

Centre vortex geometry at finite temperature



Derek Leinweber

In collaboration with:

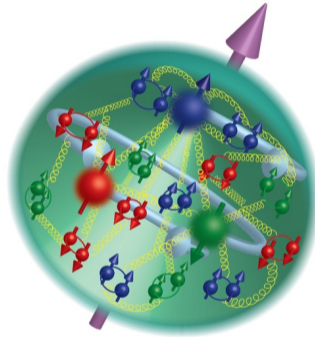
Jackson Mickley

James Biddle & Waseem Kamleh



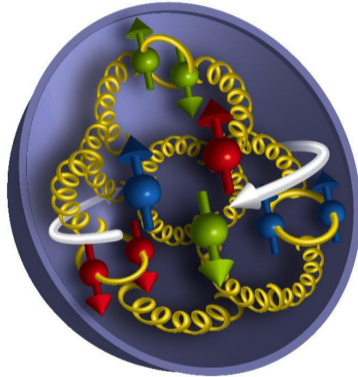
THE UNIVERSITY
of ADELAIDE

Images of the proton



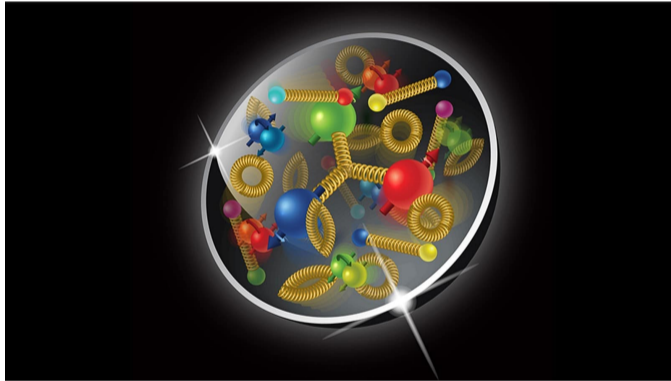
- Artistic rendering of proton structure revealing its intricate and dynamic system of quarks and gluons. (Image by Argonne National Laboratory.)

Images of the proton



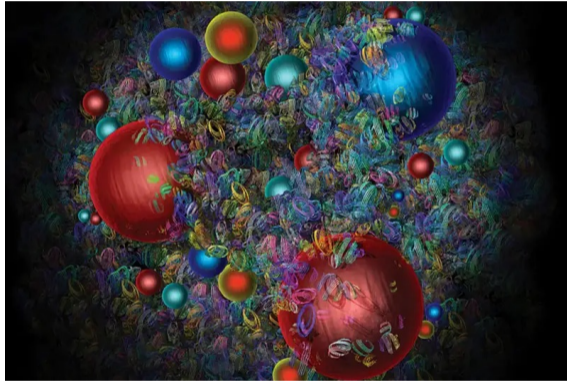
- The internal structure of a proton, with quarks, gluons, and quark spin shown. The nuclear force acts like a spring, . . . Brookhaven National Laboratory.

Images of the proton



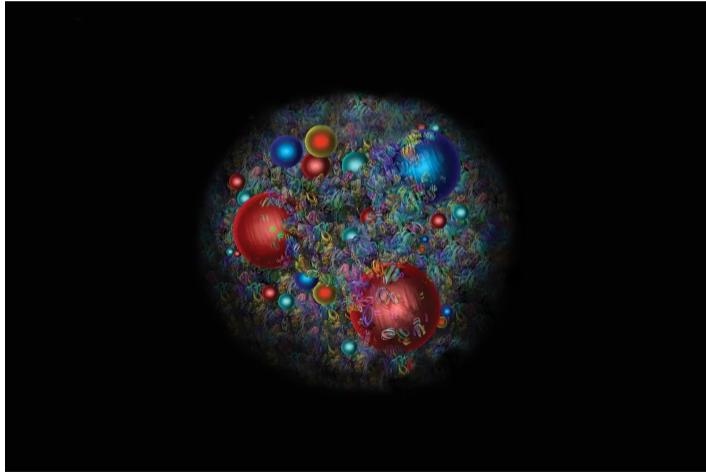
- Quark and gluon sea: in this illustration of the proton the large spheres represent the three valence quarks, the small spheres other quarks and the springs the gluons holding them together. (Courtesy: Brookhaven National Laboratory.)

Images of the proton

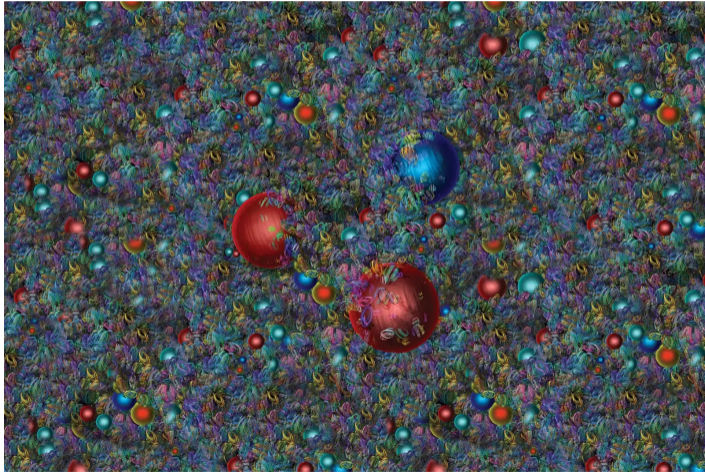


- **Glorious complexity** An artist's impression of the mayhem of quarks and gluons inside the proton. Credit: Daniel Dominguez (CERN).

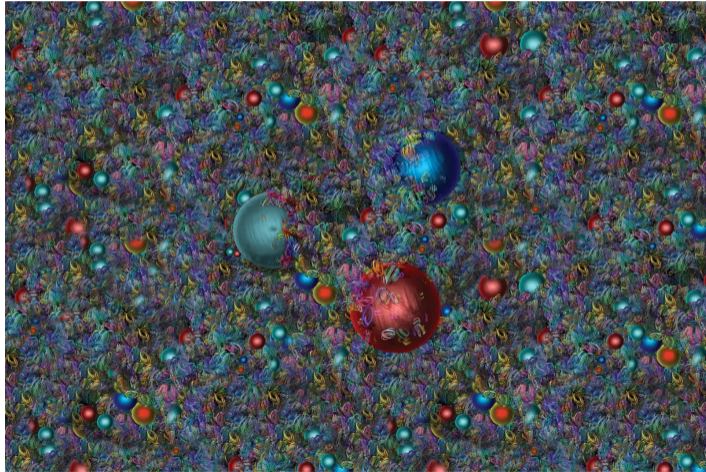
Images of the proton



The reality of a lattice QCD calculation



Valence quarks embedded in ground-state QCD-vacuum fields



Introduction

- What aspect of the nontrivial ground-state field structure confines quarks?
- Is there an essential aspect that captures the salient features of QCD?
 - Confinement.
 - Dynamical generation of mass via chiral symmetry breaking.

Centre-Vortices in the Ground-State QCD-Vacuum Fields

- Foundations
 - Spaghetti Vacuum, a condensate of vortices with finite thickness.
 - H. B. Nielsen and P. Olesen, Nucl. Phys. B **160** (1979), 380-396
 - Importance of the Z_N centre of the $SU(N)$ gauge group.
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- Renewed Excitement...
 - L. Del Debbio, M. Faber, J. Greensite,... Phys. Rev. D **55** (1997), 2298-2306
 - M. Faber, J. Greensite and S. Olejnik, Phys. Rev. D **57** (1998), 2603-2609
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 - R. Bertle, M. Faber, J. Greensite and S. Olejnik, JHEP **03** (1999), 019
 - M. Faber, J. Greensite, S. Olejnik and D. Yamada, JHEP **12** (1999), 012

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- Review: "The confinement problem in lattice gauge theory,"
 J. Greensite, Prog. Part. Nucl. Phys. **51** (2003), 1. 500+ citations

Centre-Vortices in the Ground-State QCD-Vacuum Fields

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- What do Centre Vortices look like?

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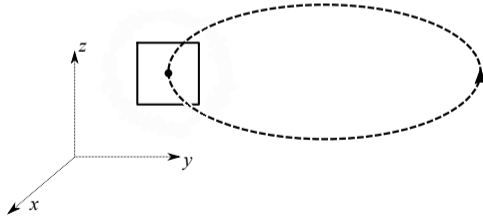
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Centre-Vortices in the Ground-State QCD-Vacuum Fields

- What are Centre Vortices and how do we locate them?
- What do Centre Vortices look like?
- How does centre-vortex geometry change at finite temperature?
- What happens to centre vortices in the process of deconfinement?

What Are Centre Vortices?

- Centre vortices in 3D are tube-like topological defects present in the QCD vacuum.
- We locate thin vortex lines on the lattice.
- The vortex line can be thought of as the 'axis of rotation' of the vortex.



A centre vortex (dashed line) intersecting a lattice plaquette (solid square).

Vortex Structure in the Colour Fields of the QCD Vacuum



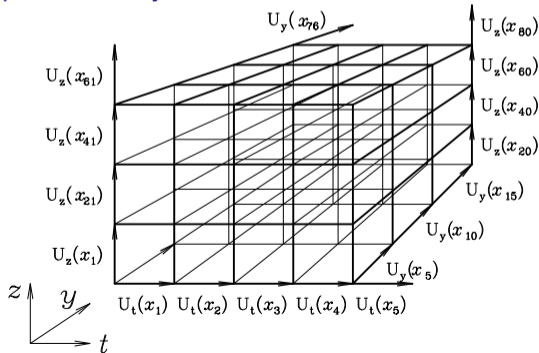
How do you find centre vortices?

Lattice Links

- On the lattice, the **gluon-field** is encoded in terms of the **link variable**

$$U_{\mu}(x) \simeq \exp(i a g A_{\mu}(x)) ,$$

a 3×3 complex special-unitary matrix.

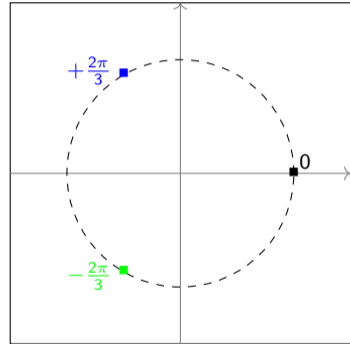


Centre Group of SU(3)

- Centre elements commute with every group element,

$$z = \exp\left(\frac{2\pi i}{3} m\right) I, \quad m \in \{-1, 0, 1\} \simeq \mathbb{Z}_3.$$

- Each of the three centre phases corresponds to a centre element of SU(3).



1. Maximal Centre Gauge

- Gauge transformations bring the links close to an element of the group centre

$$z = \exp\left(\frac{2\pi i}{3} m\right) I, \quad m \in \{-1, 0, +1\}.$$

- This is done by maximising the functional

$$R = \sum_x \sum_{\mu} |\operatorname{tr}[U_{\mu}(x)]|^2$$

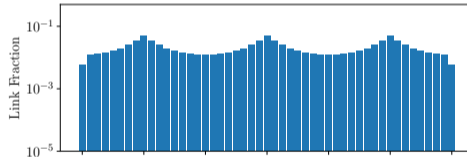
- This is called **Maximal Centre Gauge**

1. Maximal Centre Gauge

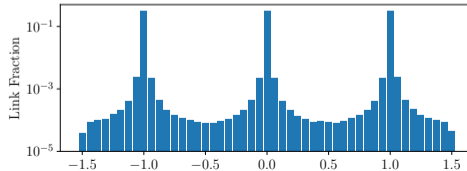
- Distribution of link phases.

$$\text{tr } U_{\mu}^{\text{MCG}}(x) = \underbrace{r_{\mu}(x)}_{\text{real}} \exp \left(\underbrace{\frac{2\pi i}{3} \phi_{\mu}(x)}_{-\pi < \text{phase} \leq \pi} \right), \quad -\frac{3}{2} < \phi_{\mu}(x) \leq \frac{3}{2}.$$

- $\phi_{\mu}(x)$ before gauge fixing.



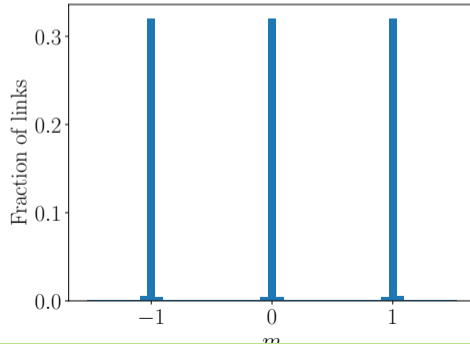
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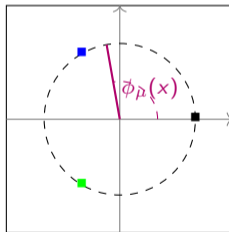
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2. Centre Projection

- Project onto $Z(3)$

$$U_{\mu}^{\text{MCG}}(x) \rightarrow Z_{\mu}(x) = \exp\left(\frac{2\pi i}{3} m_{\mu}(x)\right) I, \quad m_{\mu}(x) \in \{-1, 0, +1\}.$$



- Eight degrees of freedom are replaced by one of the three cube-roots of 1.

2. Centre Projection

- This projection allows us to define 3 sets of configurations:
 - Untouched - $U_\mu(x)$
 - Vortex Only - $Z_\mu(x)$
 - Vortex Removed - $R_\mu(x) = Z_\mu^\dagger(x) U_\mu(x)$

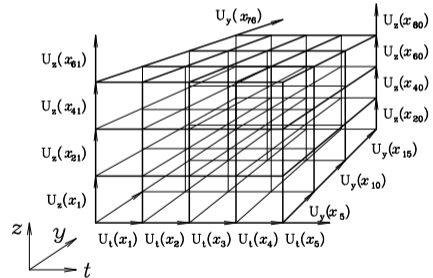
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- Untouched Ensembles
 - $32^3 \times N_t$ Iwasaki pure gauge (PG), spacing $a = 0.100$ fm
 - $32^3 \times 64$ dynamical 2 + 1 flavour, spacing $a = 0.1022$ fm, $m_\pi = 701$ MeV
 - $32^3 \times 64$ dynamical 2 + 1 flavour, spacing $a = 0.0933$ fm, $m_\pi = 156$ MeV
 - S. Aoki, *et al.* (PACS-CS), Phys. Rev. D **79**, 034503.

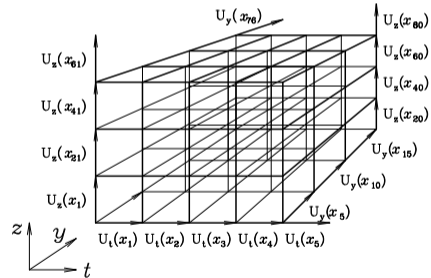
3. Identifying Vortices

- Examine the product of $Z_\mu(x)$ around each elementary square (plaquette).



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- Each plaquette takes a value from $Z(3)$.



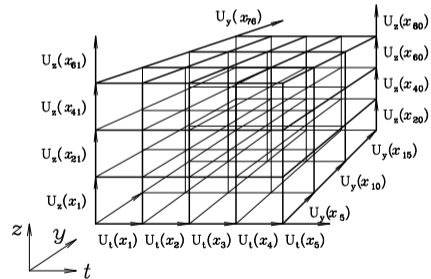
3. Identifying Vortices

- Examine the product of $Z_\mu(x)$ around each elementary square (plaquette).
- Each plaquette takes a value from $Z(3)$.

- Non-trivial plaquettes with values

$$\exp\left(\frac{2\pi i}{3} m\right) \neq 1, \quad \text{i.e. } m \in \{-1, +1\},$$

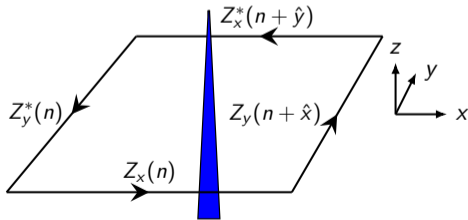
identify our thin vortices.



What do centre vortices look like?

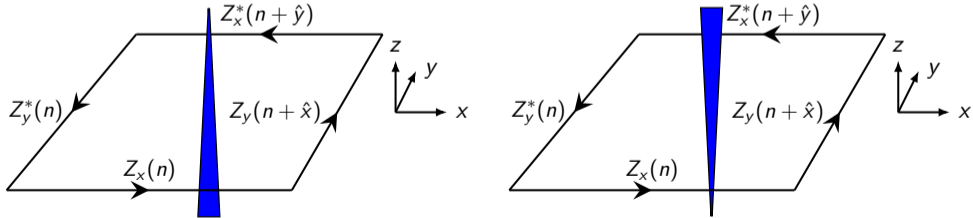
Rendering Projected Vortices

- Vortex sheets are sliced to vortex lines in a 3D slice of the 4D lattice.
- Flow of centre charge $+1$ is indicated using a right-handed coordinate system.
- For example,
 - An $m = +1$ vortex in the x - y plane is plotted in the $+\hat{z}$ direction.

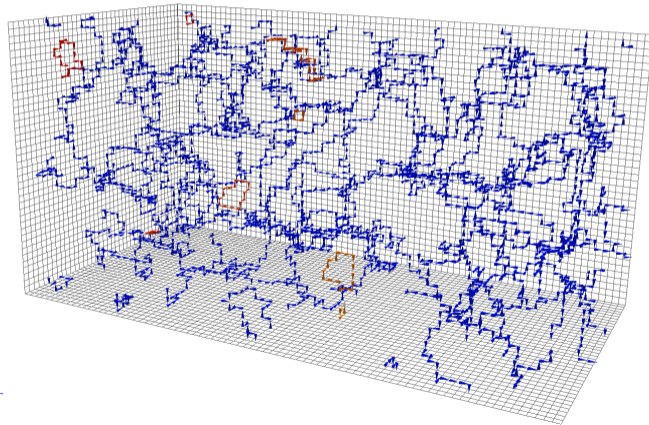


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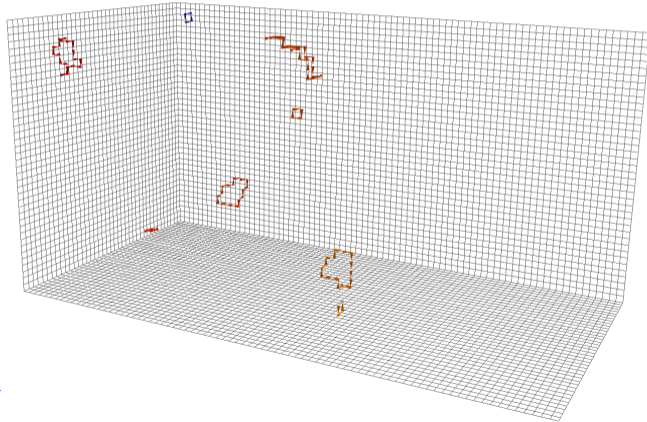
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 - An $m = -1$ vortex in the x - y plane is plotted in the $-\hat{z}$ direction.



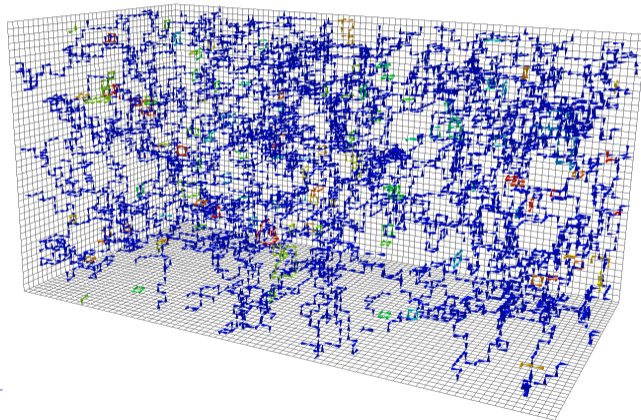
Vortices on a Pure-Gauge $32^3 \times 64$ Lattice



Secondary Loops on a Pure-Gauge $32^3 \times 64$ Lattice

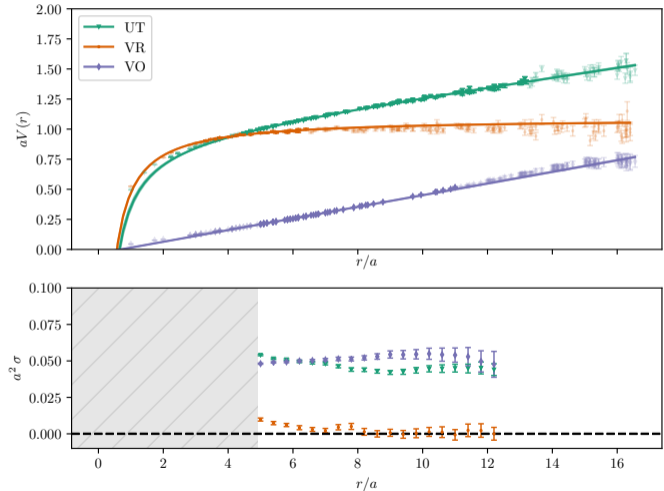


2 + 1 Flavour $32^3 \times 64$ Dynamical-Fermion Lattice $m_\pi = 156$ MeV



Static Quark Potential

- **Vortices** capture the **full** string tension.
- **Vortex removal** leaves no residual confining potential.
- Centre vortices are the origin of confinement in QCD.

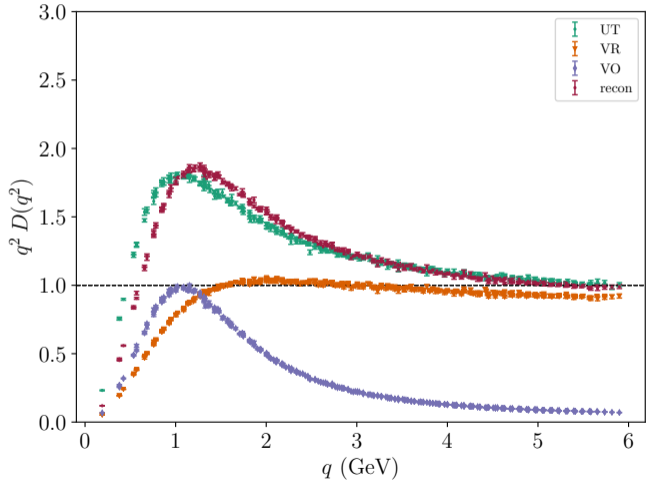


Landau-Gauge Gluon Propagator

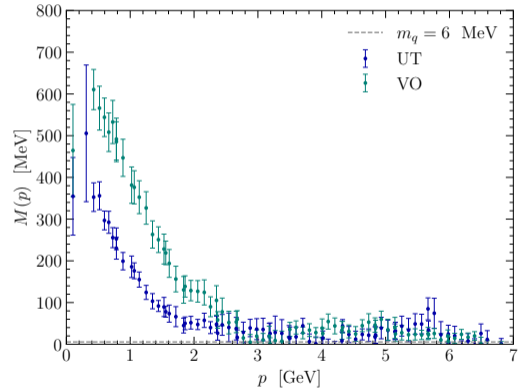
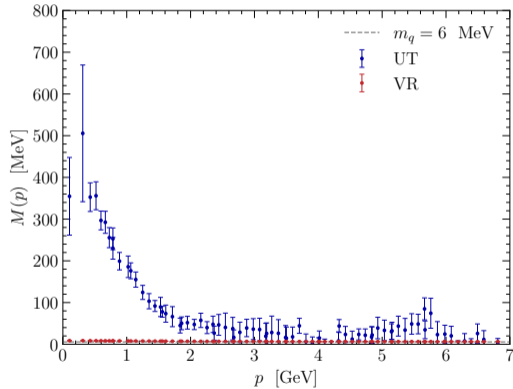
- Scalar gluon propagator

$$D(q^2) \equiv \frac{Z(q^2)}{q^2} \rightarrow \frac{1}{q^2} \text{ at tree level}$$

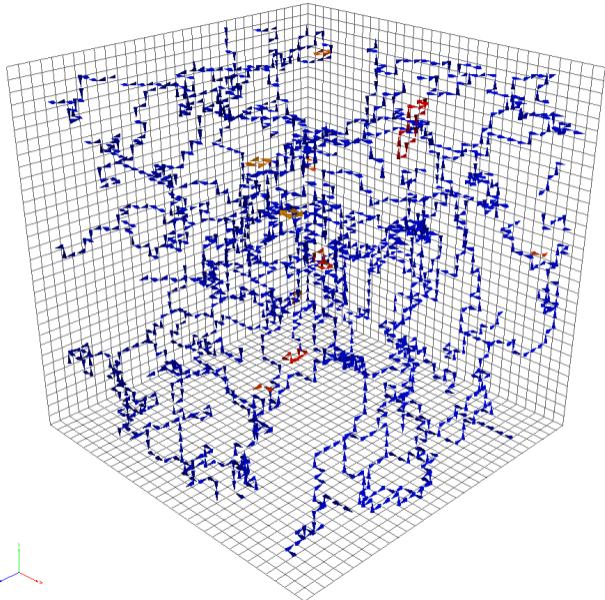
- **Vortex Removal (VR)** almost eliminates infrared enhancement.
- **Vortex-Only (VO)** configurations capture the long-distance physics.
- **Reconstructed** propagator.



Dynamical Mass Generation in the Quark Propagator



Centre Vortex Structure at Finite Temperature

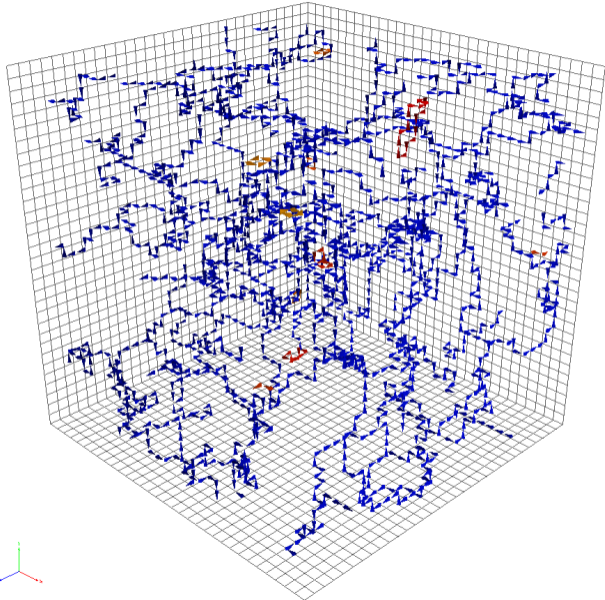


- Temporal (fixed t) slice
- Viewing spatial plaquettes
 - x - y , x - z , y - z plaquettes
- $32^3 \times N_t$ lattices

| N_t | T (MeV) | T/T_c |
|-------|-----------|---------|
| 64 | 31 | 0.11 |
| 12 | 164 | 0.61 |
| 8 | 247 | 0.91 |
| 6 | 329 | 1.22 |
| 5 | 395 | 1.46 |
| 4 | 493 | 1.83 |

$T_c = 270$ MeV



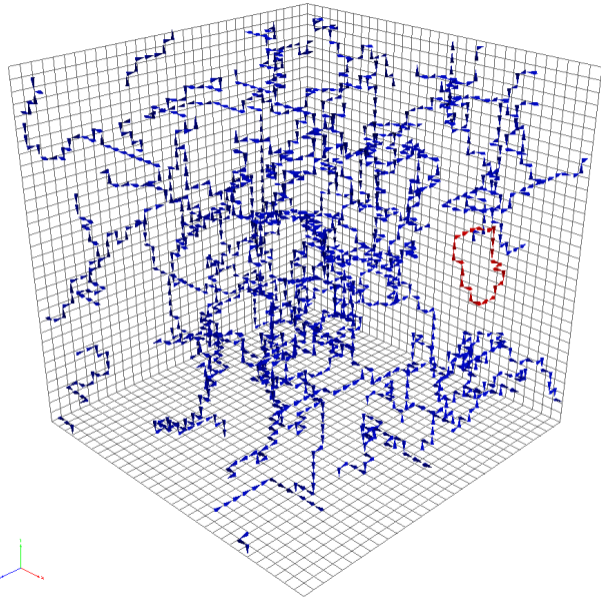


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- $N_t = 12$ lattice

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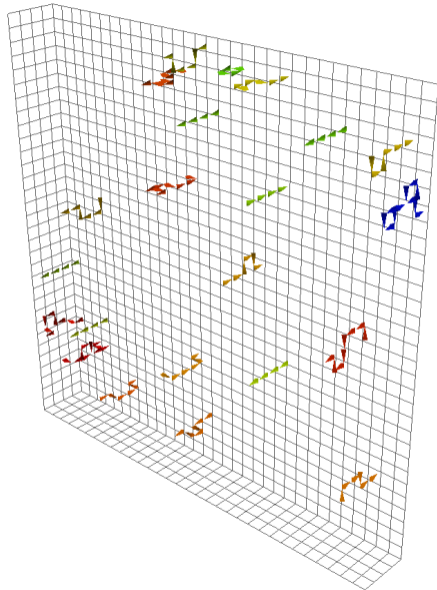
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- Spatial (fixed x) slice
- Viewing mixed plaquettes
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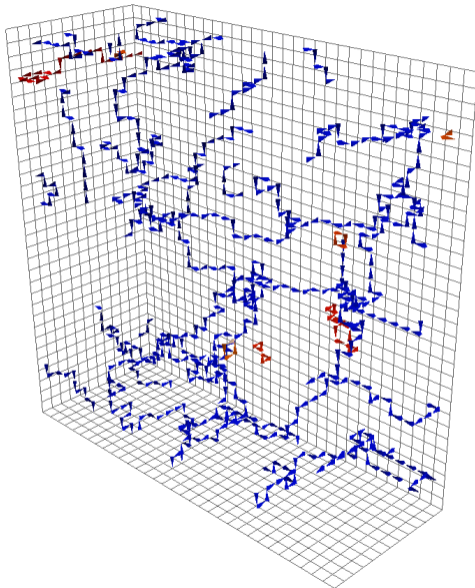
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Quantifying Centre Vortex Structure at Finite Temperature

Temporal Alignment of the Vortex Sheet

- Denote the centre charge, m , of a plaquette as

$$P_{\mu\nu}(\vec{x}, t) = \exp\left(\frac{2\pi i}{3} m_{\mu\nu}(\vec{x}, t)\right) \mathbb{I}, \quad \text{with } m_{\mu\nu}(\vec{x}, t) \in \{-1, 0, 1\}.$$

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- Identify cases where a vortex at one time remains in place at a later time.
- Define the indicator function

$$\chi_{\mu\nu}(\mathbf{x}, t; \tau) = \begin{cases} 1, & m_{\mu\nu}(\mathbf{x}, t) m_{\mu\nu}(\mathbf{x}, t + \tau) > 0 \\ 0, & \text{otherwise} \end{cases},$$

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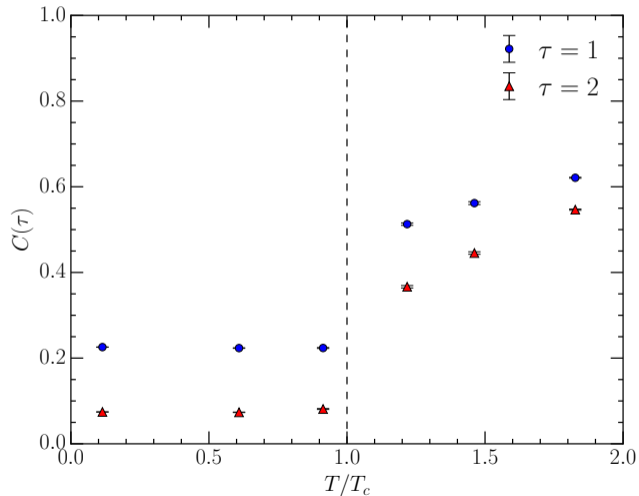
- Calculate the correlation measure

$$C(\tau) = \frac{1}{N_{\text{vor}} N_t} \sum_{\substack{\mathbf{x}, t, \\ i, j}} \chi_{ij}(\mathbf{x}, t; \tau),$$

where N_{vor} , the average number of vortices per slice, ensures $C(\tau = 0) = 1$.

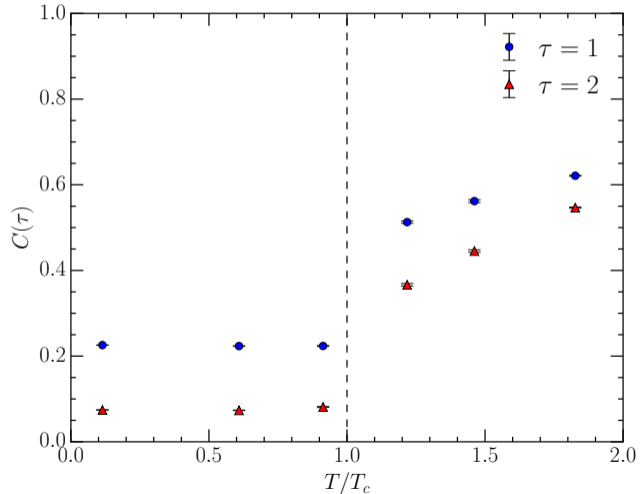
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- $C(\tau)$ changes significantly across the critical temperature, T_c .



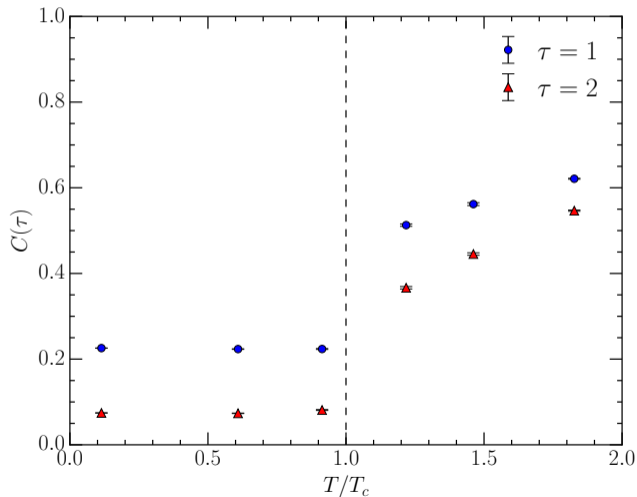
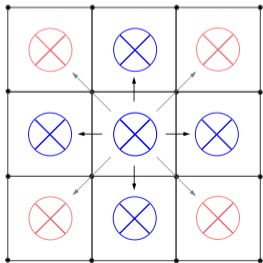
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- $C(\tau)$ changes significantly across the critical temperature, T_c .
- Rising values indicate continued alignment with temperature.



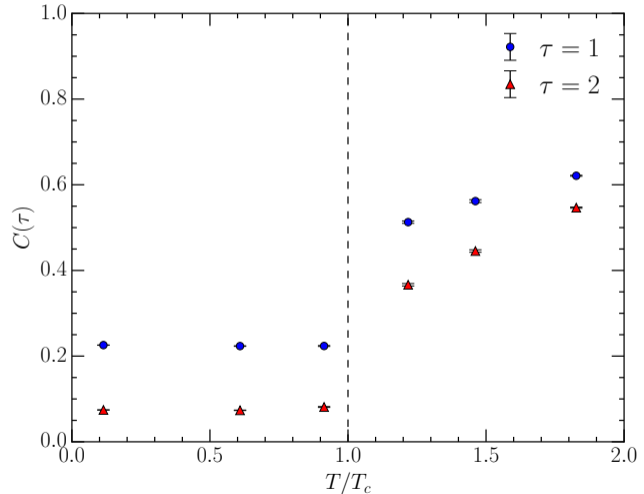
Temporal Alignment of the Vortex Sheet

- $C(\tau)$ changes significantly across the critical temperature, T_c .
- Rising values indicate continued alignment with temperature.
- $C(\tau = 1) \simeq 0.2$ below T_c .



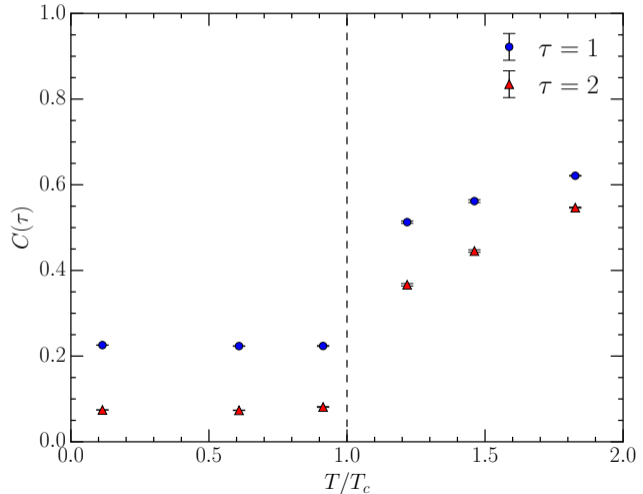
Temporal Alignment of the Vortex Sheet

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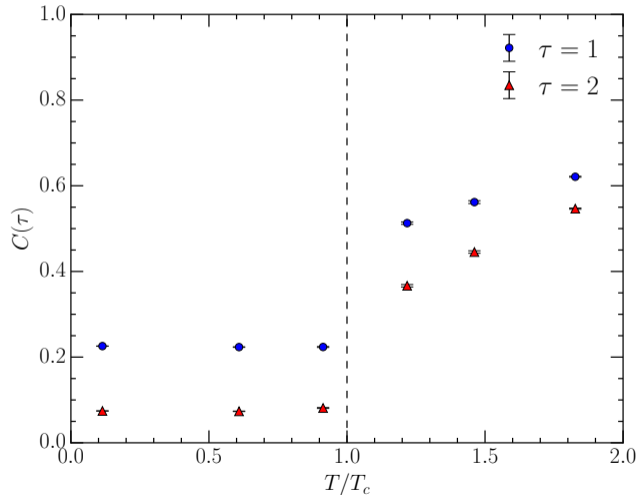
Temporal Alignment of the Vortex Sheet

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 - Halving the lattice spacing, and
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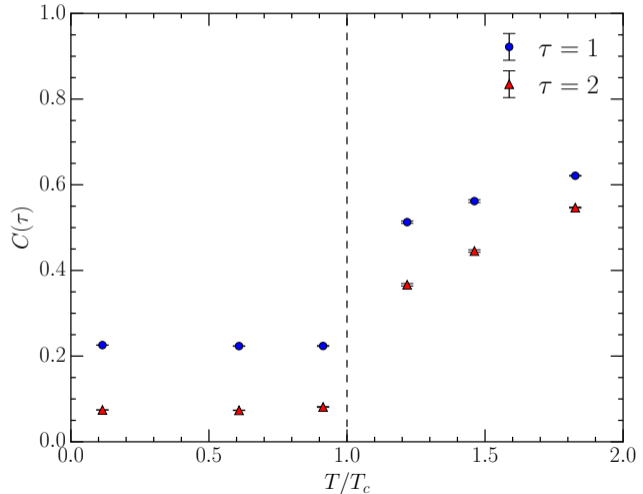
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Temporal Alignment of the Vortex Sheet

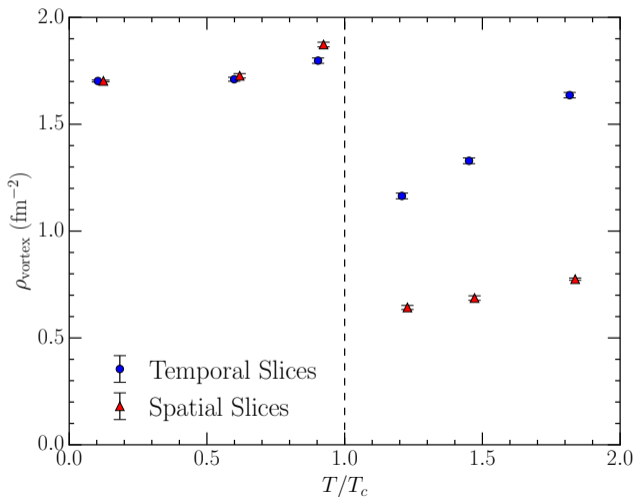
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- Considered the simple Wilson gauge action as an alternative.
- The correlations persist.



Vortex Area Density

- The vortex density of a slice is

$$\rho_{\text{vortex}} = \frac{\text{Number of Vortices}}{\text{Number of Plaquettes}} \frac{1}{a^2}.$$

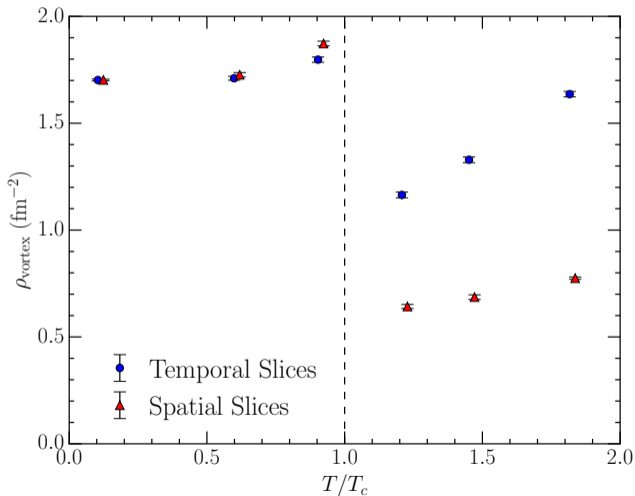


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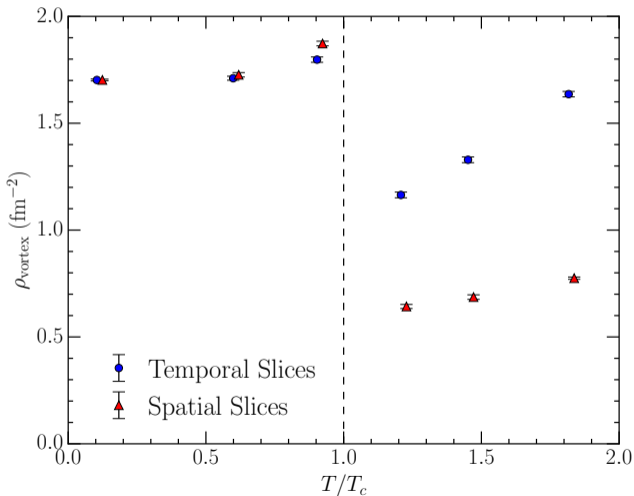


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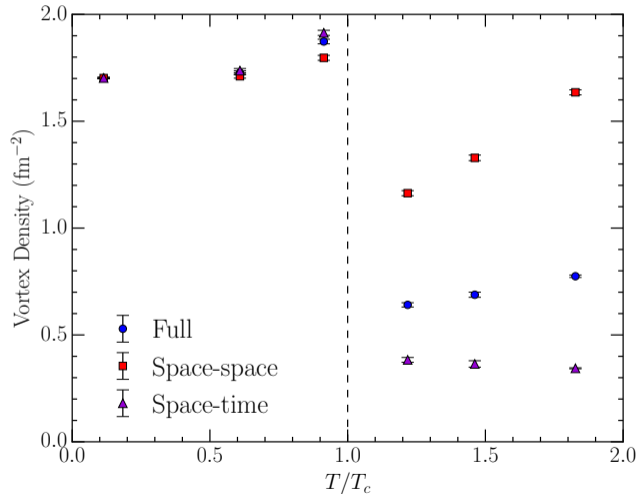
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- The sharp spatial-slice decrease as T_c is crossed coincides with the absence of a percolating cluster.



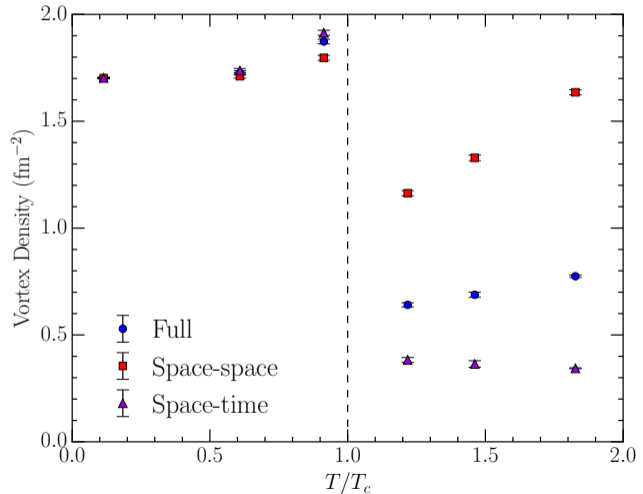
Vortex Area Density for Spatial Slices

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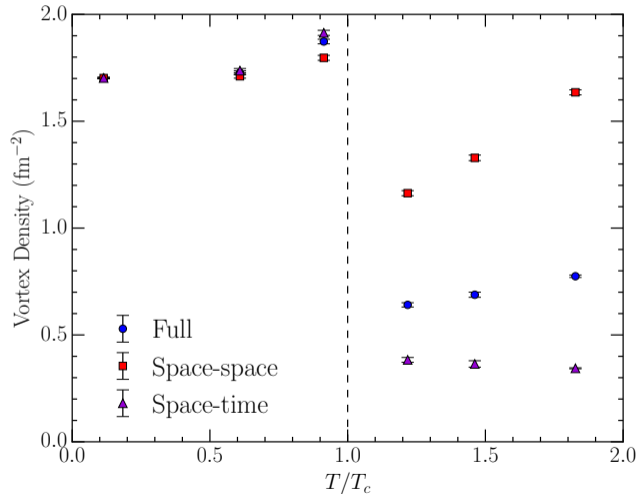
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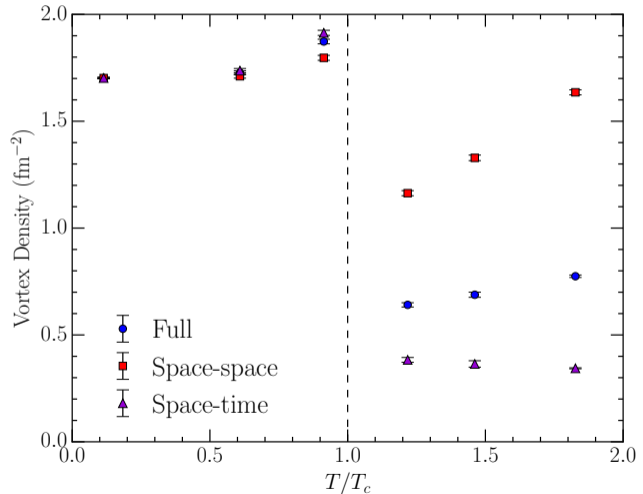
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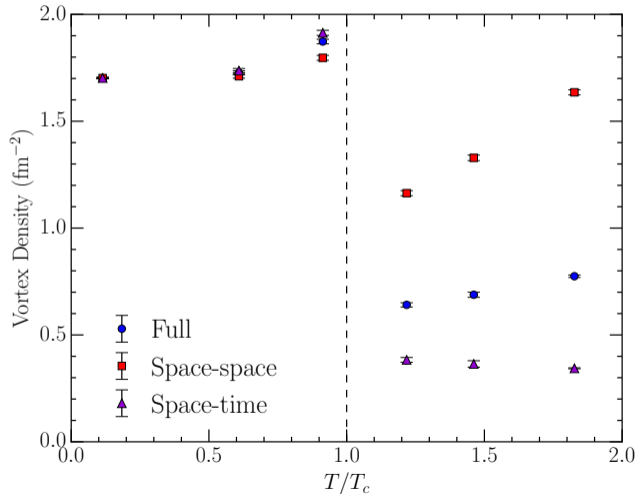
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- Consistent with continued temporal alignment of the vortex sheet.



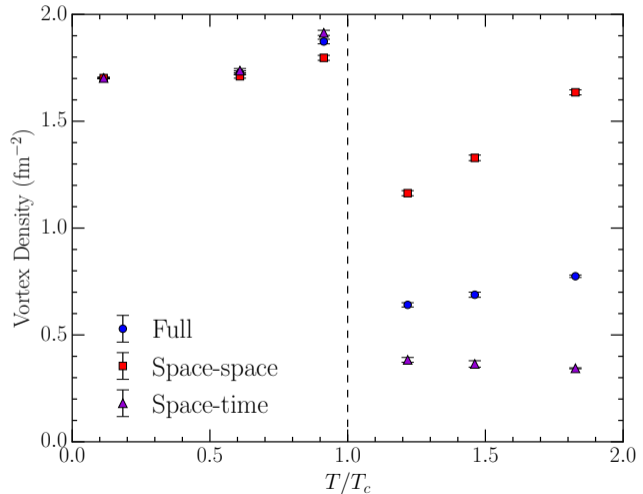
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- Consistent with continued temporal alignment of the vortex sheet.
- Mild increase in the **full** density is due to rapid increase in the **spatial** plaquettes.



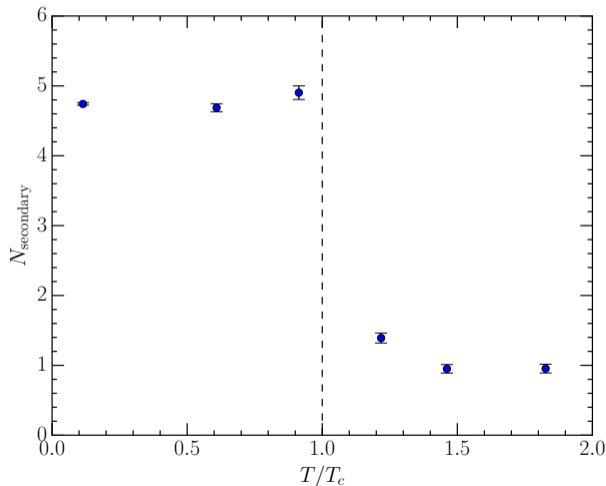
Vortex Area Density for Spatial Slices

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- Consistent with continued temporal alignment of the vortex sheet.
- Note increase in **space-time** vortex density just below T_c .



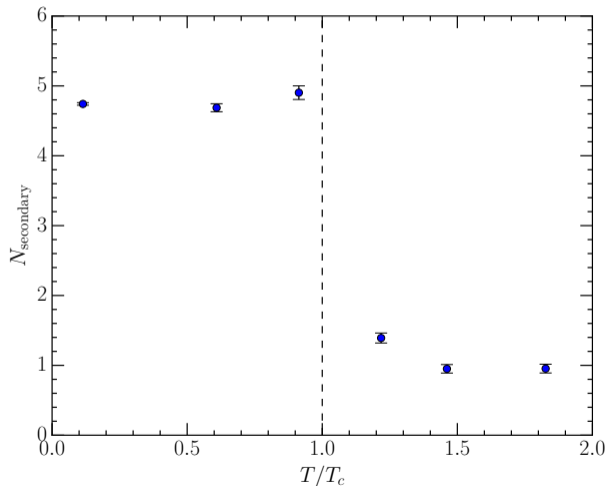
Non-percolating secondary vortex clusters in temporal slices

- Secondary clusters are smaller loops, disconnected from the main percolating cluster.



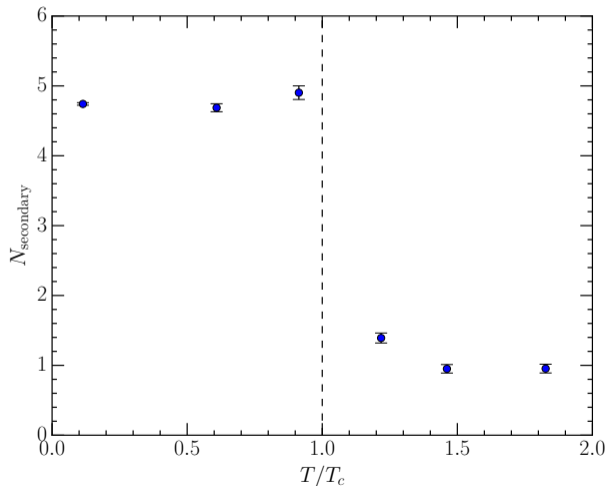
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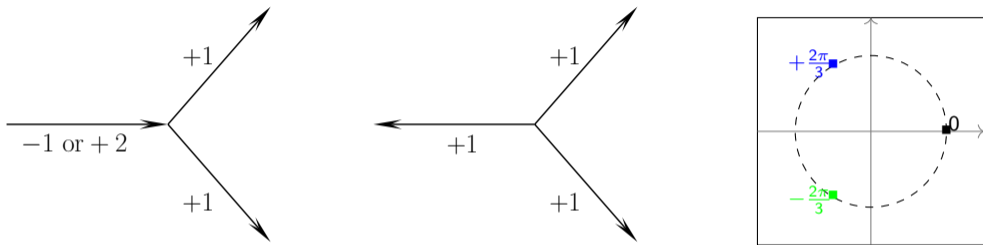
Non-percolating secondary vortex clusters in temporal slices

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- $N_{\text{secondary}}$ is the average number seen on a temporal slice.
- Pronounced drop:
 - A key aspect of vortex geometry characterising deconfinement.



Branching Point Geometry

Branching Points versus Monopoles

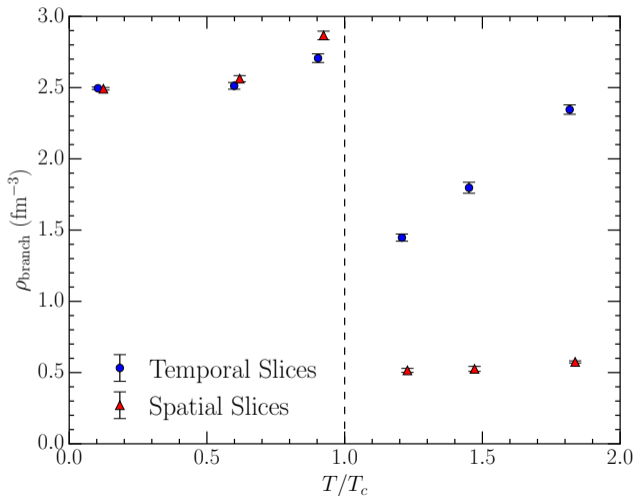


- Our convention illustrates the directed flow of charge $m = +1$.
- Arrows indicate the direction of flow for the labelled charge.
- However, a vortex monopole with charge $+1$ flowing out of the vertex (centre) is equivalent to a vortex branching point with centre charge $+2$ flowing into a vertex (left).

Branching-point geometry

- Branching-point volume density

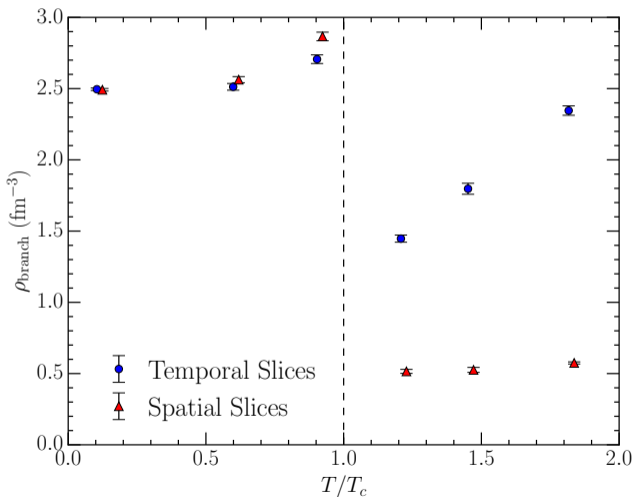
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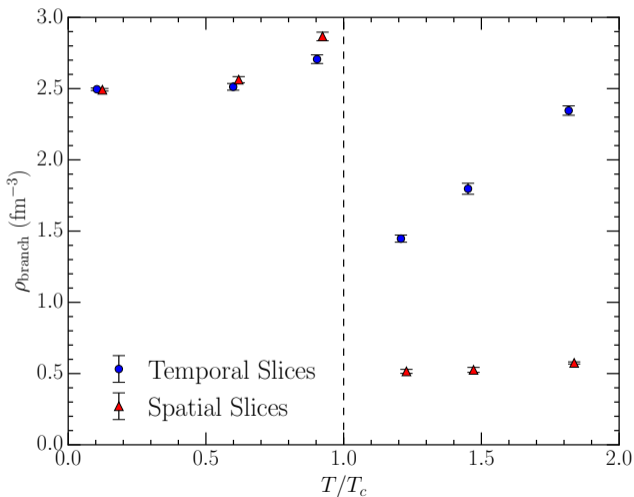
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Branching-point geometry

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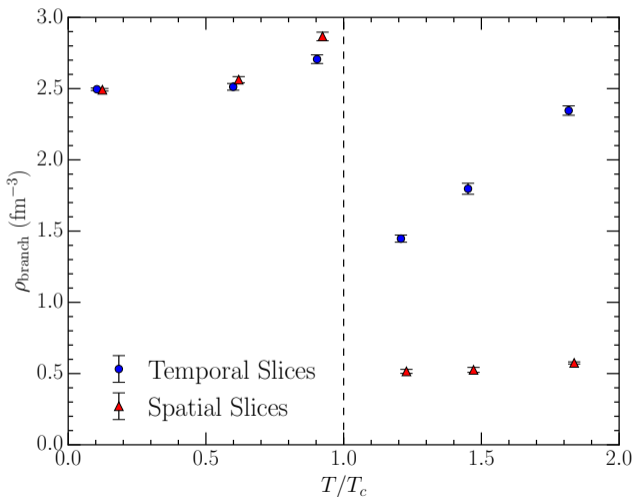
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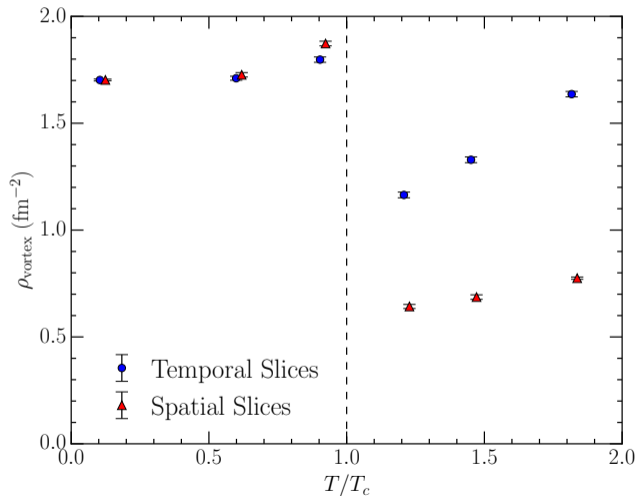
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- There are two drivers behind the rapid increase in ρ_{branch} at large T .



Vortex Area Density

- Already observed an increase in the vortex density
 - Naturally allows for more branching chances.

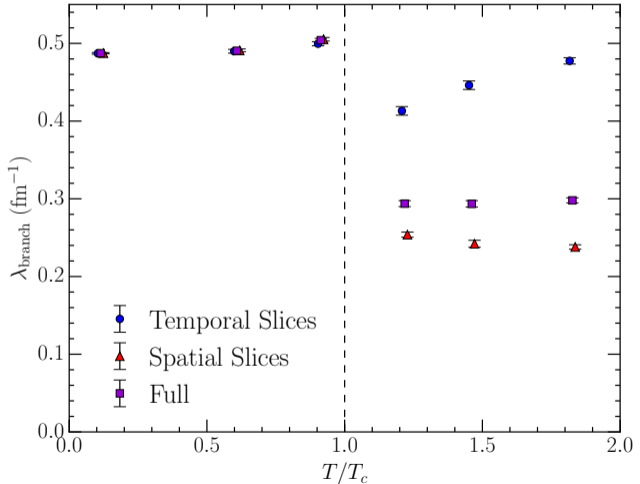


Linear Branching Density

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- Define a linear branching density

$$\sim \frac{\text{Number of branching points}}{\text{Number of vortices}} .$$

$$\lambda_{\text{branch}} = \frac{\rho_{\text{branch}}}{3 \rho_{\text{vortex}}} .$$

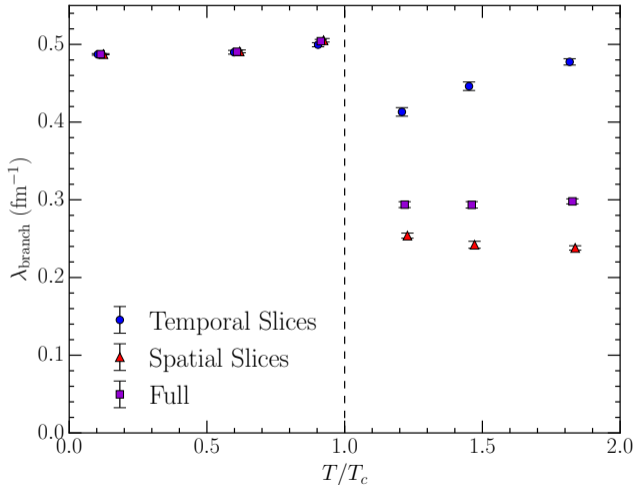


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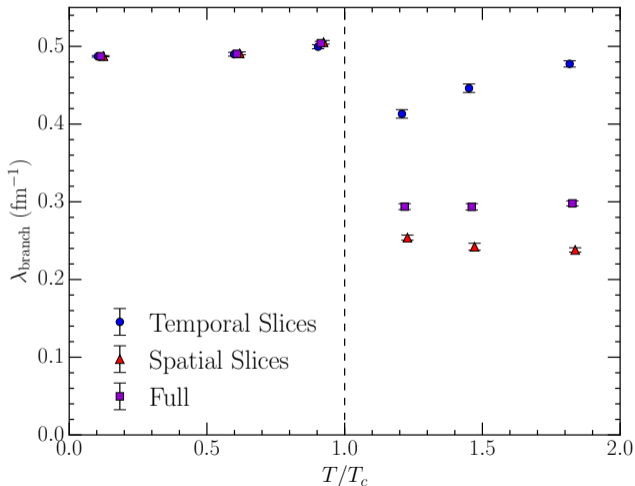
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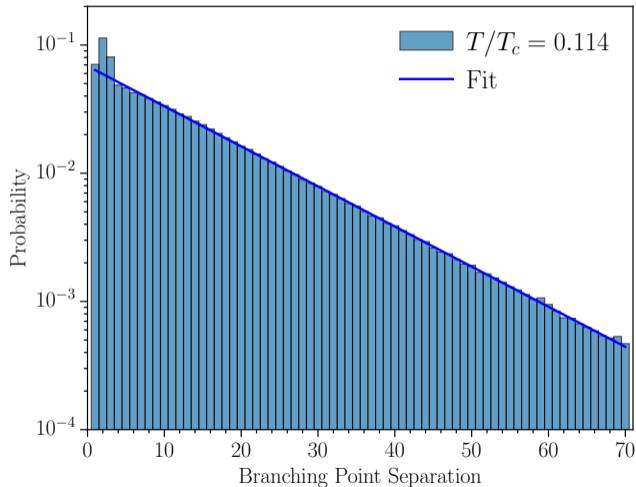
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- Describes a rate of branching (branches/fm).
- **Full** reports an average over the 4 slice directions.



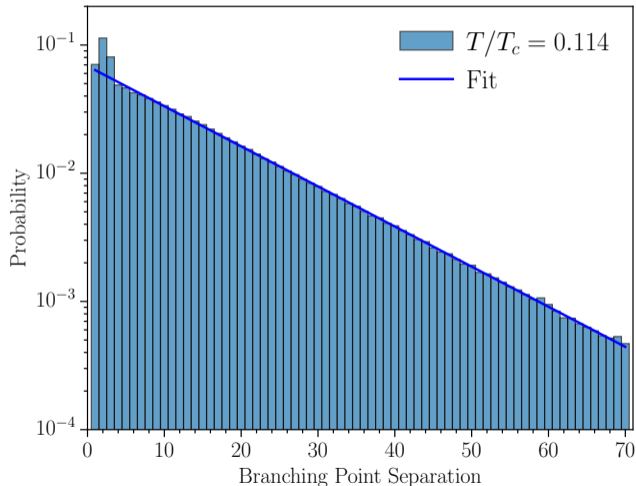
Histogram of Vortex Chain Lengths in Temporal (fixed t) Slices

- Histogram of vortex chain lengths in the percolating cluster.
- Chain length is the number of jets from one branching point to the next.
- The histograms are normalised to unit probability.

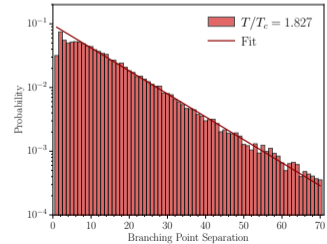
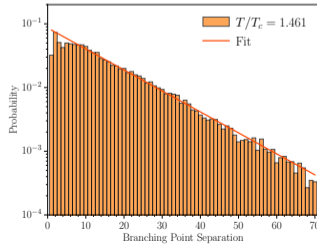
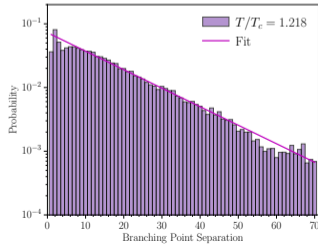
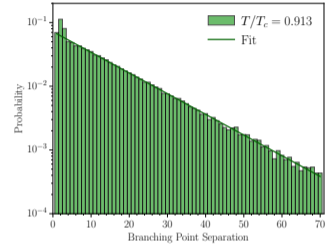
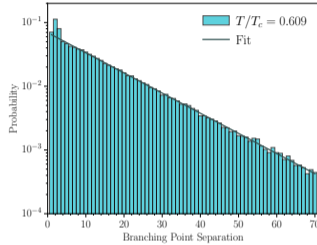
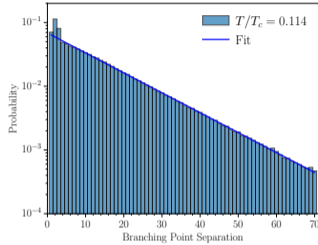


Histogram of Vortex Chain Lengths in Temporal (fixed t) Slices

- Histogram of vortex chain lengths in the percolating cluster.
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- The histograms are normalised to unit probability.
- Moderate size loops, $k > 5$, are exponentially distributed.
 - Fixed probability of branching .
 - Branching is independent of length.

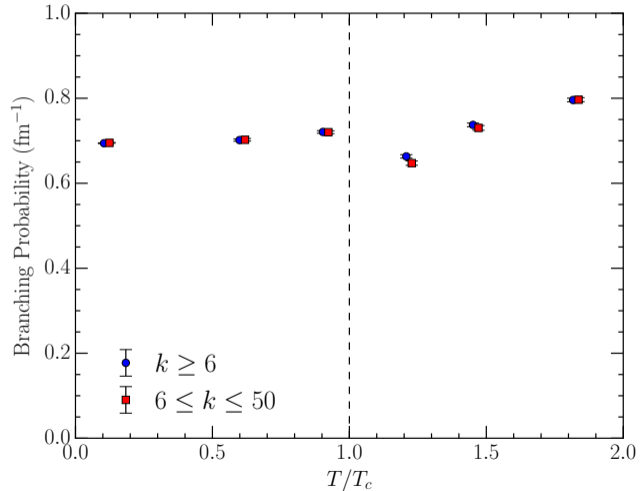


Histogram of Vortex Chain Lengths in Temporal (fixed t) Slices



Branching Probability of a Vortex Path

- Long-range ($k \geq 6$) branching probability.
- Chain lengths $k \leq 70$ and $k \leq 50$ are considered.
- Note the mild drop across the phase transition.
- Probability grows with temperature to large values.



Selected References

- “Centre vortex geometry at finite temperature,”
J. A. Mickley, W. Kamleh and DBL, Accepted Phys. Rev. D (2024) [arXiv:2405.10670 [hep-lat]].
- “Centre vortex structure in the presence of dynamical fermions,”
J. C. Biddle, W. Kamleh and DBL, Phys. Rev. D **107** (2023) no.9, 094507 [arXiv:2302.05897 [hep-lat]].
- “Impact of dynamical fermions on the center vortex gluon propagator,”
J. Biddle, W. Kamleh and DBL, Phys. Rev. D **106** (2022) 014506 [arXiv:2206.02320 [hep-lat]].
- “Static quark potential from centre vortices in the presence of dynamical fermions,”
J. Biddle, W. Kamleh and DBL, Phys. Rev. D **106** (2022) 054505 [arXiv:2206.00844 [hep-lat]].
- “Dynamical fermions, centre vortices, and emergent phenomena,” DBL, J. Biddle, W. Kamleh and A. Virgili, EPJ Web Conf. **274** (2022), 01002 [arXiv:2211.13421 [hep-lat]].
Plenary presentation summary to appear in the proceedings of the XVth Quark Confinement and the Hadron Spectrum conference, 1st-6th August 2022, Stavanger, Norway.

Conclusions

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- Above T_c :
 - The vortex density, branching-point density, and branching probability drop across T_c .
 - Branching points are suppressed at short distances.
 - Smaller secondary clusters are suppressed.
- Observe a rapid rise in the branching point density with high T .
 - The branching probability and branching rate increase with T .
 - The vortex density increases \rightarrow more opportunities to branch.

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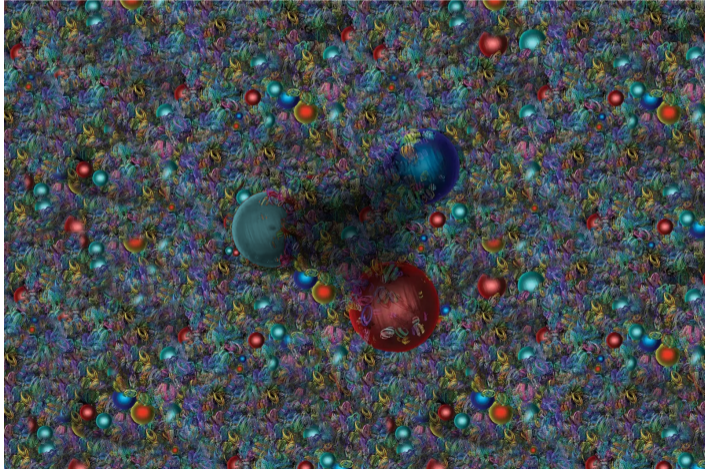
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 - A commensurate condensation of branching points on temporal slices.
- An exchange that leaves the linear branching probability averaged over the four slice dimensions invariant.

Additional Information

Ironic incongruity



- Static valence quarks suppress vacuum fields as they induce flux-tube tunnels.

Restoration of Chiral Symmetry

- If vortices are responsible for $D\chi SB$, then their removal should restore chiral symmetry

$$SU 2_L \times SU 2_R \times U(1)_A$$

- Expect hadrons related by chiral transformations to become degenerate

$$\begin{array}{ccc} \pi & \xleftrightarrow{U(1)_A} & a_0 \\ \rho & \xleftrightarrow{SU 2_L \times SU 2_R} & a_1 \\ N & \xleftrightarrow{SU 2_L \times SU 2_R} & \Delta \end{array}$$

- At light quark masses, all symmetries are observed to be restored.
- A. Trewartha, W. Kamleh and DBL, J. Phys. G **44** (2017) 125002 [arXiv:1708.06789 [hep-lat]].

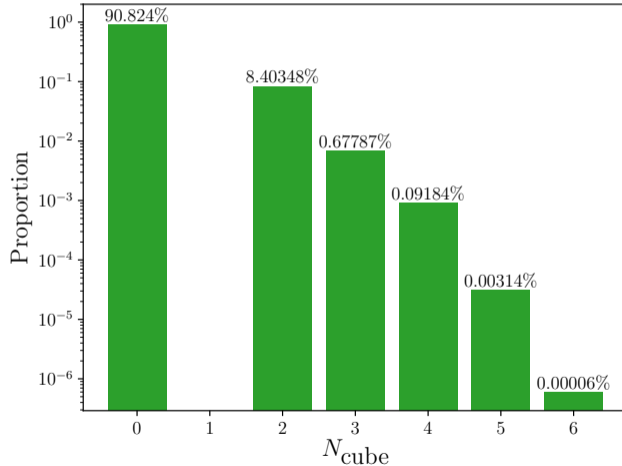
Visualising Centre Vortices

- Consider the number of vortices entering a 3D cube on the dual lattice.

| $N_{\text{cube}}(\tilde{\chi})$ | Interpretation |
|---------------------------------|--|
| 0 | No vortices present. |
| 1 | Terminating vortex, forbidden by Bianchi*. |
| 2 | Vortex line flowing through the cube. |
| 3 | Simple three-way vortex monopole. |
| 4 | Vortex intersection. |
| 5 | Complex five-way monopole path. |
| 6 | Vortex intersections or double monopoles. |

*Bianchi identity implies a continuous flow of centre vortex flux through a spatial cube.

Visualising Centre Vortices



Space-Time Oriented Vortices

Rendering Space-Time Oriented Projected Vortices

- Every link in the spatial volume has a forward and backward time-oriented plaquette associated with it.

Rendering Space-Time Oriented Projected Vortices

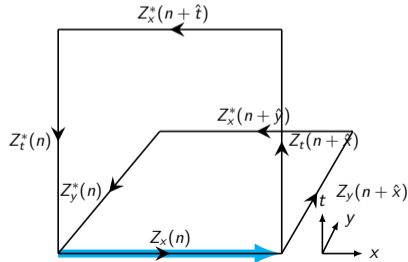
- Every link in the spatial volume has a forward and backward time-oriented plaquette associated with it.
- The three jets associated with the spatial x - y , y - z and z - x plaquettes, are complemented by
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- See “Visualization of center vortex structure,” to link vortices to topological charge.
J. C. Biddle, W. Kamleh and DBL, Phys. Rev. D **102** (2020) 034504 [arXiv:1912.09531 [hep-lat]].

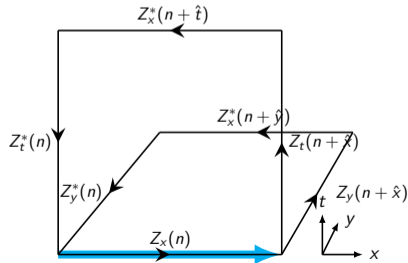
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- If a spatial link belongs to a vortex in a space-time plaquette then:
 - The link is rendered in cyan for an $m = +1$ vortex.



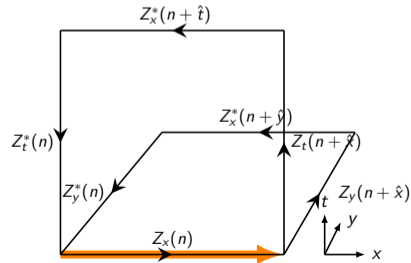
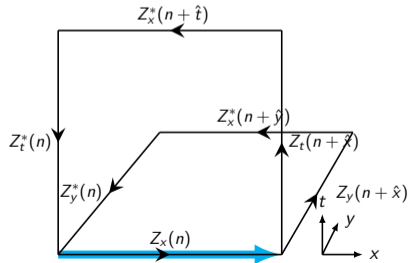
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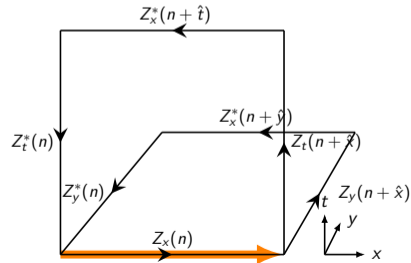
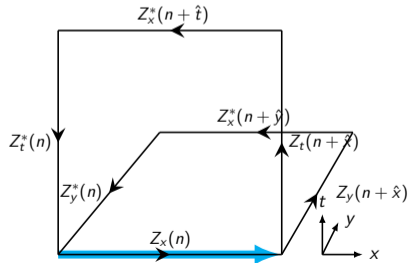
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 - The link is rendered as a positively-directed arrow for forward space-time plaquettes.



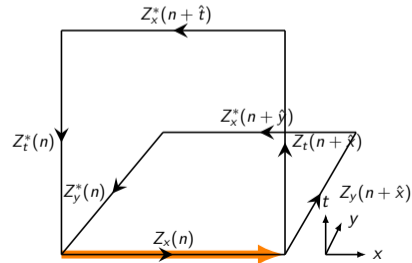
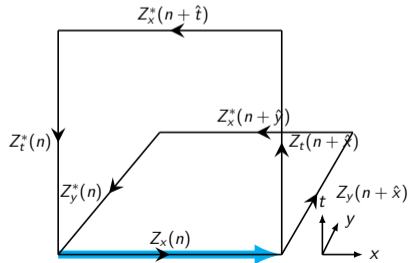
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- If a spatial link belongs to a vortex in a space-time plaquette then:
 - The link is rendered in **cyan** for an $m = +1$ vortex, and in **orange** for $m = -1$.
 - The link is rendered as a positively-directed arrow for forward space-time plaquettes.
 - The link is rendered as a negatively-directed arrow for backward space-time plaquettes.

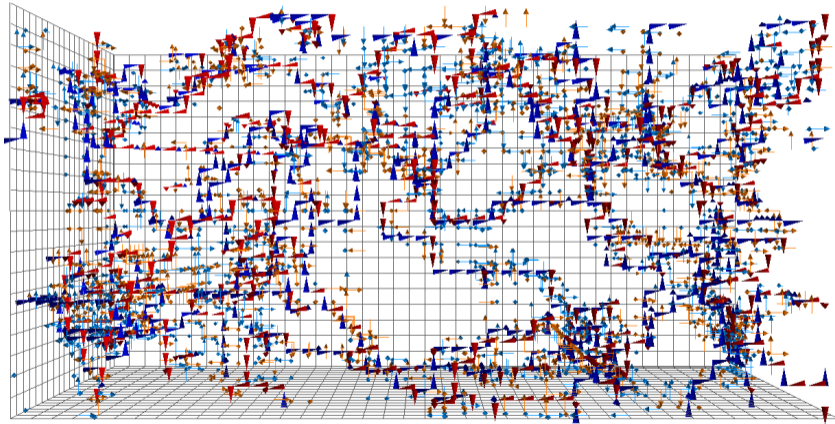


Rendering Space-Time Oriented Projected Vortices

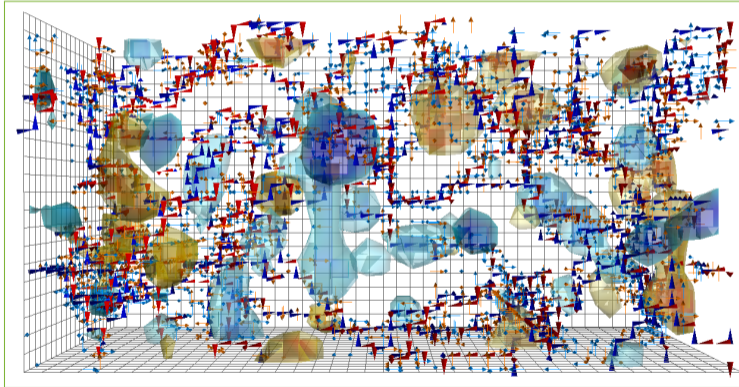
- If a spatial link belongs to a vortex in a space-time plaquette then:
 - The link is rendered in **cyan** for an $m = +1$ vortex
 - The link is rendered as a positively-directed arrow for forward space-time plaquettes.
 - The link is rendered as a negatively-directed arrow for backward space-time plaquettes.
- As one steps forwards in time, positively-directed links become negatively-directed.



Time slice $t = 1$



Animation of Centre Vortex Structure [Google: YouTube CSSM Visualisations](#)



Interactive 3D Visualisation Techniques

- Rendered in AVS Express Visualisation Edition.
<http://www.avs.com/solutions/express/>
- Exported in VRML.
- Converted to U3D format via pdf3d ReportGen.
<https://www.pdf3d.com/products/pdf3d-reportgen/>
- Imported into \LaTeX via media9 package.
- Viewed in Adobe acroread (Linux, use 9.4.1 when 3D support was maintained).
<ftp://ftp.adobe.com/pub/adobe/reader/unix/9.x/9.4.1/>