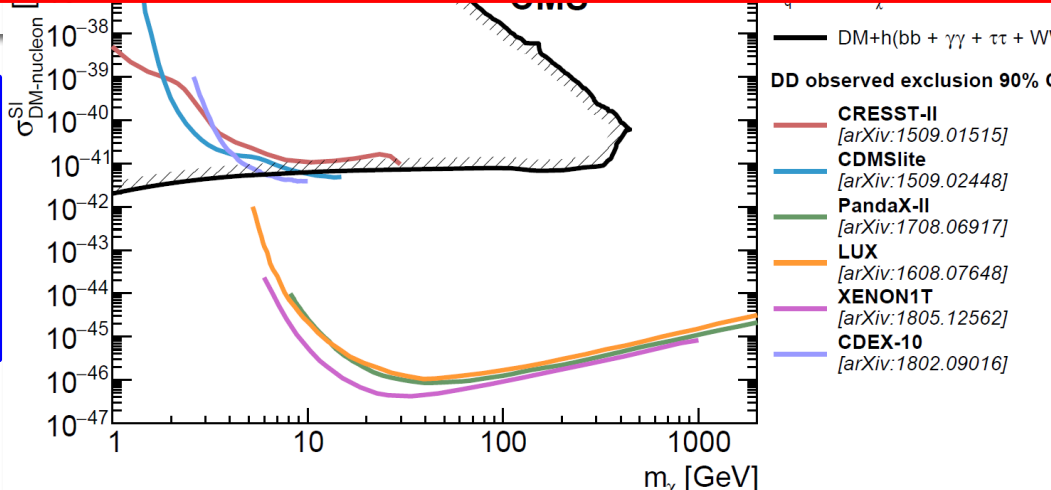
Current Search Status of DM (WIMP)

1



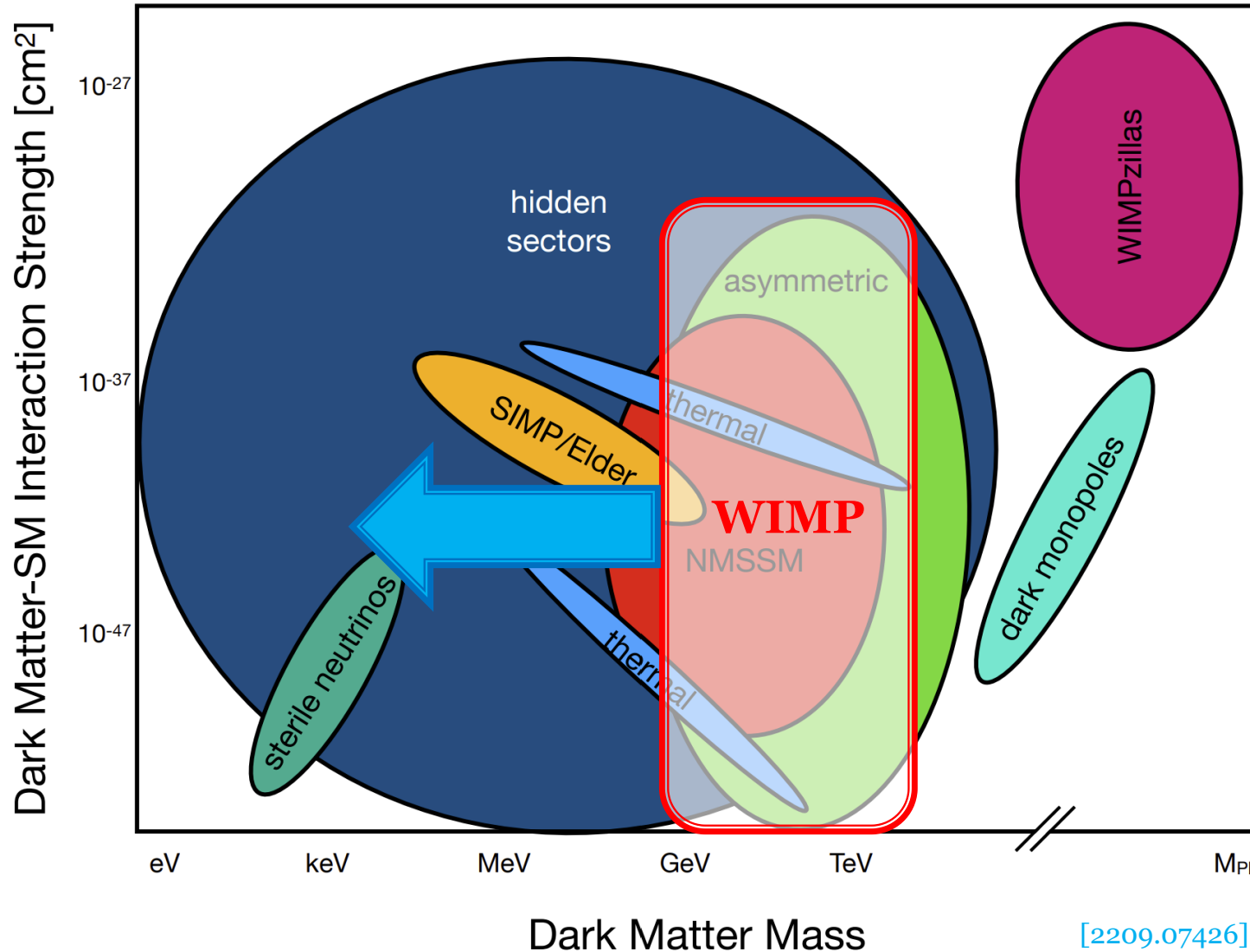
Lack of a conclusive WIMP signal
→ increasingly motivating alternatives

❖ **No solid observation**
→ Stringent constraints
 on WIMP



Particle DM: Larger Space beyond WIMP

Report of the Topical Group on Particle Dark Matter for Snowmass 2021



- ❖ Light DM models involving new light mediators: growing attention.
- ❖ The experimental detection of such light DM candidates: **challenging** as the conventional search efforts are designed and optimized for WIMP.

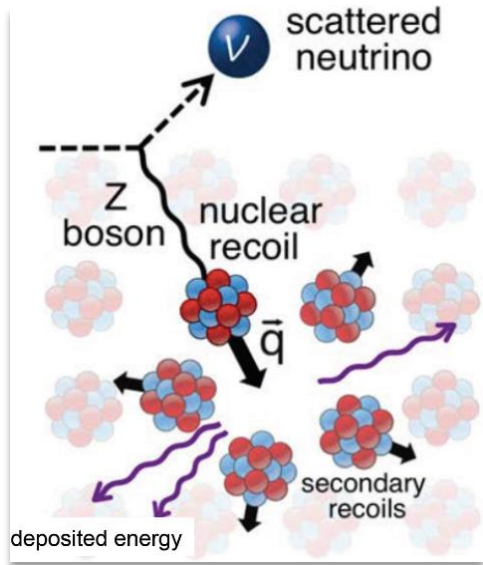
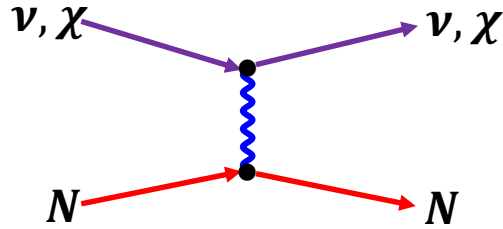
[2209.07426]



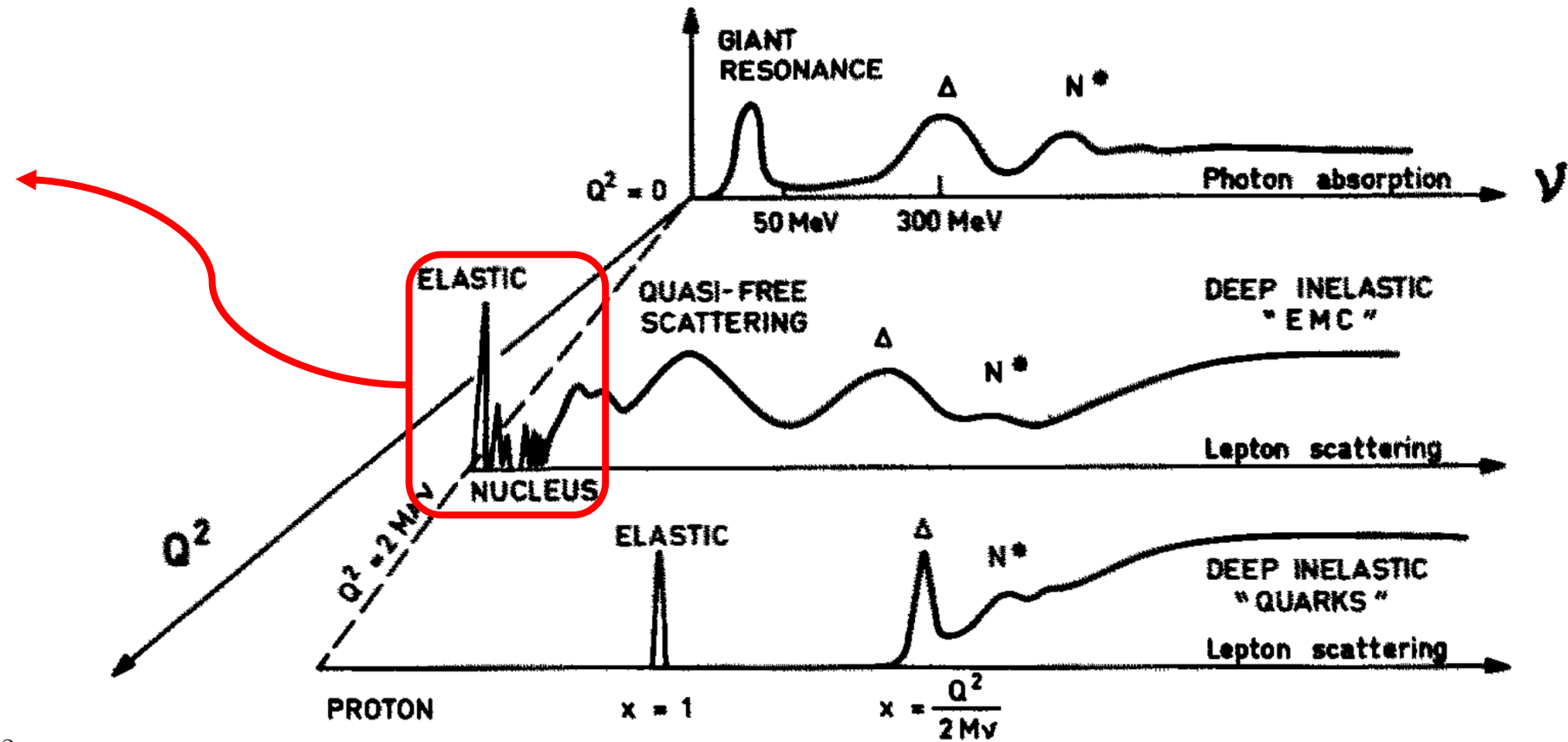
Nuclear Scattering

Nuclear Scattering: Elastic

❖ Coherent elastic nuclear scattering



NUCLEAR RESPONSE FUNCTION



$$\sigma \propto Q_W^2 \propto (N - (1 - 4 \sin^2 \theta_W)Z)^2$$

$$\Rightarrow \sigma \propto N^2$$

B. Frois, Nuclear Structure (1985)

Nuclear Scattering: Elastic Neutrino

❖ Coherent Elastic ν Nucleus Scattering (CE ν NS)

PRD (1974)

Coherent effects of a weak neutral current

Daniel Z. Freedman†

National Accelerator Laboratory, Batavia, Illinois 60510

and Institute for Theoretical Physics, State University of New York, Stony Brook, New York 11790

(Received 15 October 1973; revised manuscript received 19 November 1973)

If there is a weak neutral current, then the elastic scattering process $\nu + A \rightarrow \nu + A$ should have a sharp coherent forward peak just as $e + A \rightarrow e + A$ does. Experiments to observe this peak can give important information on the isospin structure of the neutral current. The experiments are very difficult, although the estimated cross sections (about 10^{-38} cm² on carbon) are favorable. The coherent cross sections (in contrast to incoherent) are almost energy-independent. Therefore, energies as low as 100 MeV may be suitable. Quasi-

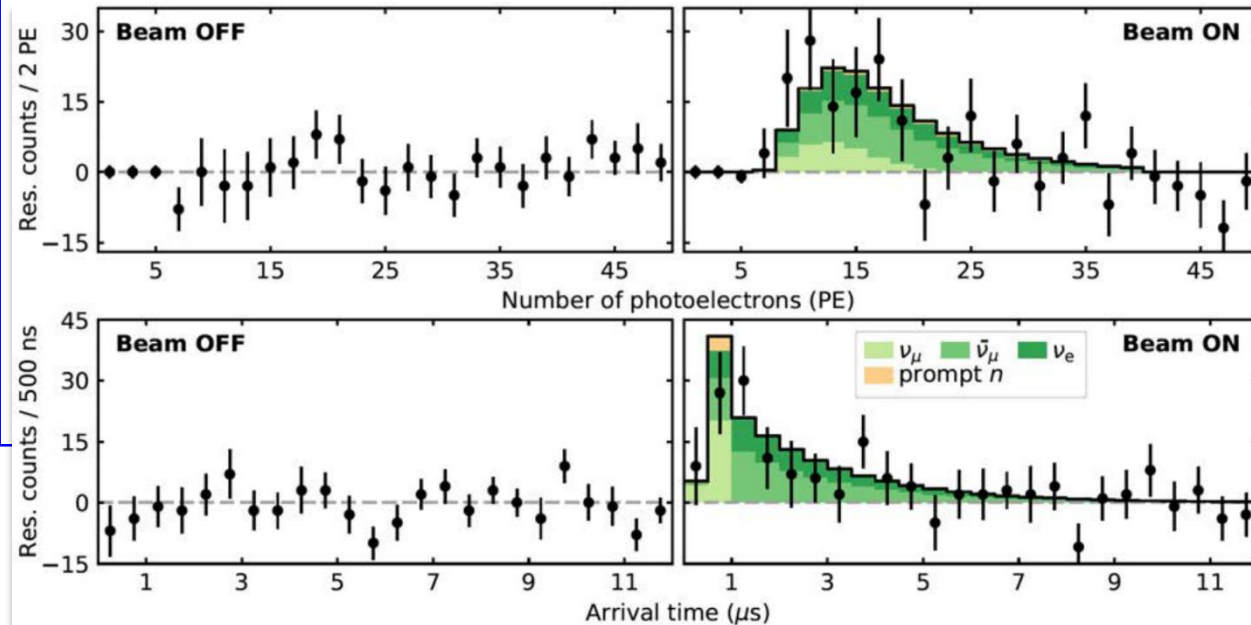
PRD (1984) Principles and applications of a neutral-current detector for neutrino physics and astronomy

A. Drukier and L. Stodolsky

Max-Planck-Institut für Physik und Astrophysik, Werner-Heisenberg-Institut für Physik, Munich, Federal Republic of Germany

(Received 21 November 1983)

We study detection of MeV-range neutrinos through elastic scattering on nuclei and identification of the recoil energy. The very large value of the neutral-current cross section due to coherence indicates a detector would be relatively light and suggests the possibility of a true “neutrino observatory.” The recoil energy which must be detected is very small (10 – 10^3 eV), however. We examine a realization in terms of the superconducting-grain idea, which appears, in principle, to be feasible through extension and extrapolation of currently known techniques. Such a detector could permit determination of the neutrino energy spectrum and should be insensitive to neutrino oscillations since it detects all neutrino types. Various applications and tests are discussed, including spallation sources, reactors, supernovas, and solar and terrestrial neutrinos. A preliminary estimate of the most difficult backgrounds is attempted.



134 ± 22 events observed! \rightarrow The absence of CE ν NS events is rejected at 6.7σ .

COHERENT [1708.01294]

Nuclear Scattering: Elastic DM

❖ Coherent Elastic **DM** Nucleus Scattering → DM Direct Detection

Detectability of certain dark-matter candidates

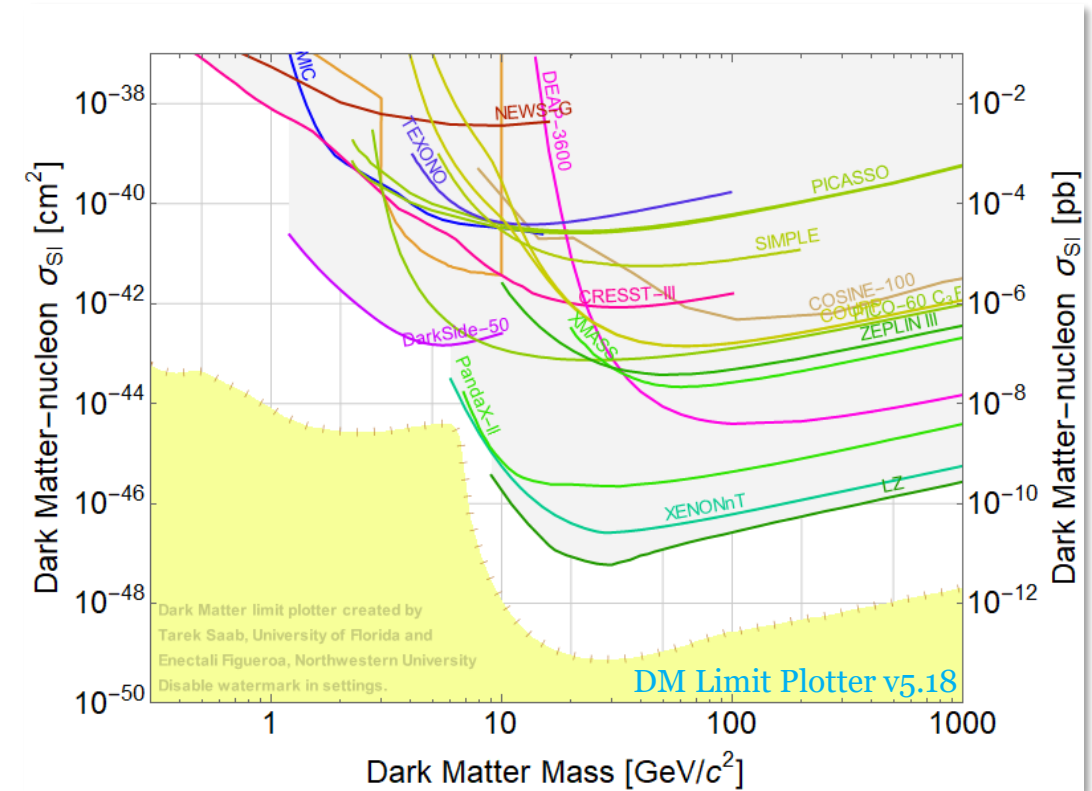
Mark W. Goodman and Edward Witten

Joseph Henry Laboratories, Princeton University, Princeton, New Jersey 08544

(Received 7 January 1985)

We consider the possibility that the neutral-current neutrino detector recently proposed by Drukier and Stodolsky could be used to detect some possible candidates for the dark matter in galactic halos. This may be feasible if the galactic halos are made of particles with coherent weak interactions and masses $1-10^6$ GeV; particles with spin-dependent interactions of typical weak strength and masses $1-10^2$ GeV; or strongly interacting particles of masses $1-10^{13}$ GeV.

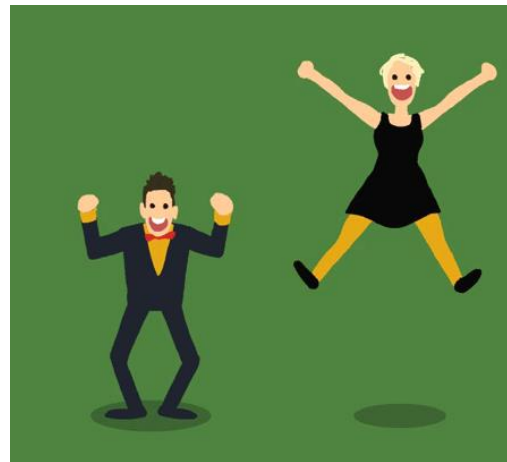
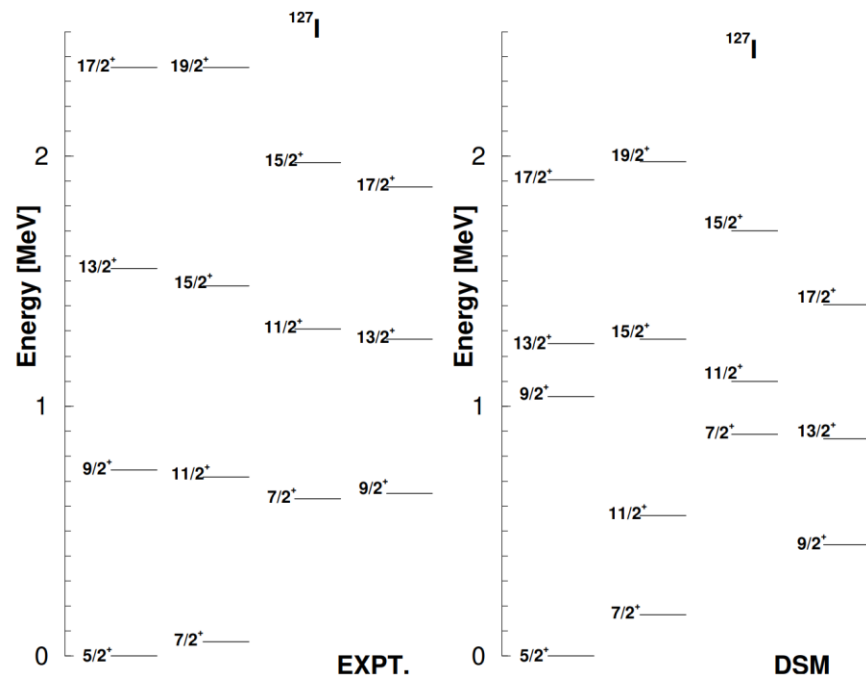
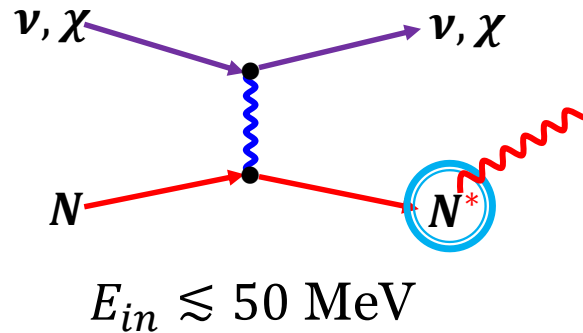
PRD (1985)



No solid observation yet.

Nuclear Scattering: *Inelastic*

❖ Why **inelastic** channel?



➤ Signatures

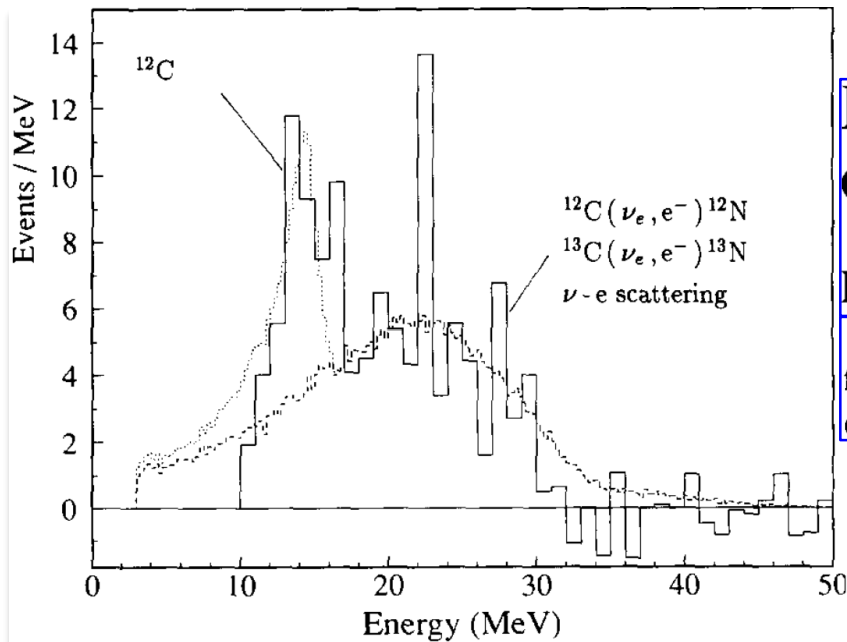
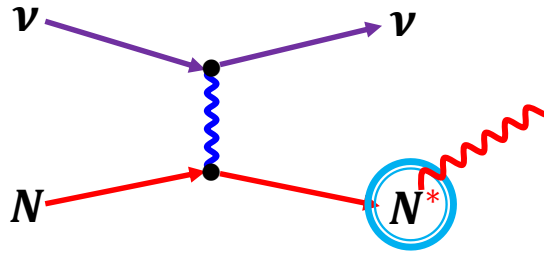
- ✓ **Elastic:** low energy nuclear recoil
- ✓ **Inelastic:** γ cascade ($\Delta E \lesssim 10 \text{ MeV}$),
 γ cascade + nucleons ($\Delta E \gtrsim 10 \text{ MeV}$)

➤ Motivation

- ✓ A new channel to study
- ✓ Larger energy $\sim O(1 - 10) \text{ MeV}$
- ✓ Better S/B ratio

Nuclear Scattering: *Inelastic Neutrino*

❖ Inelastic ν Nucleus Scattering (I ν NS)



TESTING THE STRUCTURE OF WEAK NEUTRAL CURRENTS BY INELASTIC NEUTRINO SCATTERING FROM NUCLEI*

T.W. DONNELLY* and R.D. PECCEI

*Institute of Theoretical Physics, Department of Physics,
Stanford University, Stanford, California 94305*

Received 13 June 1976

Revised manuscript received 17 September 1976

By using selected nuclear transitions specific pieces of the weak neutral current may be greatly enhanced, leading to widely different results for different models of the weak neutral current. Predictions for low-energy inelastic neutrino scattering from ^{12}C are examined within the framework of a variety of $\text{SU}(2)_W \times \text{U}(1)$ gauge theory models.

PLB (1976)

First observation
of the neutral current nuclear excitation $^{12}\text{C}(\nu, \nu')^{12}\text{C}^*(1^+, 1)$

KARMEN Collaboration

PLB (1991)

The neutral current nuclear excitation $^{12}\text{C}(\nu, \nu')^{12}\text{C}^*(1^+, 1; 15.1 \text{ MeV})$ has been observed for the first time. For ν_e and $\bar{\nu}_\mu$ from μ^+ -decay at rest the flux averaged cross section was determined to be $\langle \sigma_{\text{NC}}(\nu_e + \bar{\nu}_\mu) \rangle = [10.8 \pm 5.1 (\text{stat.}) \pm 1.1 (\text{syst.})] \times 10^{-42} \text{ cm}^2$.

➤ Recent

✓ Inclusion of multiple excited states

improvements

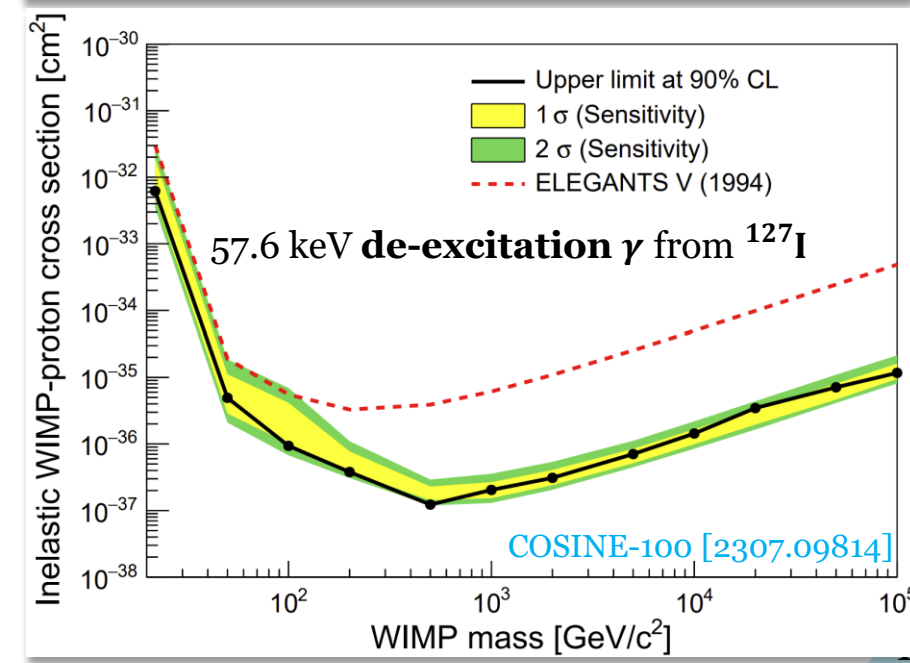
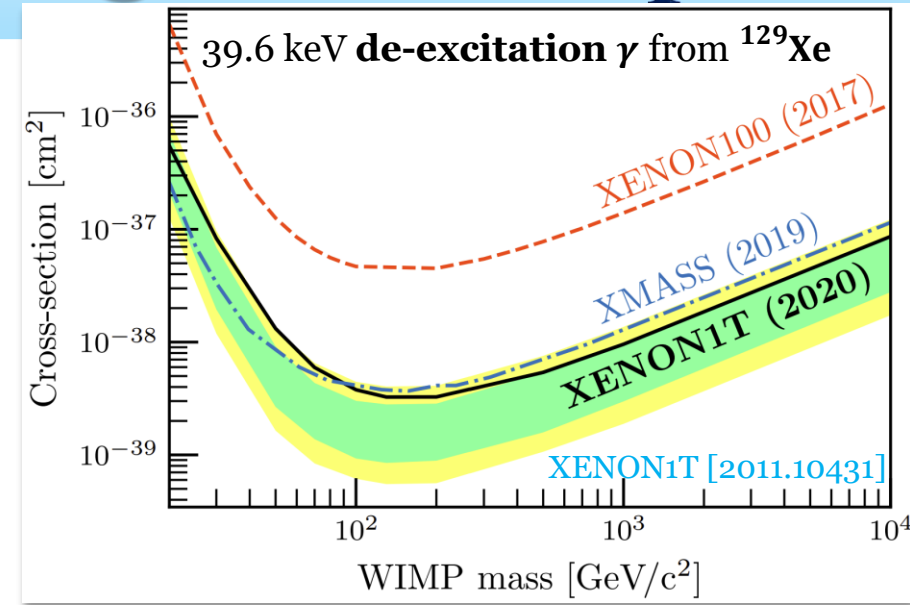
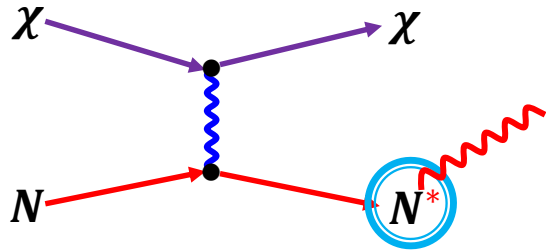
✓ Consistent handling of hadronic currents

Dutta, Newstead et al.,
[2206.08590]

✓ Exclusive cross sections for each state

Nuclear Scattering: *Inelastic* DM @ Direct Exps.

❖ Inelastic DM Nucleus Scattering



RATES FOR INELASTIC NUCLEAR EXCITATION BY DARK MATTER PARTICLES

John ELLIS ^{a)}, R.A. FLORES ^{a),1} and J.D. LEWIN ^{b)}

^{a)} CERN, CH-1211 Geneva 23, Switzerland

^{b)} Rutherford-Appleton Laboratory, Chilton, Didcot, OX11 0QX, UK

Received 29 June 1988

We calculate rates for the inelastic scattering of dark matter particles X on nuclei to produce low-lying excited nuclear states. Assuming a Maxwellian velocity distribution for the dark matter particles X , the inelastic two-body phase space suppresses all rates by factors > 10 for $m_X \leq 100$ GeV unless the excitation energy $\Delta E < 100$ keV. We catalogue all stable nuclei with excited states in this range, and we estimate inelastic scattering matrix elements by relating them to M1 transitions, after correcting for internal conversions. Our calculated rates are typically $\lesssim 10^{-4}$ events/kg·day, with the least unfavourable rates being for ^{169}Tm and ^{187}Os . Problems of natural radioactivity and expense disfavour these and many other materials, leaving ^{127}I , ^{183}W and ^{201}Hg as the least unpromising isotopes.

PLB (1988)

Is it possible to search for *inelastic nuclear scatterings* by **DM** in **neutrino detectors**?

Kinematically impossible for the ambient DM to leave *detectable MeV-range signatures* through the inelastic channel: $E_R \sim \mu_{\chi N} v^2 < \text{MeV}$

Is it possible to search for *inelastic nuclear scatterings* by **DM** in **neutrino detectors**?

Kinematically impossible for the ambient DM to leave detectable MeV-range signatures through the inelastic channel: $E_R \sim \mu_{\chi N} v^2 < \text{MeV}$





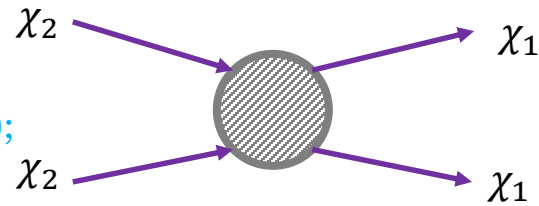
Boosted Dark Matter (BDM)

DM Boosting Mechanisms: Dark Sector



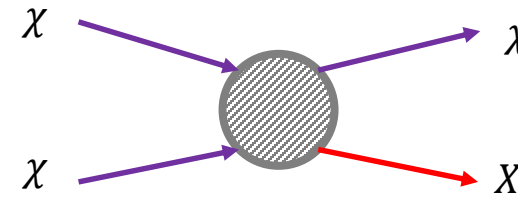
Boosted DM (BDM) coming from the Universe

[Belanger & JCP, JCAP (2012);
Agashe et al., JCAP (2014);
Kong, Mohlabeng, JCP, PLB (2015);
Berger et al., JCAP (2015);
Kim, JCP, Shin, PRL (2017);
more]



✓ Multi-component model

$$m_2 \gg m_1$$



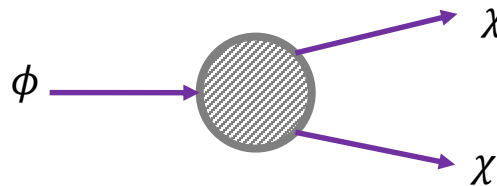
✓ Semi-annihilation model

$$m_\chi \gg m_X$$

[D'Eramo & Thaler, JHEP (2010);
Berger et al., JCAP (2015)]

Large E_k^{DM} (monochromatic) due to mass gap

❖ Relic component DM:
Non-relativistic!



✓ Decaying multi-component DM

$$m_\phi \gg m_\chi$$

[Bhattacharya et al., JCAP (2015);
Kopp et al., JHEP (2015);
Cline et al., PRD (2019);
Heurtier, Kim, JCP, Shin, PRD (2019);
more]

DM Boosting Mechanisms: Cosmic-Rays (CRs)

Cosmic-Ray-Induced BDM

- ❖ **Charged CRs:** [Bringmann & Pospelov, PRL (2019); Ema et al., PRL (2019); Cappiello & Beacom, PRD (2019); Dent & Dutta et al., PRD (2020); Jho, JCP, Park & Tseng, PLB (2020); Cho et al., PRD (2020); more]
- ❖ **CR ν (ν BDM):** [Jho, JCP, Park & Tseng, 2101.11262; Das & Sen, 2104.00027; Chao, Li, Liao, 2108.05608; Lin, Wu, Wu, Wong, 2206.06864; more]

- ❖ **BDM from astrophysical processes:**
 - Solar evaporation - Kouvaris, PRD (2015)
 - Dark cosmic rays - Hu +, PLB (2017)
 - Solar reflection - An +, PRL (2018)
 - Solar acceleration - Emken +, PRD (2018)
 - Atmospheric collider - Alvey+, PRL (2019)
 - PBH evaporation - Calabrese +, PRD (2022)
 - Blazar jets - Wang +, PRL (2022)
 - Supernova shocks - Cappiello
 - more

- ❖ **Energetic cosmic-ray-induced BDM:** energetic cosmic-rays kick DM (large $E_{e^\pm, p^\pm, \text{He}, \nu, \dots}$ \rightarrow large E_χ)
 \rightarrow **Efficient for Light DM**



$e^\pm, p^\pm, \text{He}, \nu, \dots$

$e^\pm, p^\pm, \text{He}, \nu, \dots$

$e^\pm, p^\pm, \text{He}, \nu, \dots$

χ

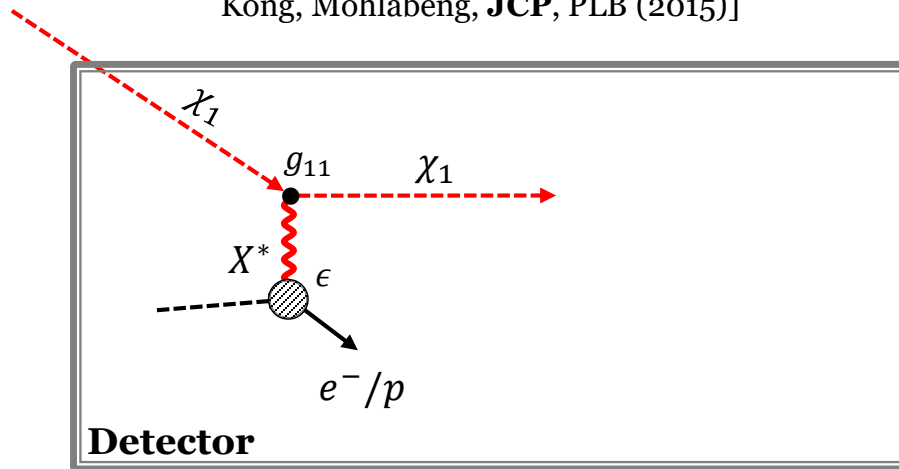
χ

BDM Signatures: Elastic or Inelastic DM

❖ BDM signal: detectable at **large volume detectors**

elastic scattering (eBDM)

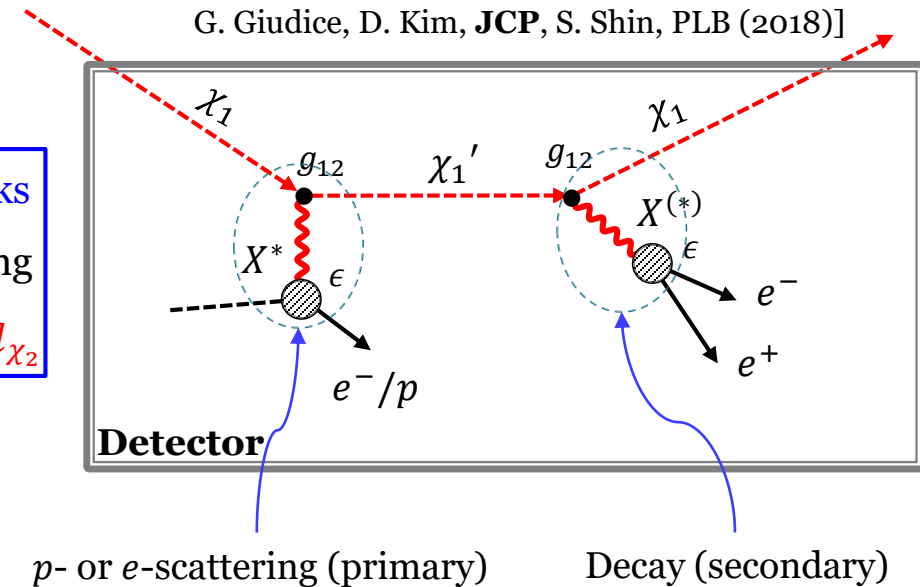
[Agashe, Cui, Necib, Thaler, JCAP (2014);
Kong, Mohlabeng, JCP, PLB (2015)]



inelastic scattering (iBDM)

[D. Kim, JCP, S. Shin, PRL (2017);
G. Giudice, D. Kim, JCP, S. Shin, PLB (2018)]

1~3 tracks
depending
on E_{th} & l_{χ_2}



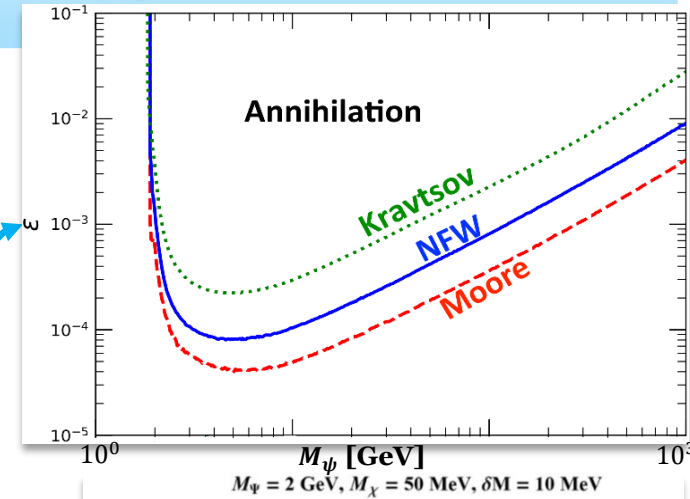
BDM Searches @ Neutrino Experiments

***Boosted DM (BDM) scenarios:
Receiving rising attention***

PHYSICAL REVIEW LETTERS **120**, 221301 (2018)

Editors' Suggestion

Search for Boosted Dark Matter Interacting with Electrons in Super-Kamiokande



Eur. Phys. J. C (2021) 81:322
<https://doi.org/10.1140/epjc/s10052-021-09007-w>

THE EUROPEAN
PHYSICAL JOURNAL C



Regular Article - Experimental Physics

**Prospects for beyond the Standard Model physics searches at the
Deep Underground Neutrino Experiment**

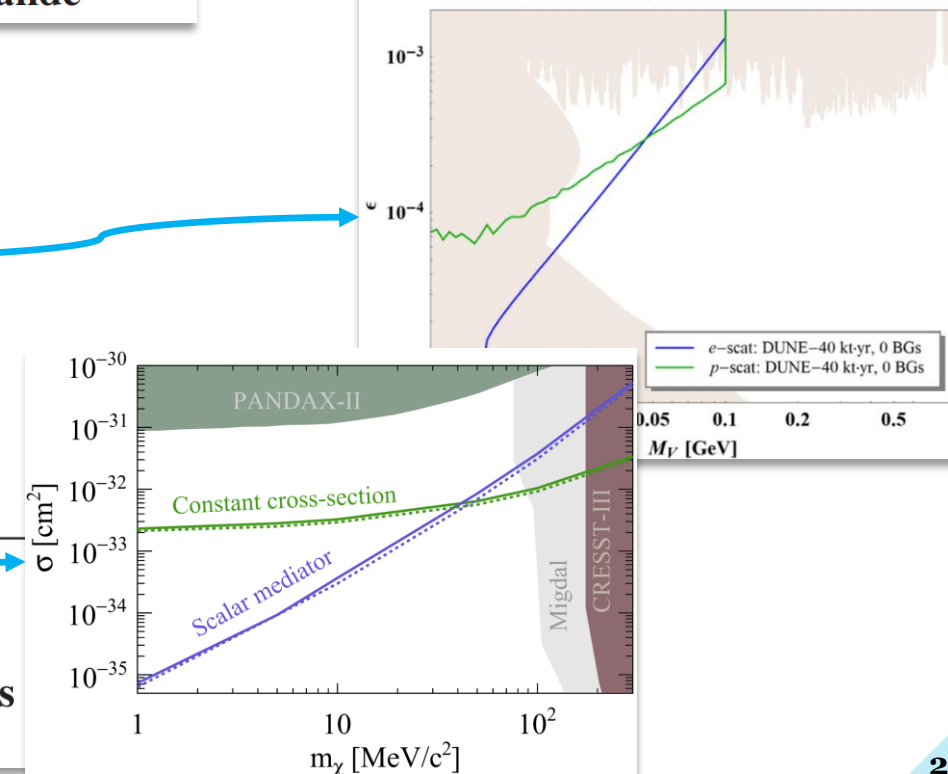
DUNE Collaboration

PHYSICAL REVIEW LETTERS **130**, 031802 (2023)

Editors' Suggestion

Featured in Physics

**Search for Cosmic-Ray Boosted Sub-GeV Dark Matter Using Recoil Protons
at Super-Kamiokande**



BDM Searches @ DM Experiments

PHYSICAL REVIEW LETTERS **122**, 131802 (2019)

Editors' Suggestion

First Direct Search for Inelastic Boosted Dark Matter with COSINE-100

PHYSICAL REVIEW LETTERS **131**, 201802 (2023)

Search for Boosted Dark Matter in COSINE-100

PHYSICAL REVIEW LETTERS **128**, 171801 (2022)

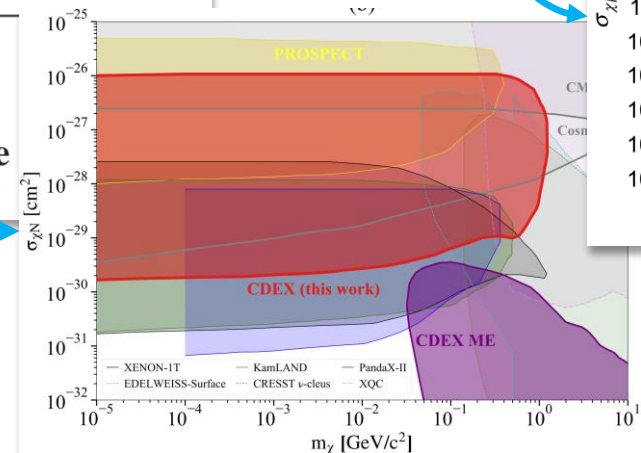
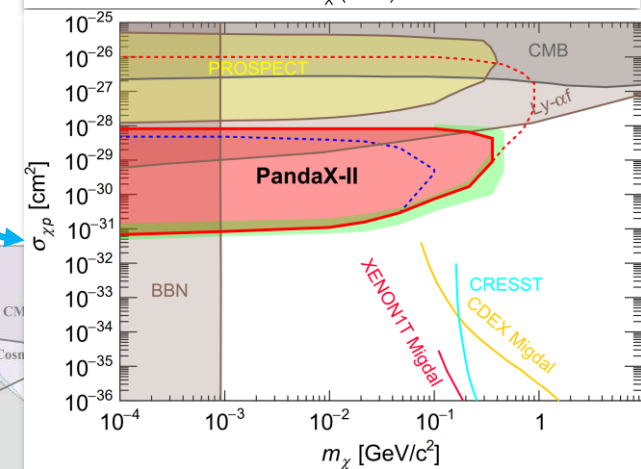
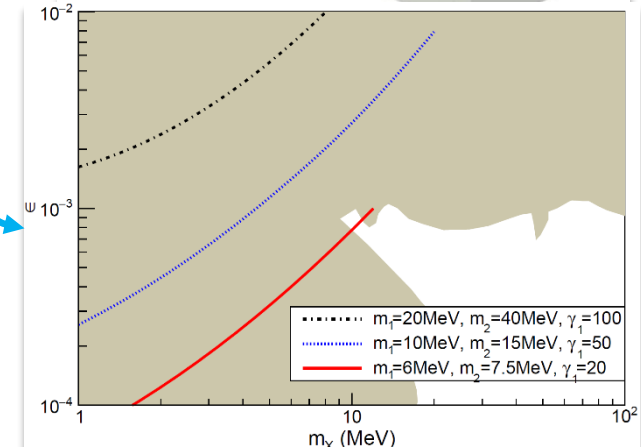
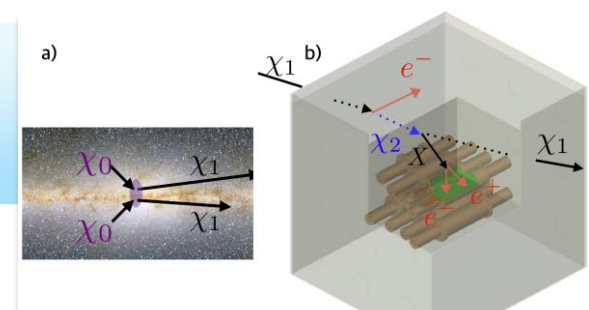
Editors' Suggestion

Search for Cosmic-Ray Boosted Sub-GeV Dark Matter at the PandaX-II Experiment

PHYSICAL REVIEW D **106**, 052008 (2022)

Constraints on sub-GeV dark matter boosted by cosmic rays from the CDEX-10 experiment at the China Jinping Underground Laboratory

✓ Searches for $e/p/N$ **elastic** scattering signals





Exciting Prospects for DM

Benchmark: BDM by Cosmic Rays

❖ **Focus:** the interaction between DM & quark

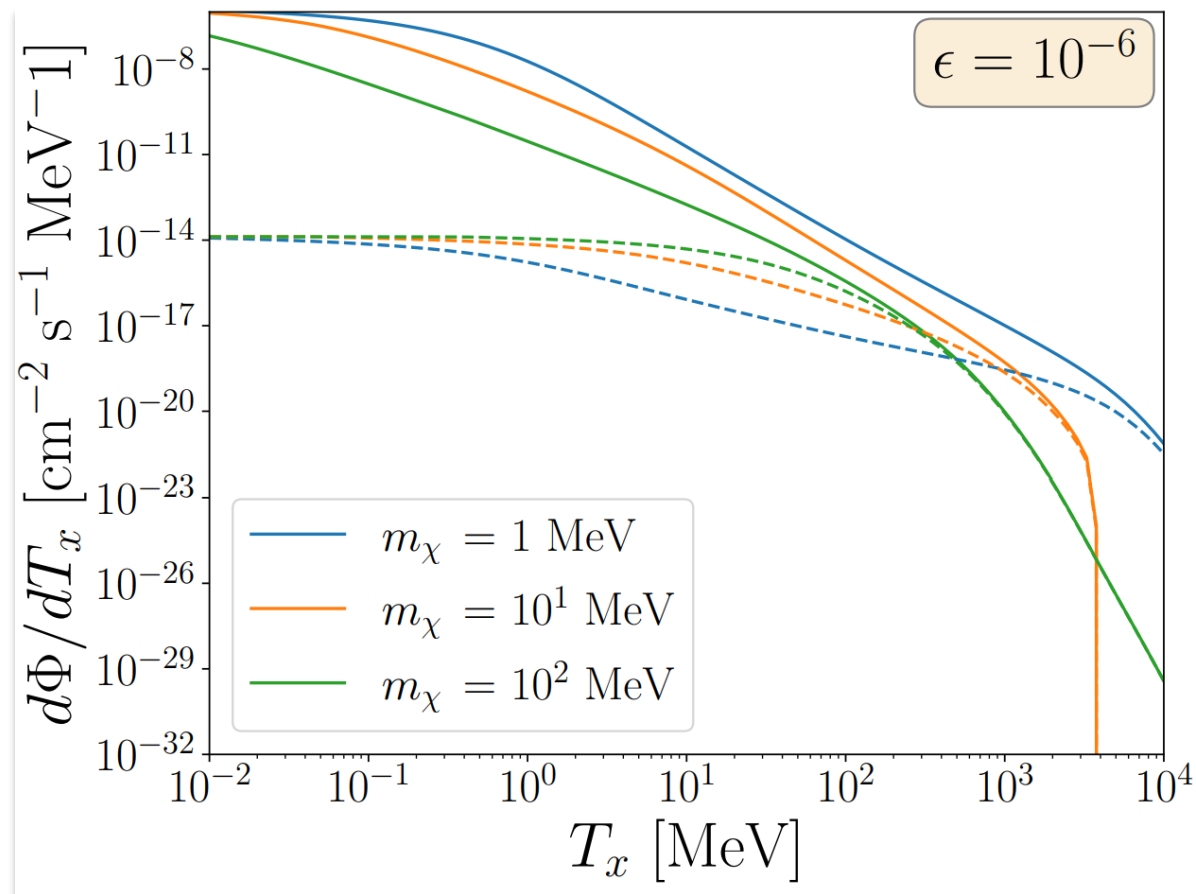
$$\mathcal{L} \supset g_D A'_\mu \bar{\chi} \gamma^\mu \chi + \epsilon Q_b A'_\mu \bar{q} \gamma^\mu q$$

➔ **DM** can be **boosted by energetic cosmic rays**
(p, He) throughout the galaxy.

❖ **CR-BDM flux**

$$\begin{aligned} \frac{d\Phi_\chi}{dT_\chi} &= \int_V dV \int_{T_i^{\min}} dT_i \frac{d^2\Gamma_{i \rightarrow \chi}}{dT_i dT_\chi} \\ &= D_{\text{eff}} \frac{\rho_\chi}{m_\chi} \sum_i \int_{T_i^{\min}} dT_i \frac{d\sigma_{\chi i}^{\text{el}}}{dT_\chi} \frac{d\Phi_i^{\text{LIS}}}{dT_i} \end{aligned}$$

[Bringmann & Pospelov, PRL (2019)]



✓ $m_{A'} = 1 \text{ MeV}$ (solid) & 100 MeV (dashed)

Inelastic Nuclear Scattering of DM

- ❖ **Gamow-Teller (GT) transitions** are the dominant contribution to the inelastic cross section.

$$\frac{d\sigma_{\chi N}^{\text{inel}}}{d\cos\theta} = \frac{2\epsilon^2 g_D^2 E'_\chi p'_\chi}{(2m_T E_R + m_{A'}^2 - \Delta E^2)^2} \frac{1}{2\pi} \frac{4\pi}{2J+1}$$

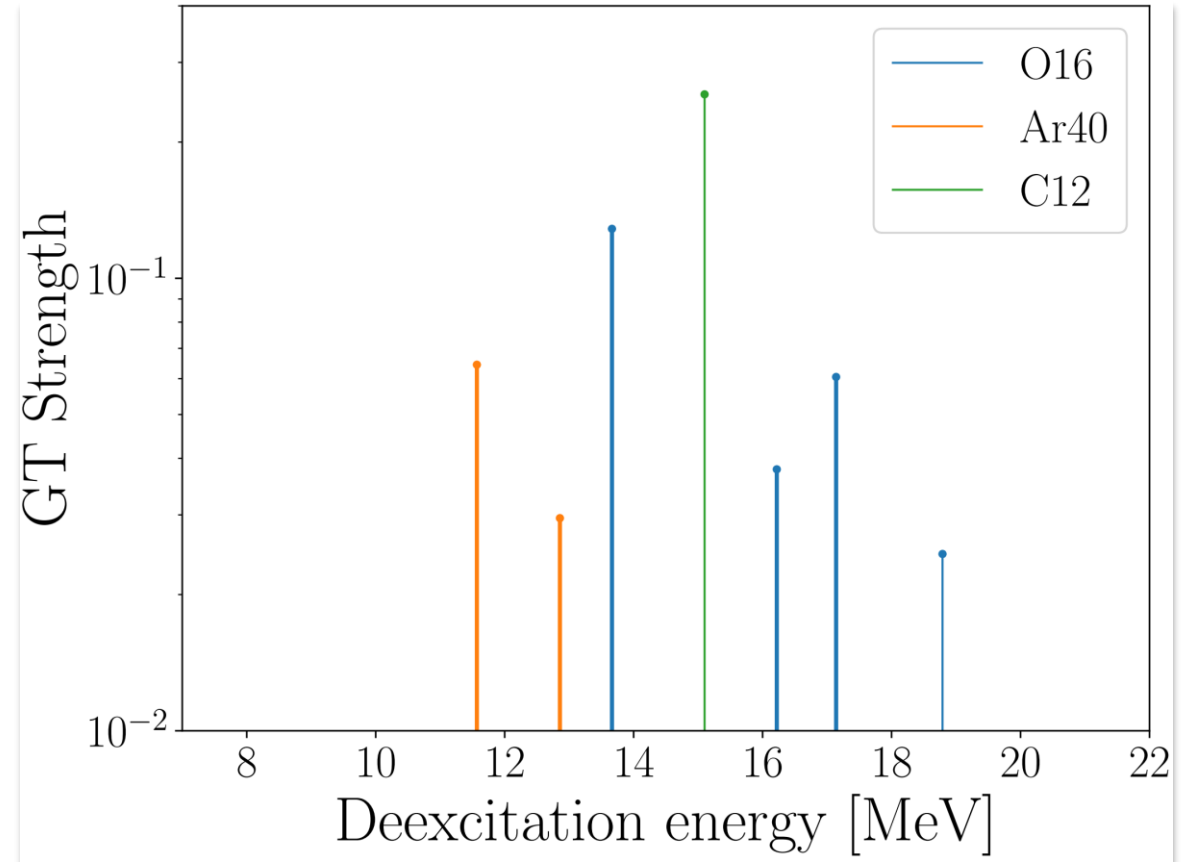
$$\times \sum_{s_i, s_f} \vec{l} \cdot \vec{l}^* \frac{g_A^2}{12\pi} \left| \langle J_f || \sum_{i=1}^A \frac{1}{2} \hat{\sigma}_i \hat{\tau}_0 || J_i \rangle \right|^2,$$

$$\sum_{s_i, s_f} \vec{l} \cdot \vec{l}^* = 3 - \frac{1}{E_\chi E'_\chi} \left[\frac{1}{2} (p_\chi^2 + p'_\chi^2 - 2m_T E_R) + \frac{3m_\chi^2}{4} \right]$$

* For more details, See e.g. Dutta et al., [2206.08590].

- ❖ The expected # of signal events

$$N_\chi = N_T \Delta t \int \sigma_{\chi N}^{\text{inel}}(E_\chi) \frac{d\Phi_\chi}{dE_\chi} dE_\chi \cdot \frac{\Gamma_{N^* \rightarrow N\gamma}}{\Gamma_{\text{total}}}$$



The **GT strengths** are derived from experimental results & the large-scale shell model code BIGSTICK.

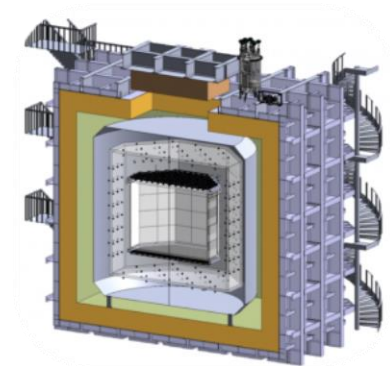
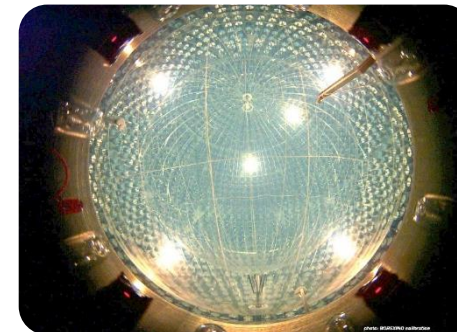
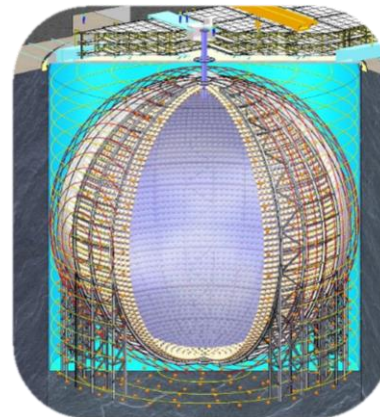
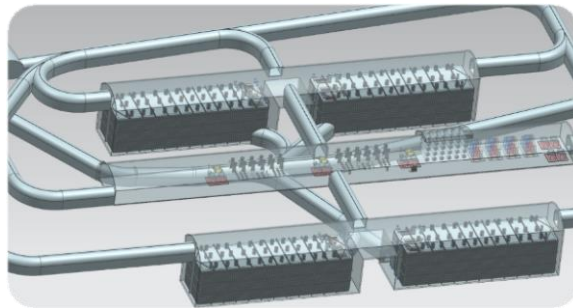
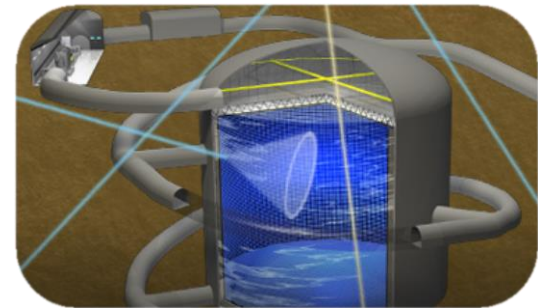
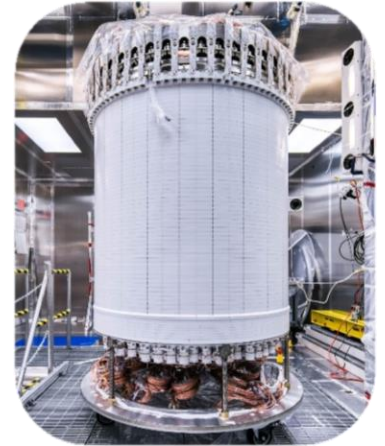
Benchmark Detectors

❖ **Key specifications** of the benchmark detectors

inelastic

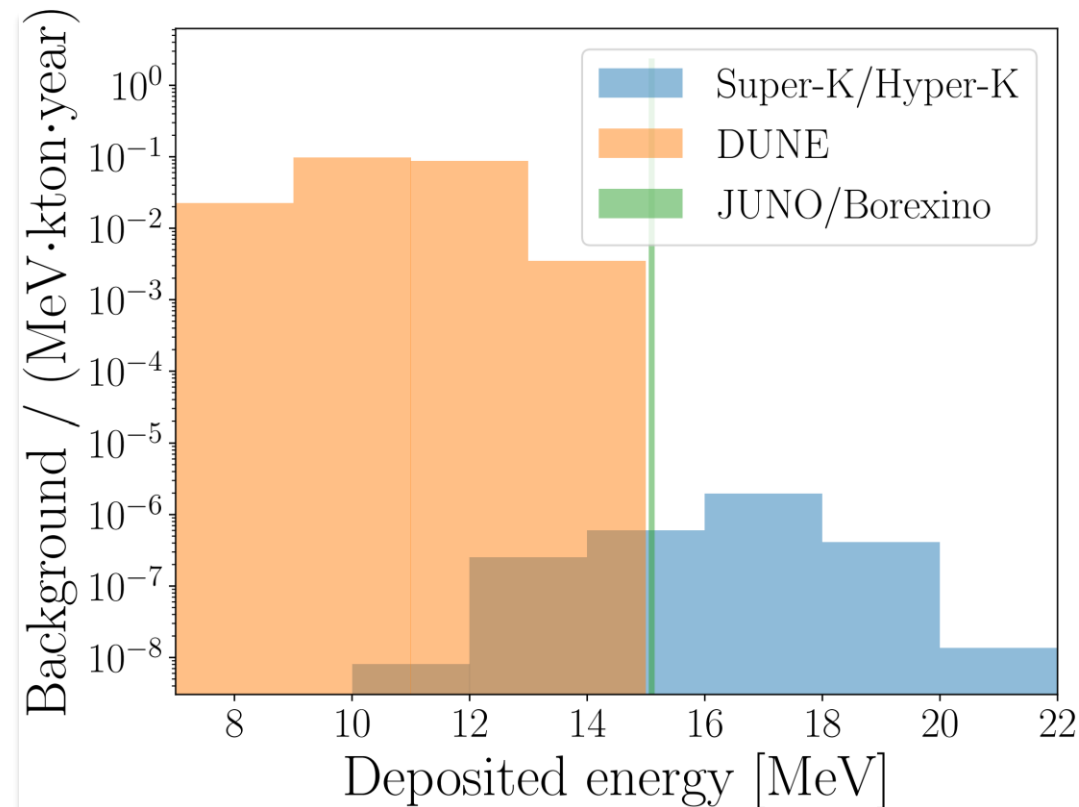
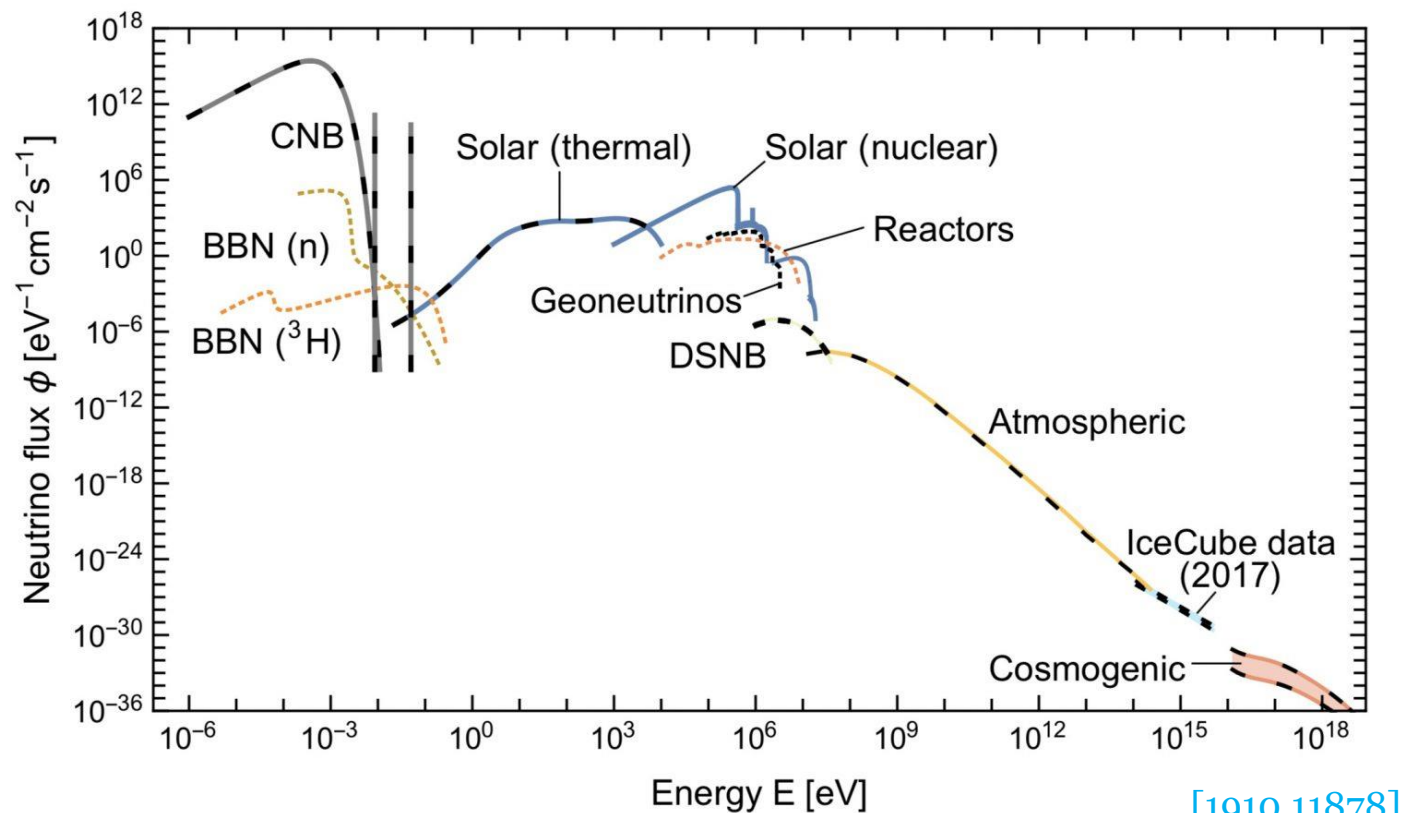
elastic

Experiment	Target	Mass (kton)	Time (year)	E_{res} (MeV)	E_{range} (MeV)	Background
Super-K	water	22.5	20	2	10 - 22	2.9×10^{-3}
Hyper-K	water	188	10	2	10 - 22	1.2×10^{-2}
DUNE	Ar	40	10	2	7 - 15	167.5
JUNO	CH ₂	20	10	-	15.1	47
Borexino	CH ₂	0.1	5.6	-	15.1	0.13
DUNE	p,n	40	10	-	50 - 1270	550
JUNO	p	20	10	-	15 - 100	1420
Borexino	p	0.1	1.22	-	12.3 - 24.6	84
LZ	Xe	0.0055	0.167	-	$(1 - 68) \cdot 10^{-3}$	1.4
DarkSide	Ar	0.02	5	-	0.03 - 0.1	1.6

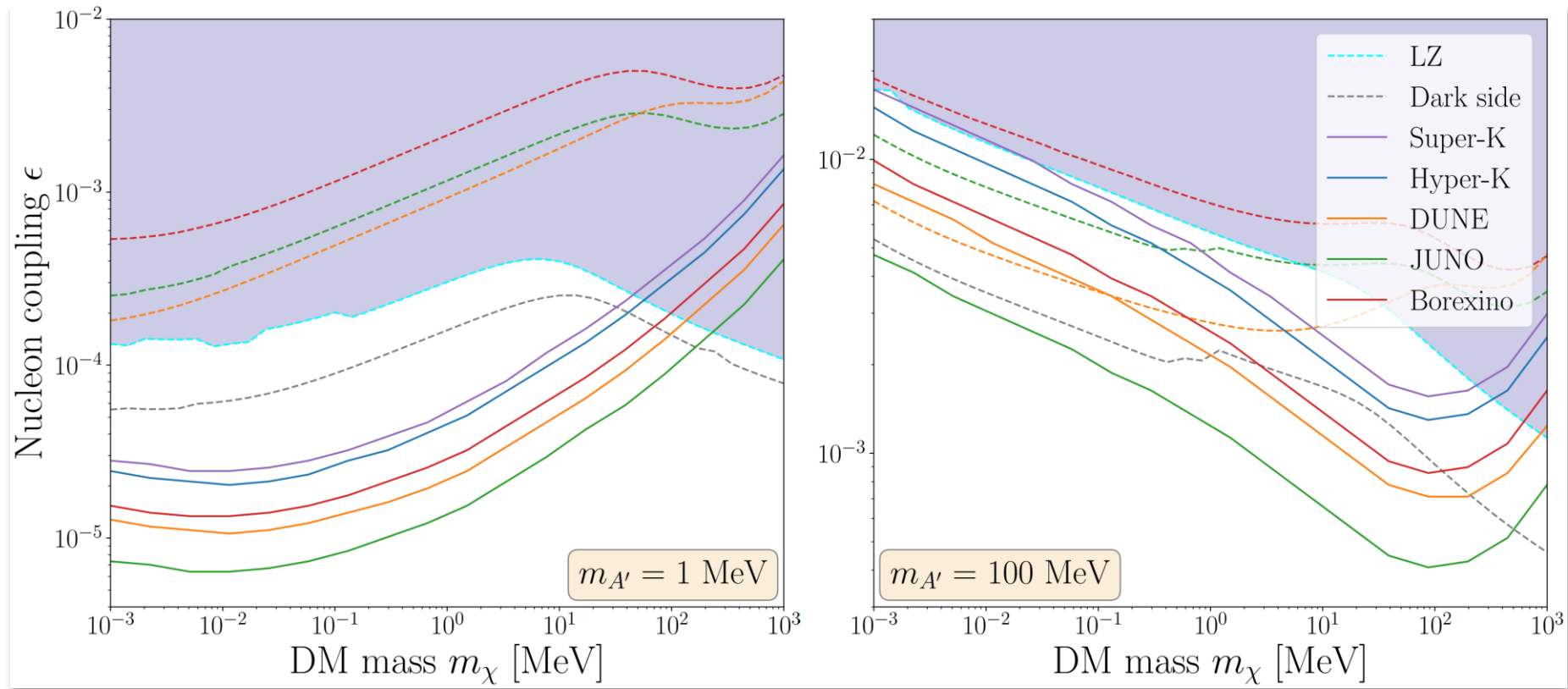


Estimated Background Rates

- ❖ The main **irreducible background**: the elastic & inelastic neutral-current scattering of **solar & atmospheric neutrinos**.

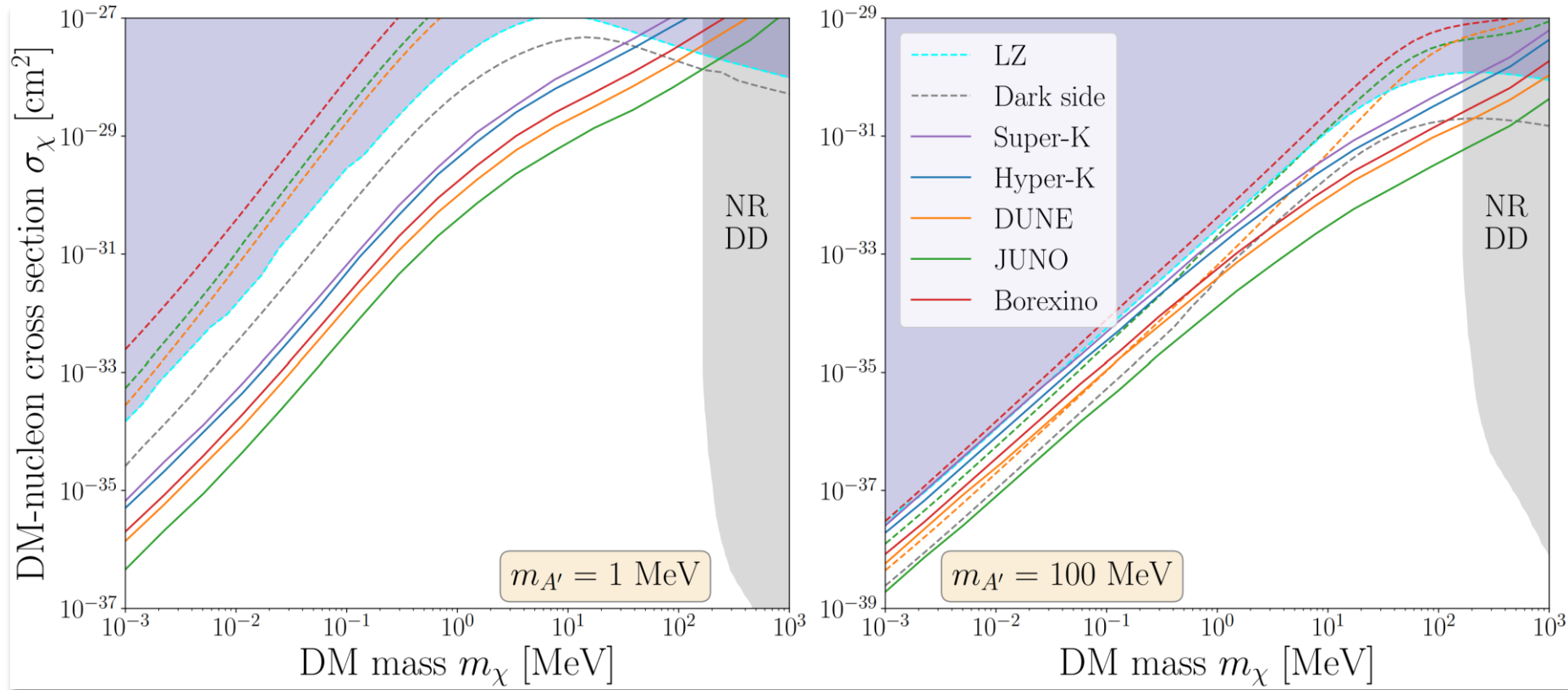


Projected/Attainable Sensitivity Reaches



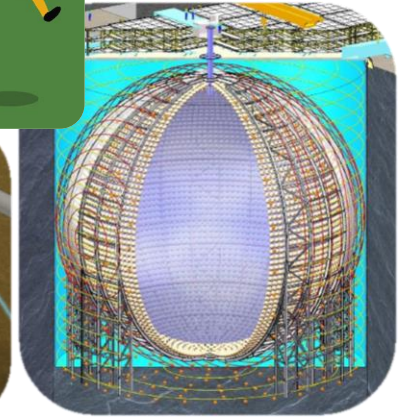
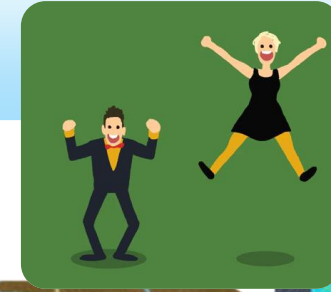
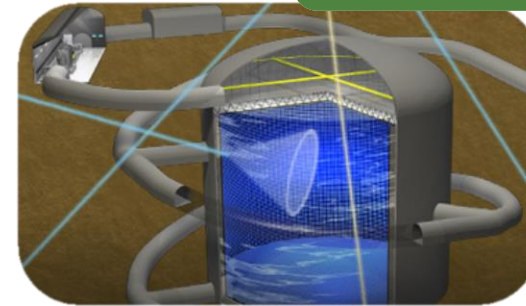
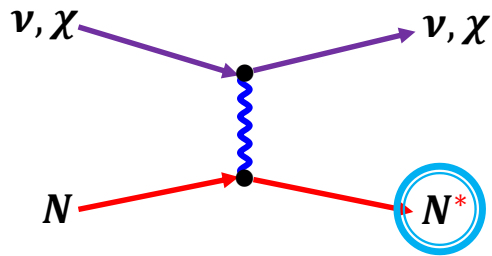
❖ Inelastic channels (solid) better than elastic channels (dashed: LZ & Borexino – data, others - future)

Projected/Attainable Sensitivity Reaches - σ_χ



❖ Inelastic channels (solid) better than elastic channels (dashed: LZ & Borexino – data, others - future)

Summary



❖ **DM-quark interactions** → **cosmic-ray boosting** of DM & **DM-nuclear scattering**.

❖ **Inelastic nuclear scattering:**

✓ **Nuclear line:** $E \sim \mathcal{O}(10)$ MeV → **No need for low threshold detectors.**

✓ **Lower background** → **Statistical advantage.**

❖ **Large volume neutrino detectors:** **wonderful to probe light DM.**

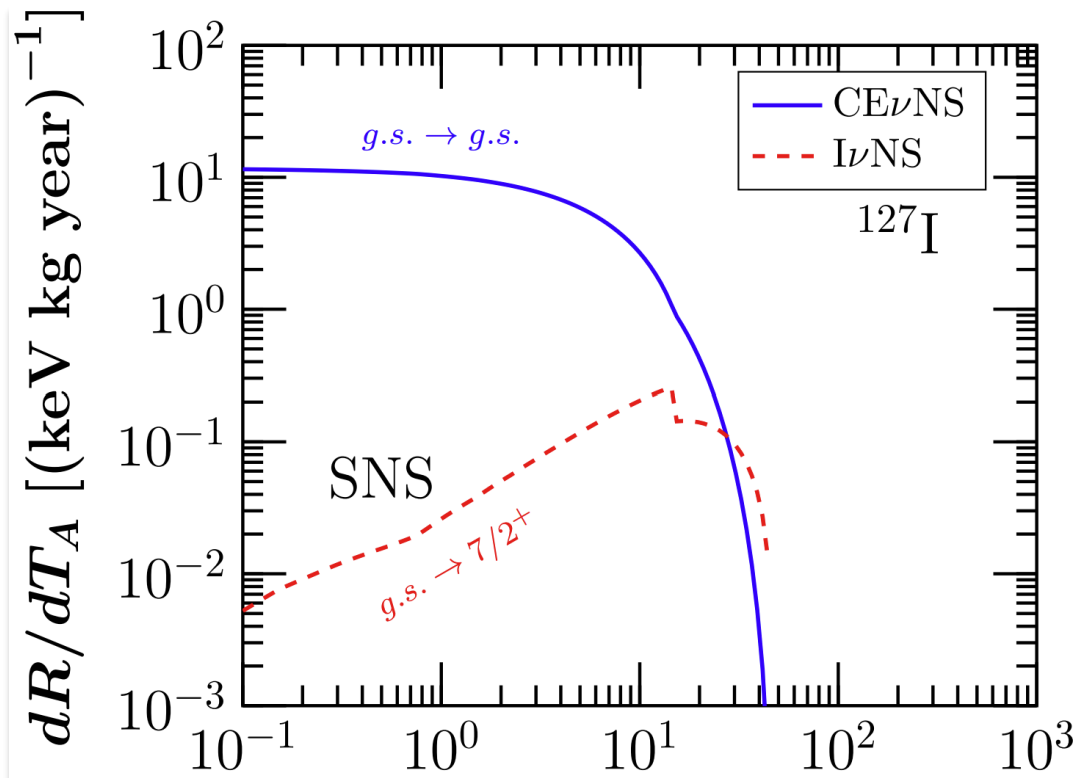
❖ **Inelastic channels:** better than elastic channels in keV – GeV range.

Thank you

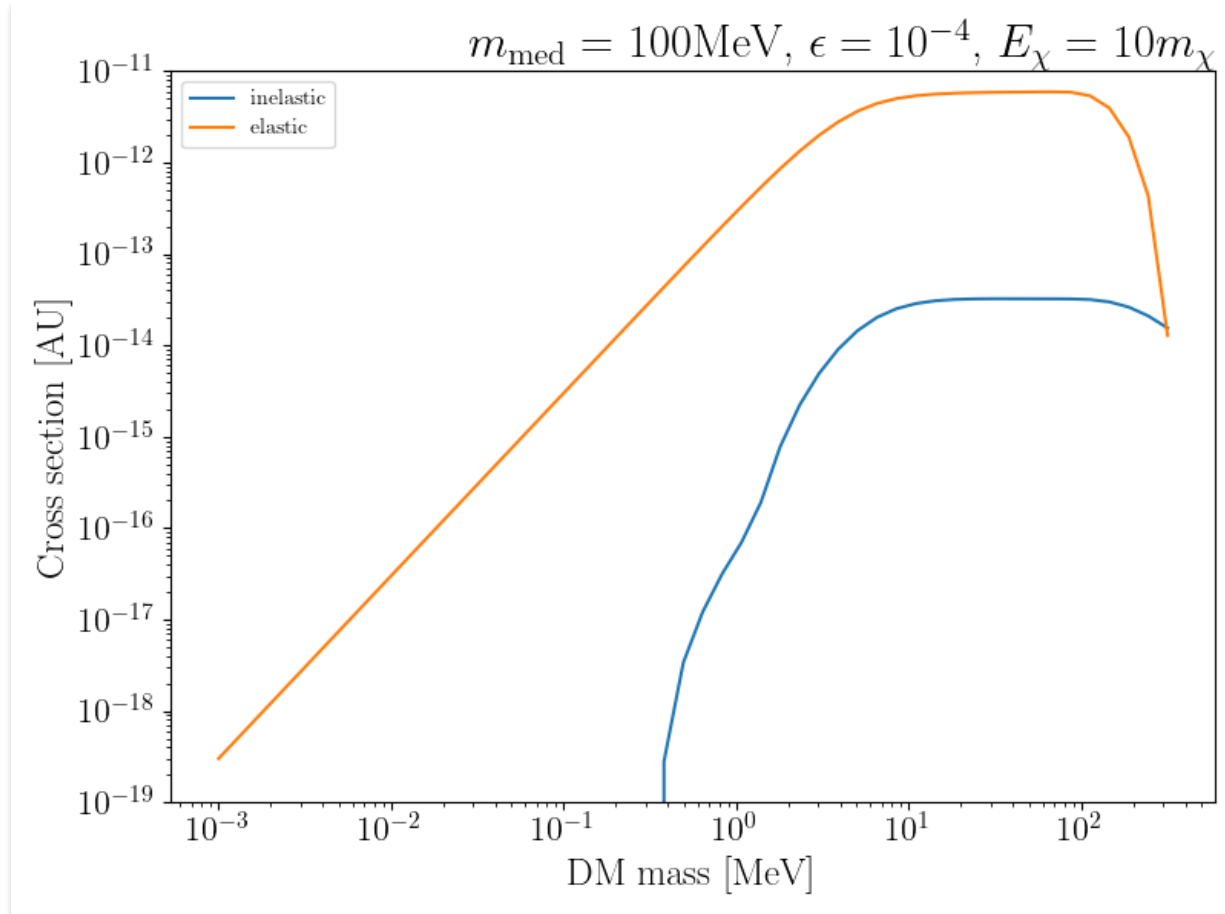


Supplemental

Elastic vs. Inelastic

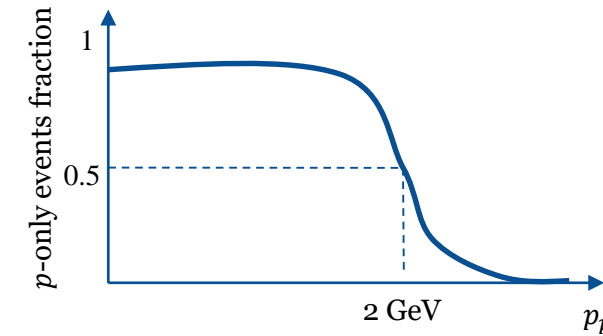
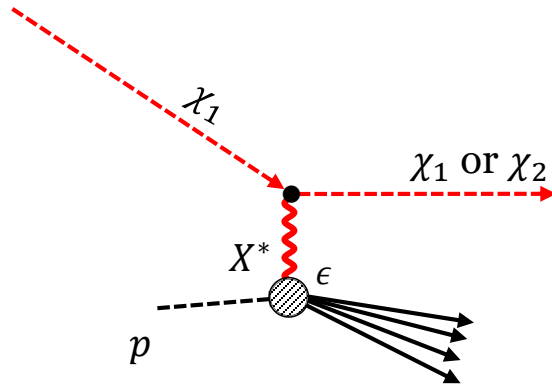


Sahu et al., [2004.04055]



p -Scattering vs. DIS

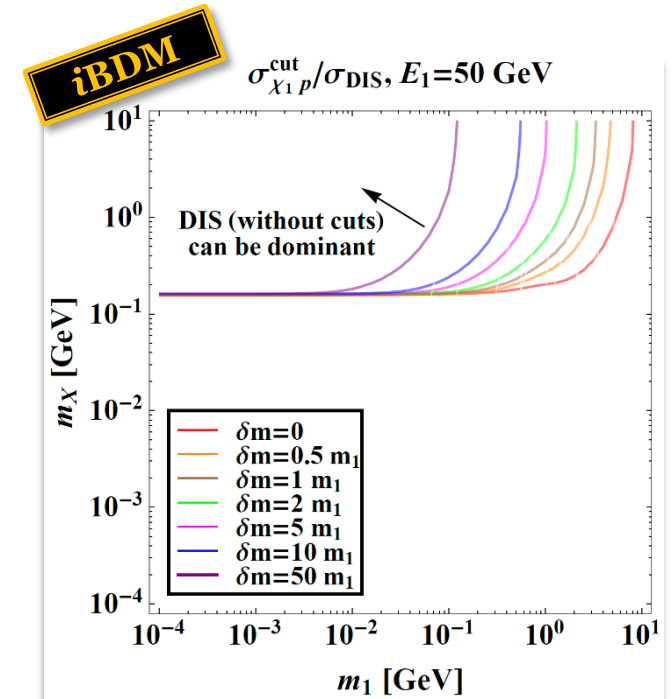
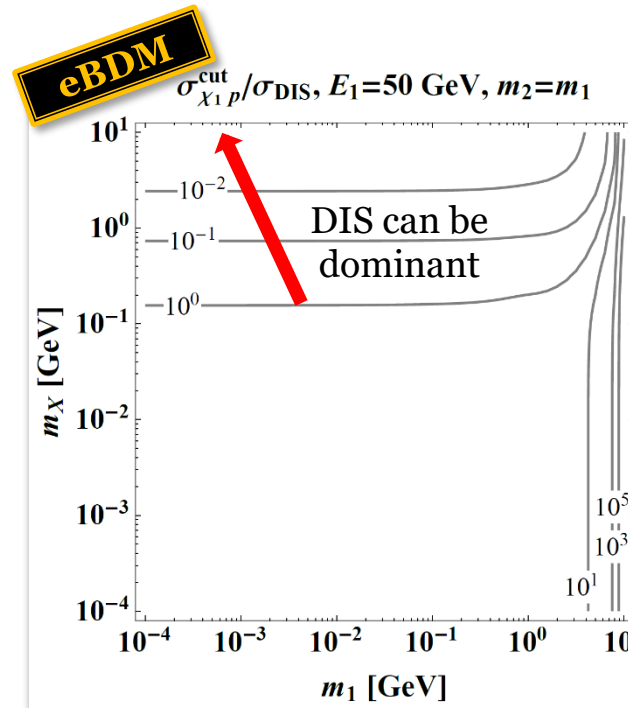
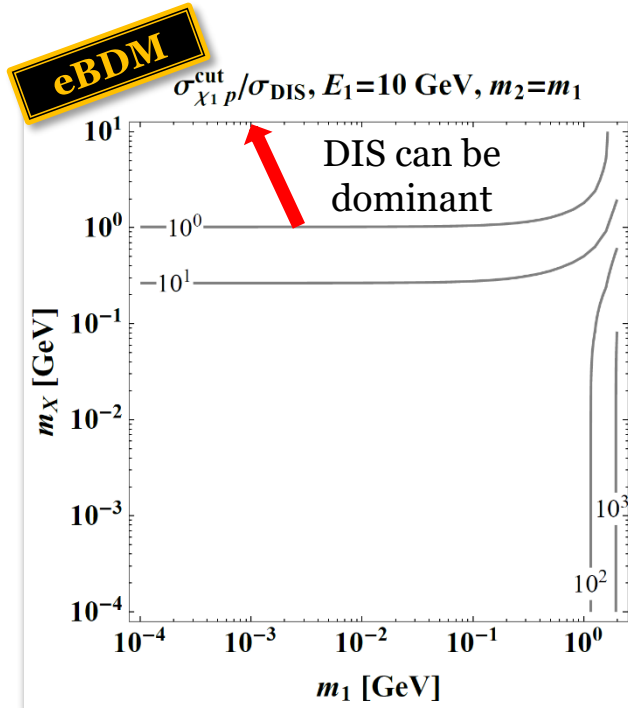
[P. Machado, D. Kim, **JCP** & S. Shin, JHEP (2020)]



- ❖ If a momentum transfer is too large, a **proton may break apart**.
- ❖ What is large? → A SK simulation study [Fechner et al, PRD (2009)] showed about 50 % events accompany (at least) a pion or a secondary particle for $p_p \approx 2$ GeV.
- ❖ We categorize any event with $p_p < 2$ GeV as the p -scattering (i.e., simplified step-function-like transition).

p -Scattering vs. DIS: Numerical Study

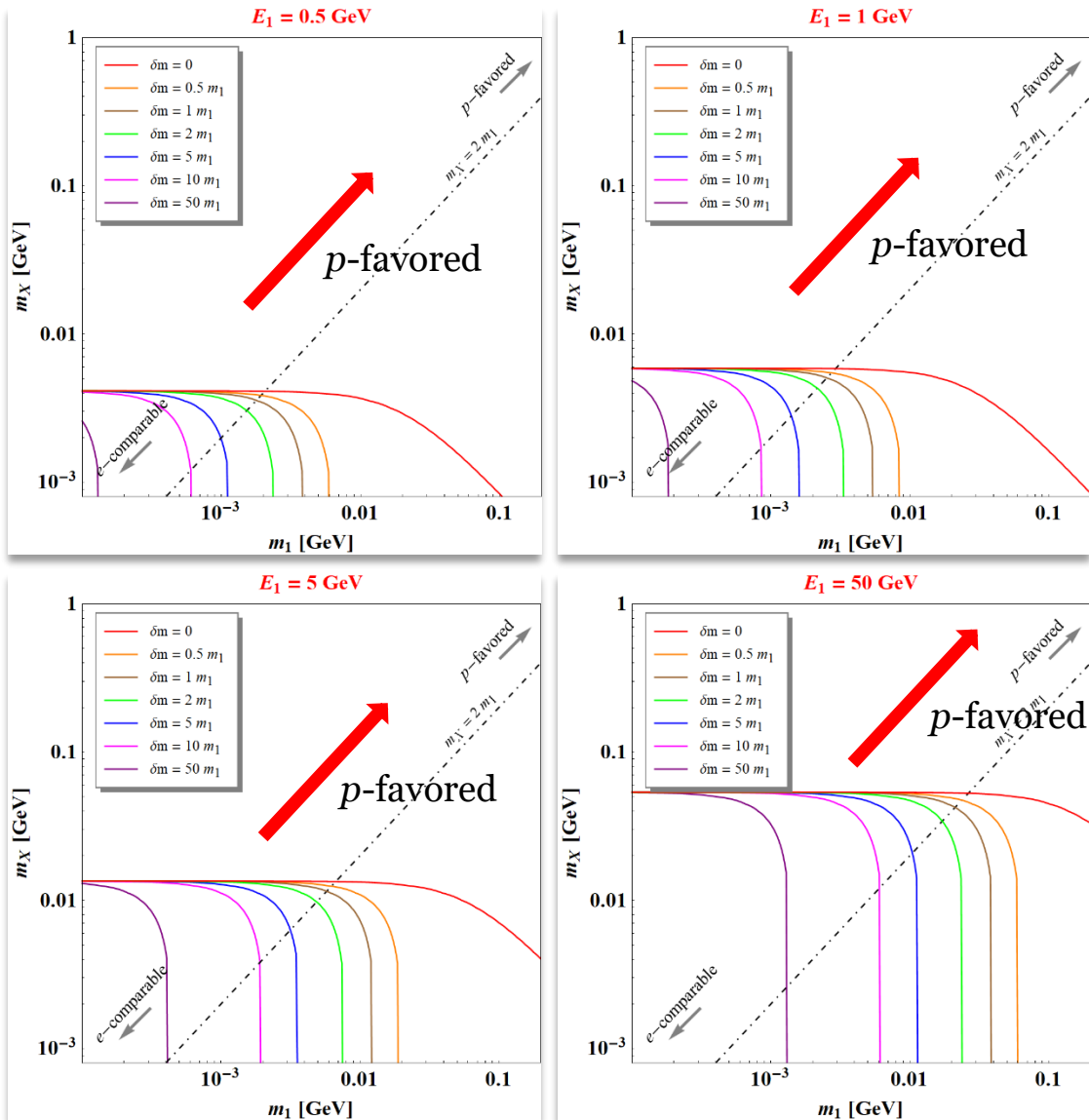
[P. Machado, D. Kim, JCP & S. Shin, JHEP (2020)]



- ✓ We study $\sigma_{\chi_1 p}^{\text{cut}}/\sigma_{\text{DIS}}$ where $200 \text{ MeV} < p_p < 2 \text{ GeV}$ is applied to $\sigma_{\chi_1 p}$ while no cuts are imposed to σ_{DIS} .
- ✓ p -scattering dominates over DIS for $m_X < O(\text{GeV})$ (cf. ν scattering via W, Z).
- ✓ As the process becomes more “inelastic”, p -scattering dominates over DIS for a given E_1 .
- ✓ DIS-preferred region expands in increasing E_1 .

p -Scattering vs. e -Scattering

[P. Machado, D. Kim, JCP & S. Shin, JHEP (2020)]



- ✓ If a BDM search hypothesizes a **heavy dark photon** (say, sub-GeV range), the p -channel may expedite discovery.
- ✓ If a model conceiving iBDM signals allows **for large mass gaps** between χ_1 and χ_2 , the p -channel is more advantageous.
- ✓ The e -channel becomes comparable in probing the parameter regions **with smaller m_1 and m_X** .
- ✓ As the boosted χ_1 comes **with more energy**, more parameter space where the e -channel is comparable opens up.
- ✓ **With cuts**, more e -channel favored region.

Many More Well-Motivated Exps.

[P. Machado, D. Kim, JCP & S. Shin, JHEP (2020)]

Dark Matter Experiments	Target	Volume [t]		Depth	E_{th}	Resolution			PID	Run Time	Refs.
	Material	Active	Fiducial	[m]	[keV]	Position [cm]	Angular [°]	Energy [%]			
DarkSide-50	LAr	46.4	36.9	3,800	$\mathcal{O}(1)$	$\sim 0.1 - 1$	-	$\lesssim 10$	-	2013-	[112]
DarkSide-20k	LAr	23	20	3,800	$\mathcal{O}(1)$	$\sim 0.1 - 1$	-	$\lesssim 10$	-	goal: 2021-	[79]
XENON1T	LXe	2.0	1.3	3,600	$\mathcal{O}(1)$	$\sim 0.1 - 1$	-	-	-	2016-2018	[113, 114]
XENONnT	LXe	5.9	~ 4	3,600	$\mathcal{O}(1)$	$\sim 0.1 - 1$	-	-	-	goal: 2020-	[113]
DEAP-3600	SP LAr	3.26	2.2	2,000	$\mathcal{O}(10)$	< 10	-	$\sim 10 - 20$	-	2016-	[99-101]
DEAP-50T	SP LAr	150	50	2,000	$\mathcal{O}(10)$	15	-	-	-	-	[99]
LUX-ZEPLIN	LXe	7	5.6	1,500	$\mathcal{O}(1)$	$\sim 0.1 - 1$	-	2.5 MeV: 2	-	goal: 2020-	[115, 116]
Neutrino Experiments	Target	Volume [kt]		Depth	E_{th}	Resolution			PID	Run Time	Refs.
Material	Active	Fiducial	[m]	[MeV]	Vertex [cm]	Angular [°]	Energy [%]				
Borexino	organic LS	0.278	0.1	3,800	~ 0.2	$\sim 9-17$	-	$\frac{5}{\sqrt{E}(\text{MeV})}$	-	> 5.6 year	[117]
KamLAND	LS	1	0.2686	1,000	0.2-1	$\frac{12-13}{\sqrt{E}(\text{MeV})}$	-	$\frac{6.4-6.9}{\sqrt{E}(\text{MeV})}$	-	~ 10 year?	[118, 119]
JUNO	LS	-	20	700	goal: 0.1	$\frac{12}{\sqrt{E}(\text{MeV})}$	μ^- : $L > 5 \text{ m}: < 1,$ $L > 1 \text{ m}: < 10$	$\frac{3}{\sqrt{E}(\text{MeV})}$	μ^\pm vs π^\pm , e^\pm vs π^0 ; difficult	goal: 2021-	[120-122]
DUNE	LArTPC	Total: 17.5 $\times 4$	≥ 10 $\times 4$	1500	$e: 30,$ $p: \lesssim 1-2$ 21-50	$\leq 1-2$	$e, \mu: 1,$ $\pi^\pm, p, n: 5$	$e: 20 (E < 0.4 \text{ GeV}),$ $10 (E < 1.0 \text{ GeV}),$ $2 + \frac{8}{\sqrt{E/\text{GeV}}} (E > 1.0 \text{ GeV})$ $p: 10 (E < 1.0 \text{ GeV}),$ $5 + \frac{5}{\sqrt{E/\text{GeV}}} (E > 1.0 \text{ GeV})$	good e, μ, π^\pm, p separation	10 kt: 2026- 20 kt: 2027-	[77, 80-84]
SK	Water Cherenkov	Total: 50	22.5	1,000	$e: 5,$ $p: 485$	5 MeV: 95, 10 MeV: 55, 20 MeV: 40	10 MeV: 25, 0.1 GeV: 3, 1.33 GeV: 1.2	10 MeV: 16, 1 GeV: 2.5	$e, \mu:$ good	≥ 14 year	[123-125]
HK	Water Cherenkov	Total: 258 $\times 2$	187 $\times 2$	650	$e: < 5,$ $p: 485$	5 MeV: 75, 10 MeV: 45, 15 MeV: 40, 0.5 GeV: 28	similar to SK	better than SK	$e, \mu:$ good, π^0, π^\pm ; mild	goal: 2027-	[85-87]
Neutrino Telescopes	Target	Effective Volume [Mt]		Depth	E_{th}	Resolution			PID	Run Time	Refs.
Material	Volume [Mt]	[m]	[GeV]	Vertex [m]	Angular [°]	Energy [%]					
IceCube	Ice	100 GeV: $\sim 30,$ 200 GeV: ~ 200	1,450	~ 100	vertical: 5, horizontal: 15	μ -track: $\sim 1,$ shower: ~ 30	100 GeV: 28, 1 TeV: 16	only μ	2011- (2008)	[89, 126]	
DeepCore	Ice	10 GeV: $\sim 5,$ 100 GeV: ~ 30	2,100	~ 10	better	μ -track: $\sim 1,$ shower: ≥ 10	-	only μ	2011- (2010)	[88, 89]	
IceCube Upgrade	Ice	-	2,150	$\mathcal{O}(1)$	much better	5 GeV: $\sim 20,$ 10 GeV: ~ 15	-	only μ	goal: 2023	[127]	
PINGU	Ice	1 GeV: $\geq 1,$ 10 GeV: ~ 5	2,100	~ 1	much better	1 GeV: 25, 10 GeV: 10	1 GeV: 55, 10 GeV: 25	only μ	$>$ 2023	[128, 129]	
Gen2	Ice	~ 10 Gt	1,360	~ 50	worse	μ -track: < 1 shower: ~ 15	-	only μ	-	[130]	

- ❖ Many existing/upcoming experiments are **potentially capable of testing models conceiving BDM signals.**
- ❖ **Additional physics opportunity on top of the main mission** of each experiment.

Many More Well-Motivated Exps.

[P. Machado, D. Kim, JCP & S. Shin, JHEP (2020)]

Detectors are
complementary to
 one another **rather**
than superior to the
 other!

