

PROBING THE IN- AND OUT-OF-EQUILIBRIUM

CHIRAL MAGNETIC EFFECTS WITH

LATTICE **QCD**

XVITH QUARK CONFINEMENT AND THE HADRON SPECTRUM 2024
CAIRNS, AUSTRALIA

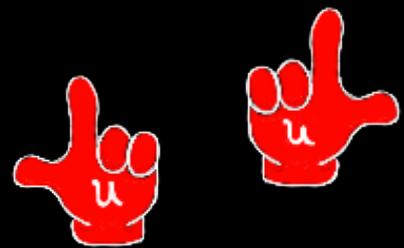
DEAN VALOIS

dvalois@physik.uni-bielefeld.de

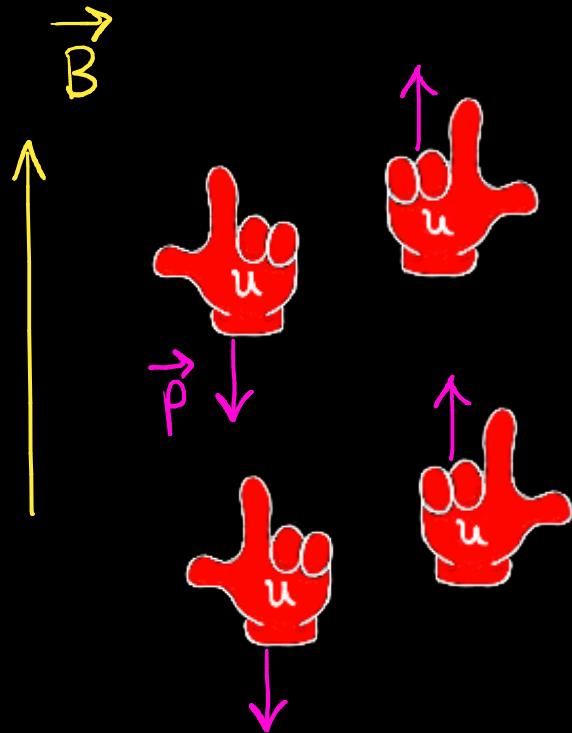
IN COLLABORATION WITH

BASTIAN BRANDT , GERGELY ENDRÓDI
EDUARDO GARNACHO , GERGELY MARKÓ

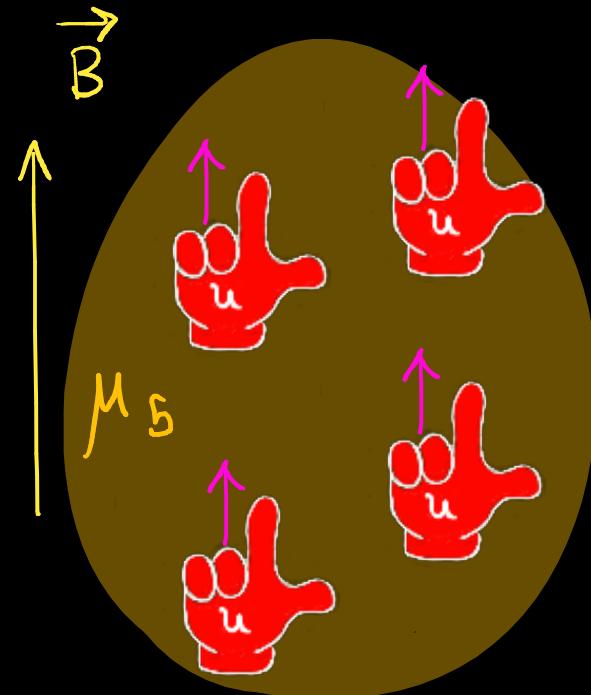
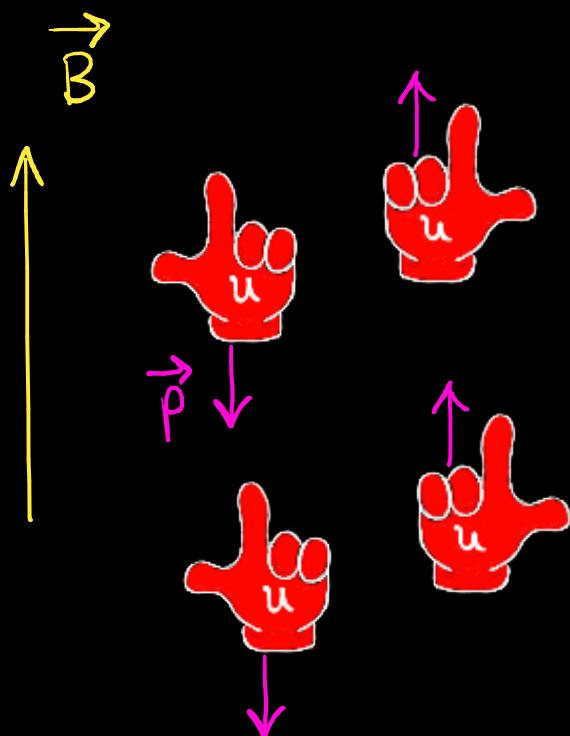
WHAT IS THE CHIRAL MAGNETIC EFFECT?



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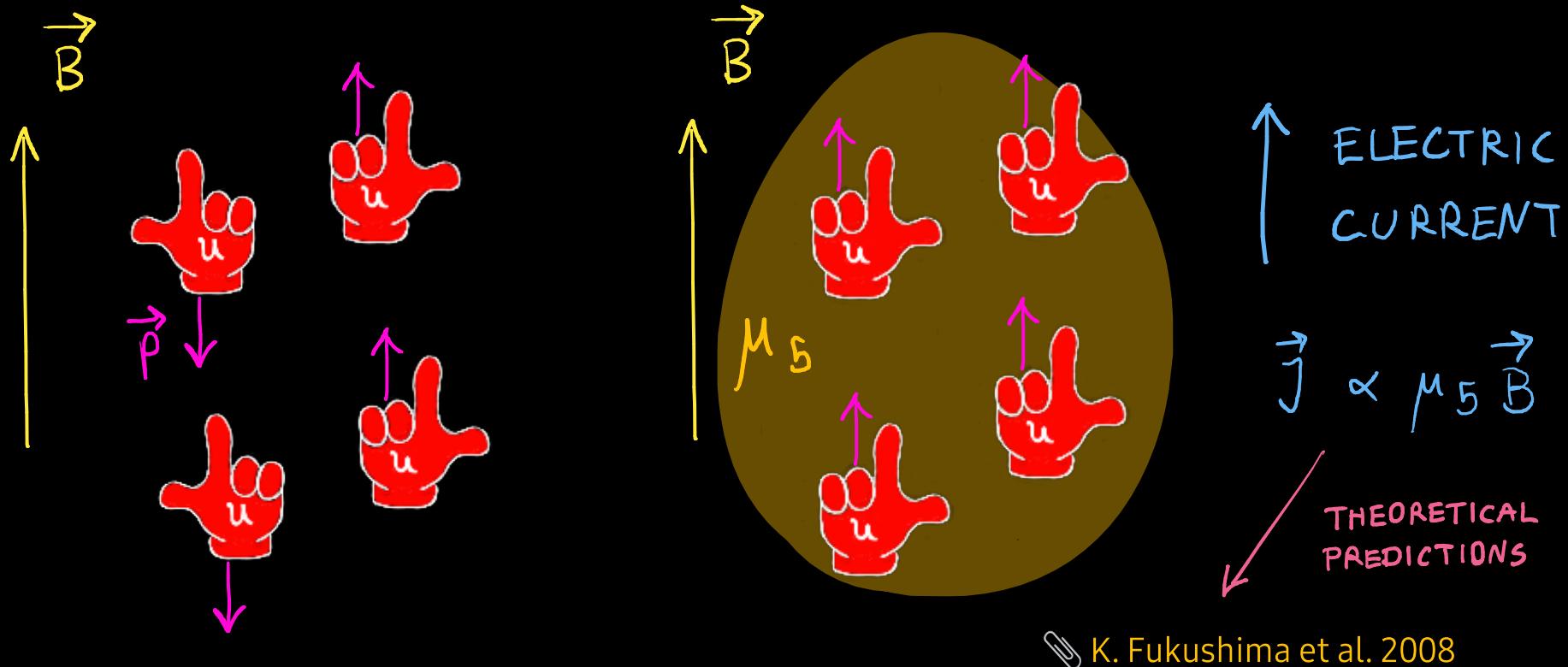


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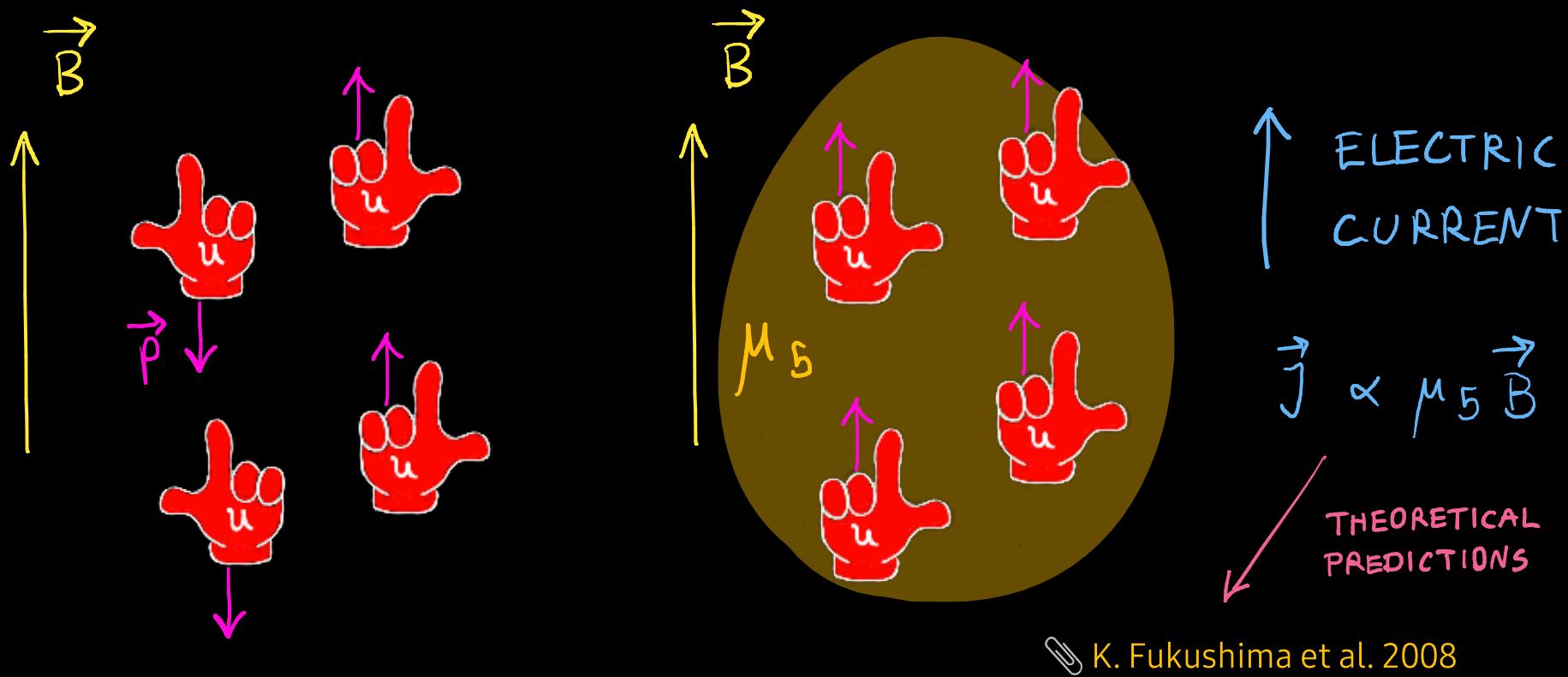


ELECTRIC CURRENT
 $\vec{j} \propto \mu_5 \vec{B}$

WHAT IS THE CHIRAL MAGNETIC EFFECT?



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IN CONDENSED MATTER :

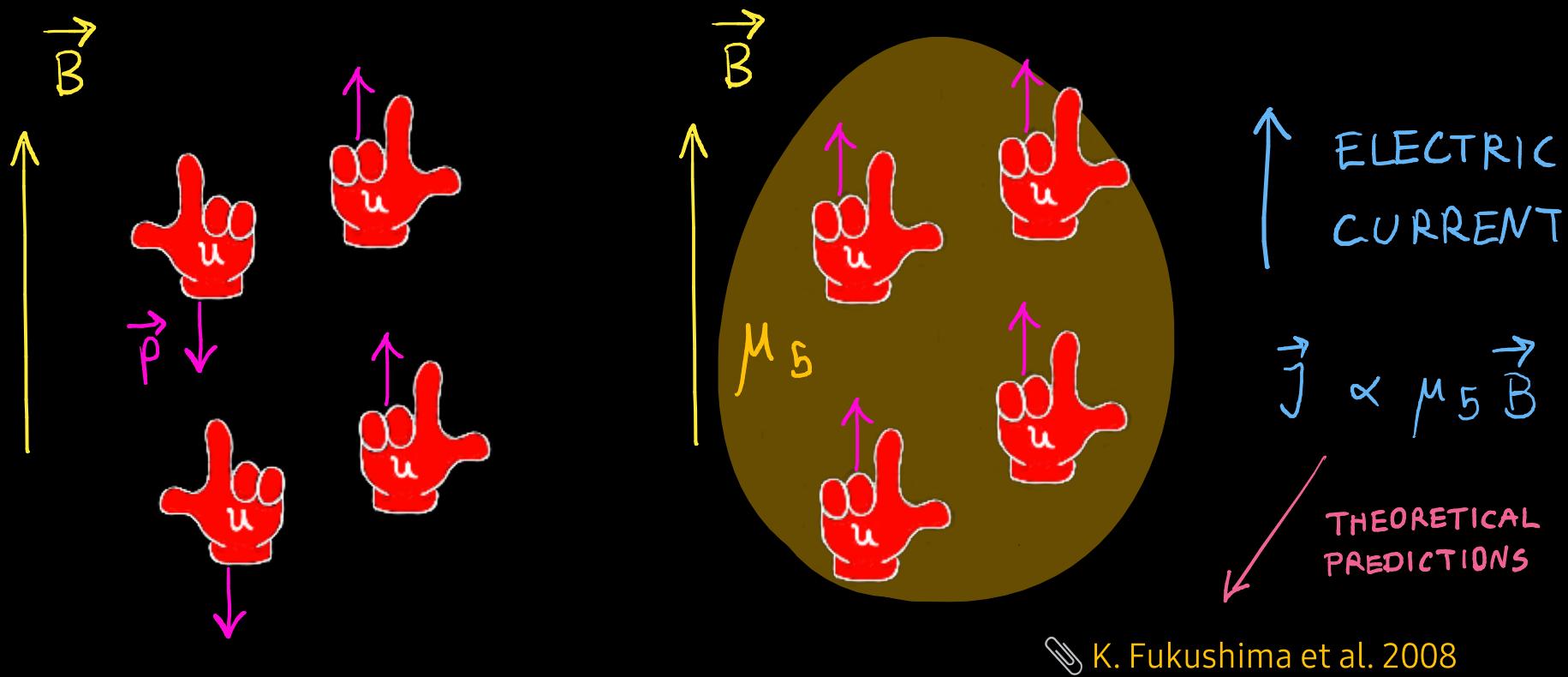
Zr_xTe₅ Q. Li, D. Kharzeev et al. 2014

Na₃Bi J. Xiong, S. K. Kushwaha et al. 2015

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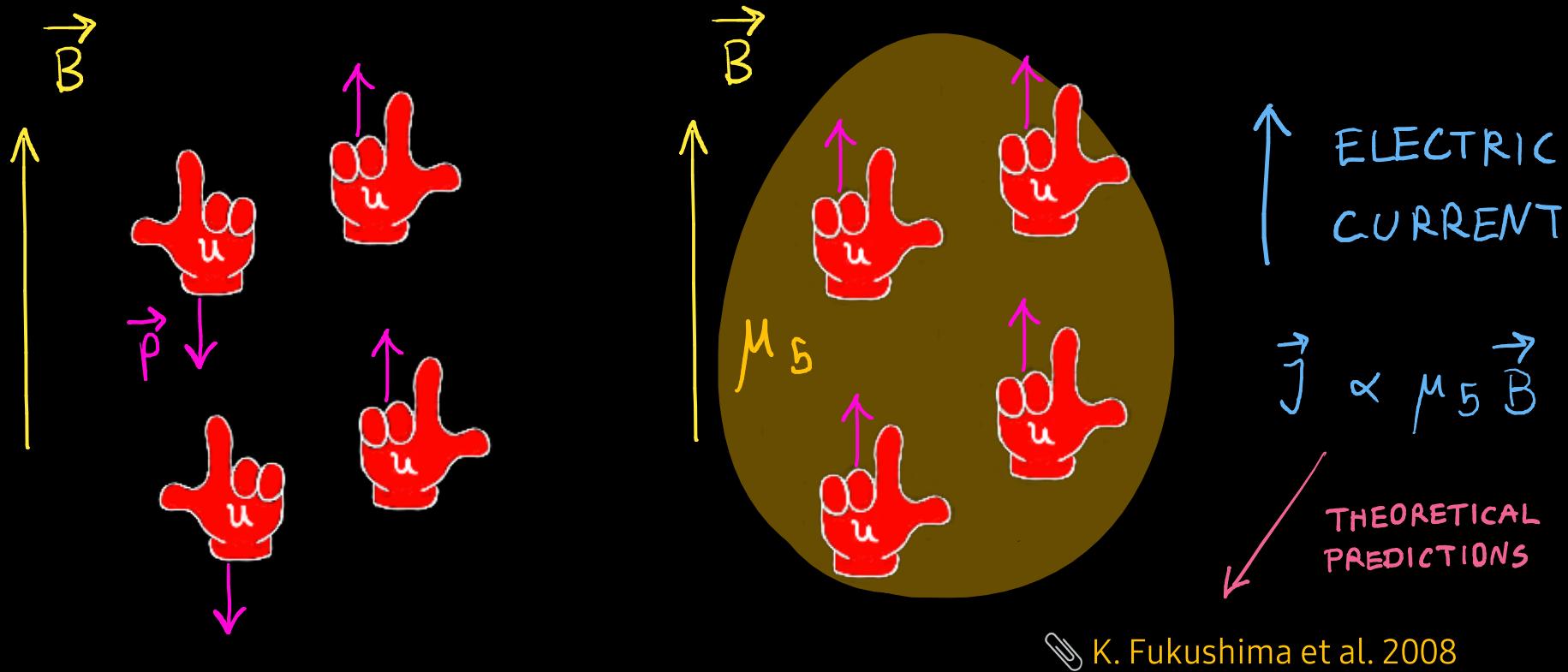
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- NON-DISSIPATIVE

WHAT IS THE CHIRAL MAGNETIC EFFECT?



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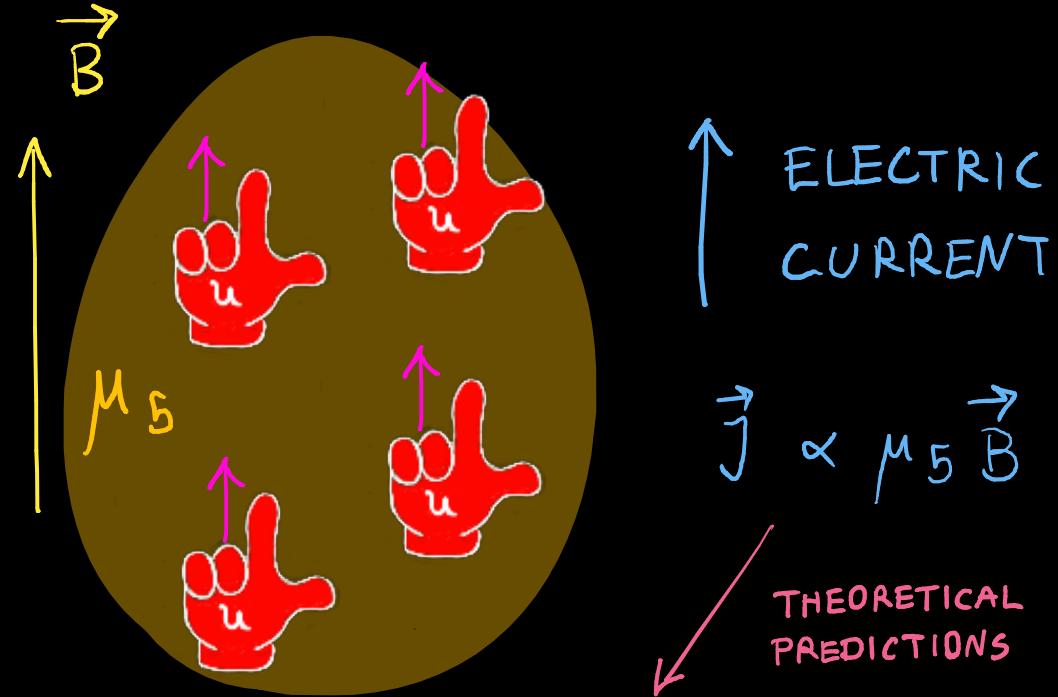
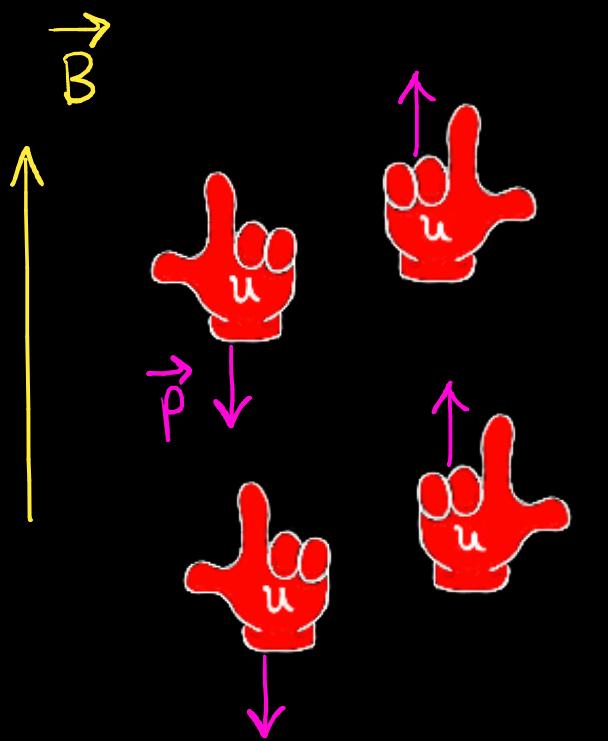
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- NON-DISSIPATIVE
- TOPOLOGICALLY PROTECTED

WHAT IS THE CHIRAL MAGNETIC EFFECT?



📎 K. Fukushima et al. 2008

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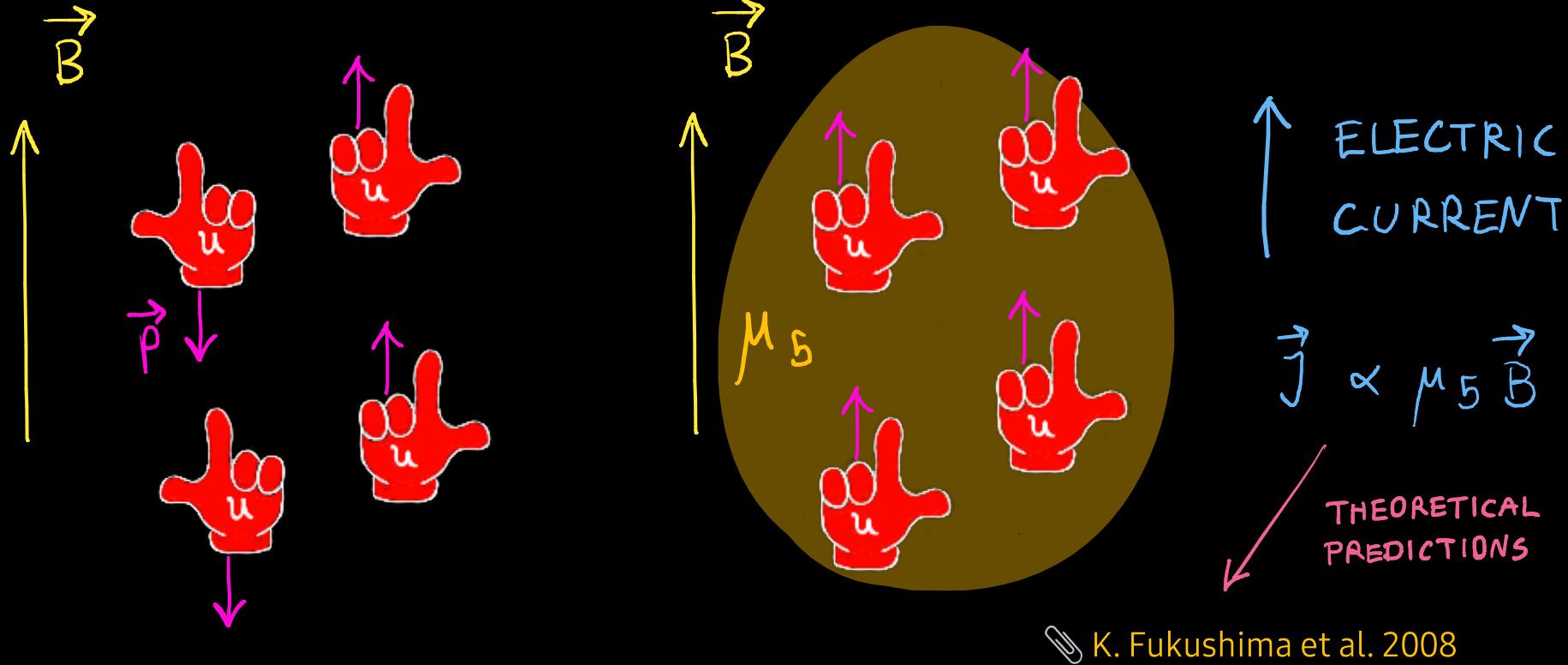
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- NON-DISSIPATIVE
- TOPOLOGICALLY PROTECTED
- IS THERE A CME IN **QCD** ?

WHAT IS THE CHIRAL MAGNETIC EFFECT?



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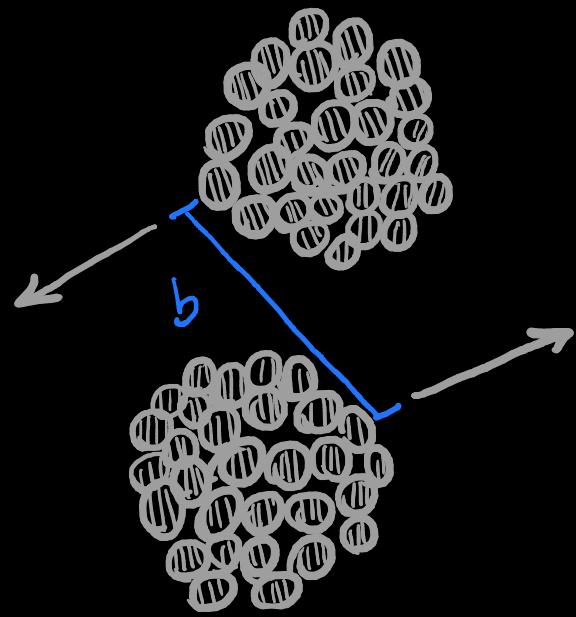
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- NON-DISSIPATIVE
- TOPOLOGICALLY PROTECTED
- IS THERE A CME IN ?

EVIDENCE FOR CME IN HIC

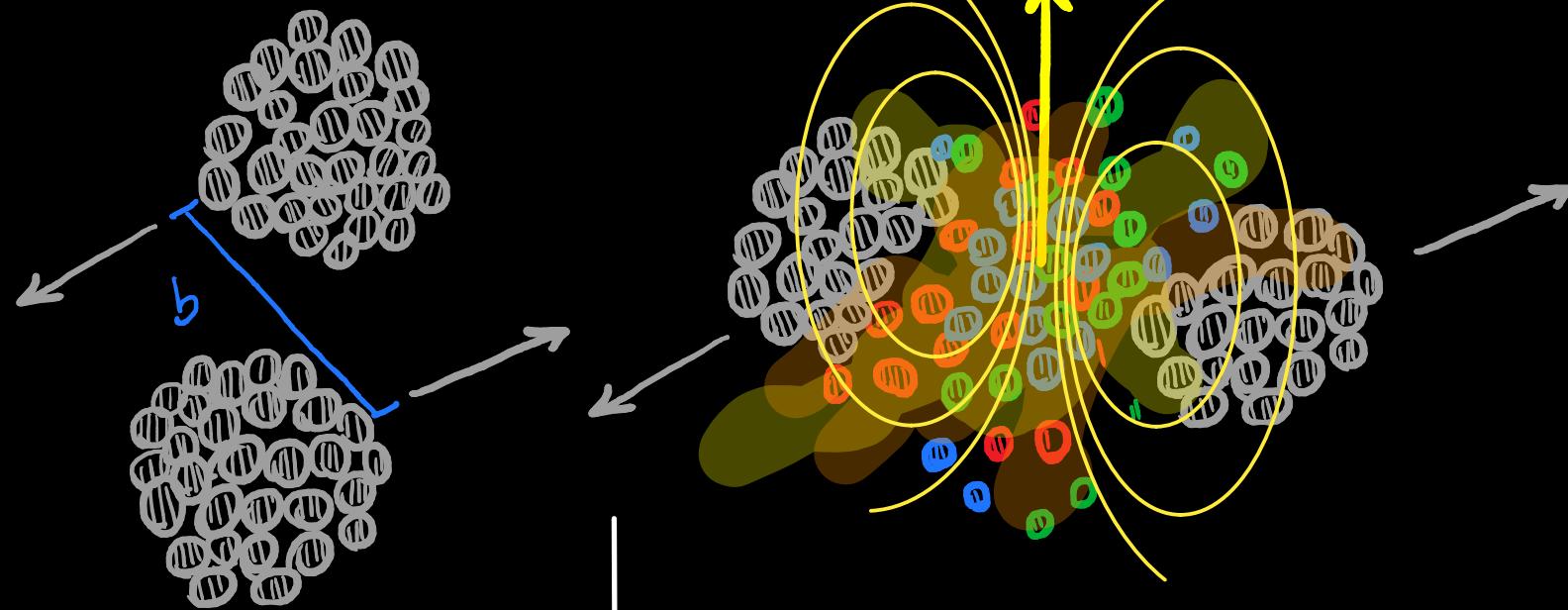
B. I. Abelev et al. (STAR Collaboration) 2009

HEAVY-ION COLLISIONS (HIC)



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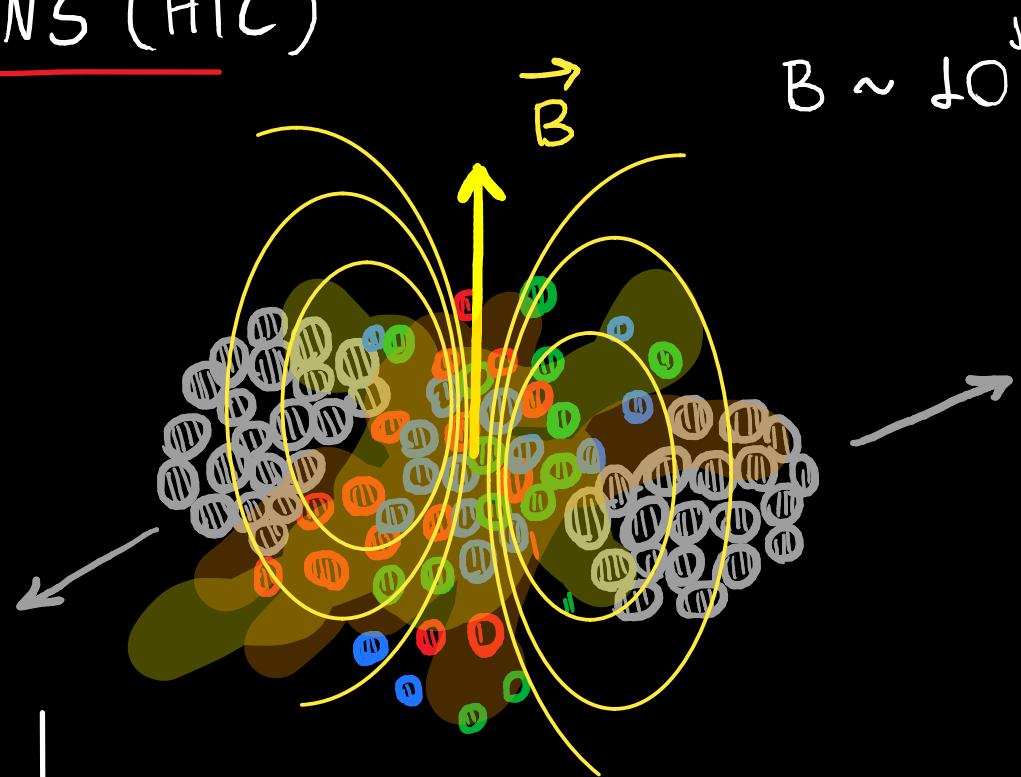
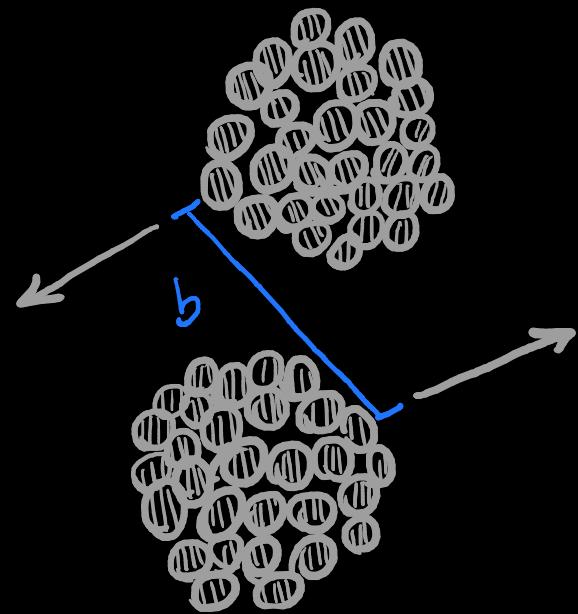
$$B \sim 10^{18} \text{ G}$$



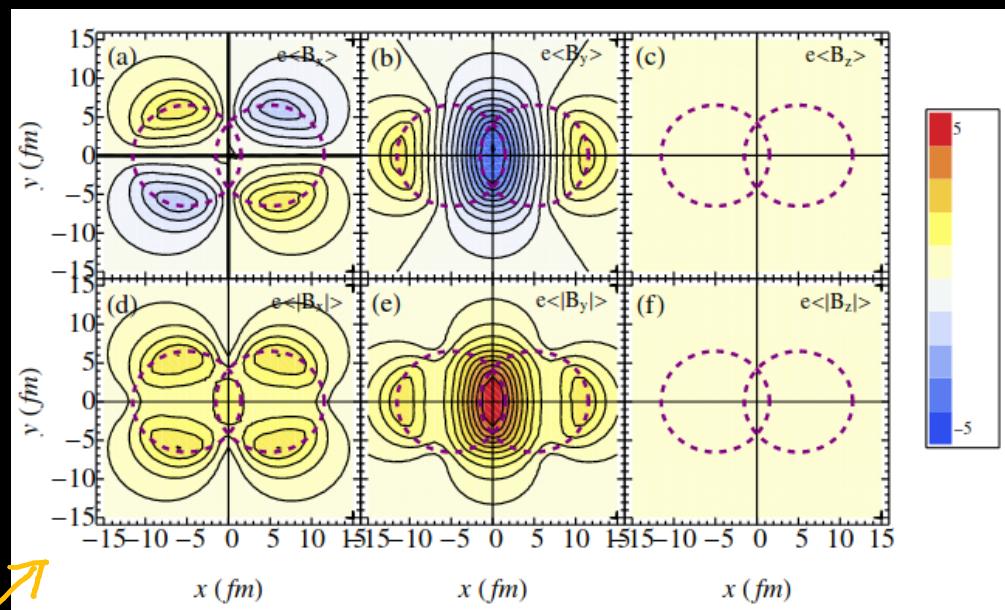
- RECENT RESULTS BY STAR

📎 M. Abdulhamid et al. 2024

HEAVY-ION COLLISIONS (HIC)



- RECENT RESULTS BY STAR
📎 M. Abdulhamid et al. 2024
- MAGNETIC FIELDS IN HIC HAVE HIGHLY NON-TRIVIAL GEOMETRIES



OUTLINE

- BRIEF OVERVIEW OF CME
- PART I - EQUILIBRIUM CME WITH NON-UNIFORM \vec{B}
- PART II - OUT-OF-EQUILIBRIUM CME WITH UNIFORM \vec{B}
- SUMMARY & CONCLUSIONS

CME IN EQUILIBRIUM

$$j_3 = C_{\text{CME}} \mu_5 B$$

CME IN EQUILIBRIUM

CME CONDUCTIVITY ↗

$$\downarrow \vec{j}_3 = C_{\text{CME}} \mu_5 B$$

VECTOR CURRENT

CME IN EQUILIBRIUM

CME CONDUCTIVITY ↗

$$\text{↗ } j_3 = C_{\text{CME}} \mu_5 B$$

VECTOR CURRENT

INITIAL EXPECTATIONS: $C_{\text{CME}} = \frac{1}{2\pi^2}$

CME IN EQUILIBRIUM

CME CONDUCTIVITY ↗

$$\text{vector current } \downarrow j_3 = C_{\text{CME}} \mu_5 B$$

VECTOR CURRENT

INITIAL EXPECTATIONS: $C_{\text{CME}} = \frac{1}{2\pi^2}$

HOWEVER: $C_{\text{CME}} = 0$

CME IN EQUILIBRIUM

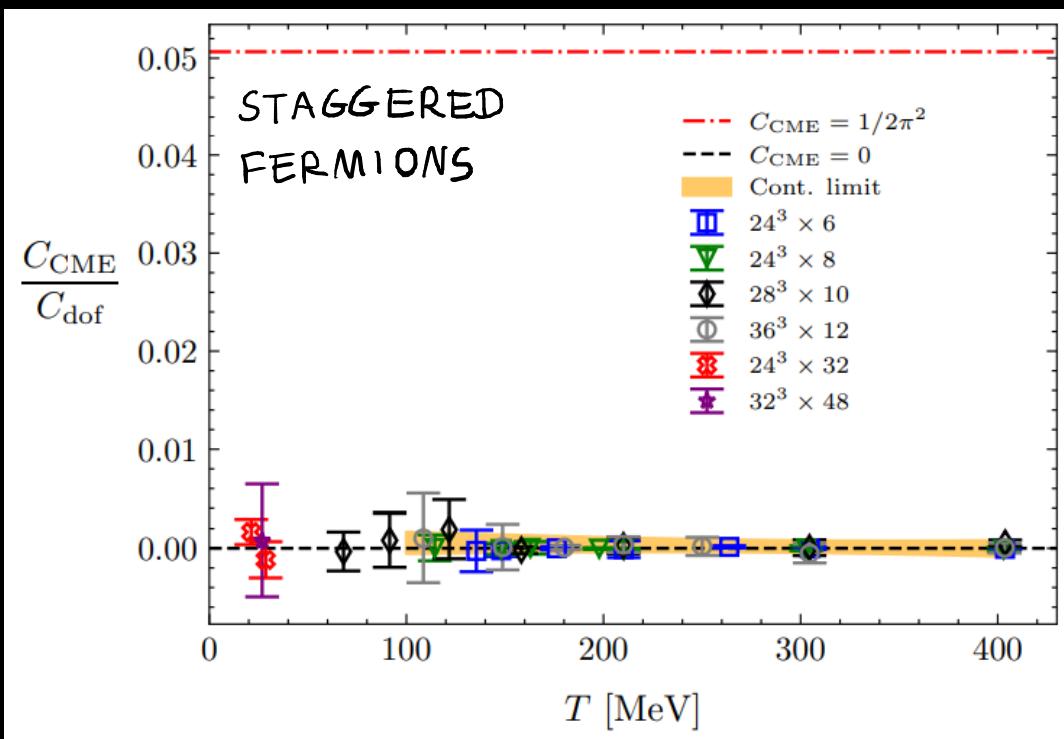
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📎 B. Brandt, G. Endrődi, E.
Garnacho-Velasco, G. Markó 2024

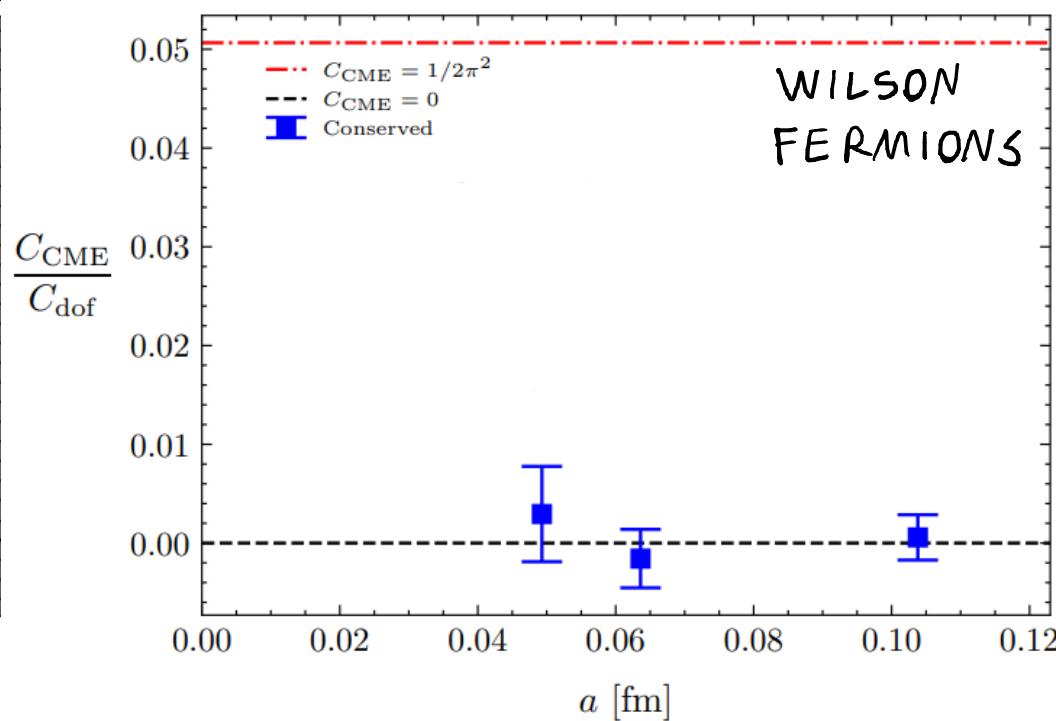
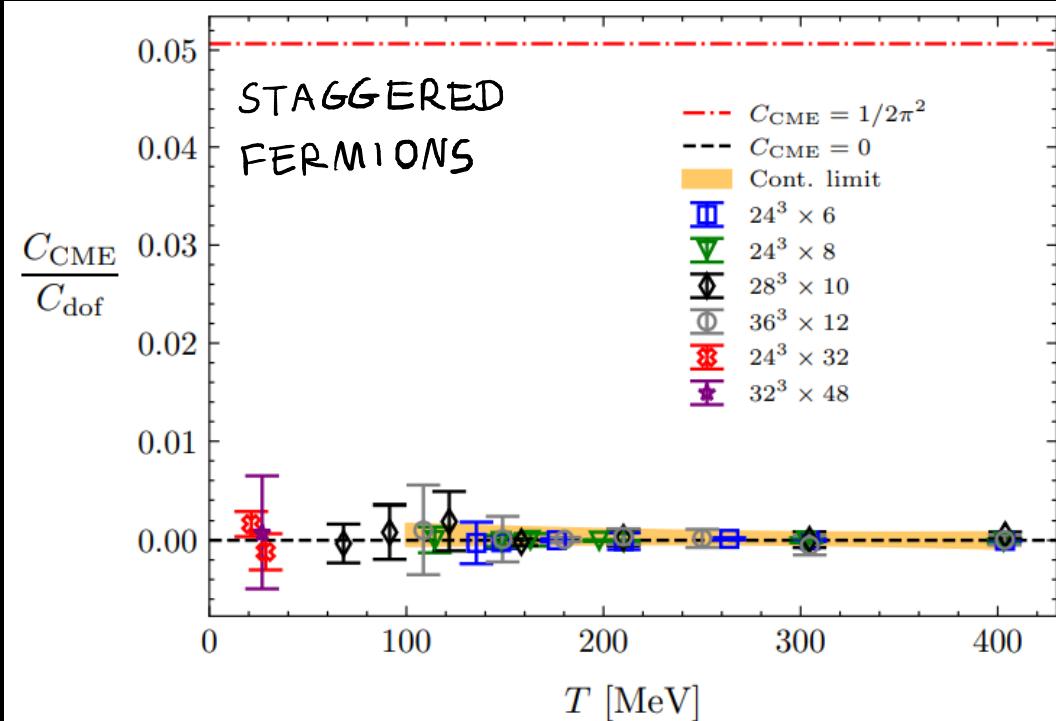
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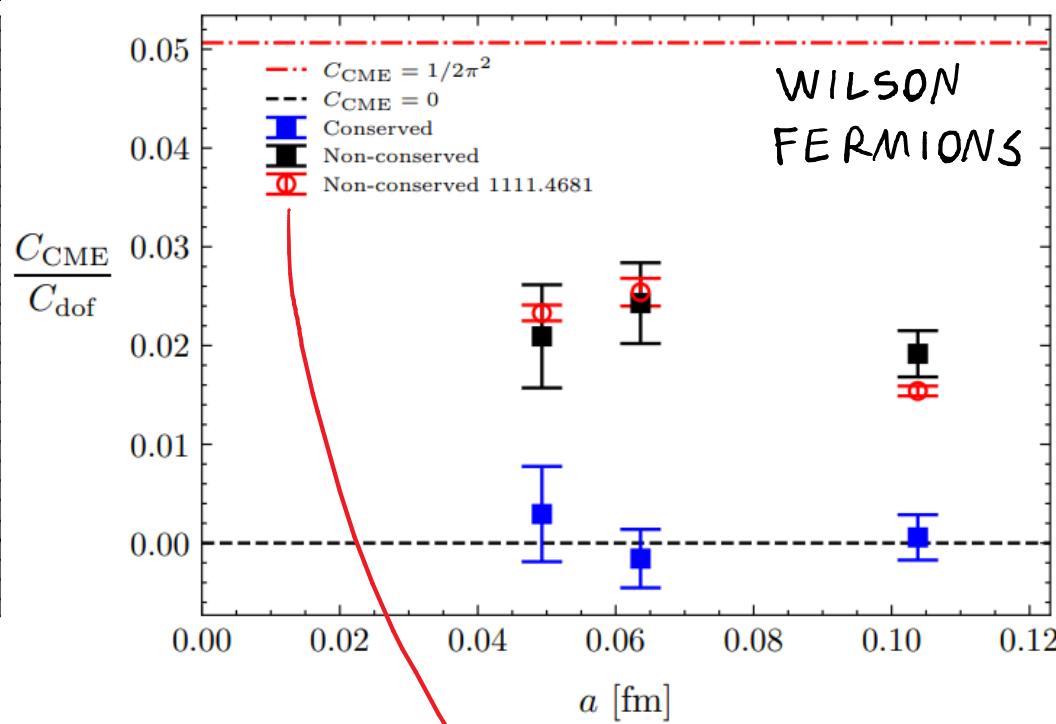
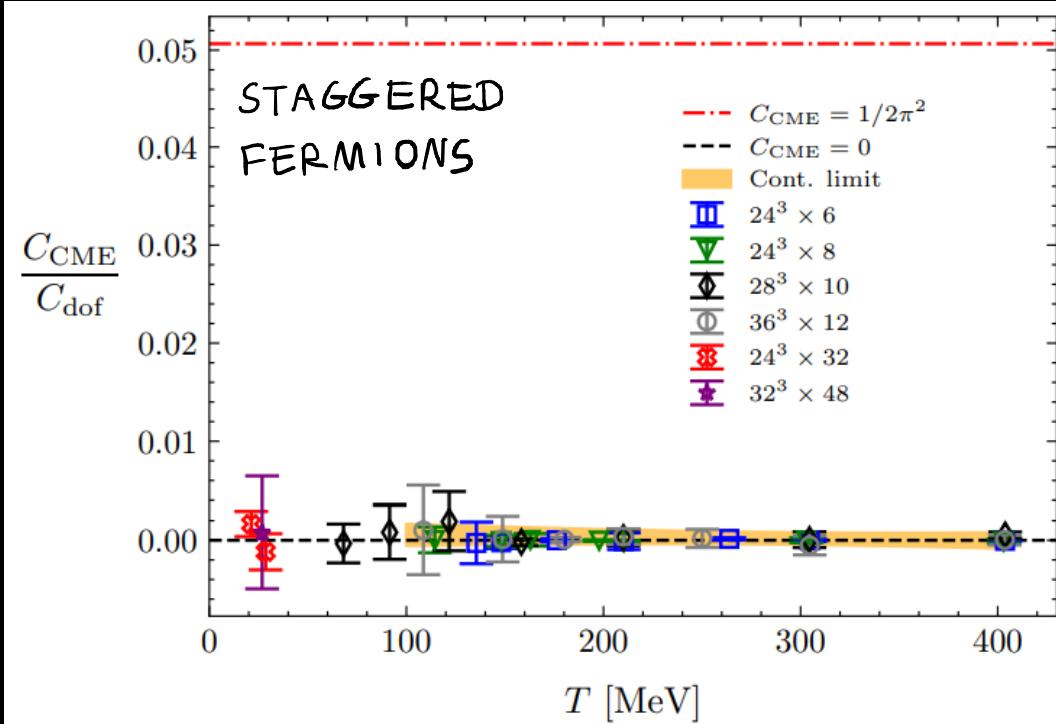
CME CONDUCTIVITY ↗

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📎 B. Brandt, G. Endrődi, E.
Garnacho-Velasco, G. Markó 2024

📎 A. Yamamoto 2011

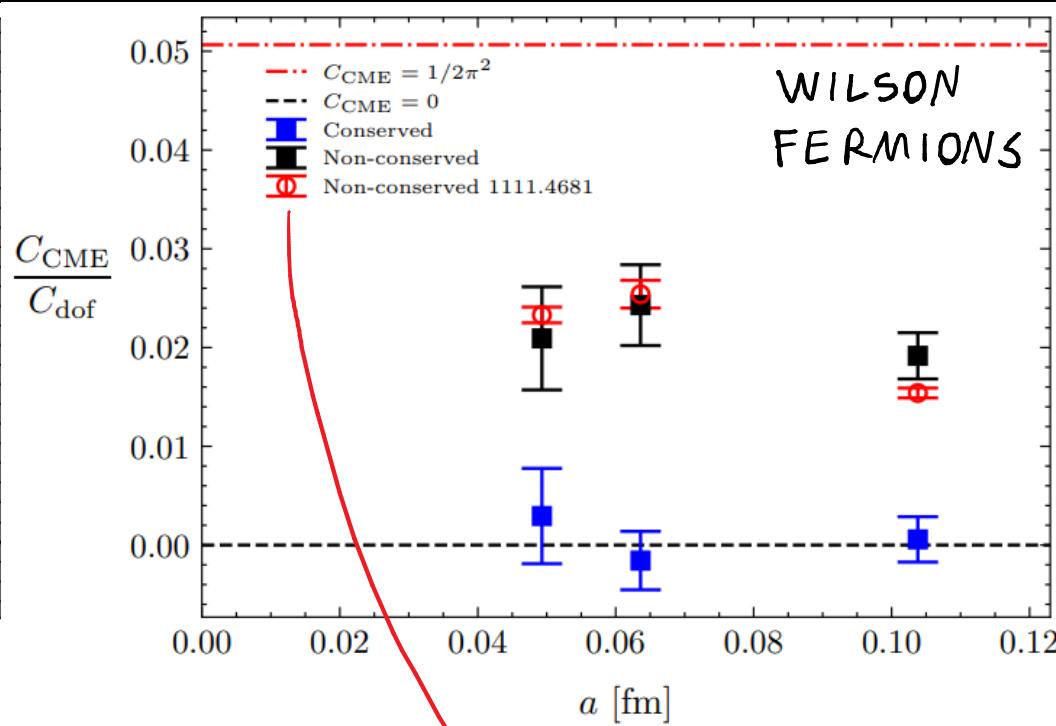
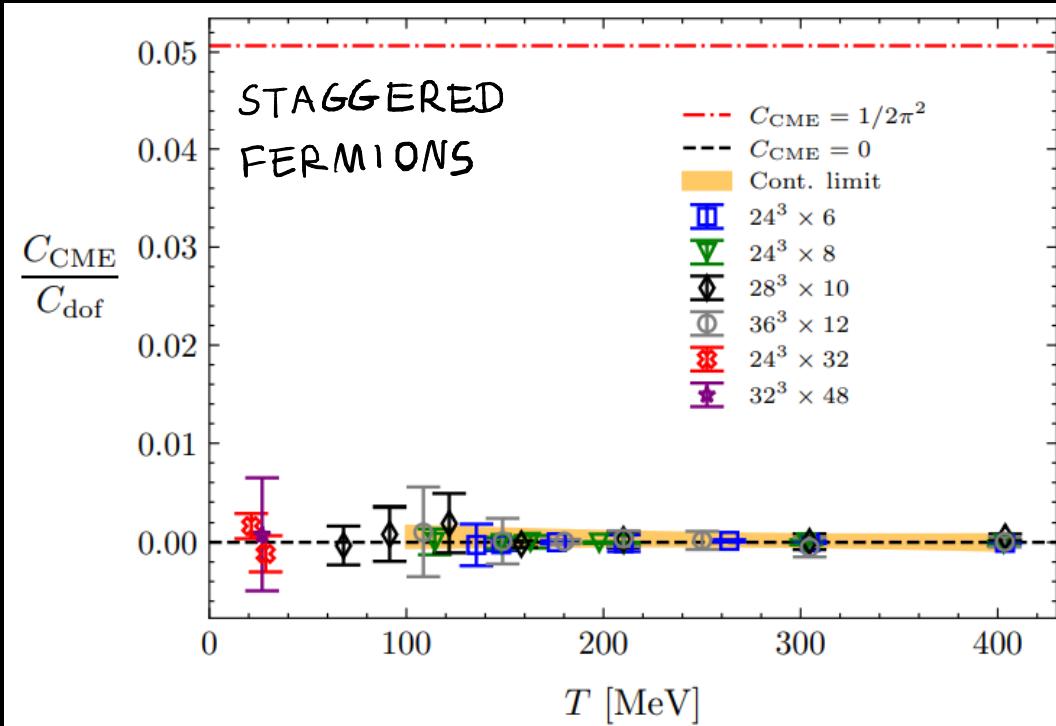
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CME CONDUCTIVITY \rightarrow

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📎 B. Brandt, G. Endrődi, E.
Garnacho-Velasco, G. Markó 2024

📎 A. Yamamoto 2011

CONCLUSION: FOR THE CORRECT CME, j_3 HAS TO BE CONSERVED
ON THE LATTICE!

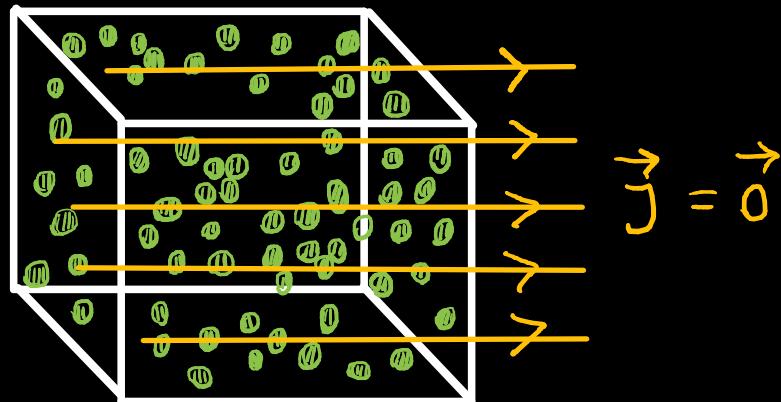
CME & BLOCH'S THEOREM

- IN QUANTUM MECHANICS :

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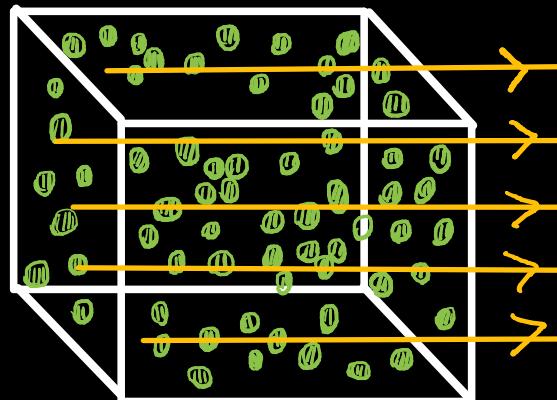
TOTAL CONSERVED CURRENT VANISHES ON THE GROUND STATE
IN EQUILIBRIUM



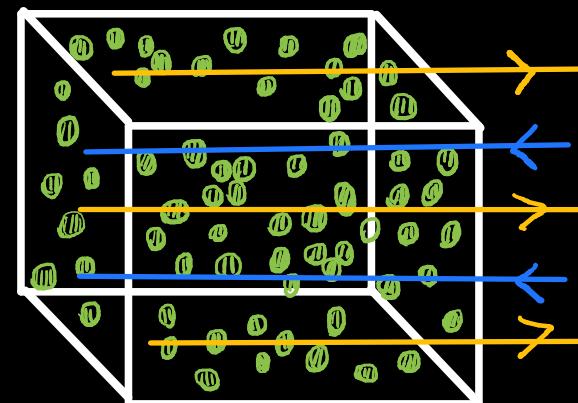
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TOTAL CONSERVED CURRENT VANISHES ON THE GROUND STATE
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$$\vec{j} = \vec{0}$$



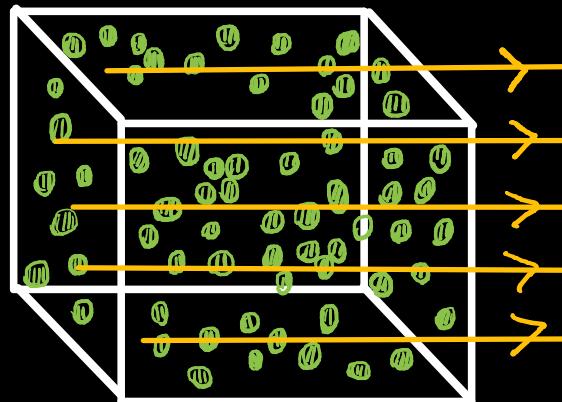
$$\vec{j}(\vec{x})$$

LOCAL CURRENTS
ARE ALLOWED

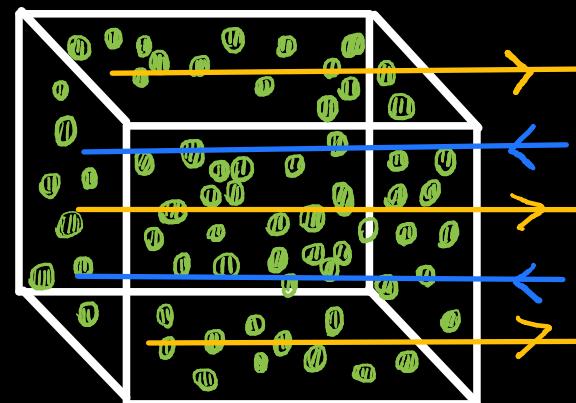
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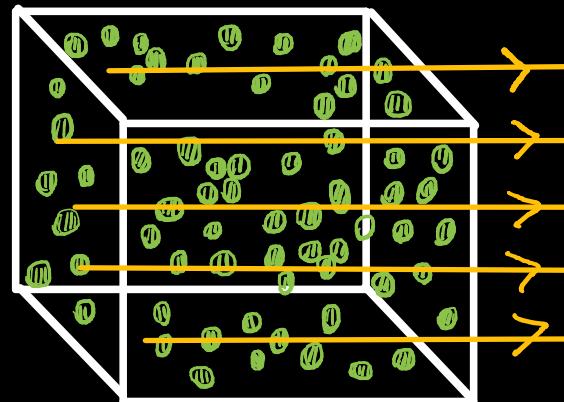
$$\int_{\vec{x}} \vec{J}(\vec{x}) = \vec{J} = \vec{0}$$

LOCAL CURRENTS
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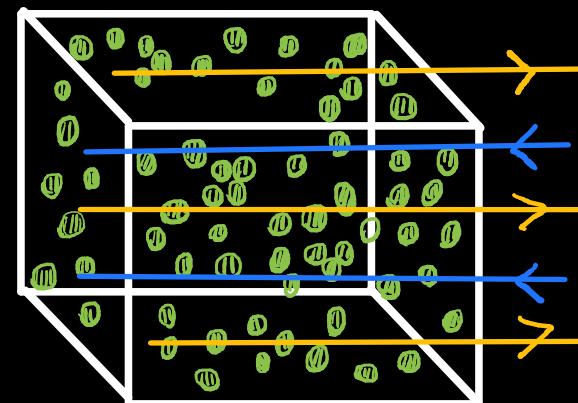
CME & BLOCH'S THEOREM

- IN QUANTUM MECHANICS :

TOTAL CONSERVED CURRENT VANISHES ON THE GROUND STATE
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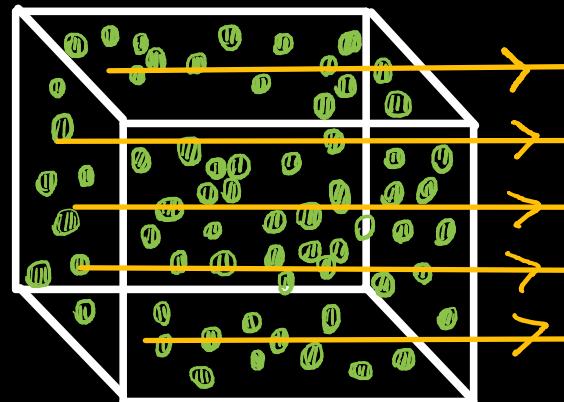
- IN QFT :

Naoki Yamamoto 2015 \Rightarrow CME = 0 IN EQUILIBRIUM

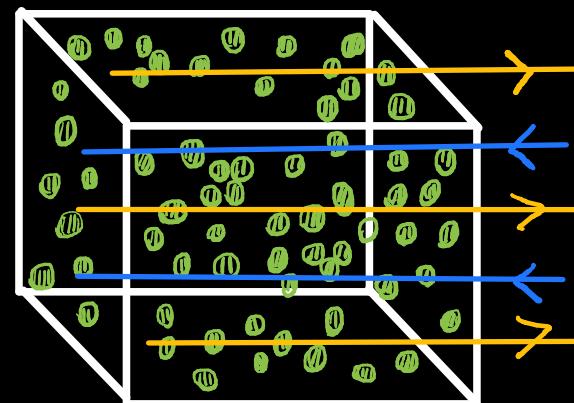
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LOCAL CURRENTS
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- IN QFT :

Naoki Yamamoto 2015 \Rightarrow CME = 0 IN EQUILIBRIUM

WHAT GENERATES A NON-TRIVIAL LOCAL $j(\vec{x})$?

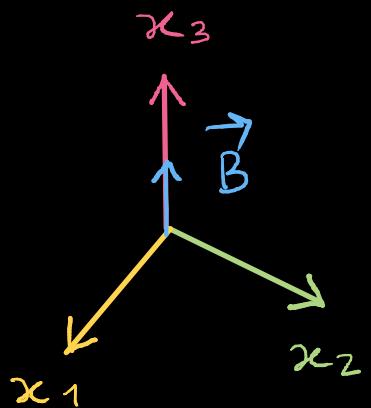
PART I -

EQUILIBRIUM CME WITH
NON-UNIFORM \vec{B}

NON - UNIFORM \vec{B} MODEL

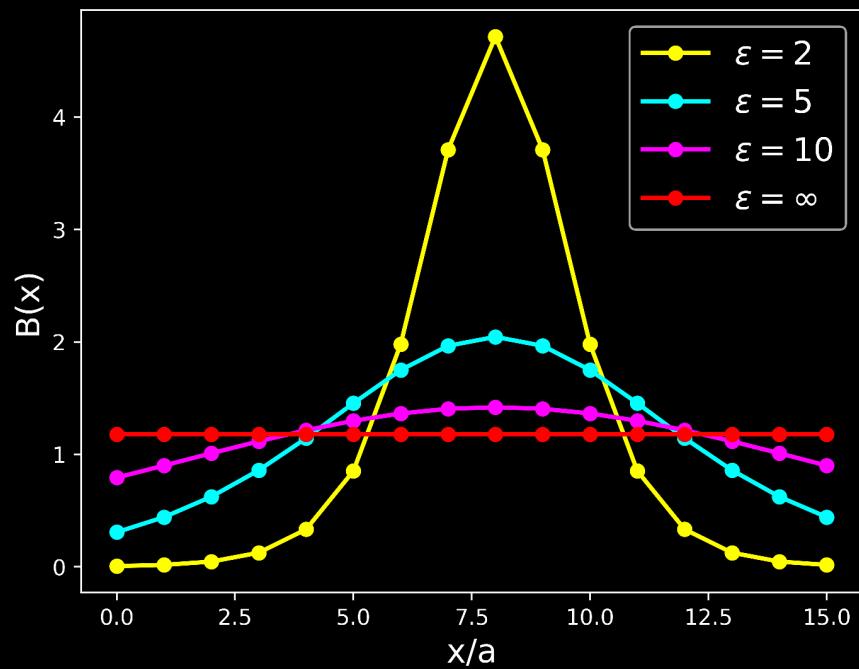
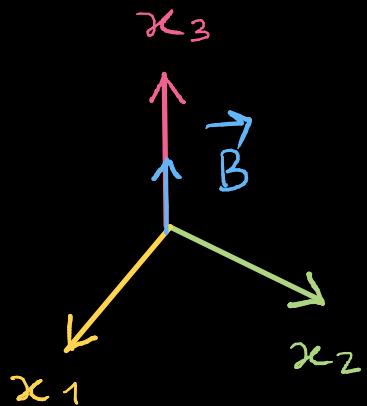
NON - UNIFORM \vec{B} MODEL

$$\vec{B}(x_1) = \frac{B}{\cosh^2\left(\frac{x_1}{\epsilon}\right)} \hat{x}_3$$



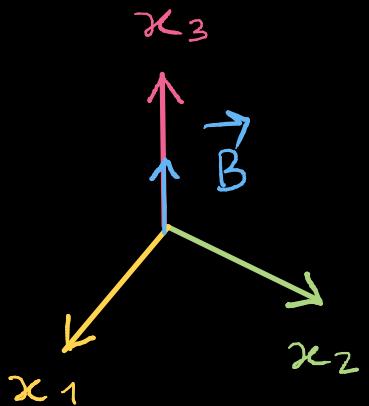
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NON-UNIFORM \vec{B} MODEL

$$\vec{B}(x_1) = \frac{\vec{B}}{\cosh\left(\frac{x_1}{\varepsilon}\right)^2} \hat{x}_3$$

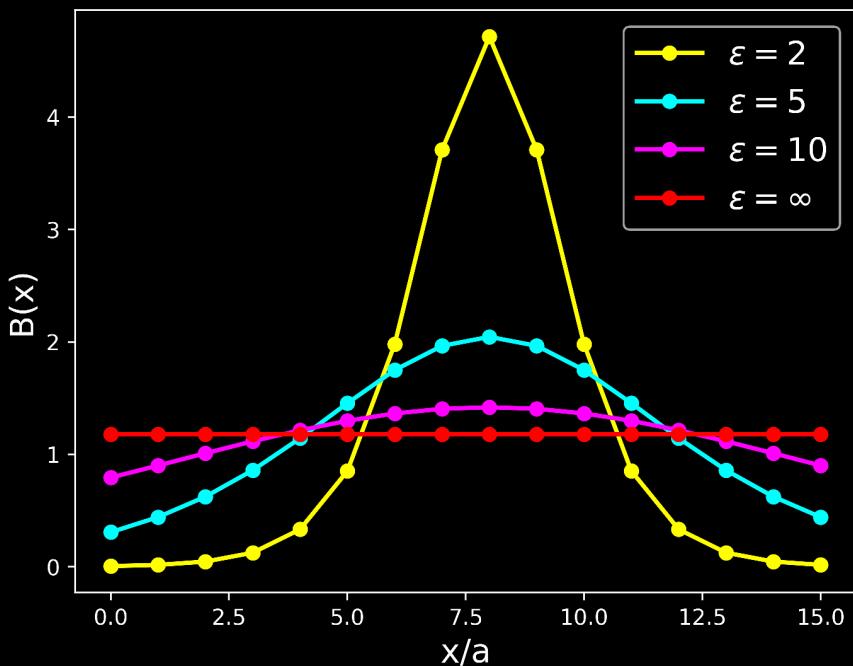


- ANALYTICALLY TREATABLE

📎 Gaoqing Cao Phys. Rev. D 97, 054021 (2018)

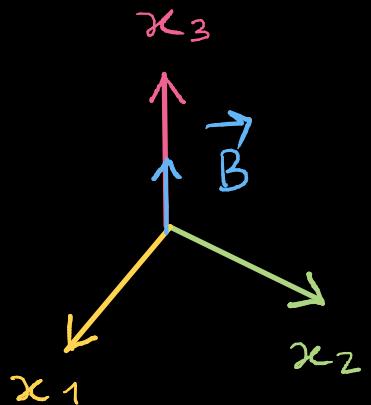
- RESEMBLES \vec{B} IN HIC

📎 W. T. Deng & X. G. Huang 2012



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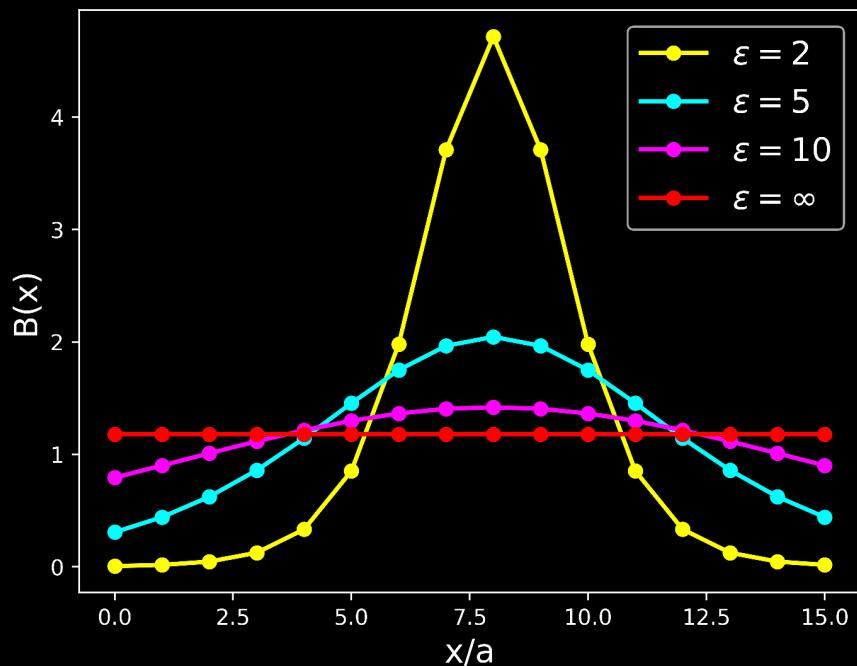


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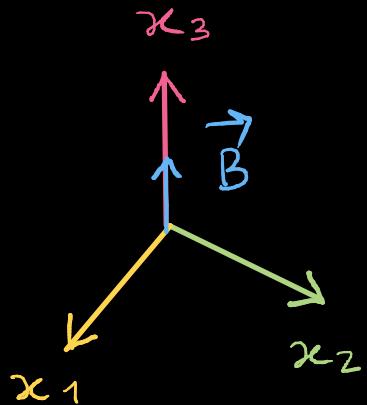
QUANTIZATION
ON A FINITE VOLUME :

$$eB = \frac{3\pi N_b}{\varepsilon L_y \tanh\left(\frac{L_x}{2\varepsilon}\right)}$$

$N_b \in \mathbb{Z}$

NON-UNIFORM \vec{B} MODEL

$$\vec{B}(x_1) = \frac{\vec{B}}{\cosh\left(\frac{x_1}{\varepsilon}\right)^2} \hat{x}_3$$

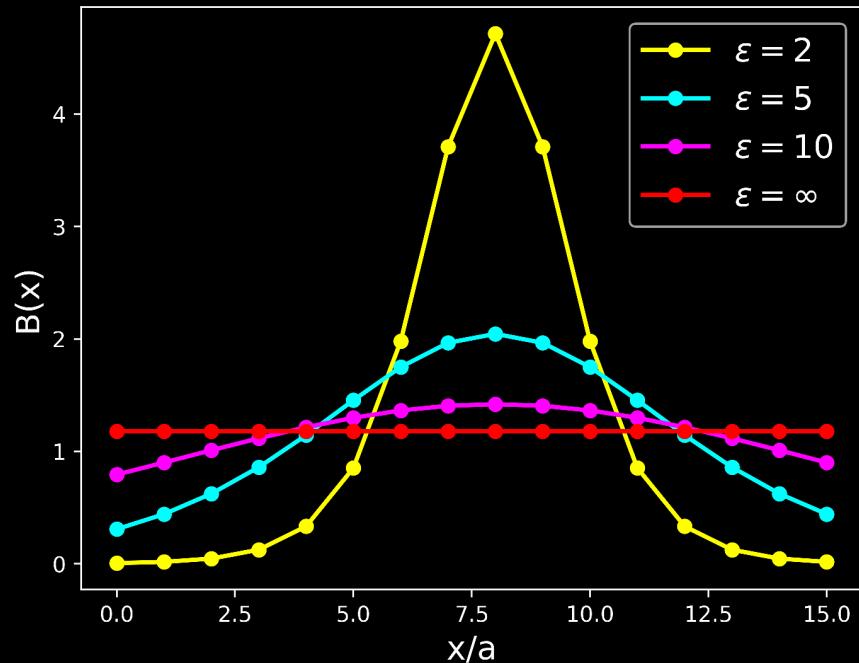


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- IMPACT ON PHASE DIAGRAM : 📎 B. Brandt et al. 2023

- USED AS A TOOL TO COMPUTE EoS : 📎 B. Brandt, G. Endrődi, G. Markó, A. D. M. Valois 2024

CME WITH FREE FERMIONS

$$j_3 = C_{\text{CME}} \mu_5 B$$

CME WITH FREE FERMIONS

$$j_3 = C_{\text{CME}} \mu_5 B$$

- OUR LATTICE OBSERVABLE: $G(x_1) \equiv \left[\frac{\partial}{L^2} \langle j_3(x_1) \rangle \right]_{\mu_5=0}$

CME WITH FREE FERMIONS

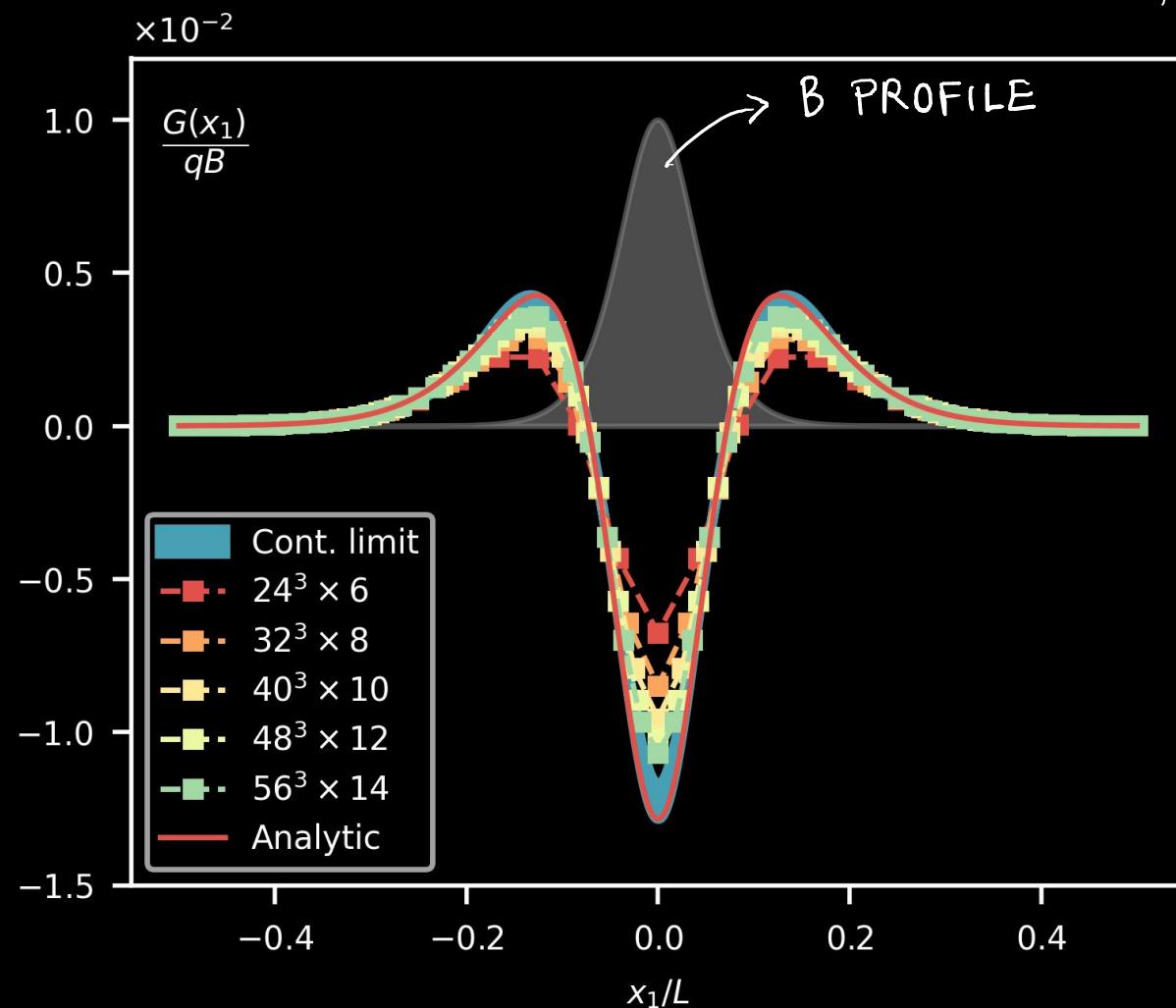
$$j_3 = C_{\text{CME}} \mu_5 B$$

• OUR LATTICE OBSERVABLE: $G(x_1) \equiv \left[\frac{1}{L^2} \frac{\partial}{\partial \mu_5} \langle j_3(x_1) \rangle \right]_{\mu_5=0} = \frac{1}{L^2} \langle j_{45} j_3(x_1) \rangle$

CME WITH FREE FERMIONS

$$j_3 = C_{\text{CME}} \mu_5 B$$

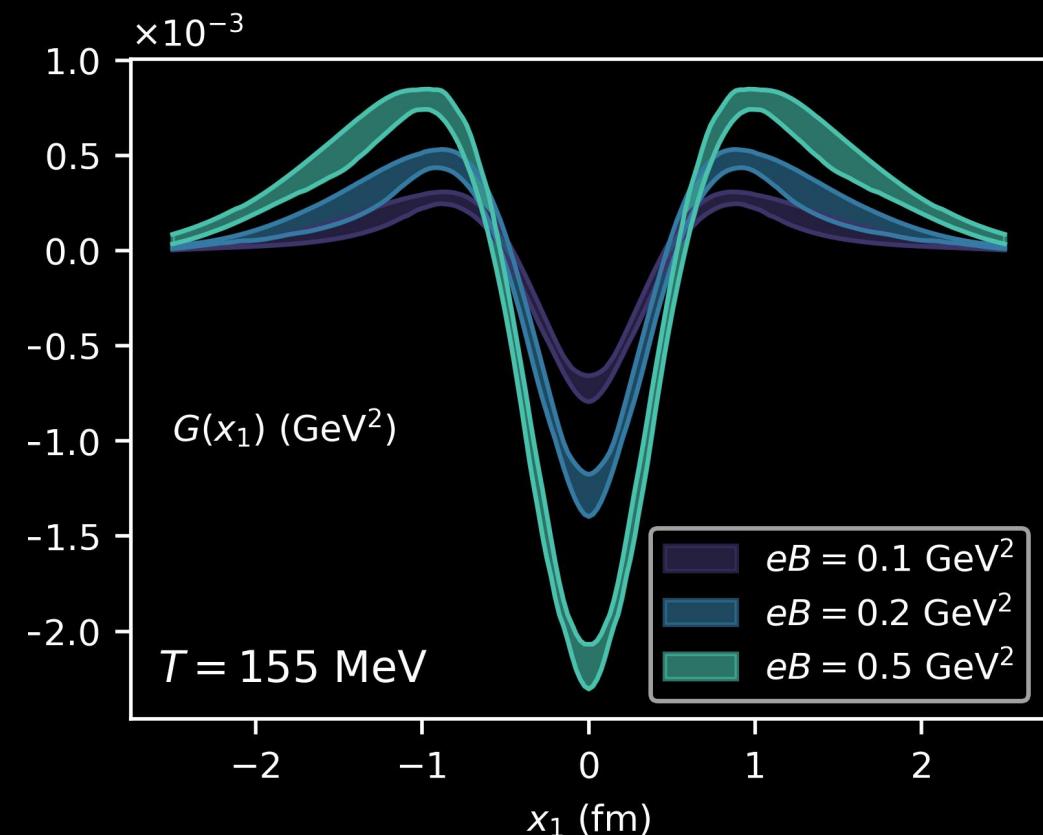
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CME IN QCD

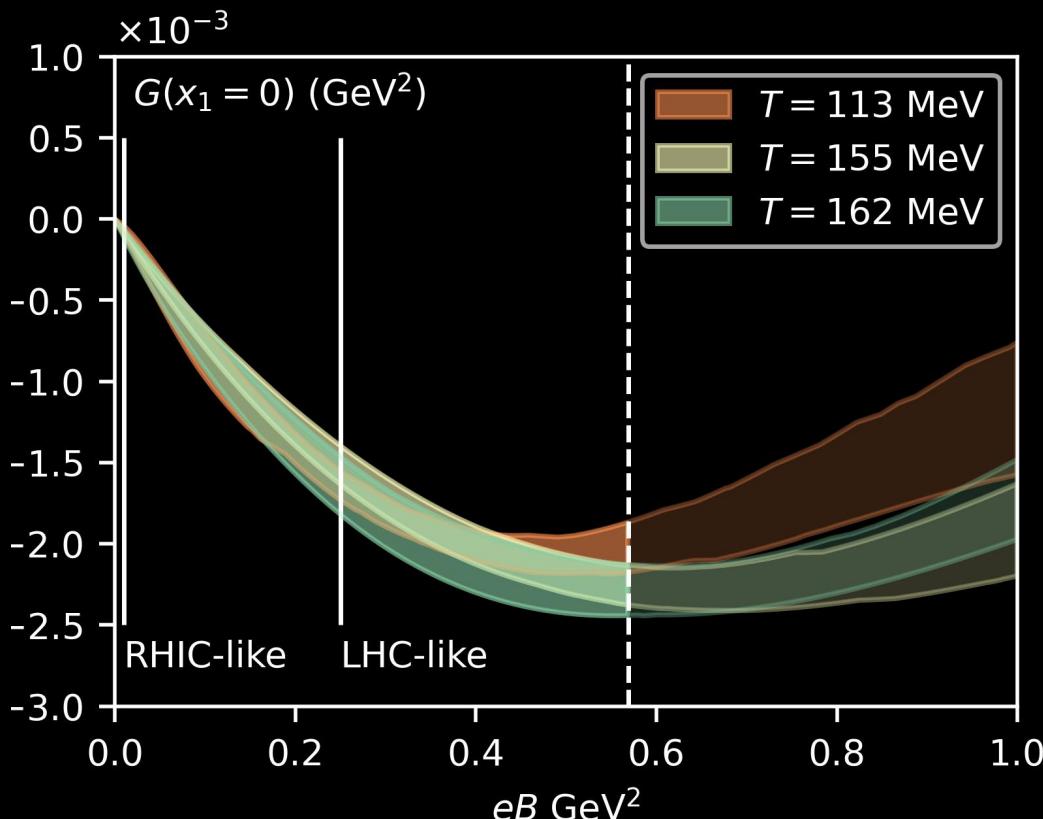
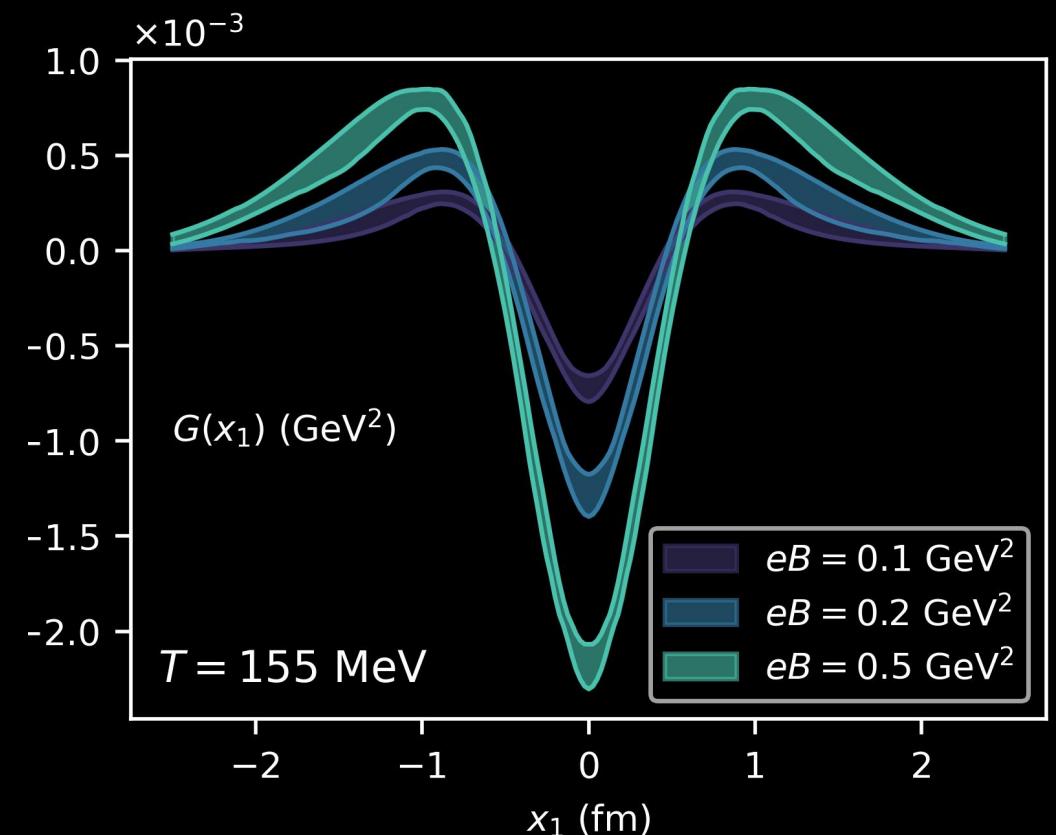
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⌚ B. Brandt, G. Endrődi, E. Garnacho-Velasco,
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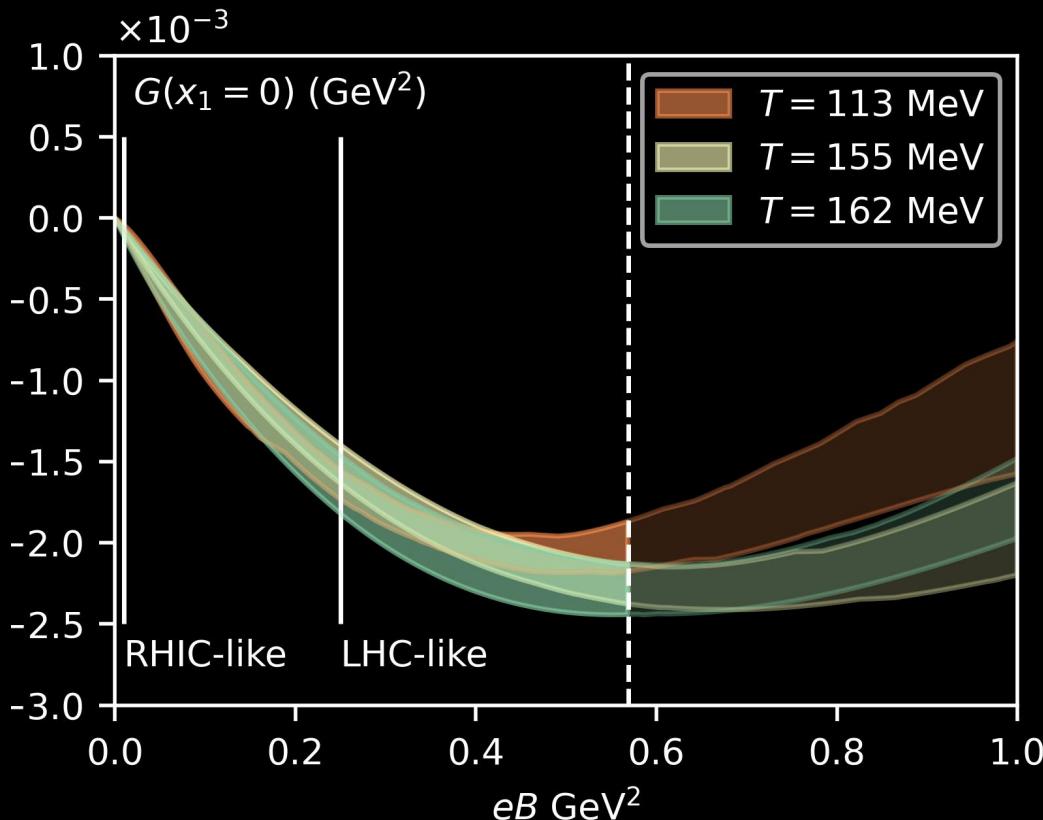
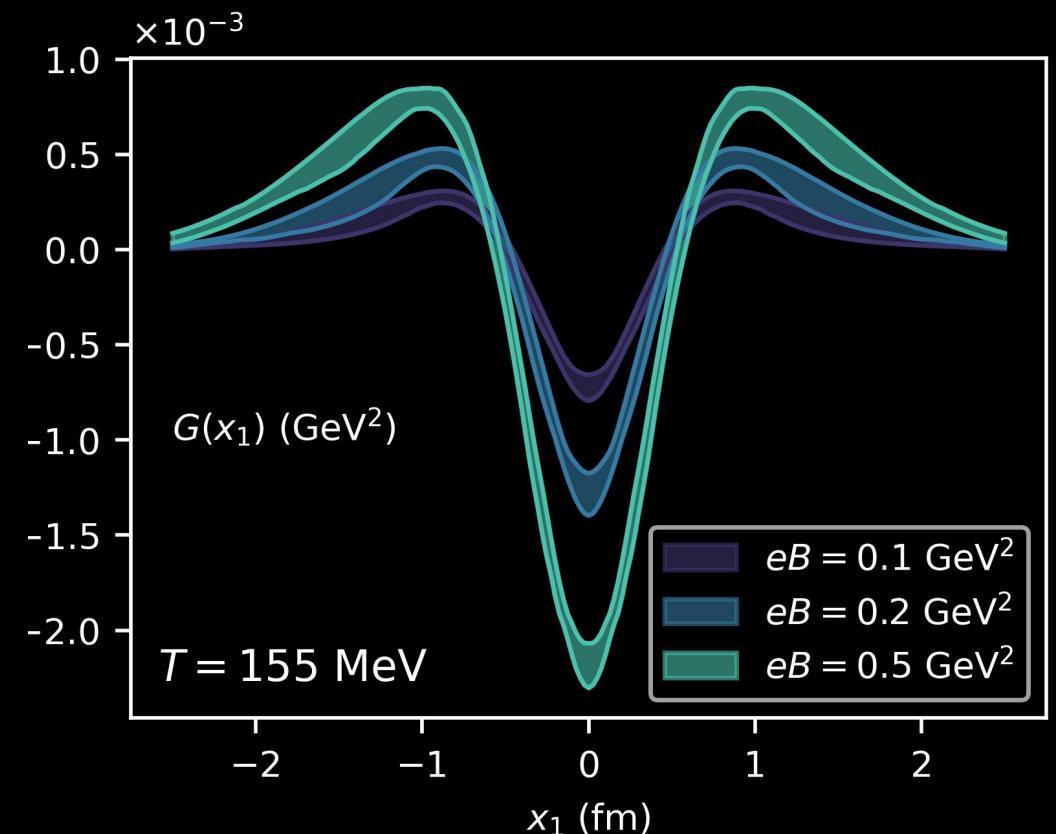
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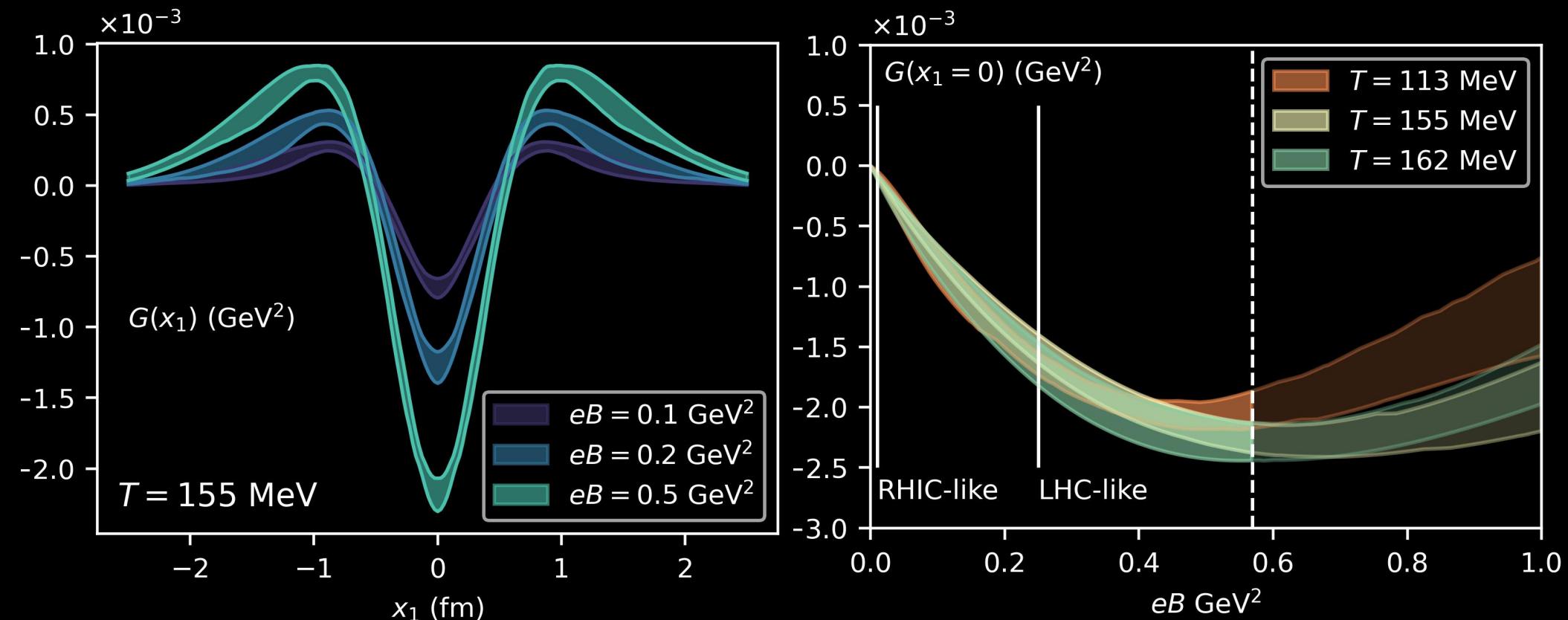
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- J_3 FLOWS ALONG $\mathcal{K}_3 \parallel \mathcal{B}$

CME IN QCD

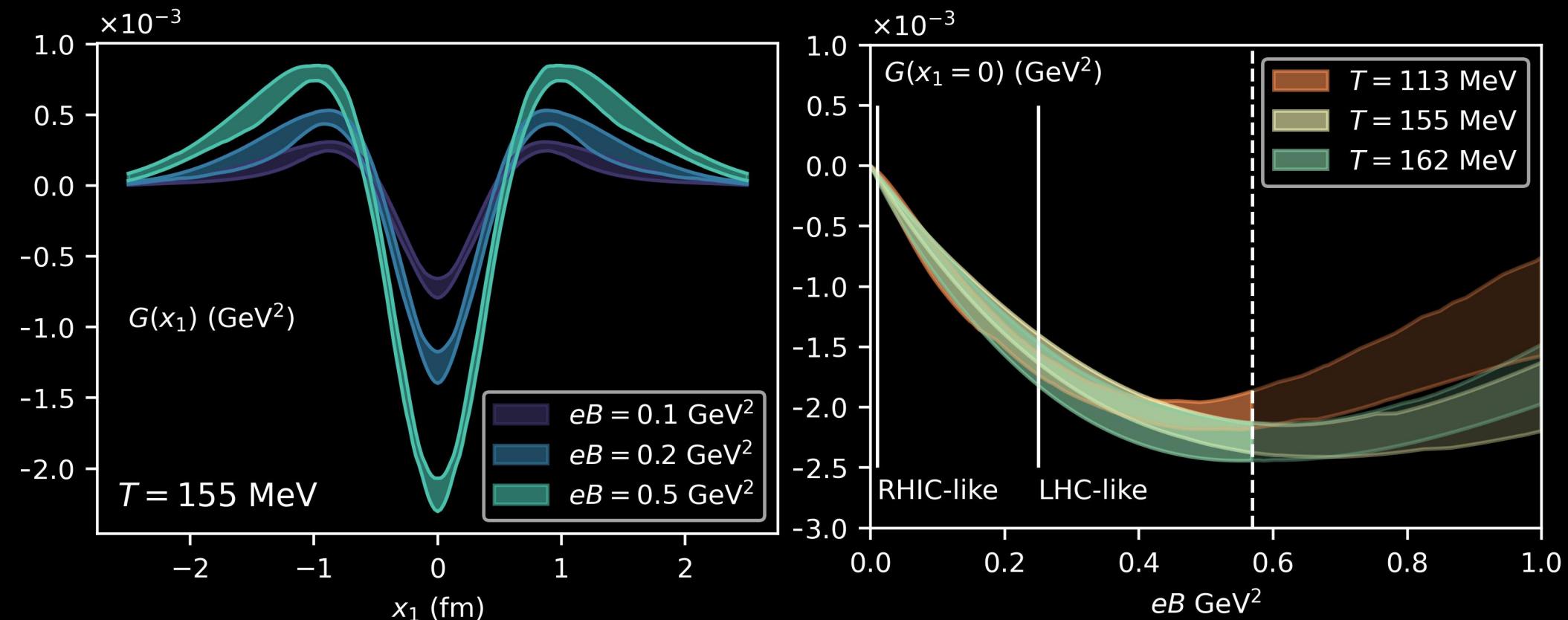
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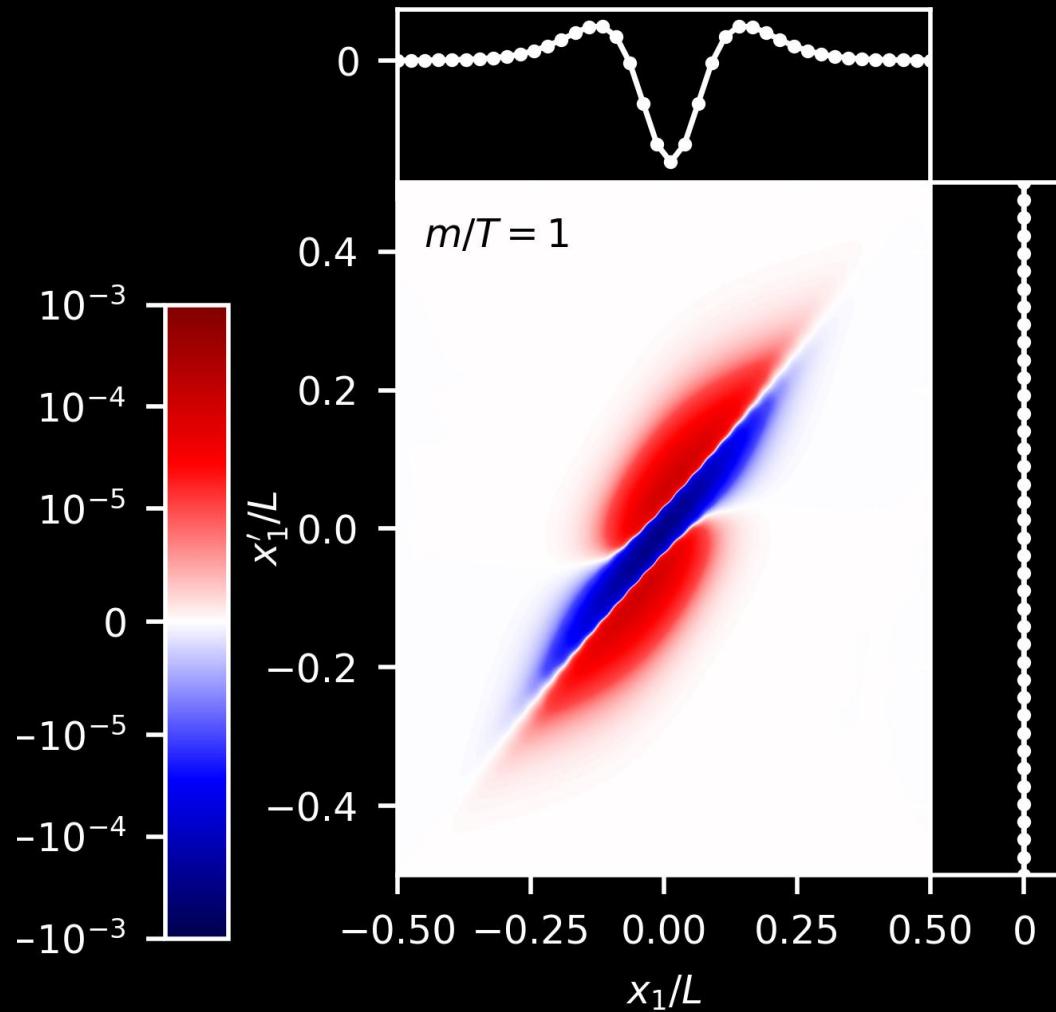
- j_3 FLOWS ALONG $x_3 \parallel \mathcal{B}$
- j_3 IS INHOMOGENEOUS ALONG x_1
- $\int dx_1 G(x_1) = 0$

CME IN QCD

$$H(x_1, x_1') = \frac{I}{L^2} \left. \frac{\delta \langle j_3(x_1) \rangle}{\delta \mu_5(x_1')} \right|_{\mu_5=0}$$

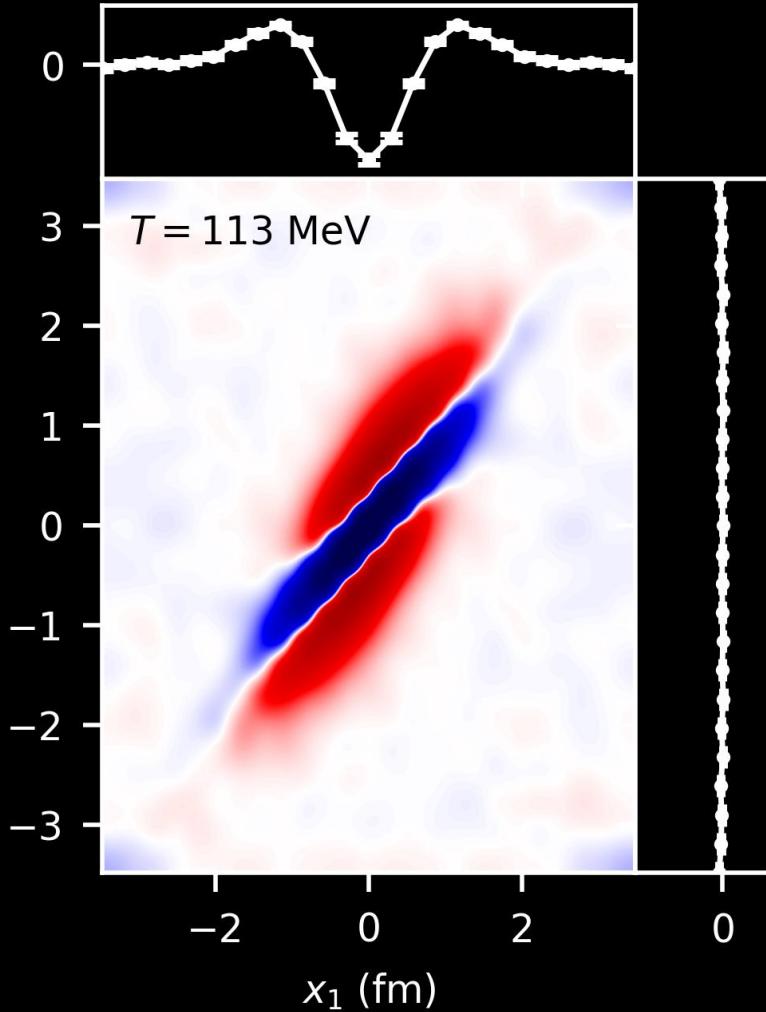
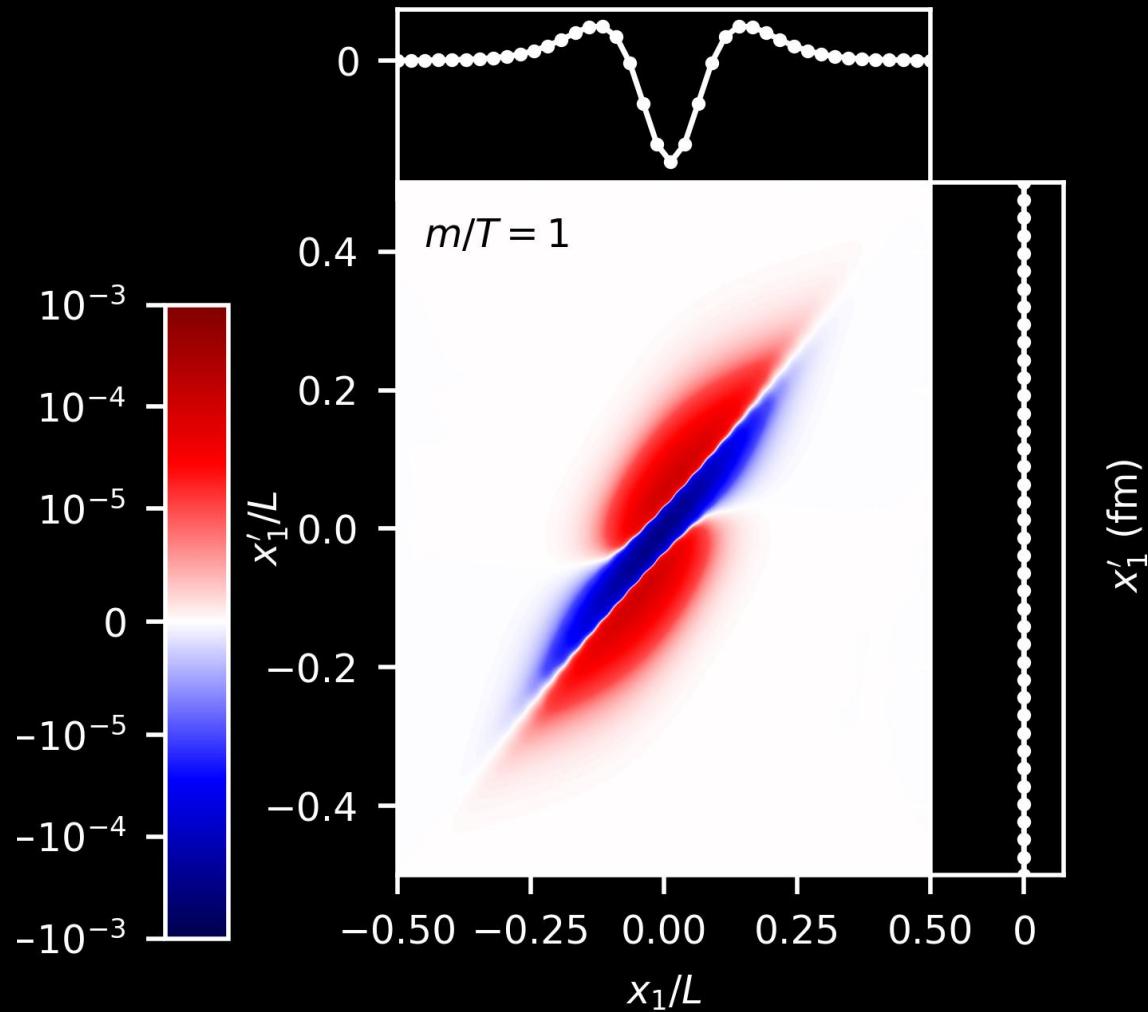
CME IN QCD

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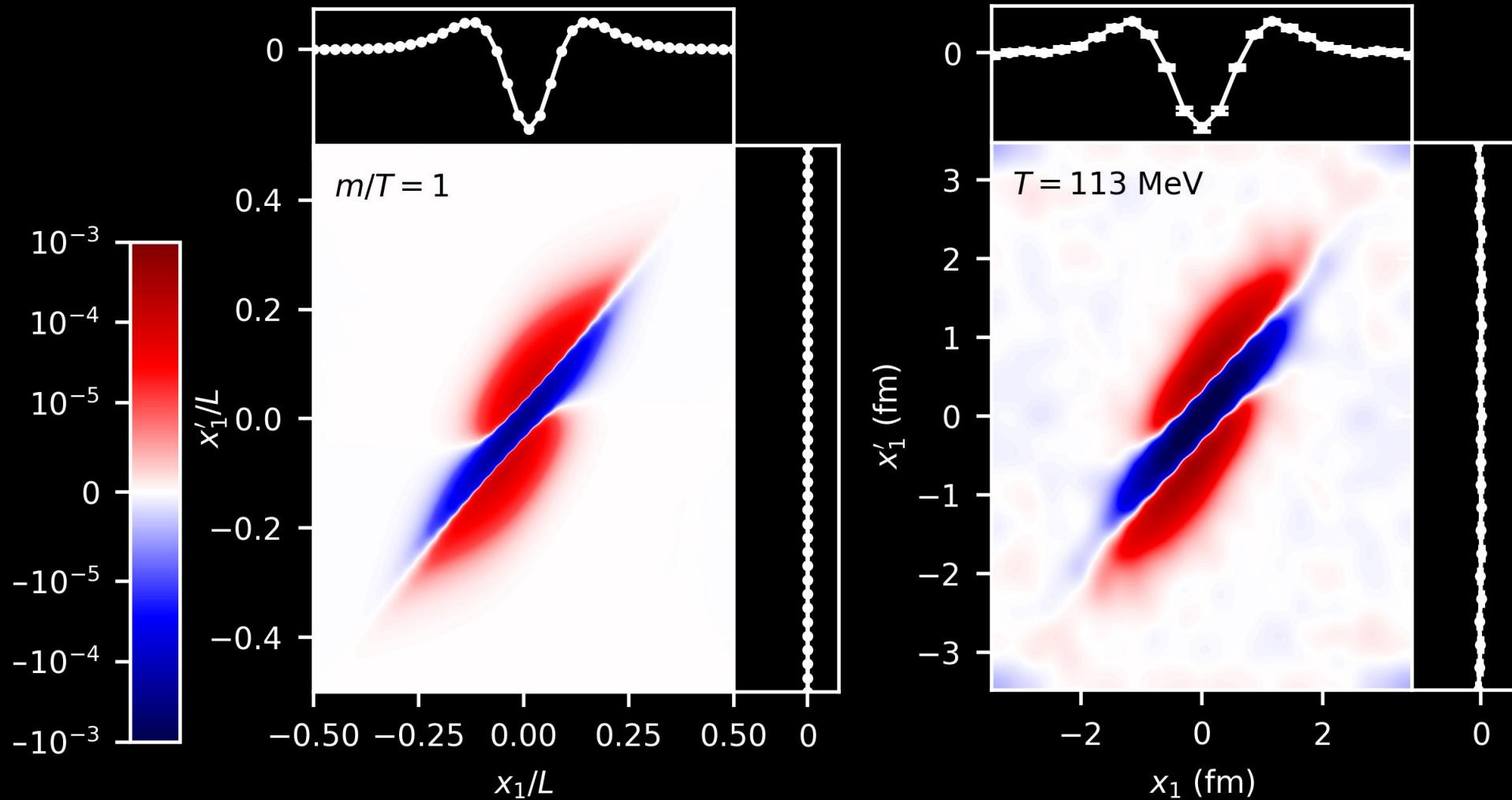
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IF $\mu_5(x_1')$ IS KNOWN \rightarrow CONVOLUTE $\mu_5(x_1')$ WITH $H(x_1, x_1')$

PART II -

OUT-OF-EQUILIBRIUM CME WITH
UNIFORM B

LINEAR RESPONSE THEORY

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$$C_{\text{CME}}^{\text{OUT}} \sim \frac{1}{eB} \lim_{\omega \rightarrow 0} \frac{\rho(\omega)}{\omega}$$

KUBO FORMULA FOR CME

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E.G.: BACKUS-GILBERT BAYESIAN MEM GAUSSIAN

MID-POINT METHOD

📎 P. V. Buividovich 2024

MID-POINT METHOD

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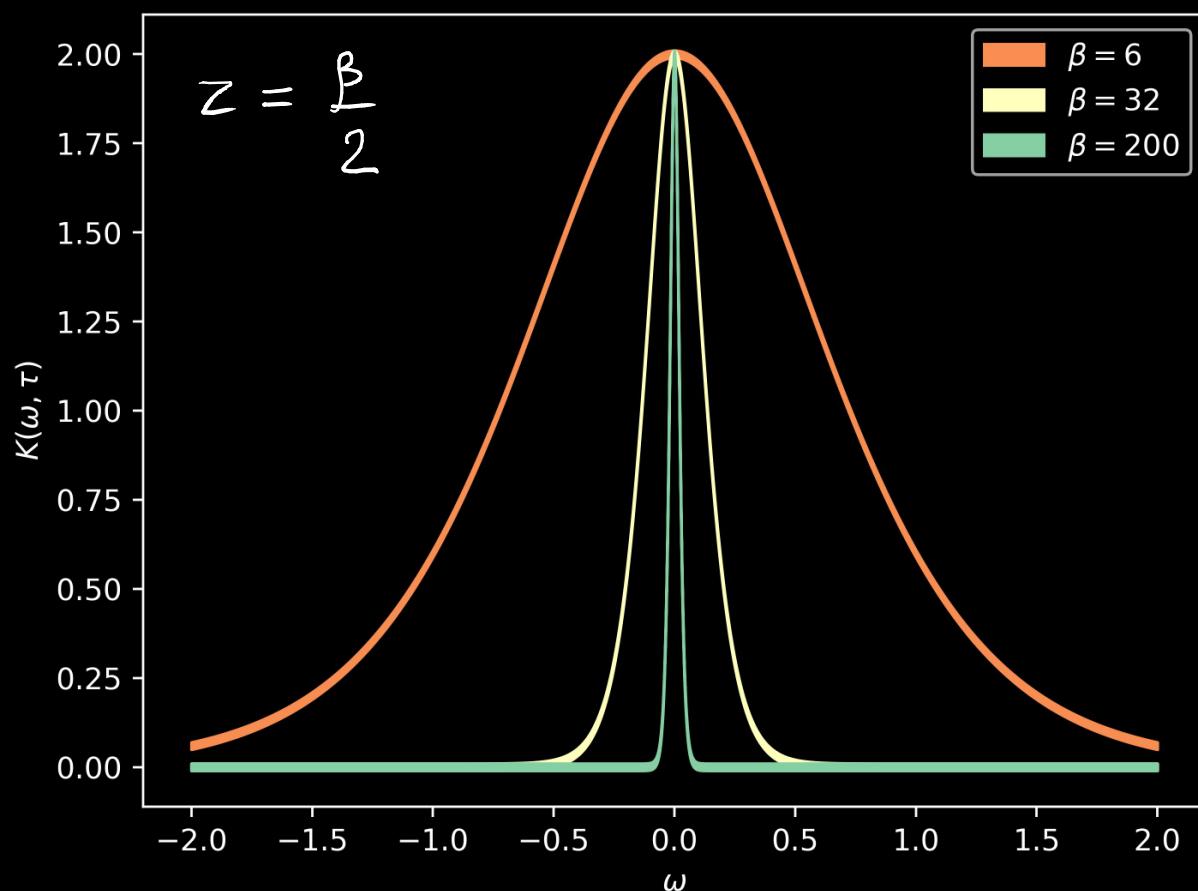
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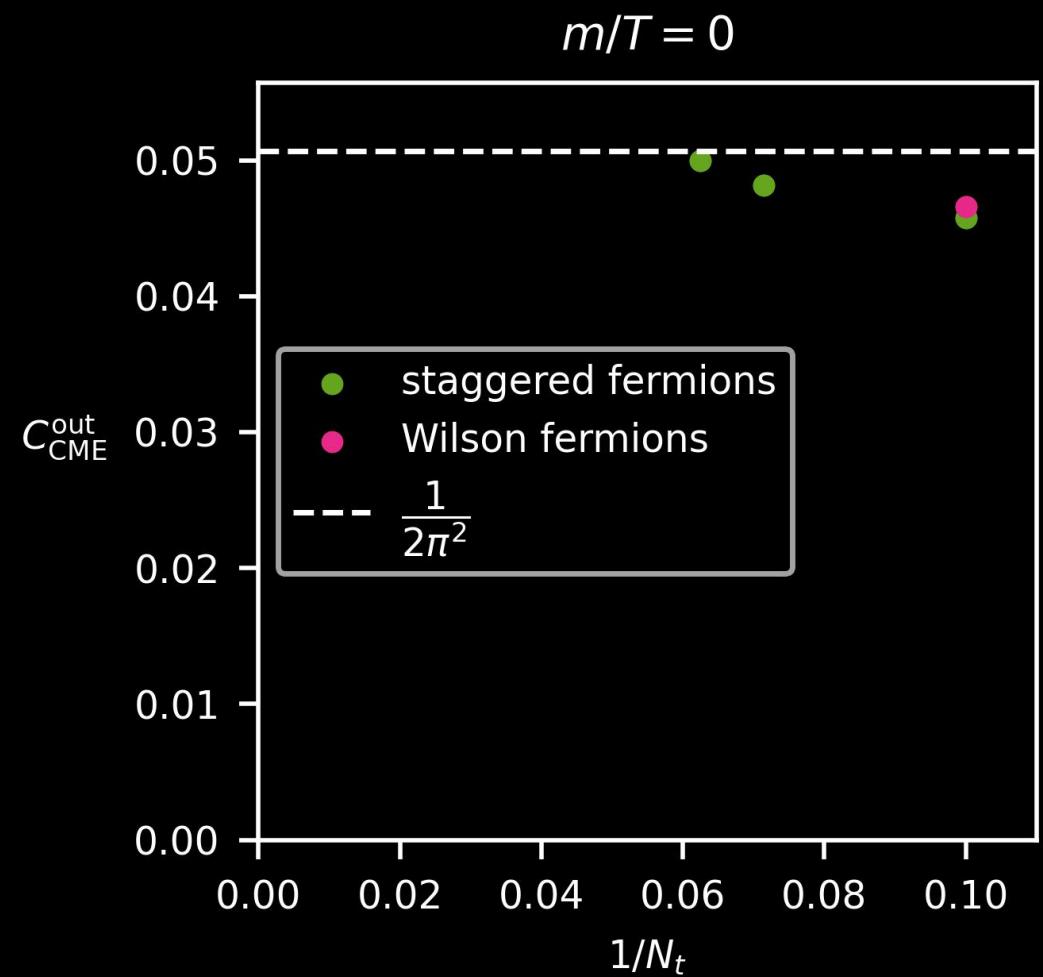
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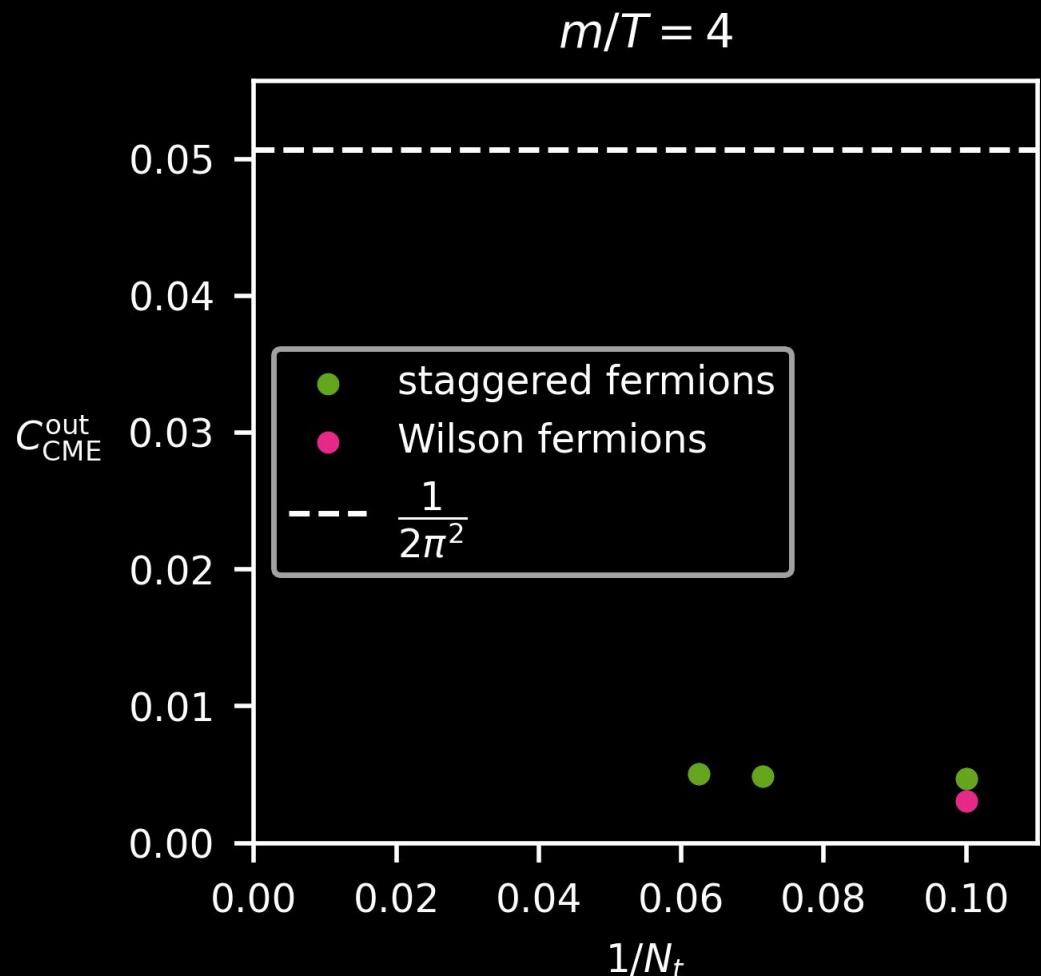
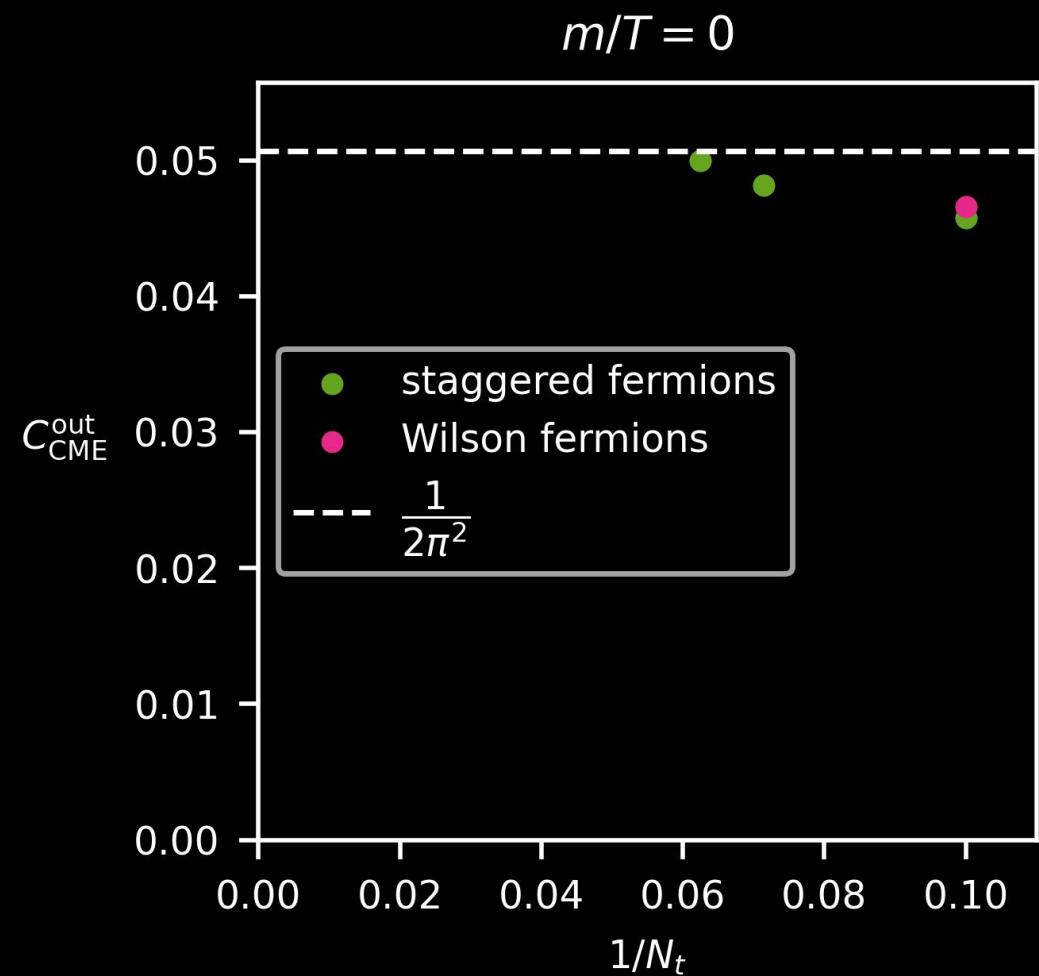
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PRELIMINARY RESULTS
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OUT-OF-EQUI.

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