

Predictions in the superheavy region from the quark-meson coupling model, QMC π -III

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Motivation

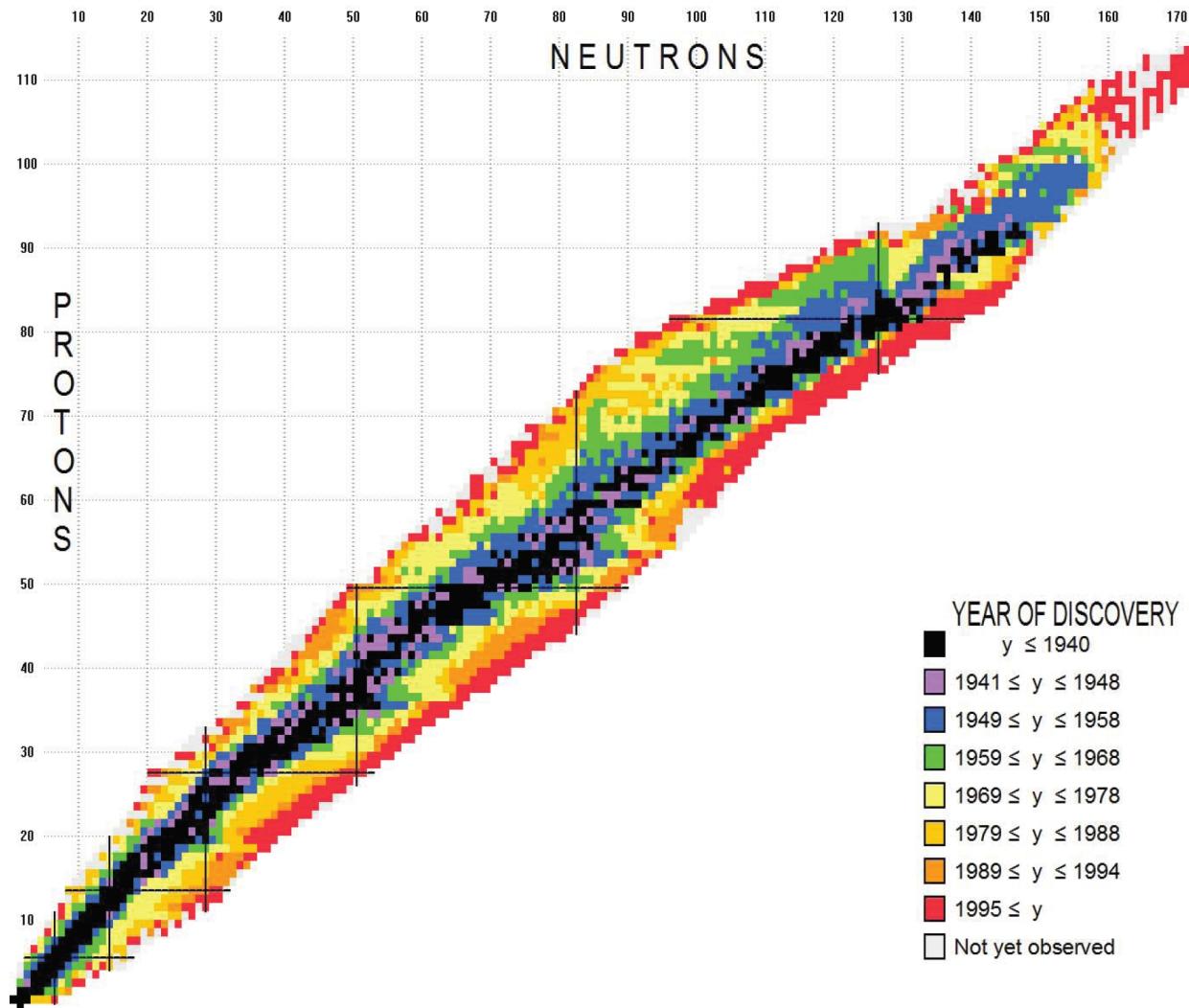
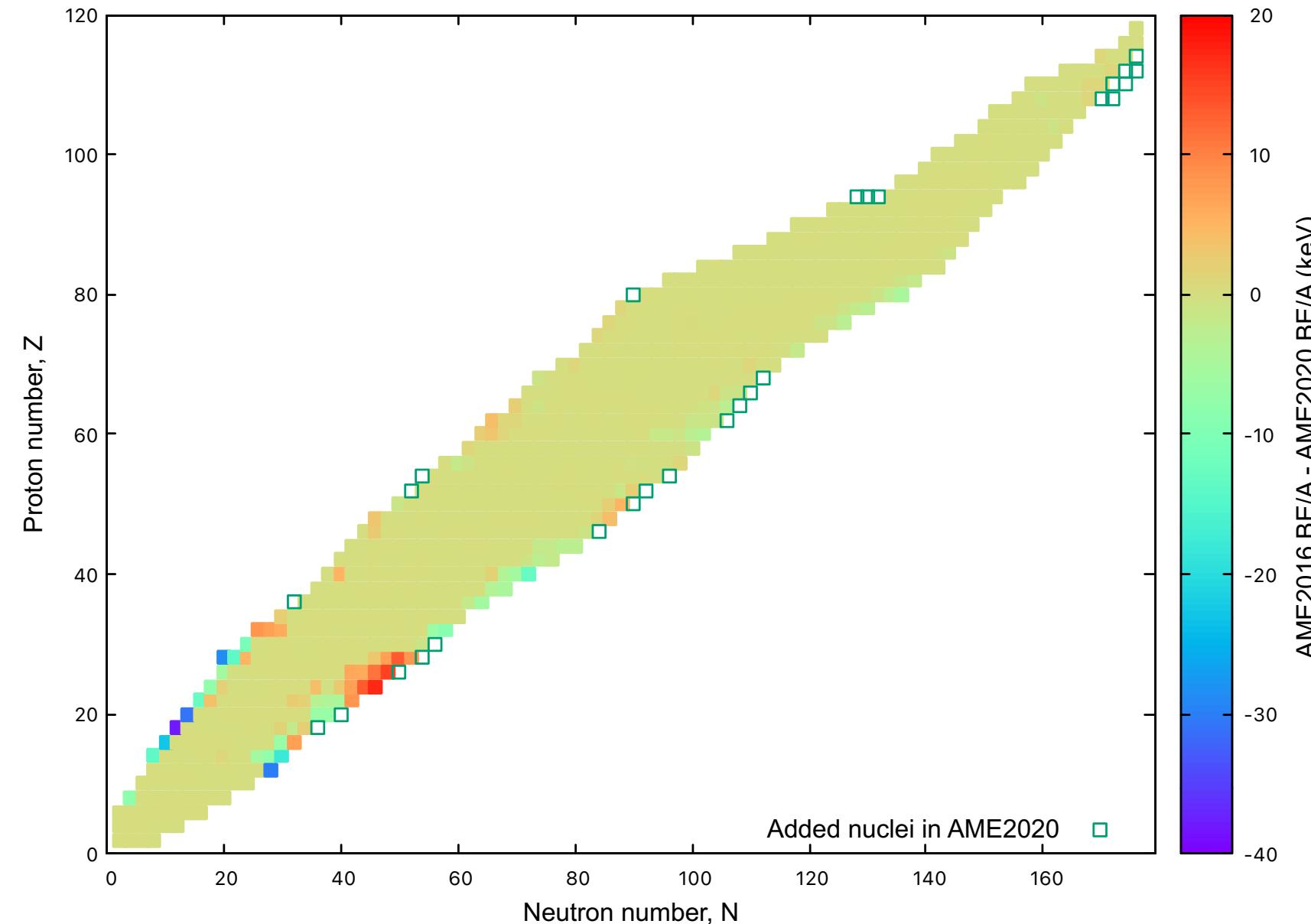


Figure from NuBASE 2017, AME 2016

Motivation



Differences in binding energy per nucleon (BE/A) between AME 2016[1] and AME 2020[2].

Motivation

- The Quark-Meson Coupling (QMC) model describes a self-consistent relationship between the dynamics of the quark structure of a nucleon and the relativistic mean fields arising within the nuclear medium.

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Thomas, A. W. (2016). QCD and a new paradigm for nuclear structure. *Heavy Ion Accelerator Symposium on Fundamental and Applied Science (HIAS 2015)* (p.01003). Canberra: EPJ Web Conf.

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Stone, J. R., Guichon, P. A., & Thomas, A. W. (2017). Super-heavy nuclei in the quark-meson coupling model. *EPJ Web Conf.*, **163** 00057.

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K. L. Martinez, A. W. Thomas, P. A. M. Guichon, and J. R. Stone. (2020). Tensor and pairing interactions within the quark-meson coupling energy-density functional. *Phys. Rev. C* **102** p. 034304.

Antić, S., Stone, J. R., Miller, J. C., Martinez, K. L., Guichon, P. A. M., & Thomas, A. W. (2020). Outer crust of a cold, nonaccreting neutron star within the quark meson-coupling model. *Phys. Rev. C* **102** 065801.

Outline

1 The QMC Model

2 Parameterization procedure

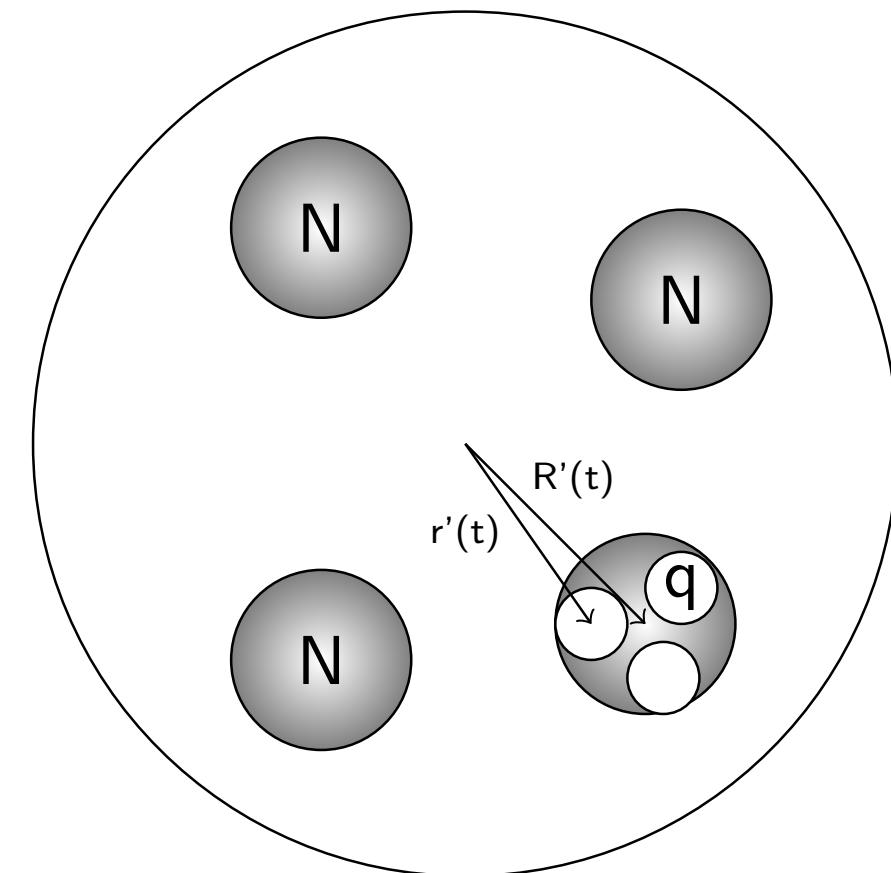
3 Results

4 Conclusion

The QMC Model

The QMC Model

- employs relativistic mean-field theory
- interactions are described through the scalar σ meson for the intermediate range attraction, vector meson ω for the short-range repulsion, and vector-isovector ρ for isospin dependence



The QMC Model

Nucleon effective mass:

$$M_B^* = M_B - g_\sigma \sigma + \frac{d}{2} (g_\sigma \sigma)^2$$

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The *classical* total energy

$$\begin{aligned} E_{QMC} &= \sum_{i=1,\dots} \sqrt{P_i^2 + M_i^2(\sigma(\vec{R}_i))} + g_\omega^i \omega(\vec{R}_i) + g_\rho \vec{l}_i \cdot \vec{B}(\vec{R}_i) \\ &\quad + E_\sigma + E_{\omega,\rho} \end{aligned}$$

The QMC Model

To get the Hamiltonian, H_{QMC} , we eliminate the meson fields by solving the equations of motion:

$$\frac{\delta H_{QMC}}{\delta \sigma(\vec{r})} = \frac{\delta H_{QMC}}{\delta \omega(\vec{r})} = \frac{\delta H_{QMC}}{\delta B(\vec{r})} = 0$$

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HF calculation

$$\mathcal{E}_{QMC} = \langle \Phi | H_{QMC} | \Phi \rangle$$

QMC model for finite nuclei

Densities

$$\rho_m(\vec{r}) = \sum_{i \in F_m} \sum_{\sigma} |\phi^i(\vec{r}, \sigma, m)|^2, \quad \rho = \rho_p + \rho_n,$$

$$\tau_m(\vec{r}) = \sum_{i \in F_m} \sum_{\sigma} \left| \vec{\nabla} \phi^{i*}(\vec{r}, \sigma, m) \right|^2, \quad \tau = \tau_p + \tau_n,$$

$$\vec{J}_m = i \sum_{i \in F_m} \sum_{\sigma \sigma'} \vec{\sigma}_{\sigma' \sigma} \times \left[\vec{\nabla} \phi^i(\vec{r}, \sigma, m) \right] \phi^{i*}(\vec{r}, \sigma', m),$$

$$\vec{J} = \vec{J}_p + \vec{J}_n$$

QMC model for finite nuclei

The contributions to the QMC Hamiltonian can be written as

$$H_{QMC} = H_0 + H_3 + H_{\text{eff}} + H_{\text{fin}} + H_{\text{so}} + H_{\pi}$$

$H_0 + H_3 \propto \rho^2$ QMC density-dependent terms

$H_{\text{eff}} \propto \rho\tau$ effective mass term

$H_{\text{fin}} \propto \nabla^2 \rho$ finite size effect

$H_{\text{so}} \propto \nabla \cdot \vec{J}$ spin-orbit

H_{π} single-pion exchange

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H_{π} single-pion exchange

Model parameters: $G_\mu = \frac{g_\mu}{m_\mu^2}$ ($\mu = \sigma, \omega, \rho$), m_σ , and λ_3

The pairing EDF

HF+BCS pairing

$$\mathcal{E}_{\text{pair},q} = \frac{1}{4} S_q(\vec{r}) \check{\rho}_q^2, \quad \chi_q(\vec{r}) = \sum_{\alpha \in q} u_\alpha v_\alpha |\phi_\alpha(\vec{r})|^2$$

Pairing strength

$$S_q(\vec{r}) = V_q \left[1 - \left(\frac{\rho(\vec{r})}{\rho_c} \right)^\alpha \right]$$

Volume pairing / Delta force (DF) pairing

Surface pairing / Density-dependent delta interaction (DDDI)

The QMC-derived pairing EDF

Pairing potential

$$V(\vec{r} - \vec{r}') = - \left(\frac{G_\sigma}{1 + d' G_\sigma \rho(\vec{r})} - G_\omega - \frac{G_\rho}{4} \right) \delta(\vec{r} - \vec{r}')$$

$$d' = d + \frac{1}{3} G_\sigma \lambda_3$$

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$$V_q = G_\sigma - G_\omega - G_\rho/4, \quad \rho_c = \frac{V_q (1 + d' G_\sigma \rho)}{d' G_\sigma^2}, \quad \alpha = 1$$

Coulomb and center-of-mass correction

$$\mathcal{E}_{\text{Coulomb}} = e^2 \frac{1}{2} \int d^3 r d^3 r' \frac{\rho_p(\vec{r}) \rho_p(\vec{r}')}{|\vec{r} - \vec{r}'|} - \frac{3}{4} e^2 \left(\frac{3}{\pi} \right)^{1/3} \int d^3 r [\rho_p]^{4/3}$$

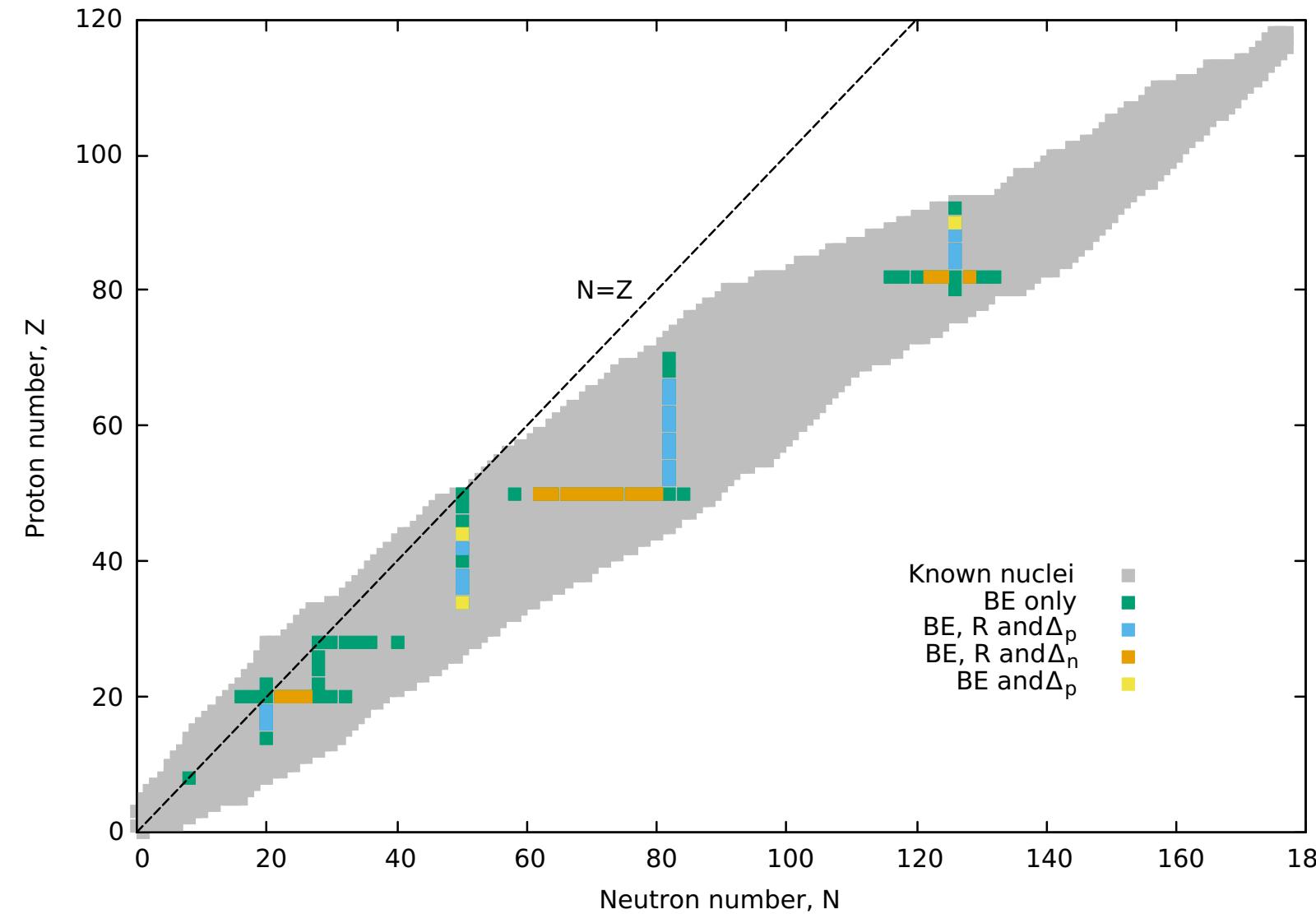
$$\mathcal{E}_{\text{c.o.m.}} = -\langle \hat{P}_{cm}^2 \rangle / (2mA)$$

QMC versions

version	pion	σ	mass	σ	self-coupling	spin-orbit	spin-tensor	pairing
QMC-I _[3]				✓				DF
QMC π -I _[4]	✓			✓				DF
QMC π -II _[2,5]	✓		bare		✓		✓	DF
QMC π -III _[6,7]	✓		bare		✓		✓	DDDI

Parameterization procedure

Experimental data set



Doubly-magic and semi-magic nuclei included in the fit. The nuclear observables and number of data points per nucleus entering the fitting procedure are indicated.

POUNDerS algorithm

Objective function $F(\hat{\mathbf{x}})$ for minimisation

$$F(\hat{\mathbf{x}}) = \sum_i^n \sum_j^o \left(\frac{\bar{s}_{ij} - s_{ij}}{w_j} \right)^2$$

where

n total number of nuclei

o total number of observables

s_{ij} experimental values

\bar{s}_{ij} fitted values

w_j effective error

$w_{BE} = 1.0$ MeV

$w_{R_{ch}} = 0.02$ fm

$w_{\Delta_{p,n}} = 0.12$ MeV

Results

QMC and pairing parameters

Parameter	QMC π -III	QMC π -II	QMC π -I	QMC-I
G_σ [fm 2]	9.62	9.66	11.16	11.85
G_ω [fm 2]	5.21	5.23	8.00	8.27
G_ρ [fm 2]	4.71	4.75	6.38	7.68
M_σ [MeV]	503	493	712	722
λ_3 [fm $^{-1}$]	0.05	0.05	-	-
V_p [MeV]	-	258	302	284
V_n [MeV]	-	237	291	326

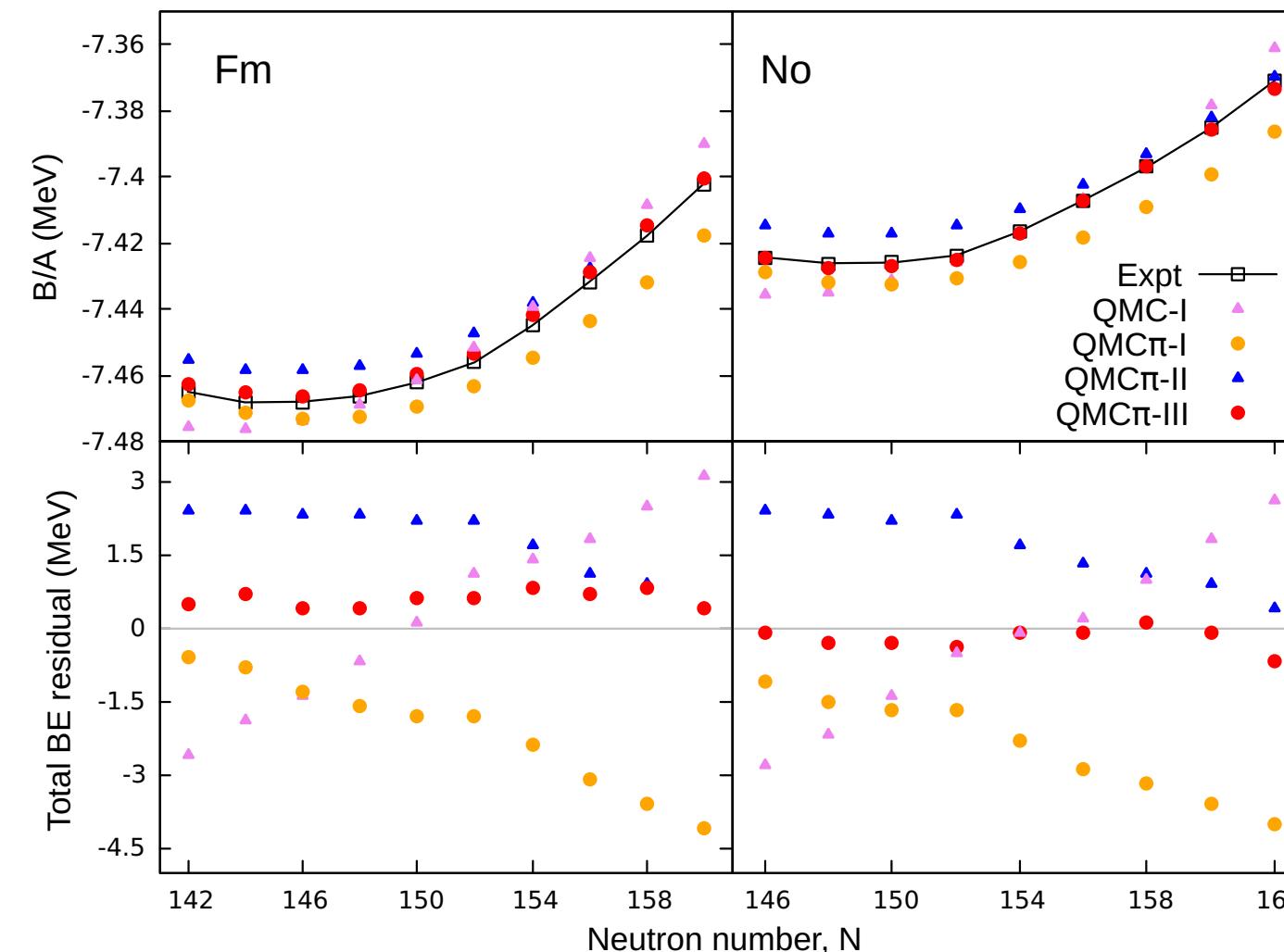
QMC nuclear matter properties

Accepted ranges:

Saturation density ρ_0	0.16 fm $^{-3}$
Saturation energy E_0	-16 MeV
Symmetry energy S_0	29–33 MeV
Slope of symmetry energy L_0	58.9 MeV
Incompressibility K_0	200–315 MeV

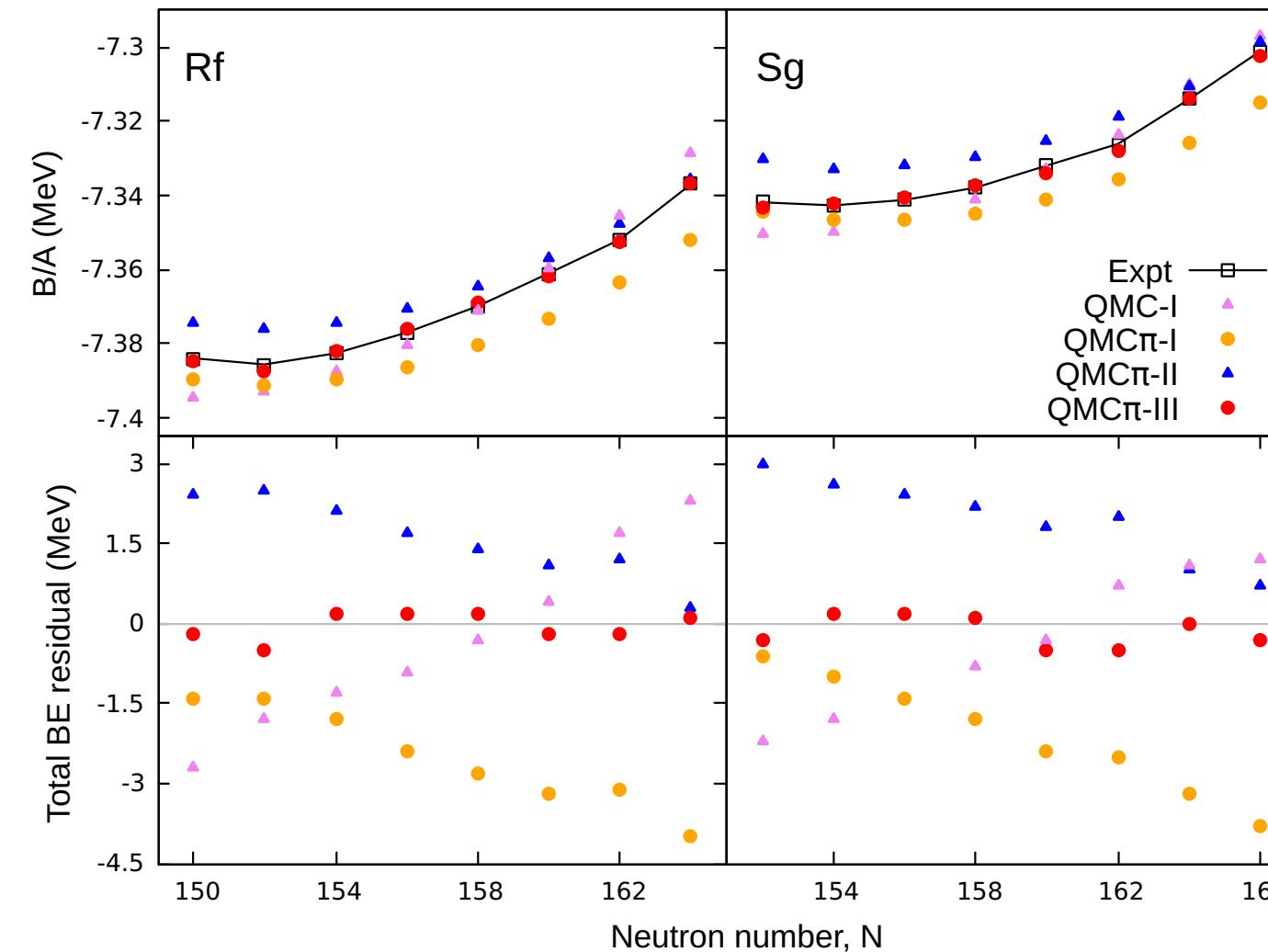
NMP	QMC π -III	QMC π -II	QMC π -I	QMC-I	SV-min
ρ_0 (fm $^{-3}$)	0.15	0.15	0.15	0.16	0.16
E_0 (MeV)	-15.7	-15.7	-15.8	-15.9	-15.9
S_0 (MeV)	29	29	30	30	31
L_0 (MeV)	43	40	17	23	93
K_0 (MeV)	233	230	319	340	222

QMC predictions for binding energies of even-even superheavy elements (SHE)



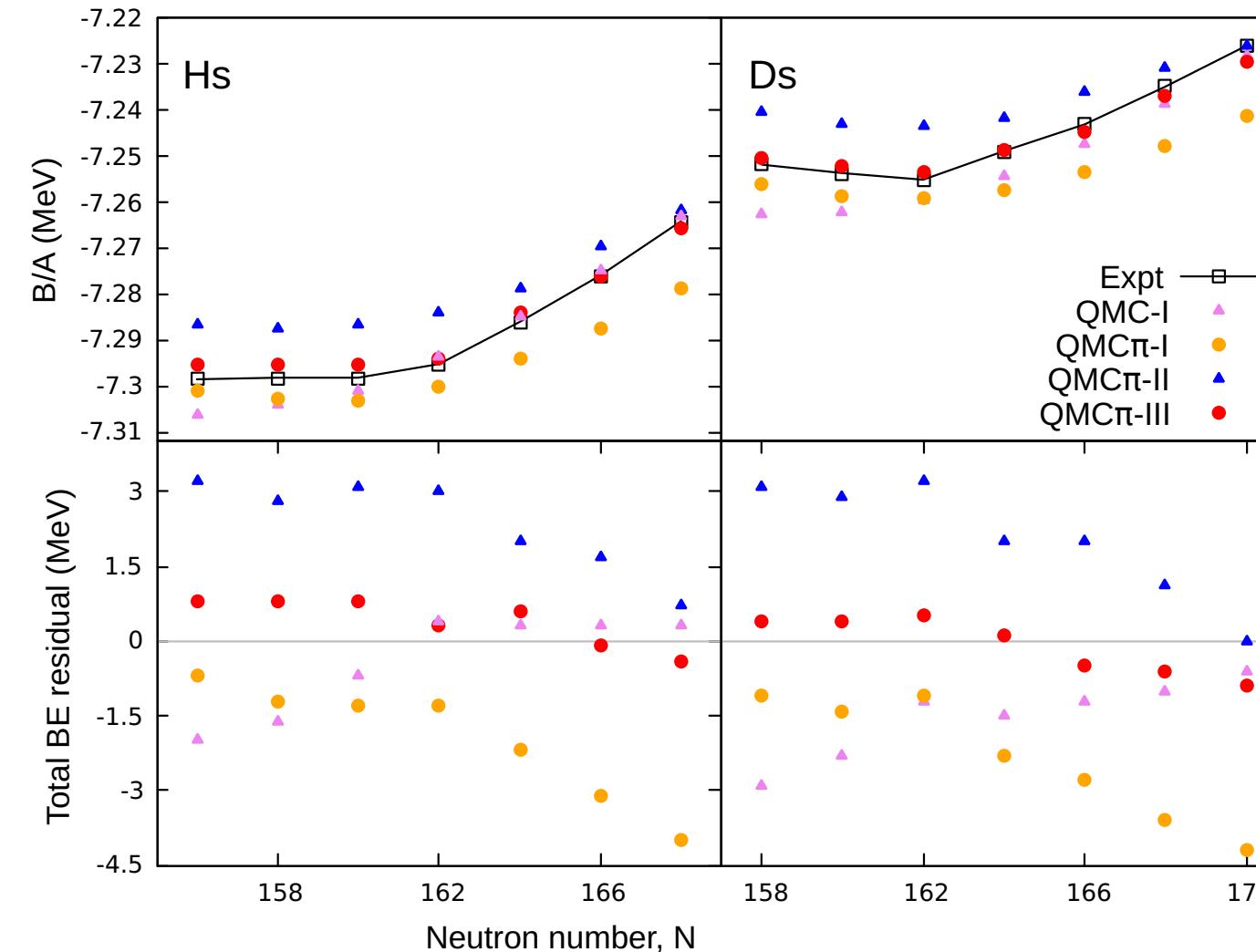
BE per nucleon and total BE residuals along Fm and No chains.

QMC predictions for binding energies of even-even superheavy elements (SHE)



BE per nucleon and total BE residuals along Rf and Sg chains.

QMC predictions for binding energies of even-even superheavy elements (SHE)



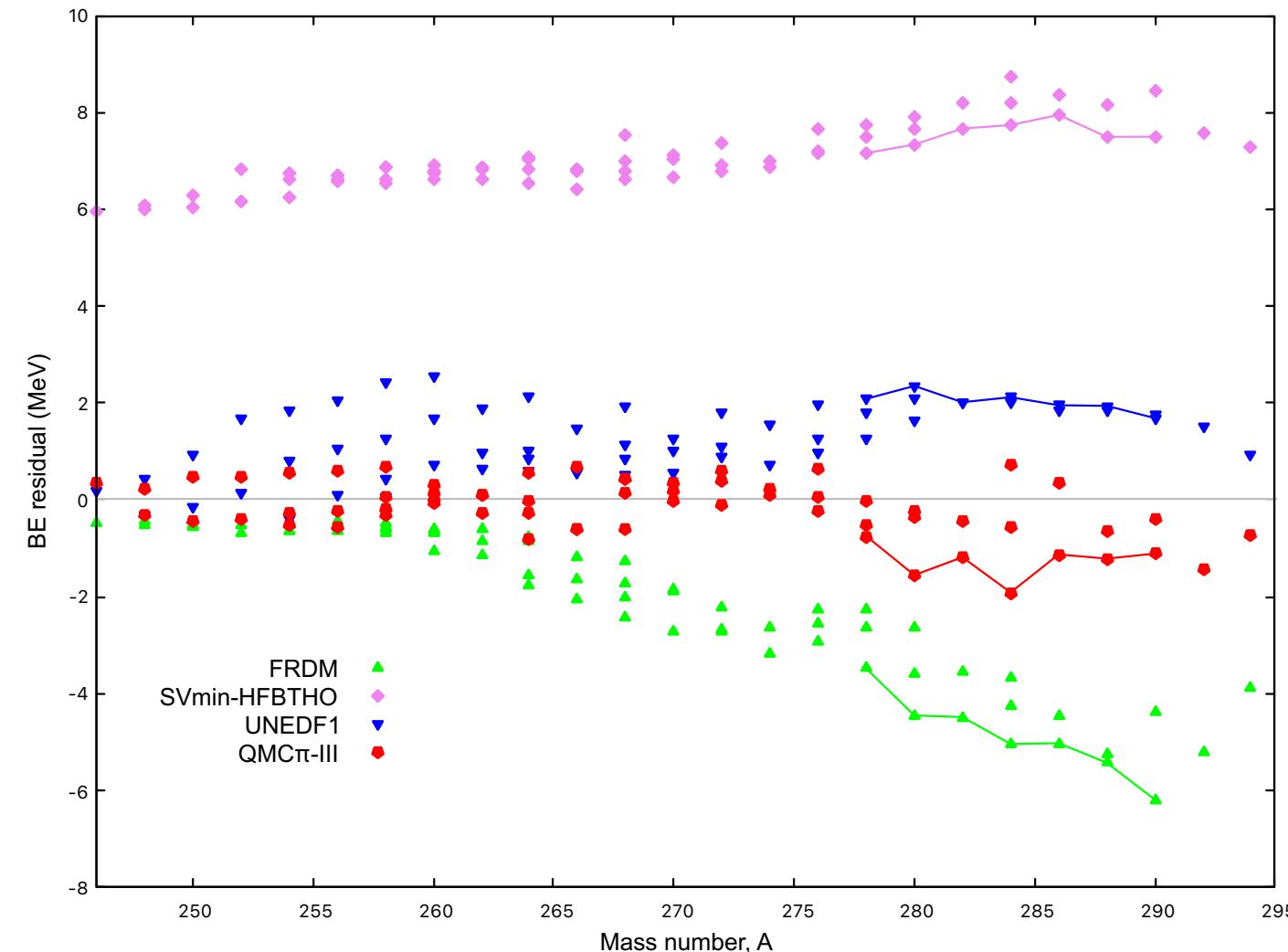
BE per nucleon and total BE residuals along Hs and Ds chains.

QMC predictions for binding energies of even-even superheavy elements (SHE)

Comparison of total BE residuals for $Z \geq 112$ from different QMC versions.

Element	Z	N	QMC-I	QMC π -I	QMC π -II	QMC π -III
Cn	112	164	-1.8	-1.2	2.7	0.2
	112	166	-1.7	-1.9	2.5	-0.8
	112	168	-1.4	-2.5	1.5	-0.6
	112	170	-0.8	-2.8	1.0	-0.6
	112	172	-0.4	-3.4	0.4	-1.1
Fl	114	170	-1.2	-1.7	1.6	0.3
	114	172	-0.5	-2.0	1.4	0.2
	114	174	0.1	-2.4	1.0	-0.6
Lv	116	174	-0.4	-1.2	2.4	-0.3
	116	176	0.3	-1.6	2.1	-1.3
Og	118	176	-0.7	-0.8	3.5	-0.8

Comparison of binding energies of even-even superheavy elements (SHE)

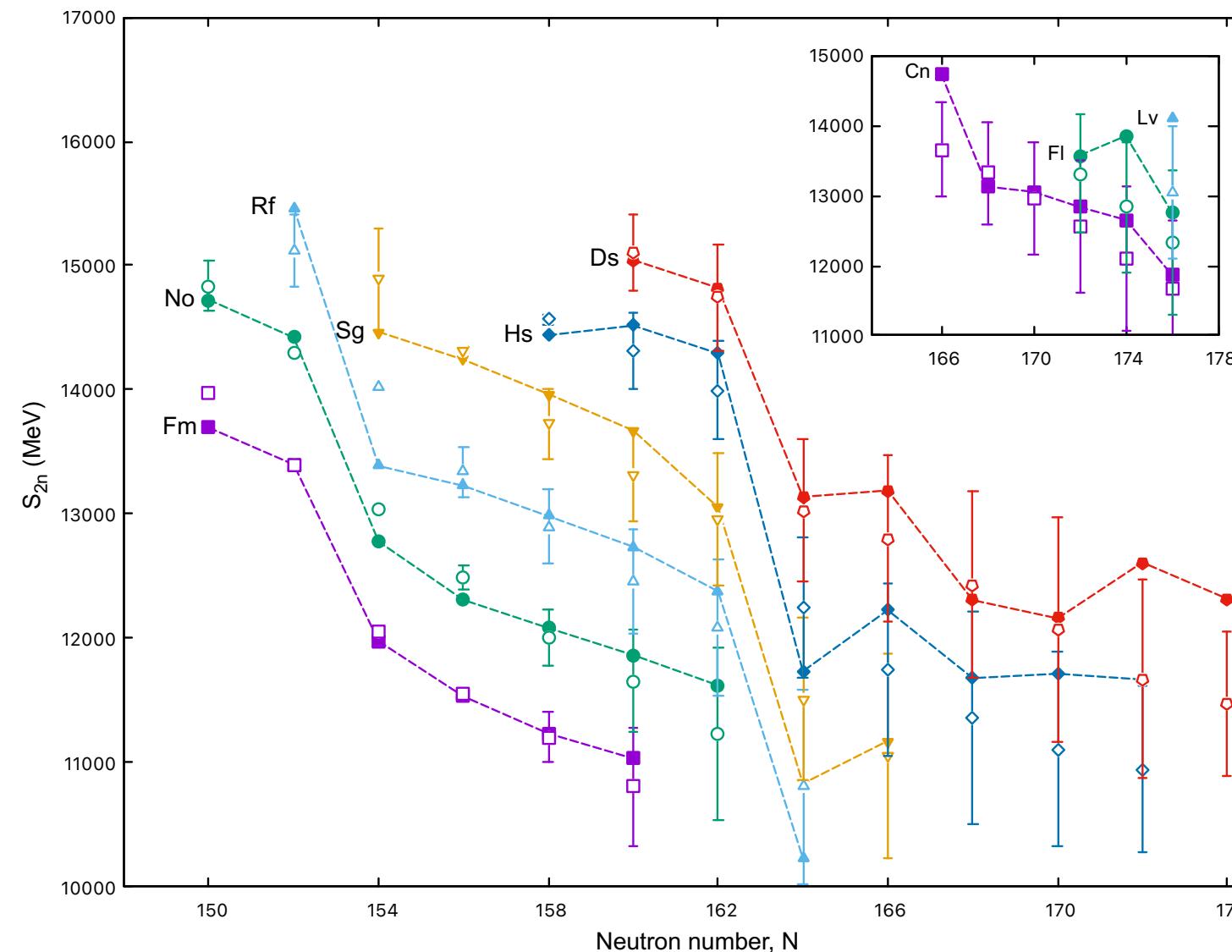


Comparison of SHE binding energy residuals from nuclear models plotted against mass number A and computed using AME 2020[2].

Comparison of *rms* percent deviations and *rms* residuals for SHE with available data

Model	<i>rms</i> % deviation	<i>rms</i> residual (MeV)
QMC π -III	0.03	0.52
QMC π -II [11]	0.11	2.04
QMC π -I [6]	0.12	2.42
QMC-I [5]	0.08	1.50
FRDM	0.11	2.25
SV-min	0.36	6.99
UNEDF1	0.07	1.31
DD-ME δ	0.12	2.28

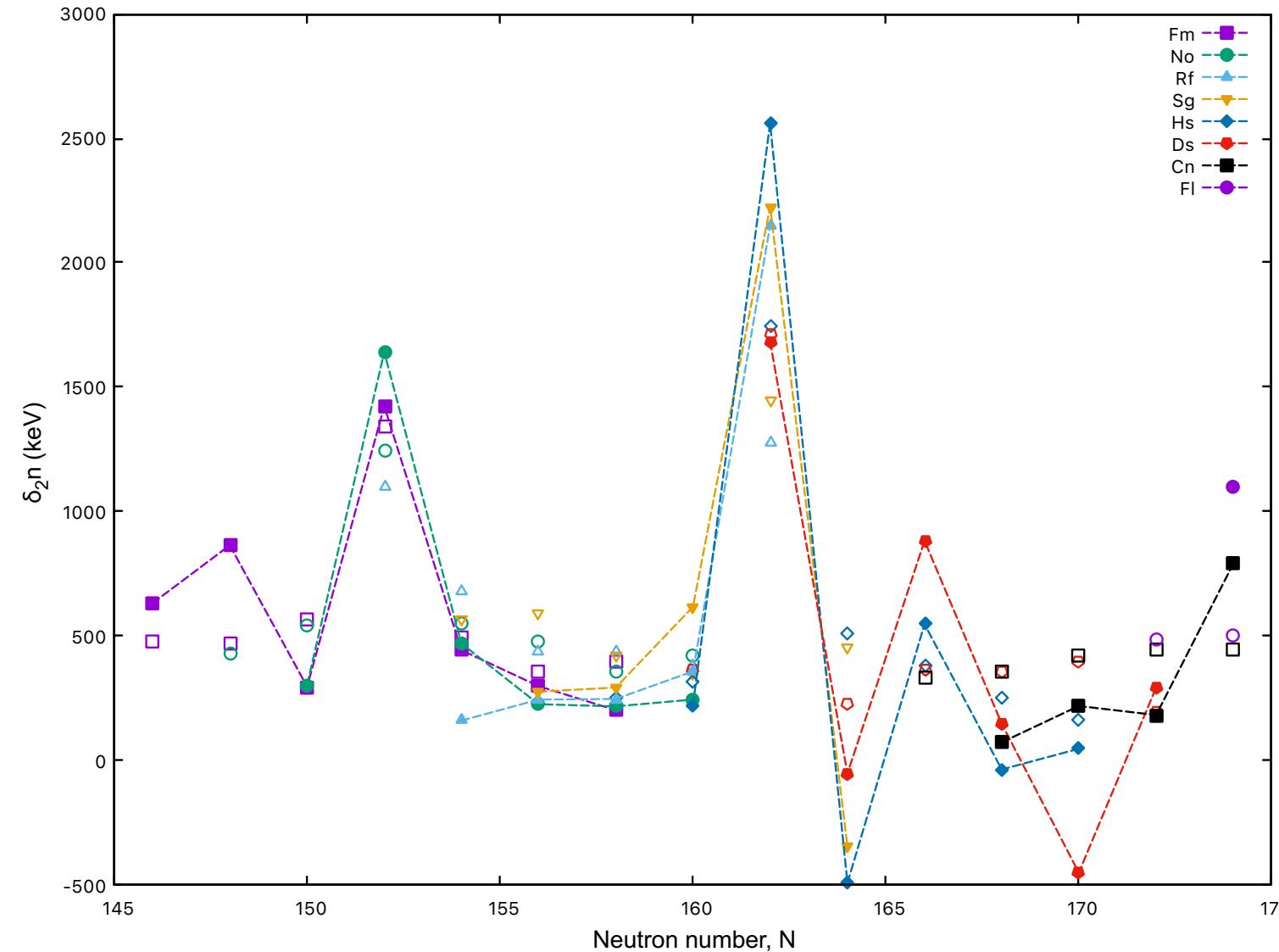
Subshell closures in the SHE region



Two-neutron separation energies, S_{2n} , plotted against neutron number.

Results from QMC π -III are shown as filled symbols and connected by lines while experimental data and errors are shown as empty symbols with vertical errorbars. Inset shows the S_{2n} values plotted against N for SHE with $Z \geq 112$.

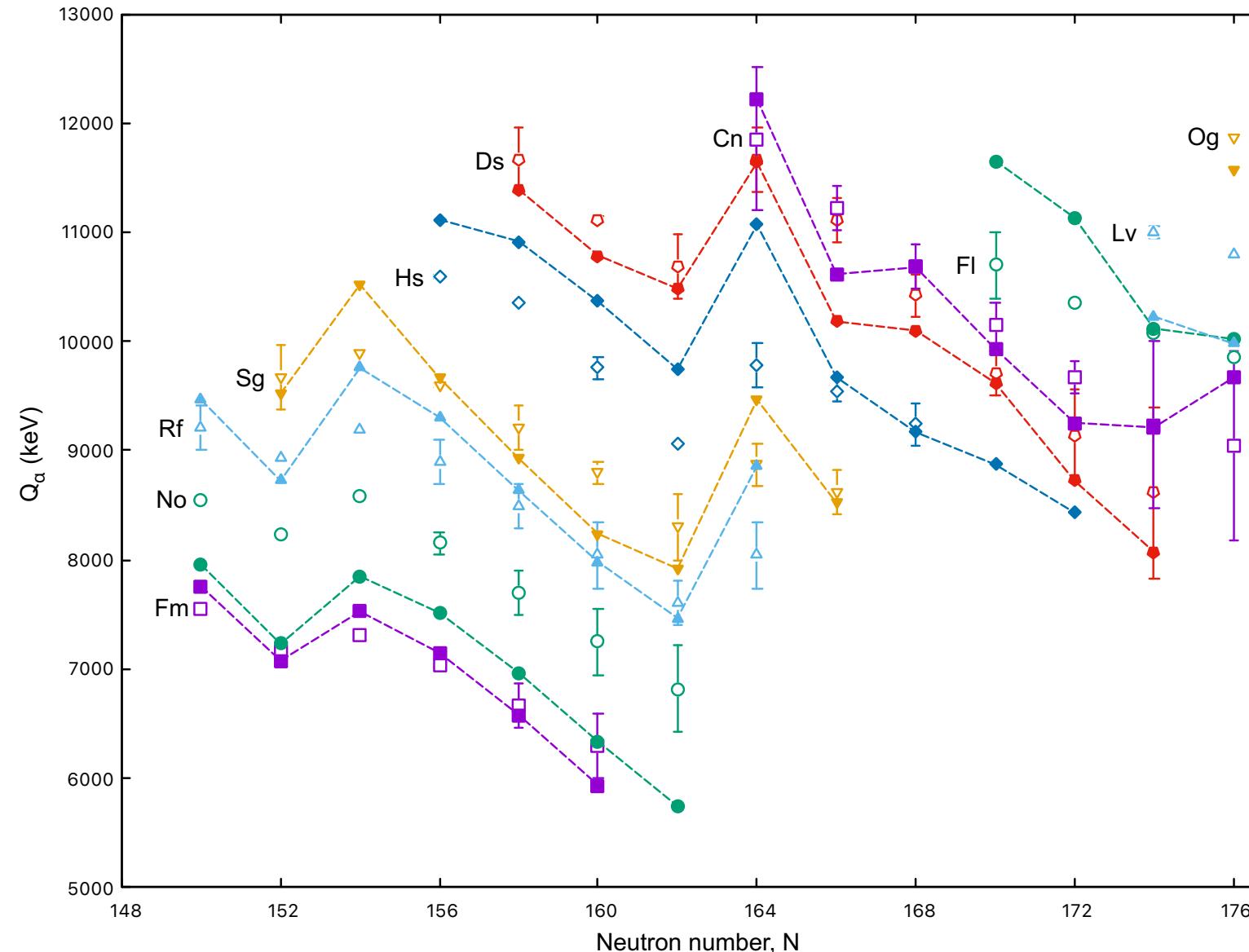
Subshell closures in the SHE region



Two-neutron shell gaps, δ_{2n} , plotted against neutron number.

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Q_α energies

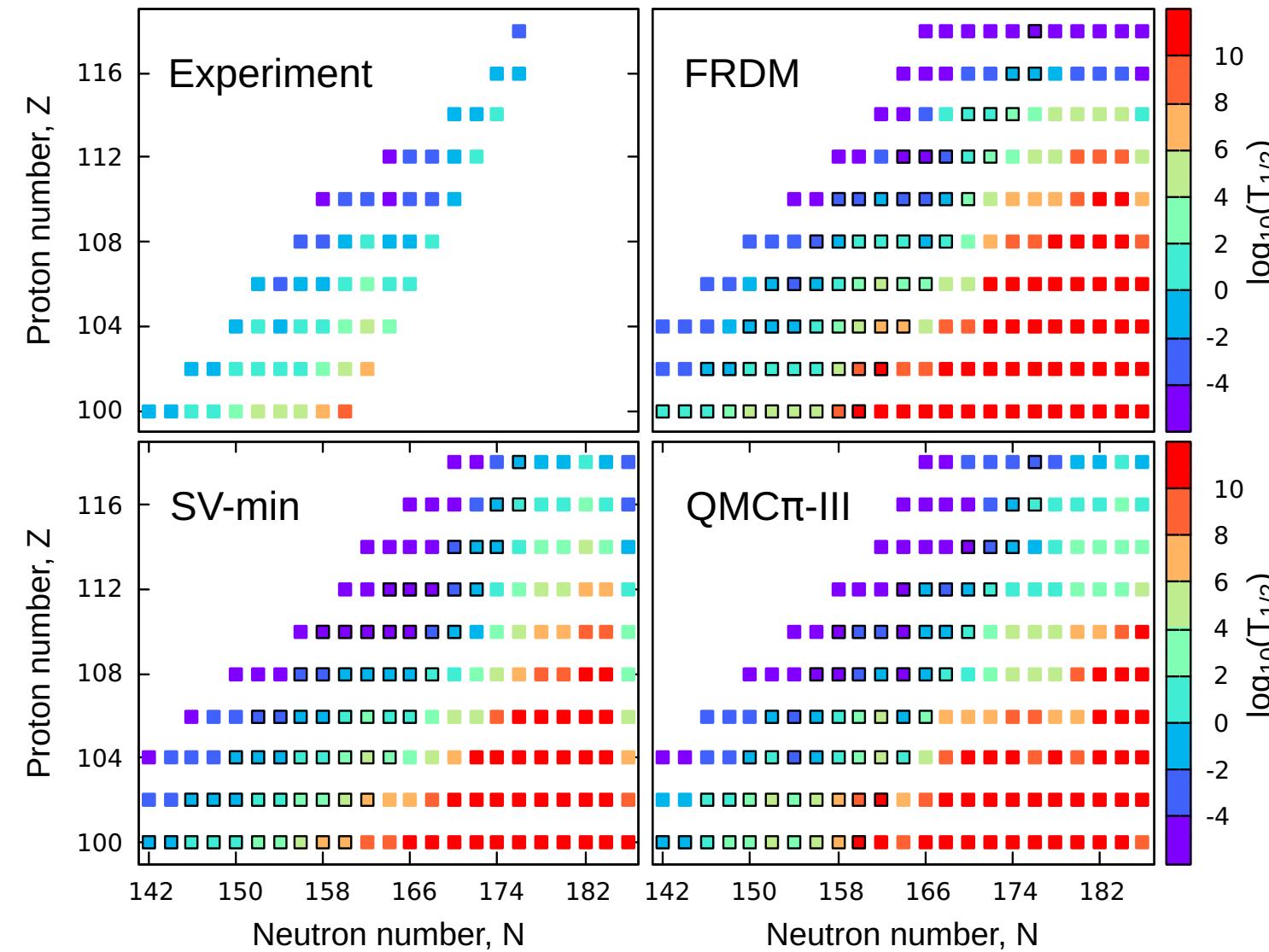


Comparison of Q_α energies from QMC π -III predictions and AME 2020 data.

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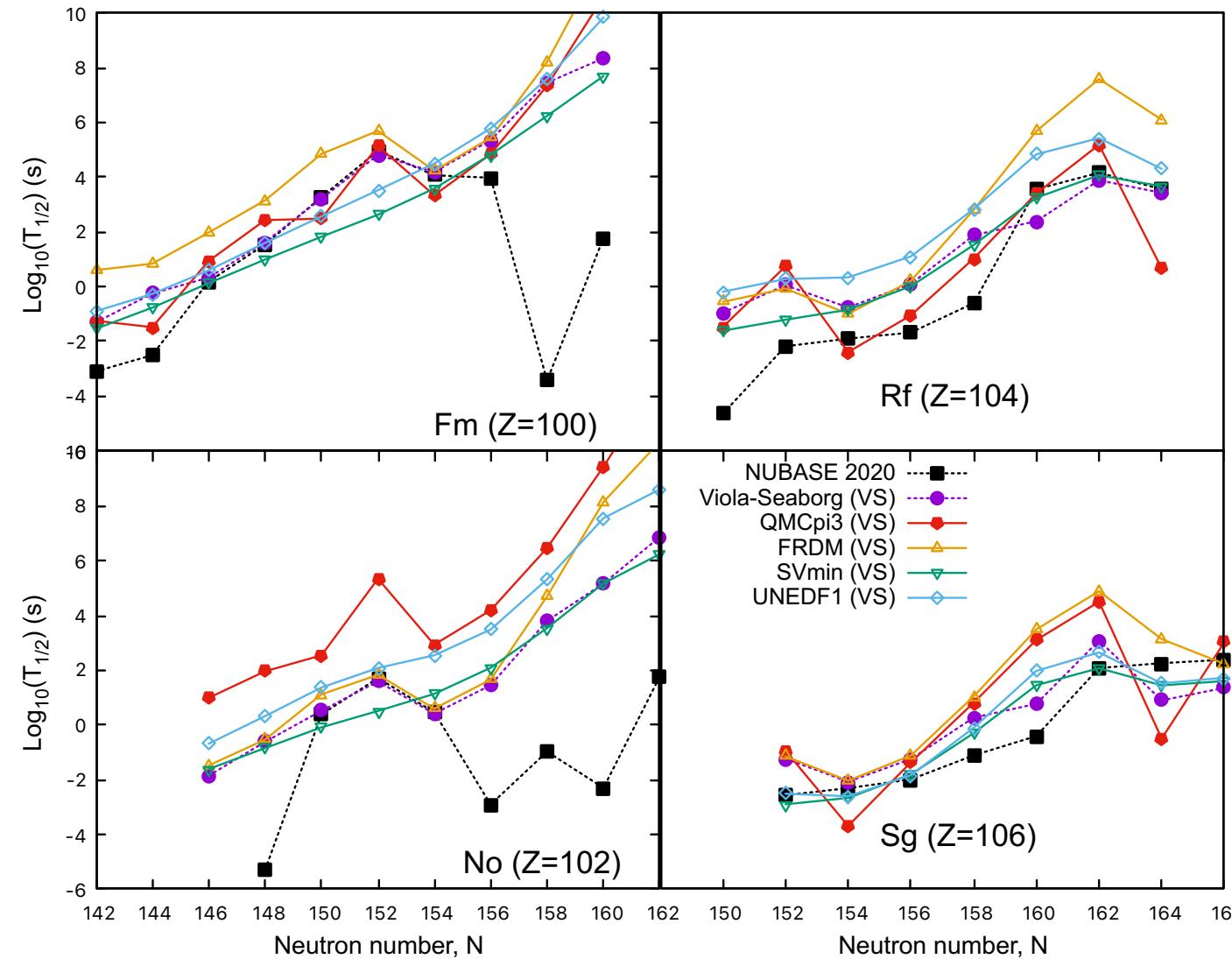
α decay half-lives

Viola-Seaborg relationship (VS): $\log_{10}(T_{1/2}) = \frac{aZ+b}{\sqrt{Q_\alpha}} + cZ + d + h_{log}$



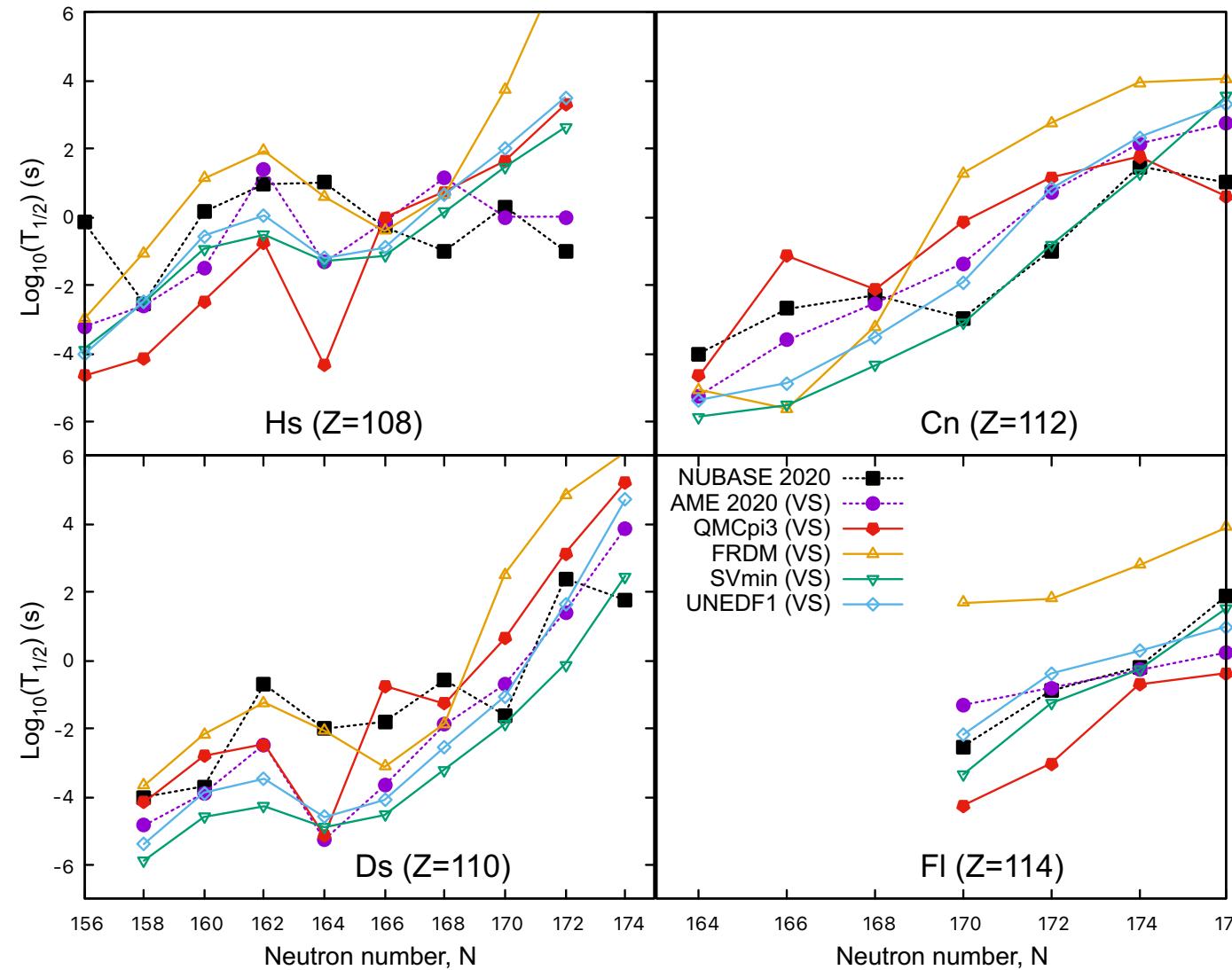
Comparison of $\log_{10}(T_{1/2})$ (s) in the SHE region.

α decay half-lives



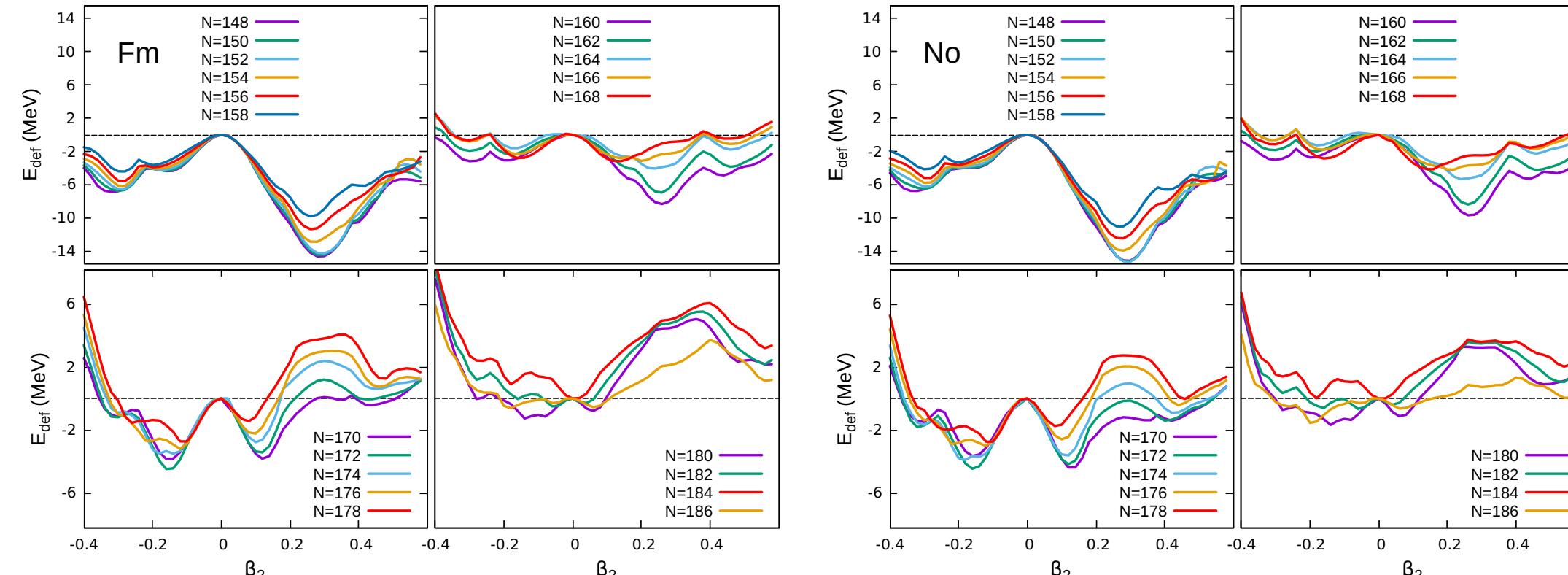
$\text{Log}_{10}(T_{1/2})$ values of Fm, No, Rf, and Sg isotopic chains from NUBASE 2020 data, AME 2020 Q_α data, predictions from QMC π -3, and other nuclear models computed using the Viola-Seaborg relationship.

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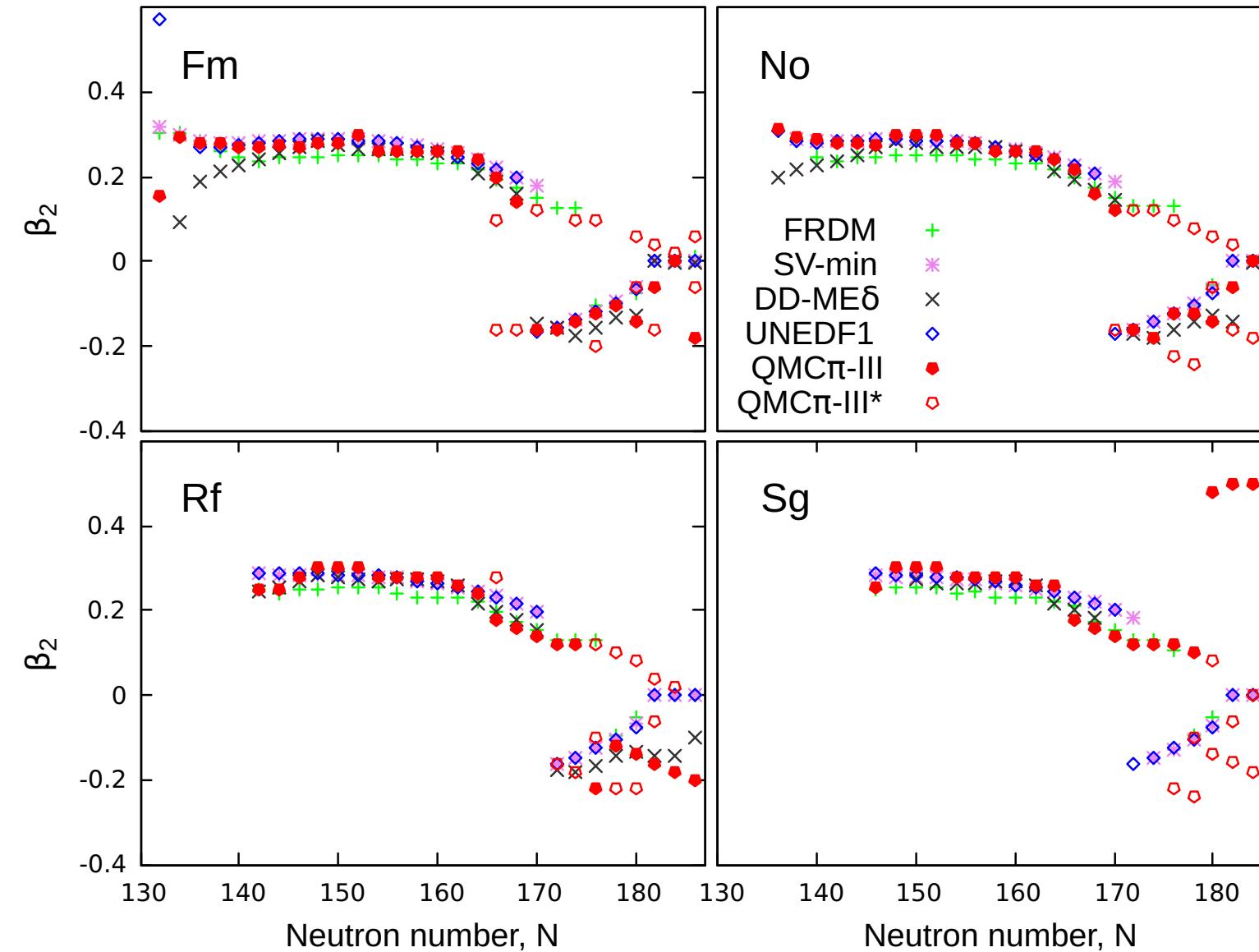
$\log_{10}(T_{1/2})$ values of Hs, Ds, Cn, and Fl isotopic chains from NUBASE 2020 data, AME 2020 Q_α data, predictions from QMC π -3, and other nuclear models computed using the Viola-Seaborg relationship.

Deformation properties



Deformation plots for even-even SHE with $100 \leq Z \leq 110$ and with N from the proton dripline up to $N \leq 186$ obtained from QMC π -3.

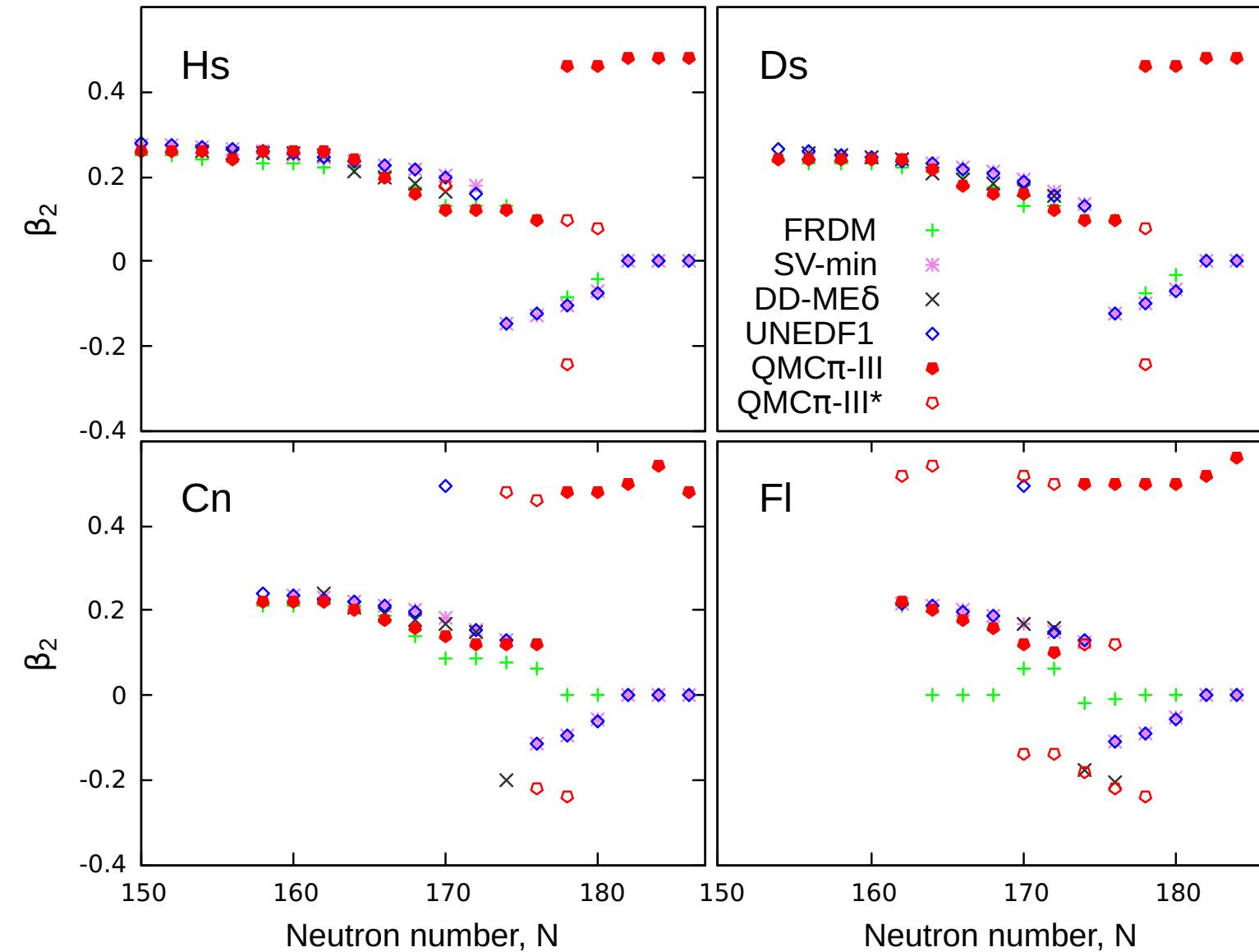
Deformation properties



Comparison of deformation parameter, β_2 , values along the Fm, No, Rf, and Sg isotopic chains from several nuclear models.

For QMC π -III, the first minima are shown as filled red symbols, while the other minima, which are very close to the first, are shown as empty red symbols and labelled 'QMC π -III*'.

Deformation properties



Comparison of deformation parameter, β_2 , values along the Hs, Ds, Cn and Fl isotopic chains from several nuclear models.

For QMC π -III, the first minima are shown as filled red symbols, while the other minima, which are very close to the first, are shown as empty red symbols and labelled 'QMC π -III*'.

Even-even nuclei across the nuclear chart

Root-mean-square residuals: $\sqrt{\frac{1}{d} \sum_i^n \sum_j^o (\bar{s}_{ij} - s_{ij})^2}$

Model	BE (MeV)	R_{ch} (fm)	# of parameters
QMC π -III	1.59	0.024	5
QMC π -II	2.34	0.029	5+2
QMC π -I	2.78	0.028	4+2
QMC-I	3.84	0.030	4+2
SV-min	3.64	0.024	11+3
UNEDF1	2.06	0.029	11+2
DD-ME δ	2.41	0.035	14+2
FRDM	0.89	-	17

Conclusion

Summary and conclusions

- The structure of finite nuclei was described using the QMC model.

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- The latest version of the QMC model performs very well in the superheavy region, yielding excellent results for binding energies and energy differences compared to available data.

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- The latest version of the QMC model performs very well in the superheavy region, yielding excellent results for binding energies and energy differences compared to available data.
- Improvement in the predictions was seen as the QMC model developed.
- In regions where data for ground-state observables are not yet available, QMC π -III showed comparable results to other nuclear models even though there are significantly fewer model parameters.

Thank you for your attention.

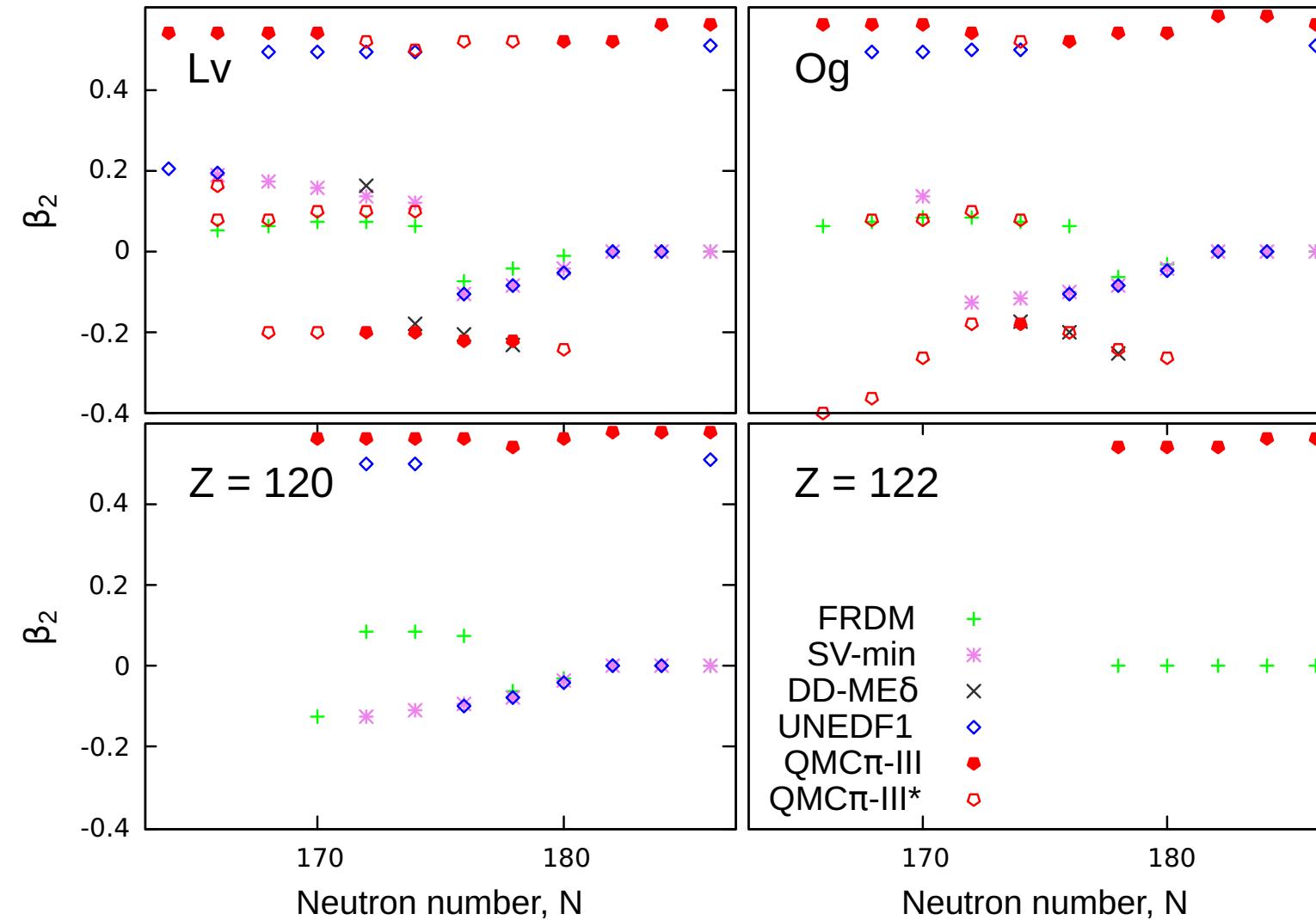
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References II

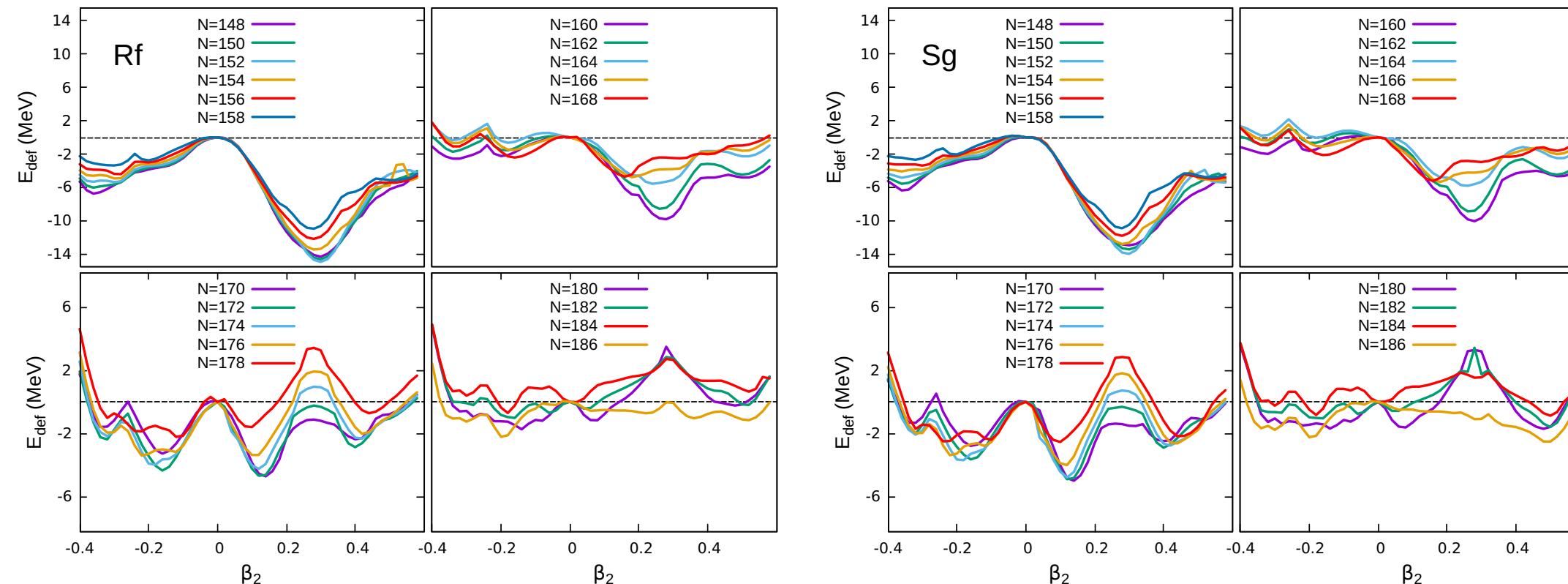
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Deformation properties



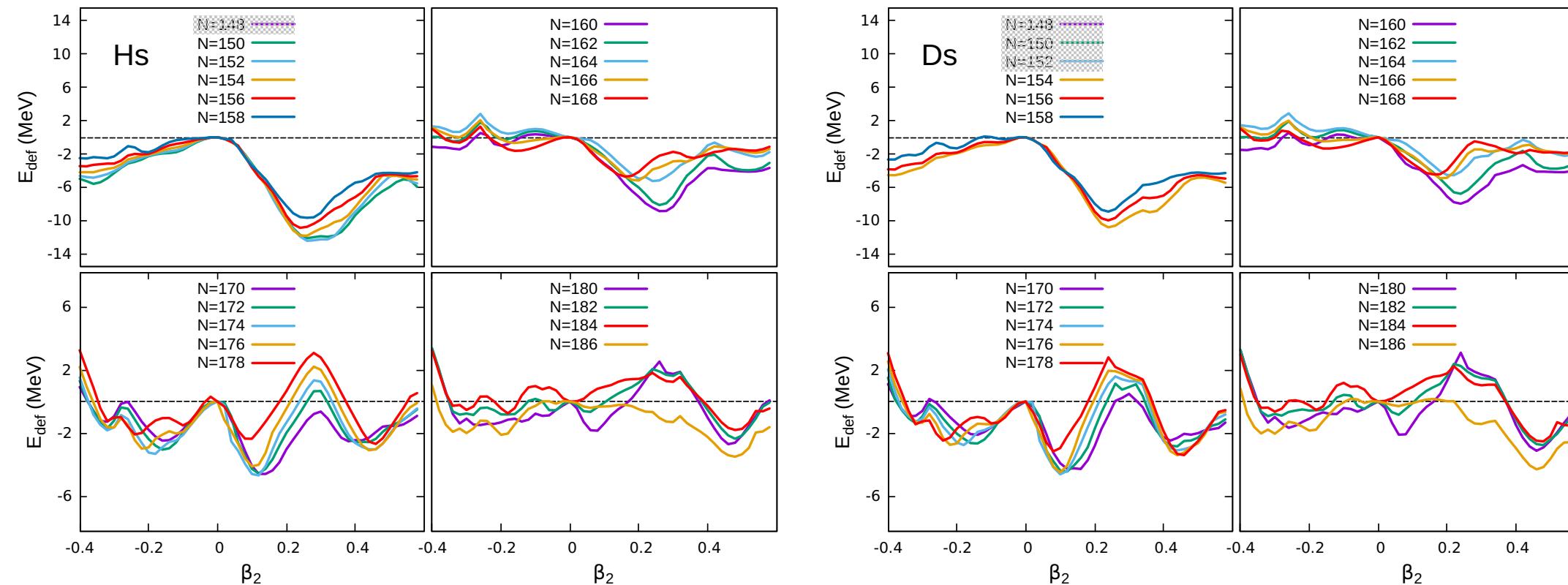
Same as in Figure 37 but for Hs, Ds, Cn and Fl isotopic chains.

Deformation properties



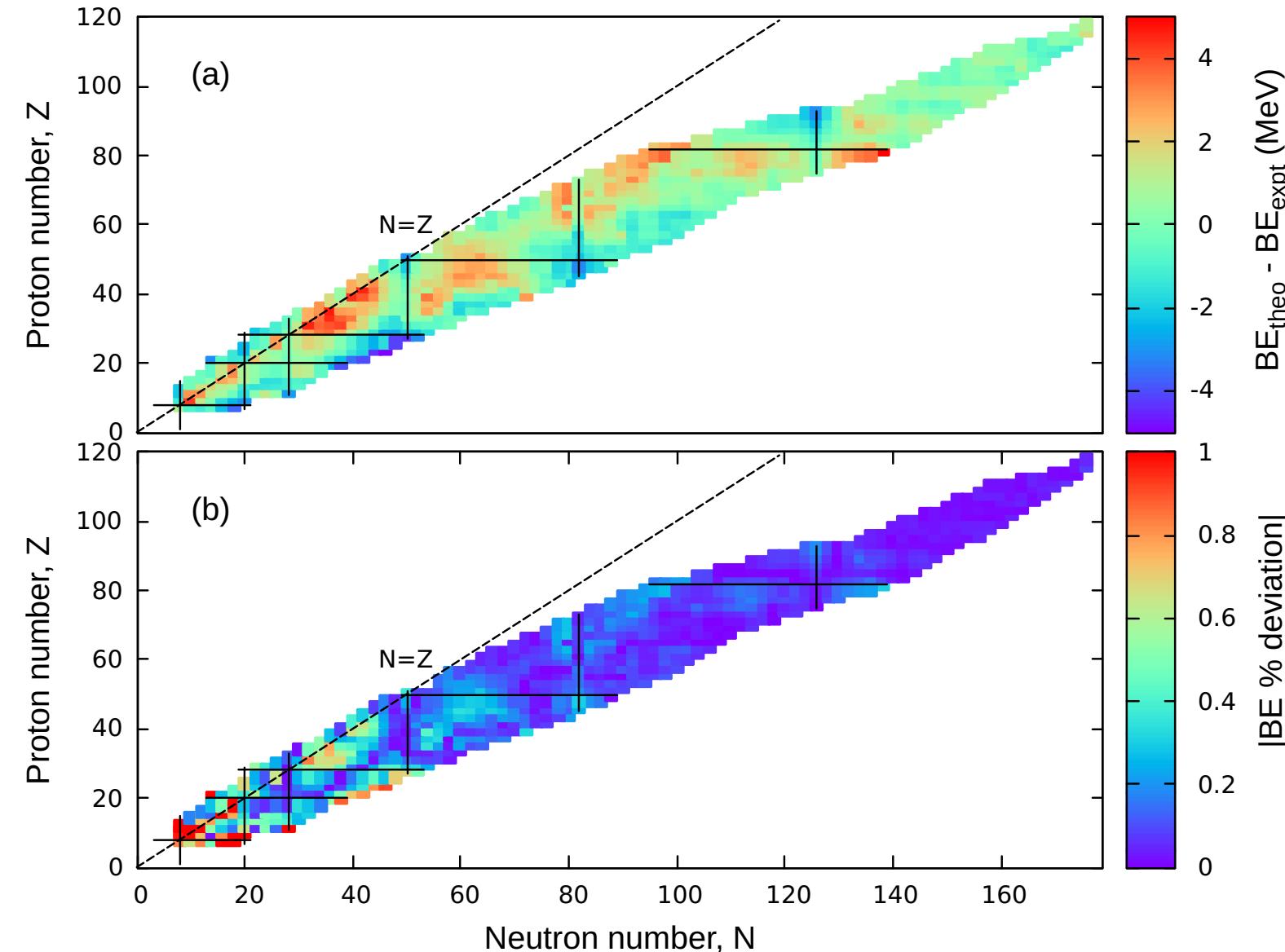
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Deformation properties



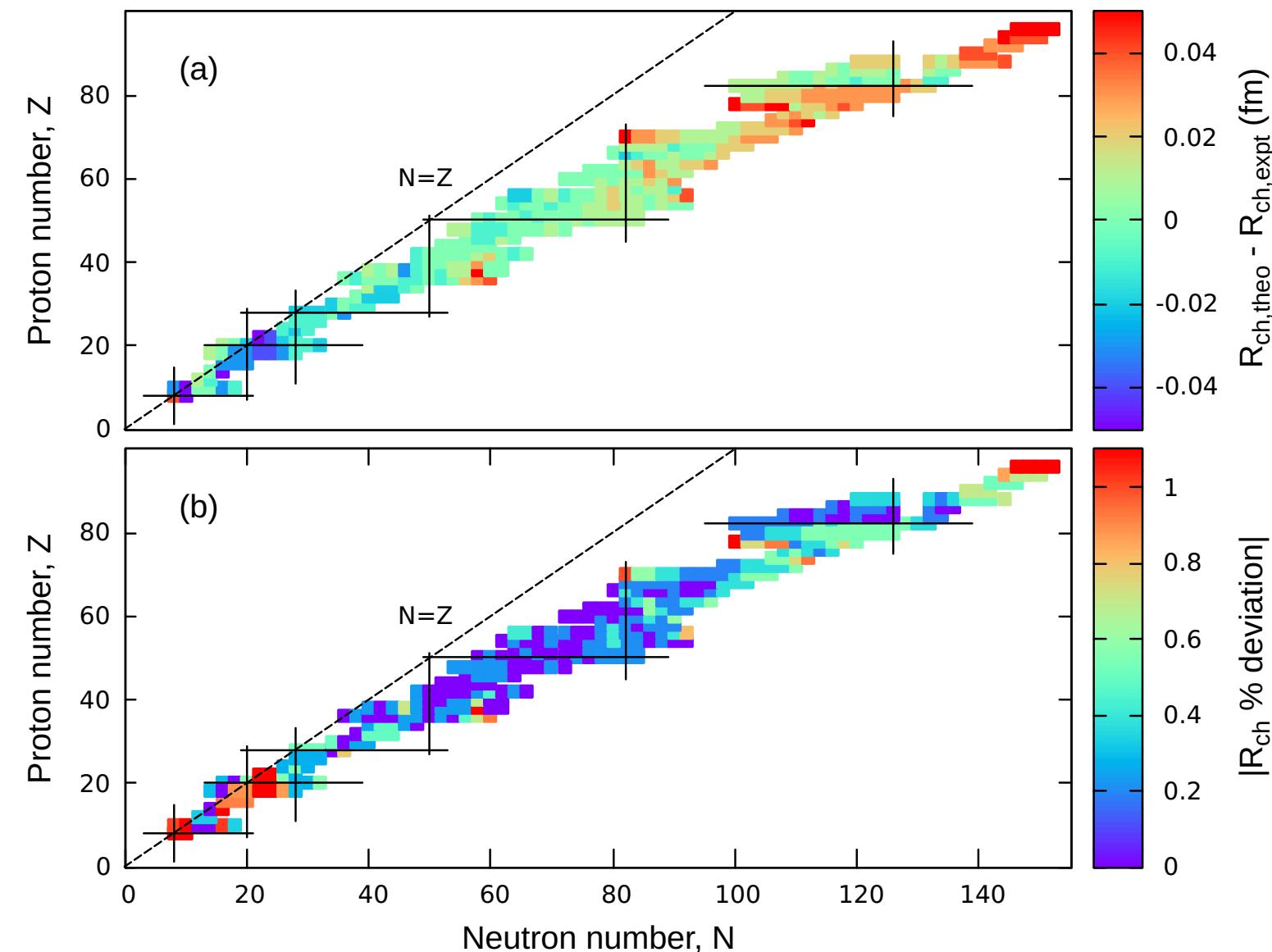
Deformation plots for even-even SHE with $100 \leq Z \leq 110$ and with N from the proton dripline up to $N \leq 186$ obtained from QMC π -3.

Even-even nuclei across the nuclear chart



QMC π -III predictions for (a) BE residuals and (b) absolute BE % deviation for even-even nuclei ($Z > 8$) with known masses.

Even-even nuclei across the nuclear chart



QMC π -III predictions for (a) R_{ch} residuals and (b) absolute R_{ch} % deviation for even-even nuclei ($Z > 8$) with known masses.