भौतिक अनुसंधान प्रयोगशाला Physical Research Laboratory

Neutrino physics : phenomenological and theoretical perspectives

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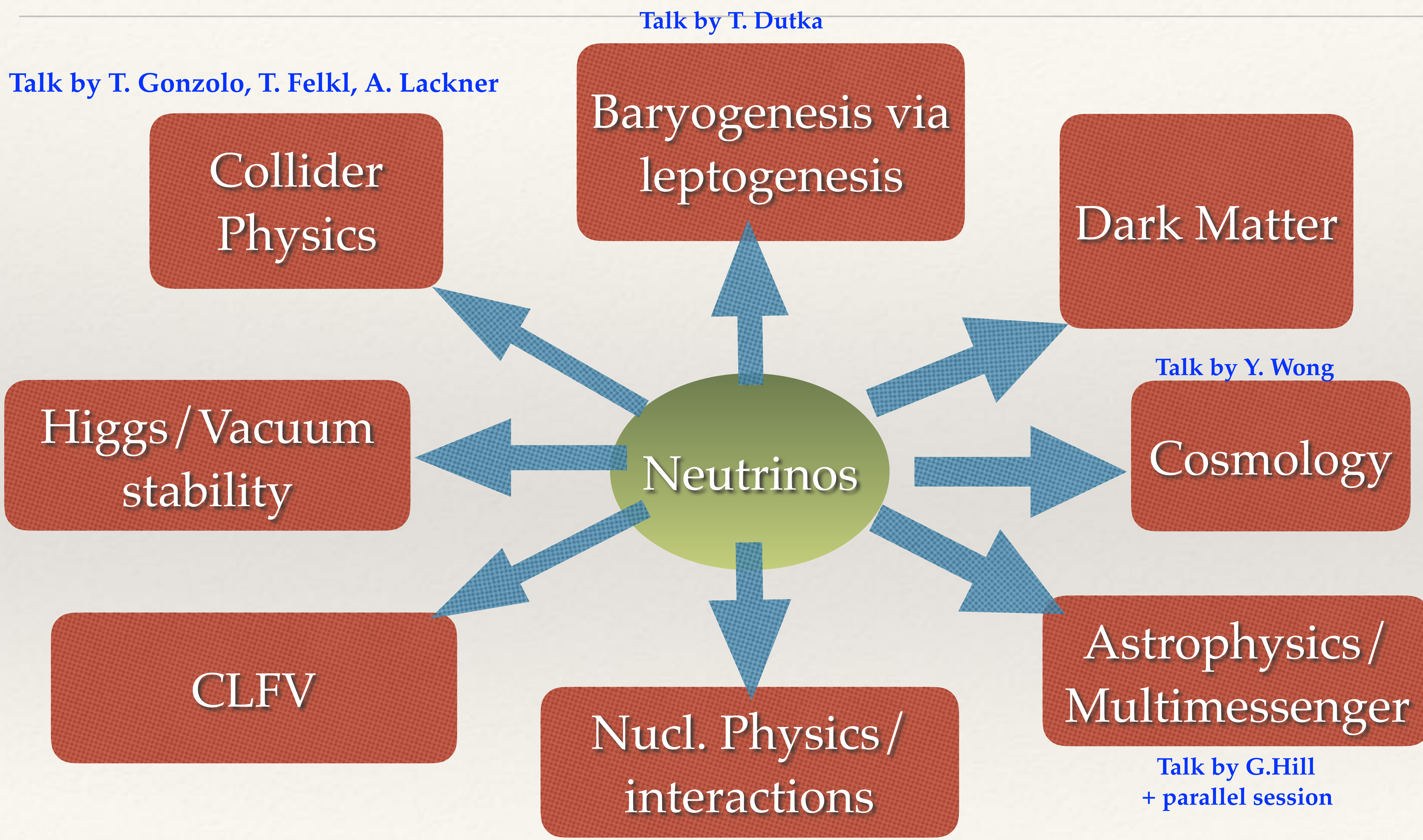
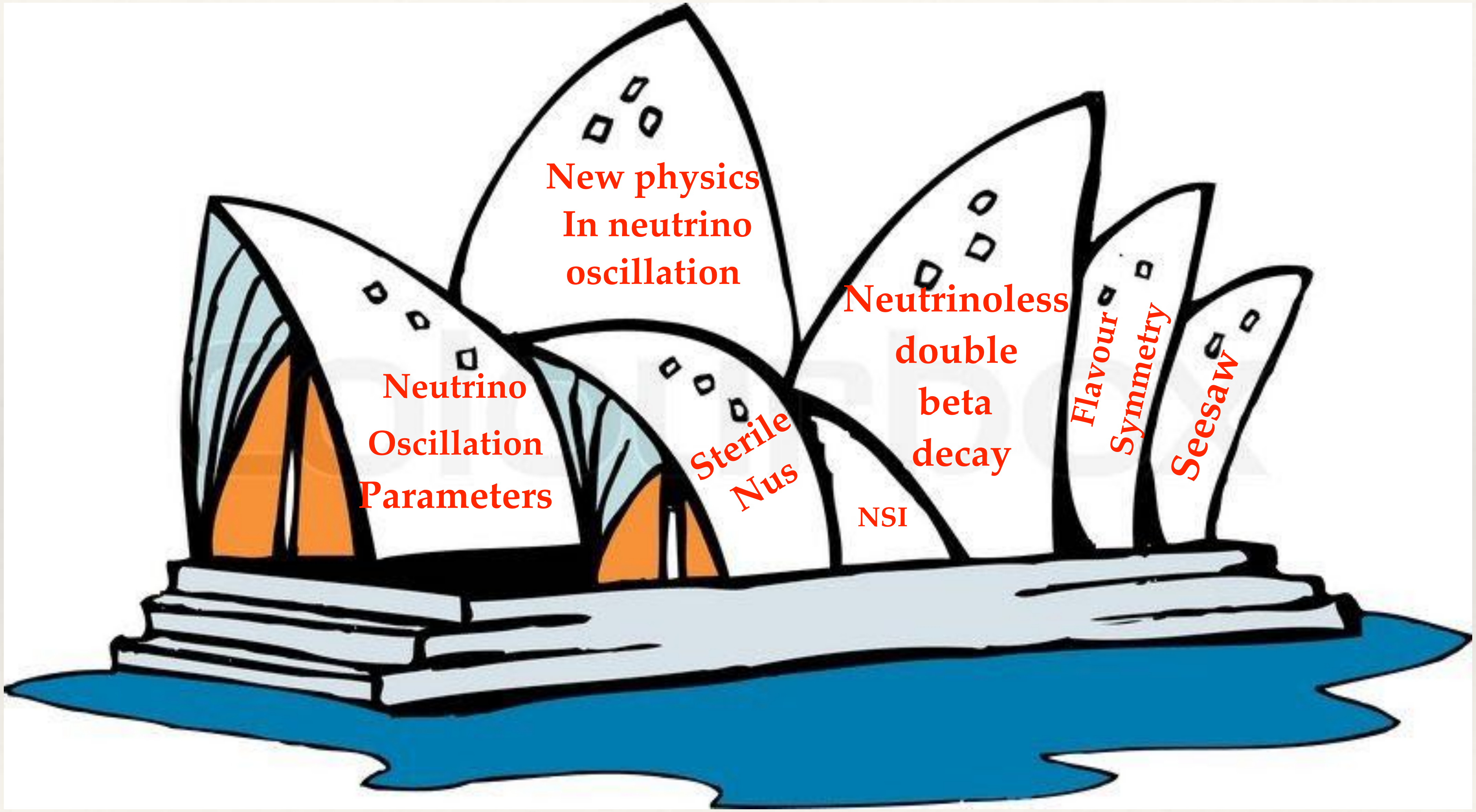
Neutrino: questions

- Unknown oscillation parameters - hierarchy, octant of 2-3 mixing angle and CP phase
- Absolute neutrino masses — beta decay, cosmology
- Nature of neutrinos - Dirac or Majorana — neutrino less double beta decay
- Are there more than three flavours — sterile neutrinos
- Origin of neutrino masses and mixing — seesaw , flavour symmetry — physics beyond Standard Model

Neutrino Connections

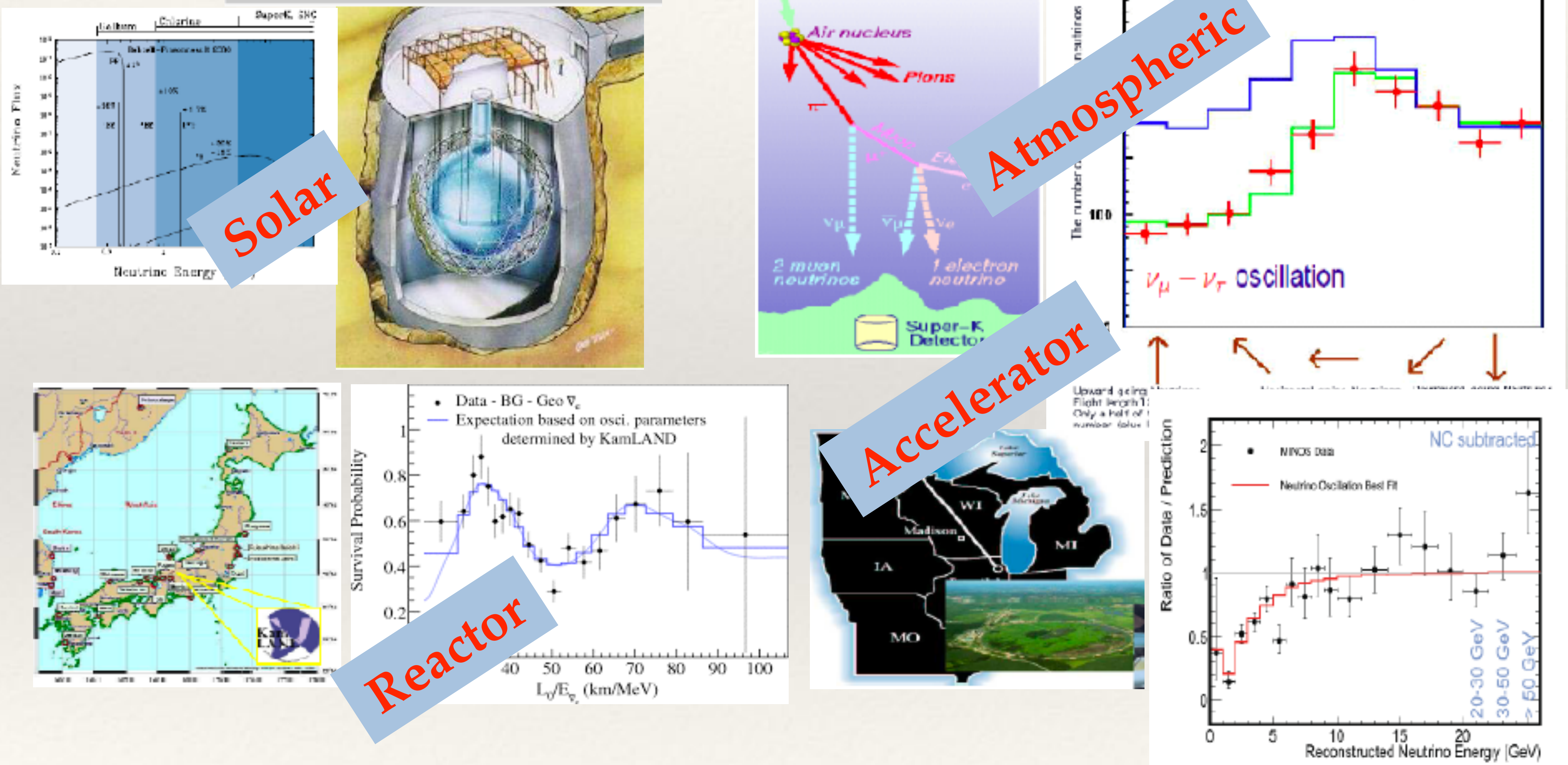


Plan of talk


New physics
In neutrino
oscillation
Neutrino
Oscillation
Parameters
Sterile
Nus
NSI
Neutrinoless
double
beta
decay
Flavour
Symmetry
Seesaw

Neutrino Oscillations

$$\frac{CC}{NC} = \frac{\nu_e}{\nu_e + \nu_\mu + \nu_\tau} < 1$$



Three Neutrino Paradigm

- Measurement of non-zero θ_{13} in reactor experiments → three neutrino picture

Atm +LBL | Sol+KL

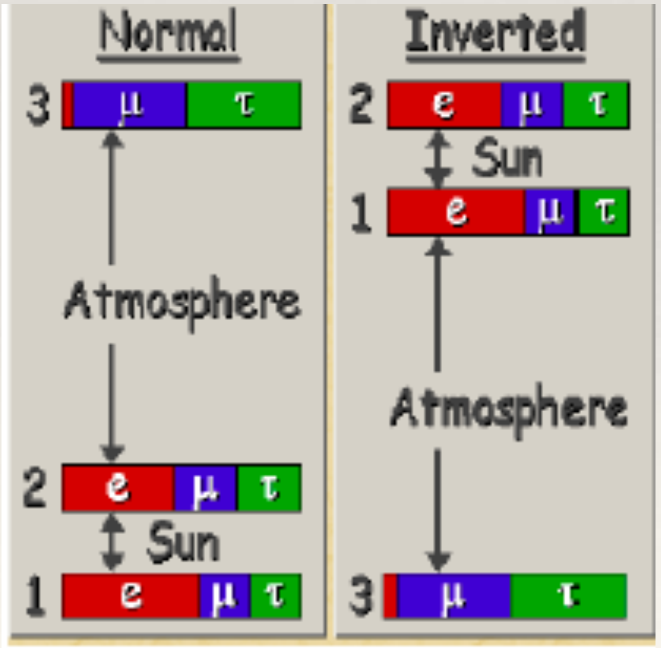
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & & \\ & c_{23} & s_{23} \\ & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & & e^{-i\delta} s_{13} \\ & 1 & \\ -e^{i\delta} s_{13} & & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & \\ -s_{12} & c_{12} & \\ & & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$c_{12} = \cos\theta_{12}$ etc., δ CP-violating phase

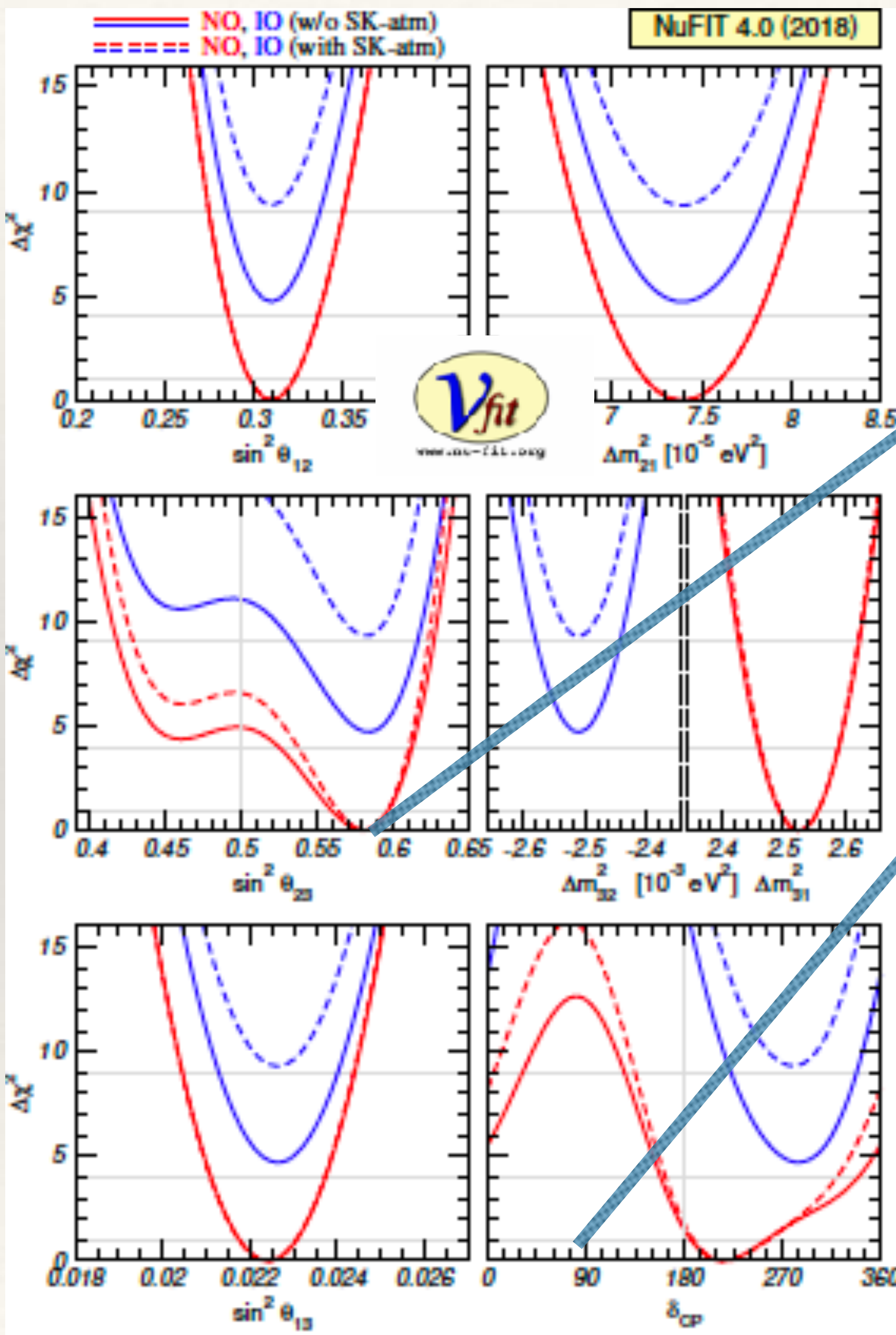
- Δm^2_{21}, θ_{12}, θ_{13} Solar + KamLND
- Δm^2_{31}, θ_{13} Reactor
- Δm^2_{31}, θ_{23}, θ_{13}, δ_{CP} Atm + LBL

↓

Interplay among different sectors because of θ_{13}



Current Status



- Best-fit θ_{23} in second octant
- Preference for NO
- $\delta_{CP} = +90^\circ$ disfavoured at more than 3σ irrespective of mass ordering
- Oscillation experiments not sensitive to Majorana phases

Similar result from Bari and Valencia groups

Degeneracy problem

- The main problem in determination of hierarchy, octant and δ_{CP} in LBL experiments is due to presence of degeneracies
- Degeneracy → different set of parameters giving the same probability → equally good fit to the data
-

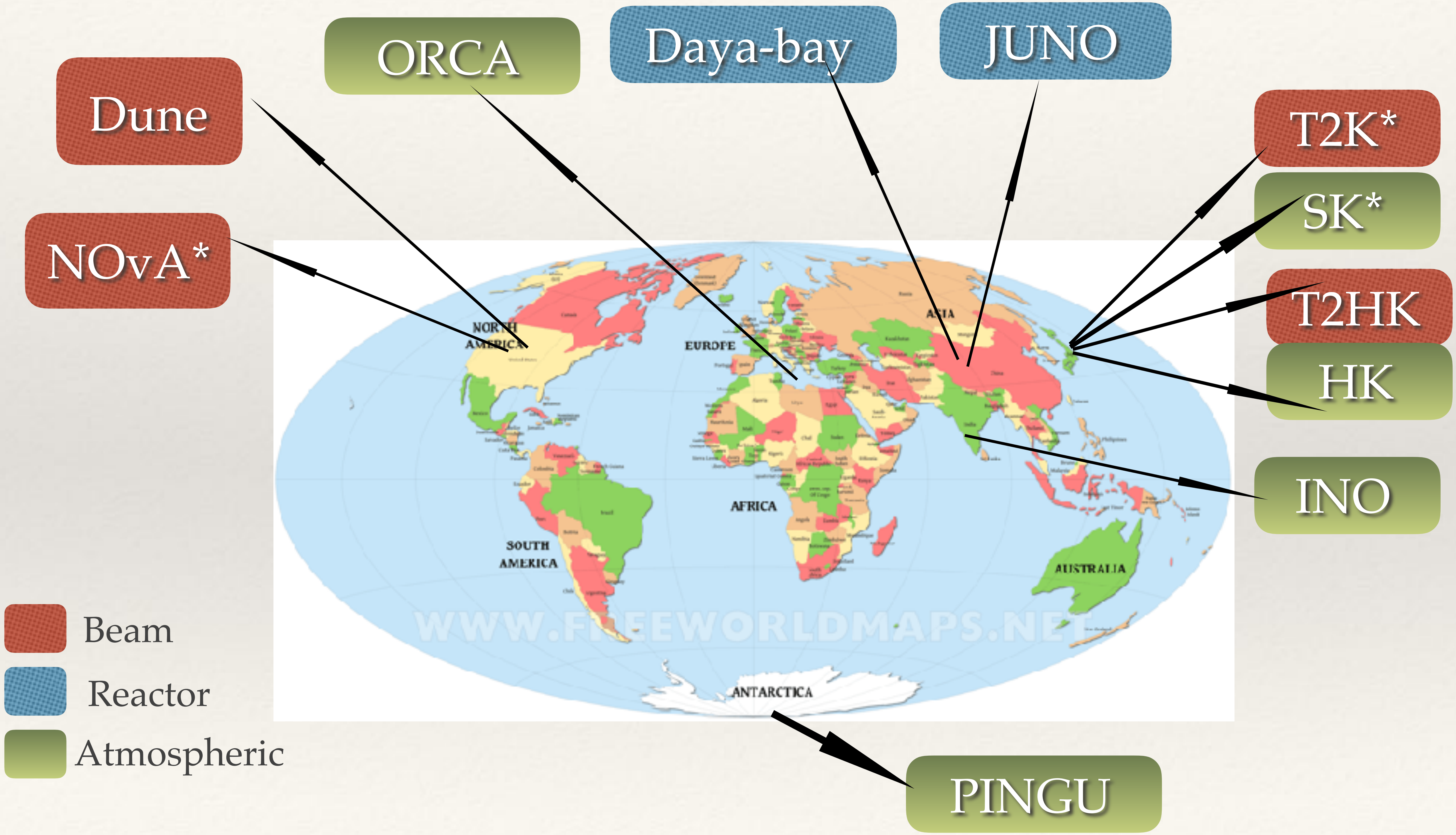
Hierarchy - δ_{CP} degeneracy	Intrinsic octant degeneracy	Octant - δ_{CP} degeneracy
$P_{\mu e}(\Delta, \delta_{CP}) = P_{\mu e}(-\Delta, \delta'_{CP})$	$P_{\mu\mu}(\theta_{23}) = P_{\mu\mu}(\theta_{23} - \pi/2 - \theta_{23})$	$P_{\mu e}(\theta_{23}, \delta_{CP}) = P_{\mu e}(\theta'_{23}, \delta'_{CP})$
Minakata, NunoKawa, 2001	Fogli and Lisi, 1996	Gandhi, Ghosal, Goswami, Shankar 2005

Comprehensive Approach

$$P_{\mu e}(\theta_{23}, \Delta, \delta_{CP}) = P_{\mu e}(\theta'_{23}, -\Delta', \delta'_{CP}) \Rightarrow \text{generalized}\,(\text{hierarchy} - \theta_{23} - \delta_{CP})\,\text{degeneracy}.$$

Coloma, Minakata, Parke, 2014
Ghosh, Ghoshal, Goswami, Nath, Raut, 2015

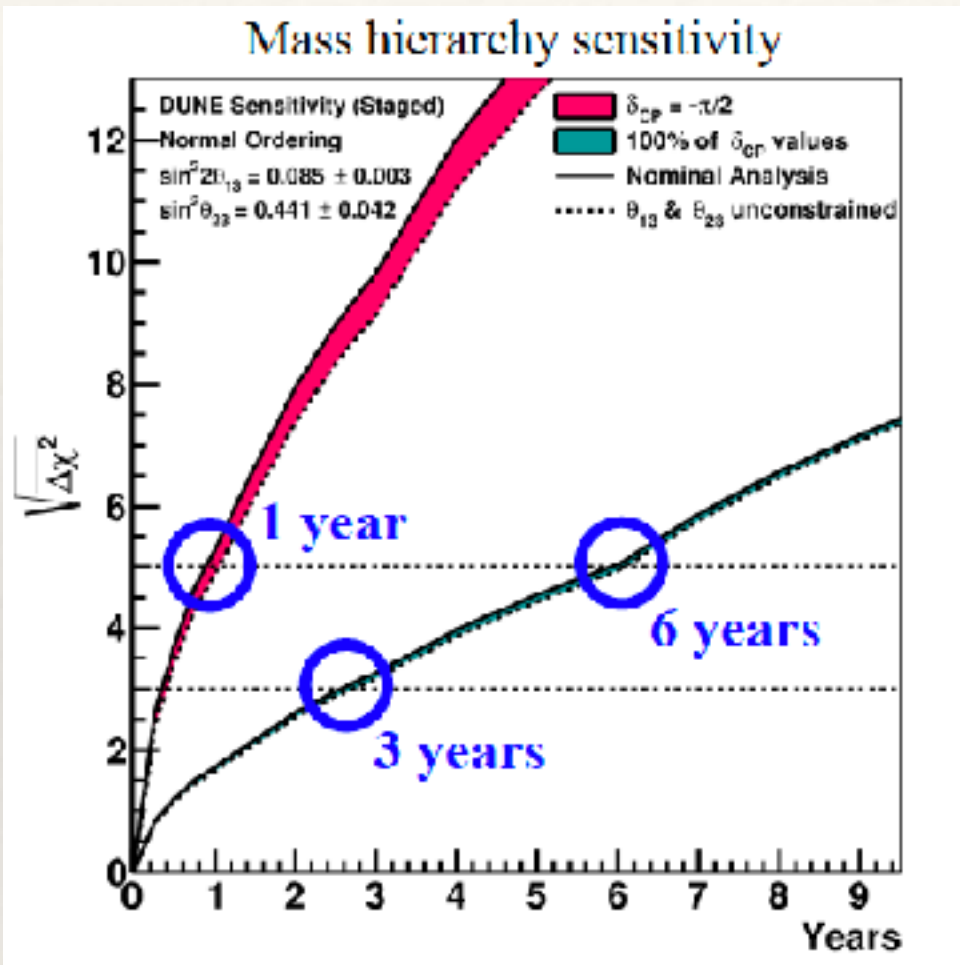
Ongoing and planned experiments



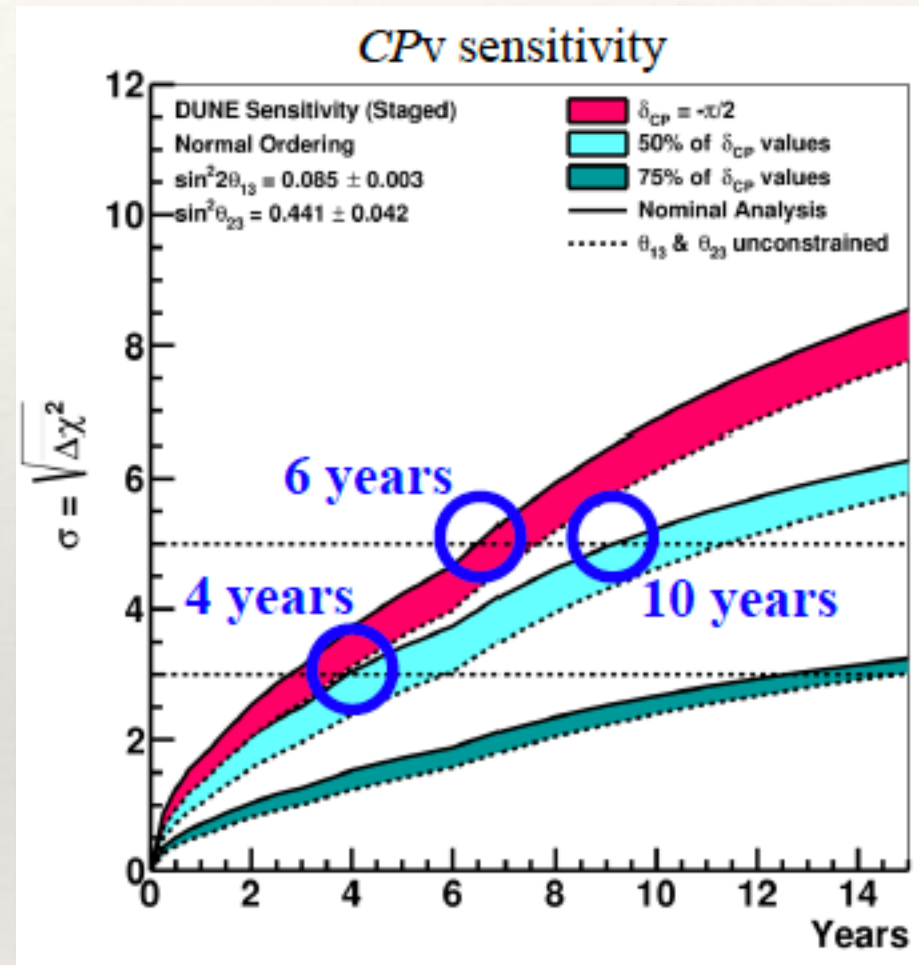
Future Goals

- Determination of hierarchy, octant and CP phase
- Probing new physics in oscillation experiments
- Testing models of flavour symmetry
- Synergy between different experiments

Mass hierarchy and CP with DUNE



Hierarchy sensitivity due to enhanced matter effects

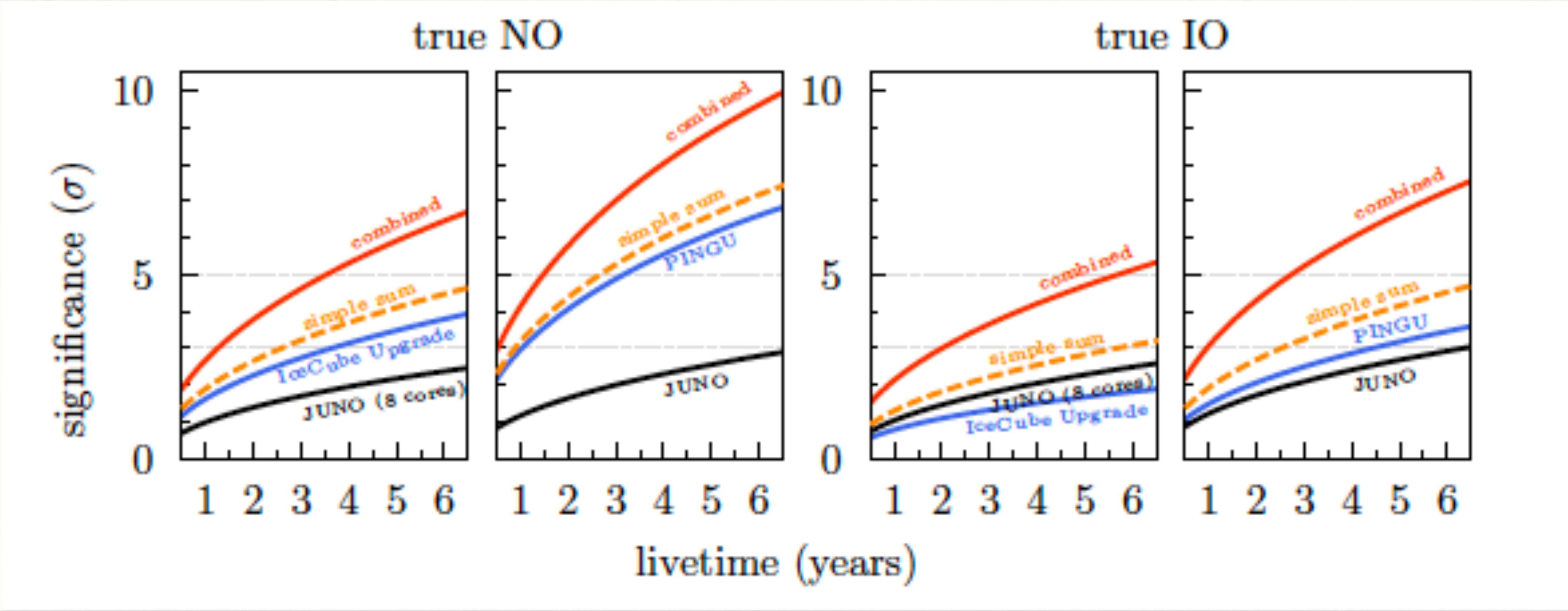


Matter effects help in removing wrong hierarchy-wrong CP solutions

From: R. Patterson's slides

Hierarchy: Juno+IceCube upgrade

8 core JUNO + IceCube upgrade/PINGU / (better efficiency for lower energy neutrinos)



5σ sensitivity in 4(6) years NO (IO)

IceCube : earth matter effect of atmospheric neutrinos
JUNO: interference effect in vacuum oscillation

} **Synergy**

hep-ex 1911.06745

New Physics

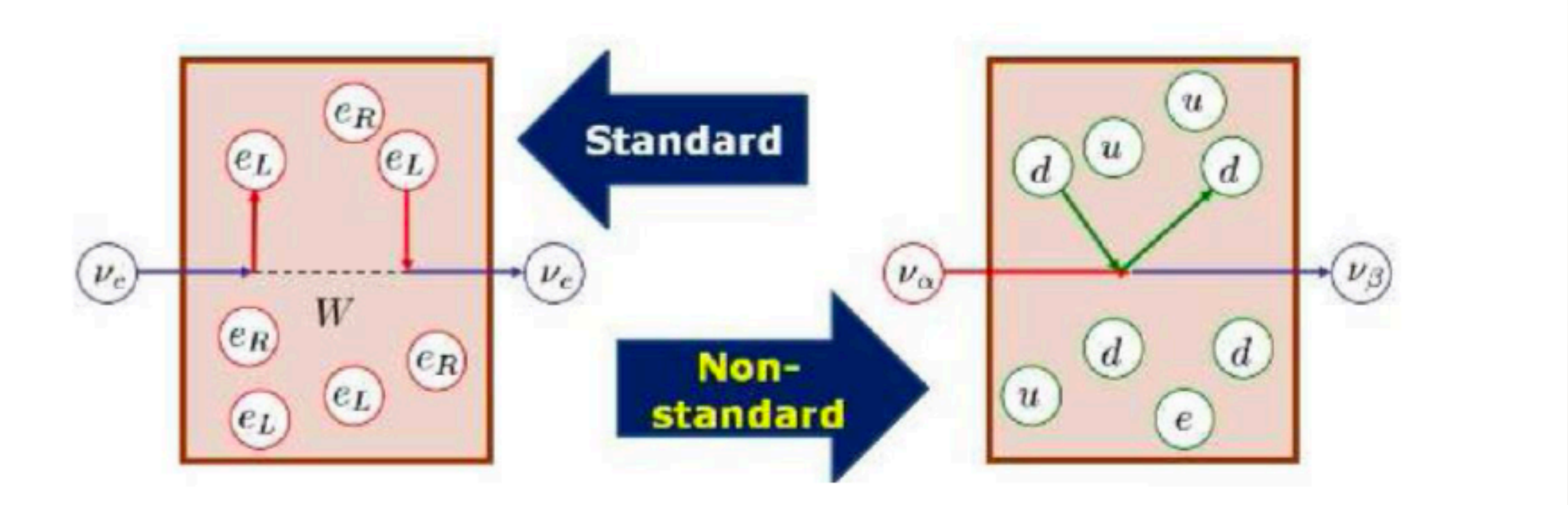
- Sterile Neutrinos
- Non-standard Interactions (NSI)
- Non-unitary mixing
- CPT and Lorentz symmetry violation
- Long range forces
- Neutrino decay

Two appoaches

Impact on the standard Three neutrino picture

Constraining new physics parameters

Non-standard interactions



$$\mathcal{L}_{\rm NSI} = -2\sqrt{2}G_F\varepsilon_{\alpha\beta}^{ff'C}\left(\overline{\nu_\alpha}\gamma^\mu P_L\nu_\beta\right)\left(\bar{f}\gamma_\mu P_C f'\right)$$

$\varepsilon_{\alpha\beta}^{ff'C}$ are NSI parameters, $\alpha, \beta = e, \mu, \tau$, $f, f' = e, u, d$ and $C = L, R$.

$f \neq f'$ → Charged Current NSI

$f = f'$ → Neutral Current NSI

Non-standard interactions

Standard-NC interaction

$$\nu_\alpha + f \rightarrow \nu_\alpha + f$$

Non-Standard NC interaction

$$\nu_\alpha + f \rightarrow \nu_\beta + f$$

$$\mathcal{L} = -G^{\alpha\beta}\epsilon^f_{\alpha\beta}\bar{\nu}_\alpha\gamma^\mu\nu_\beta\bar{f}\gamma_\mu f$$

$$\epsilon_{\alpha\beta} = \sum_{f=e,u,d}\frac{N_f}{N_e}\epsilon^f_{\alpha\beta}$$

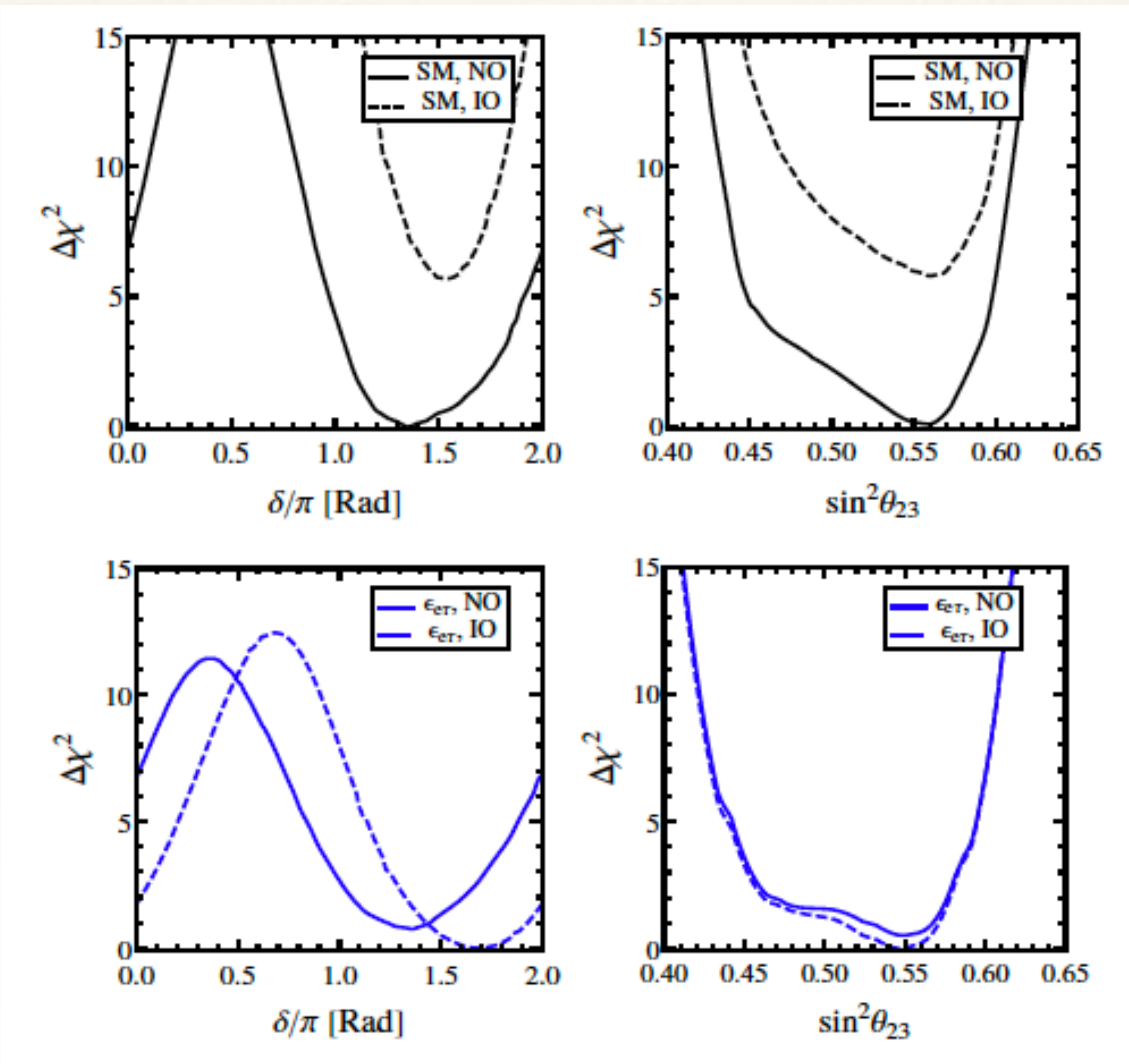
$$H = \frac{1}{2E}\left[U\text{diag}(0, \Delta m^2_{21}, \Delta m^2_{31})U^\dagger + V\right],$$

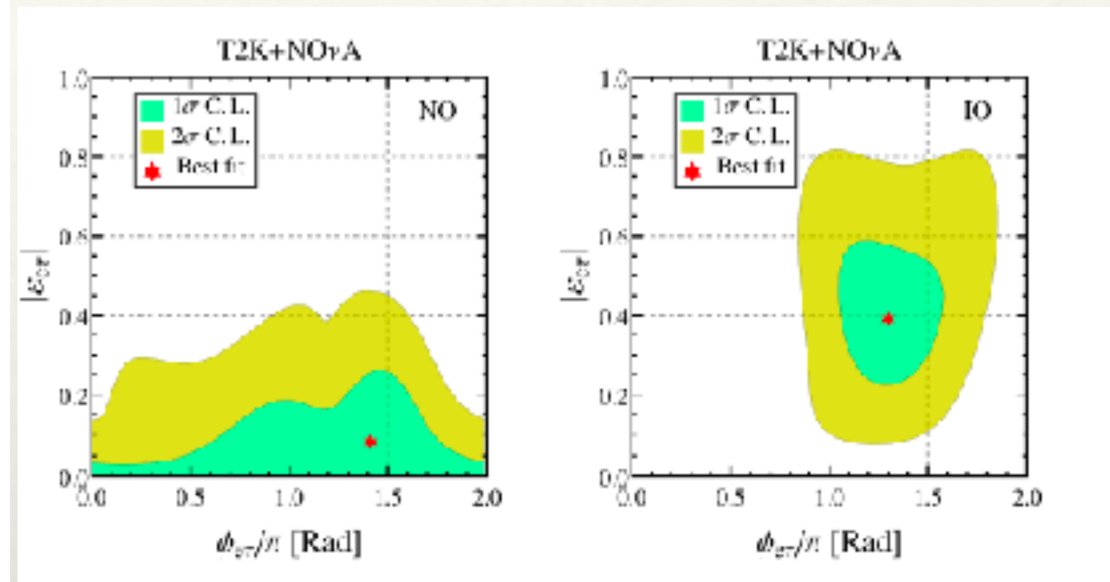
$V \Rightarrow$ matter potential in presence of NSI,

$$V = A\begin{pmatrix} 1+\epsilon_{ee} & \epsilon_{e\mu}e^{i\phi_{e\mu}} & \epsilon_{e\tau}e^{i\phi_{e\tau}} \\ \epsilon_{e\mu}e^{-i\phi_{e\mu}} & \epsilon_{\mu\mu} & \epsilon_{\mu\tau}e^{i\phi_{\mu\tau}} \\ \epsilon_{e\tau}e^{-i\phi_{e\tau}} & \epsilon_{\mu\tau}e^{-i\phi_{\mu\tau}} & \epsilon_{\tau\tau}\end{pmatrix}.$$

Here, $A \equiv 2\sqrt{2}G_F N_e E$ and $\epsilon_{\alpha\beta}e^{i\phi_{\alpha\beta}} \equiv \sum_{f,C}\epsilon^{fC}_{\alpha\beta}\frac{N_f}{N_e}$

NSI and mass hierarchy





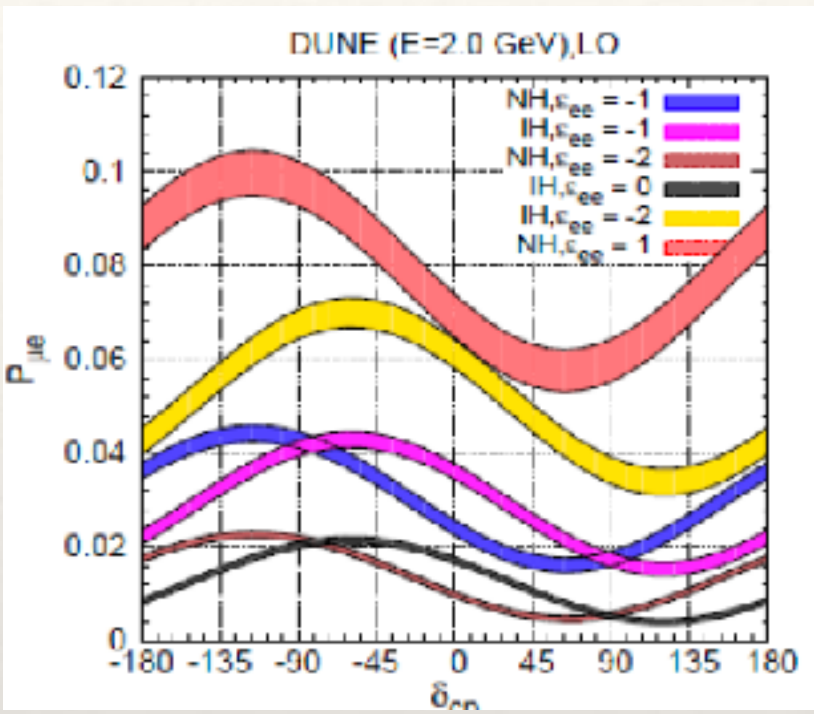
Fit to T2K and NOvA data assuming NSI

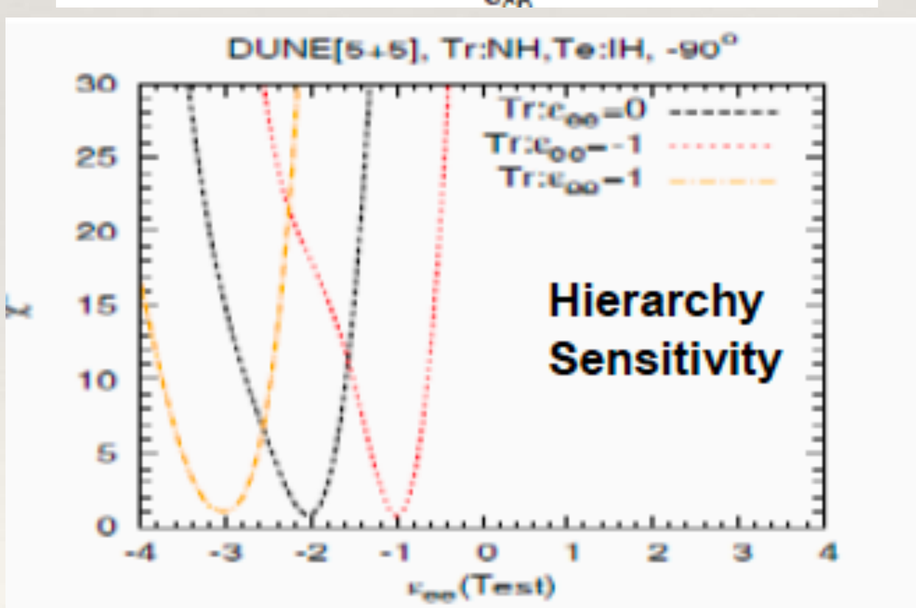
IO prefers non-zero NSI

IO no longer disfavoured

Capozzi, Chatterjee, Palazzo 1908.06992

Degeneracies due to diagonal NSI





K.N. Deepthi, S.Goswami, N. Nath, PRD 2016

$$V = 2\sqrt{2}G_F N_e(r) E \begin{pmatrix} 1+\epsilon_{ee} & 0 & 0 \\ 0 & \epsilon_{\mu\mu} & 0 \\ 0 & 0 & \epsilon_{\tau\tau} \end{pmatrix},$$

- **New degeneracies with NSI**
- $P(\epsilon_{ee}, \delta_{CP}) = P(\epsilon'_{ee}, \delta'_{CP})$

$$P(\epsilon_{ee}, \delta_{CP}) = P(-\epsilon_{ee} - 2, \delta'_{CP})$$

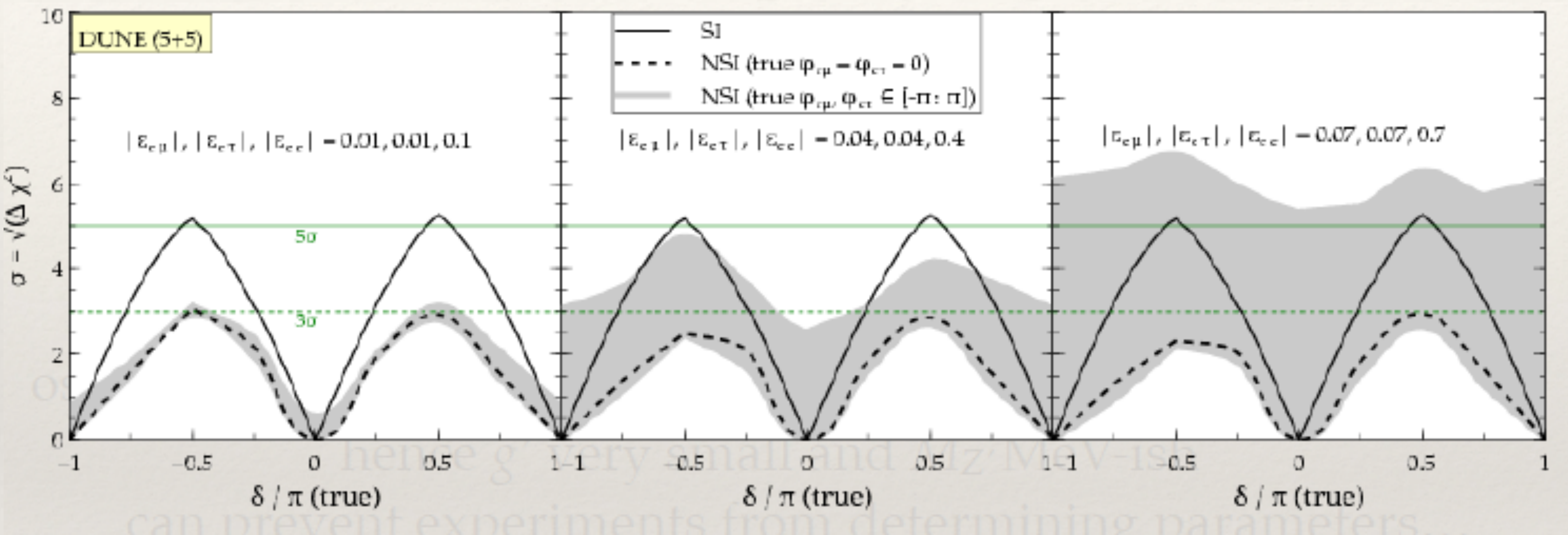
Coloma, Schwetz, PRD 2016

$$(1+\epsilon_{ee}) \rightarrow -(1+\epsilon_{ee})$$

In matter potential

Spoils the hierarchy sensitivity of Dune

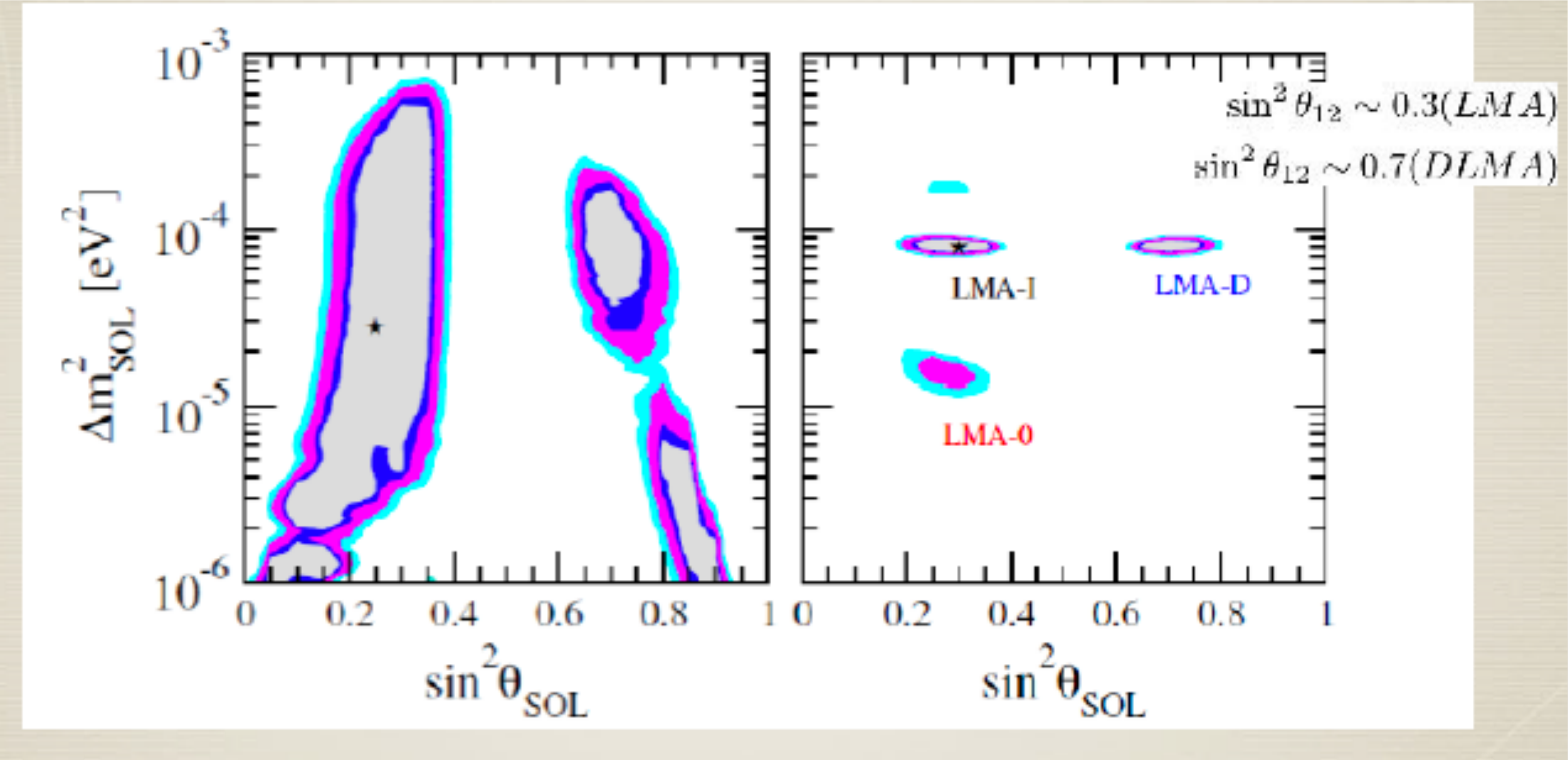
NSI and CP sensitivity



Mehta, Masood, 1603.01380

NSI can spoil CP sensitivity

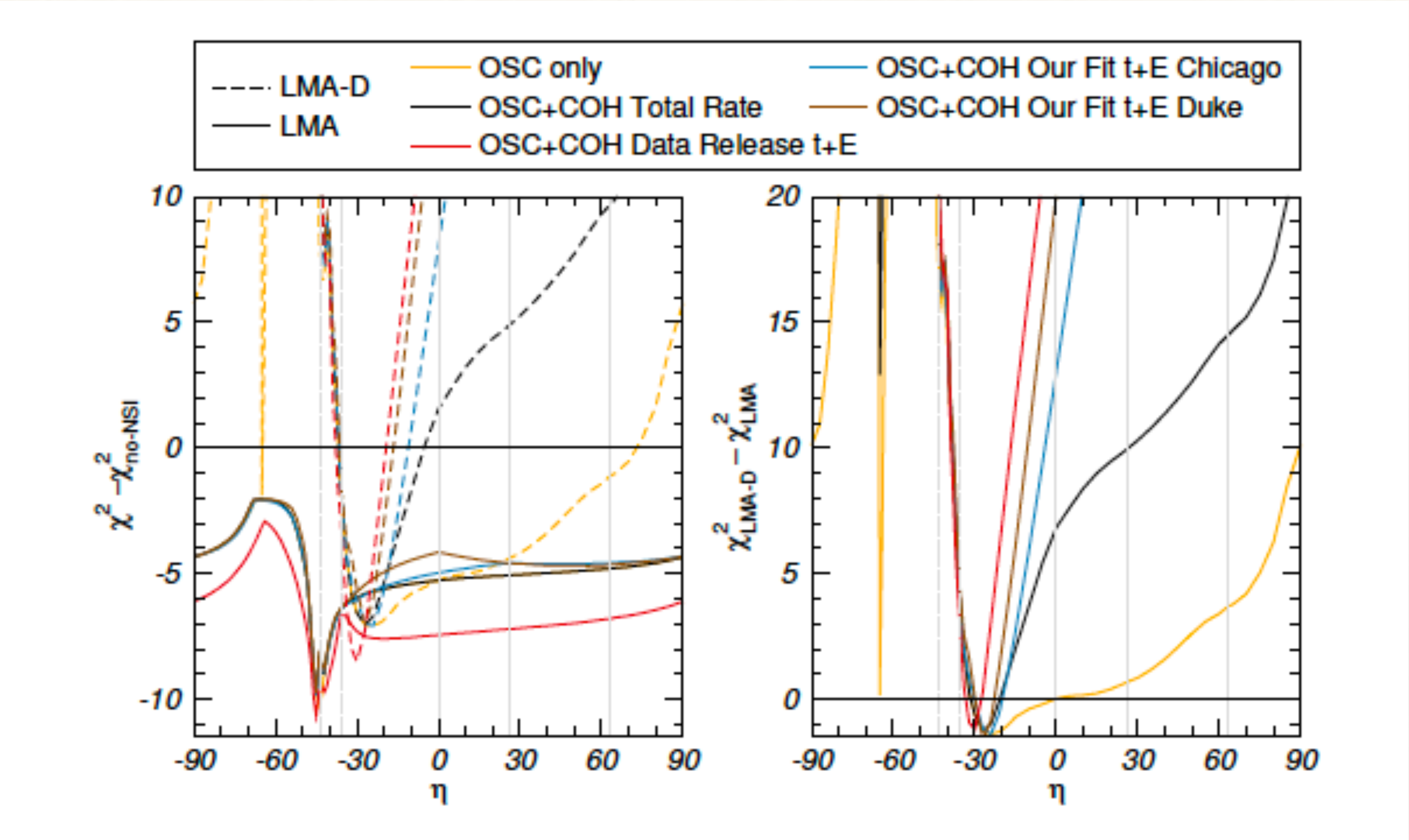
Dark-LMA solution



- Degenerate solution to the solar neutrinos problem
- $\sin^2\theta_{12} > 45^o$

Miranda, Tortola, Valle ,JHEP. 10 (2006) 008.

COHERENT constraints

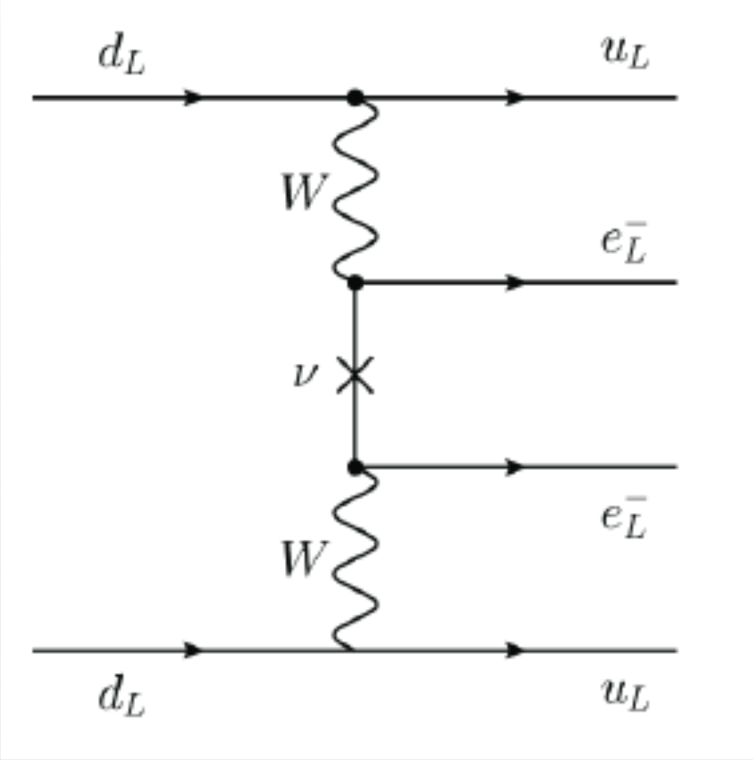


- Highly constrained by COHERENT using energy spectrum information

Coloma, Esteban,Gonzalez-Garcia, Maltoni 1911.09109

Neutrinoless double beta decay

$$(A,Z) \to (A,Z+2) + 2\,e^- \quad (0\nu\beta\beta)$$



- Standard picture $0\nu\beta\beta$ mediated by light neutrinos
- The half-life for $0\nu\beta\beta$,

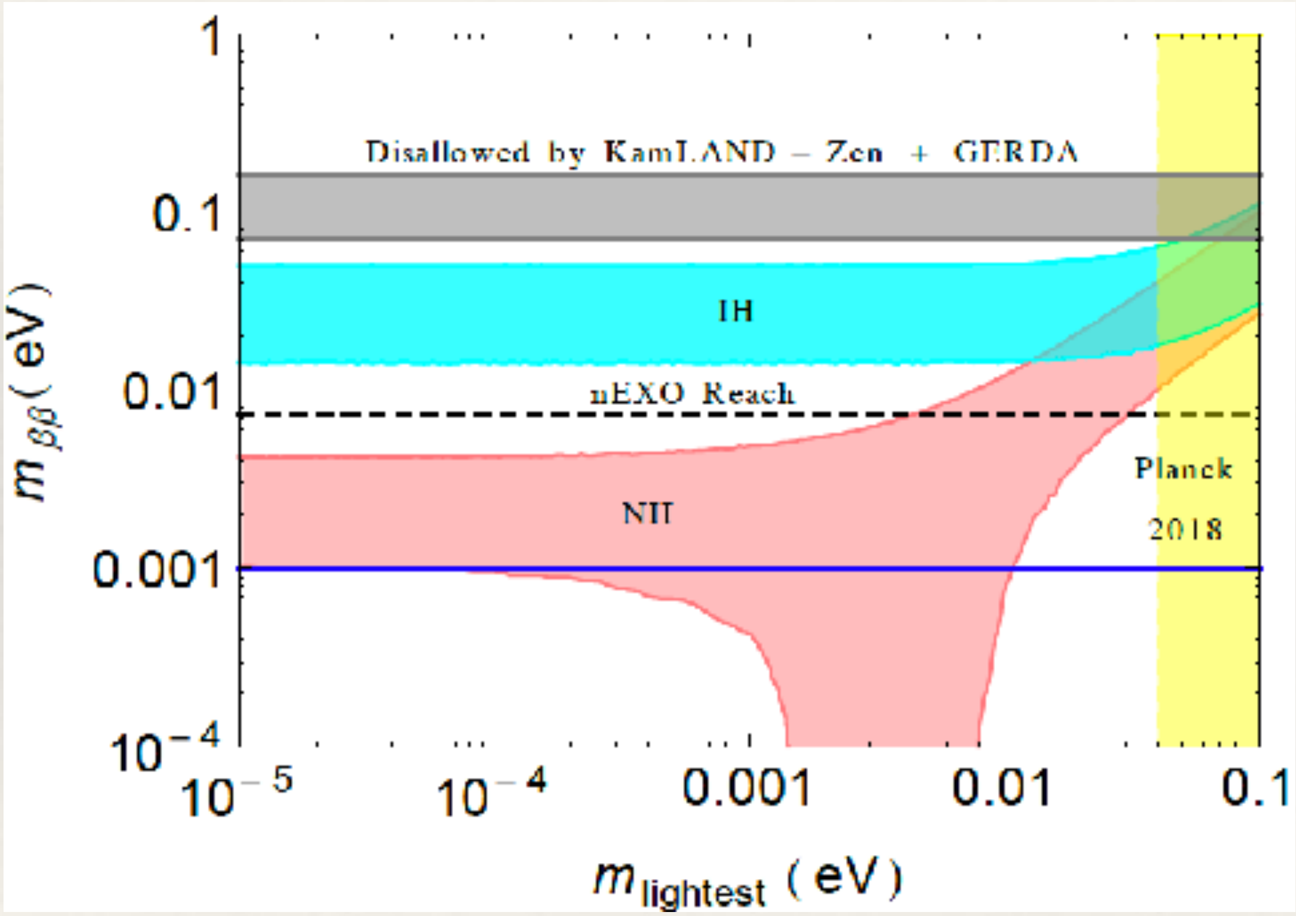
$$\frac{1}{T_{1/2}^{0\nu}} = G|\mathcal{M}_\nu|^2 \left|\frac{m^\nu_{ee}}{m_e}\right|^2 ,$$

G contains the phase space factor
M_ν is the nuclear matrix elements

- $|m_\nu^{ee}| = |U_{ei}^2\, m_i| \to$ the effective mass

Current and future sensitivity

$$|m_\nu^{ee}| = |m_1 U_{e1}^2 + m_2 U_{e2}^2 e^{2i\alpha_1} + m_3 U_{e3}^2 e^{2i\alpha_2}$$



Current Sensitivity

KamLAND-ZEN : 61-165 meV
EXO 200 : 93-286 meV
GERDA : 110-260 meV
CUORE : 110-520 meV

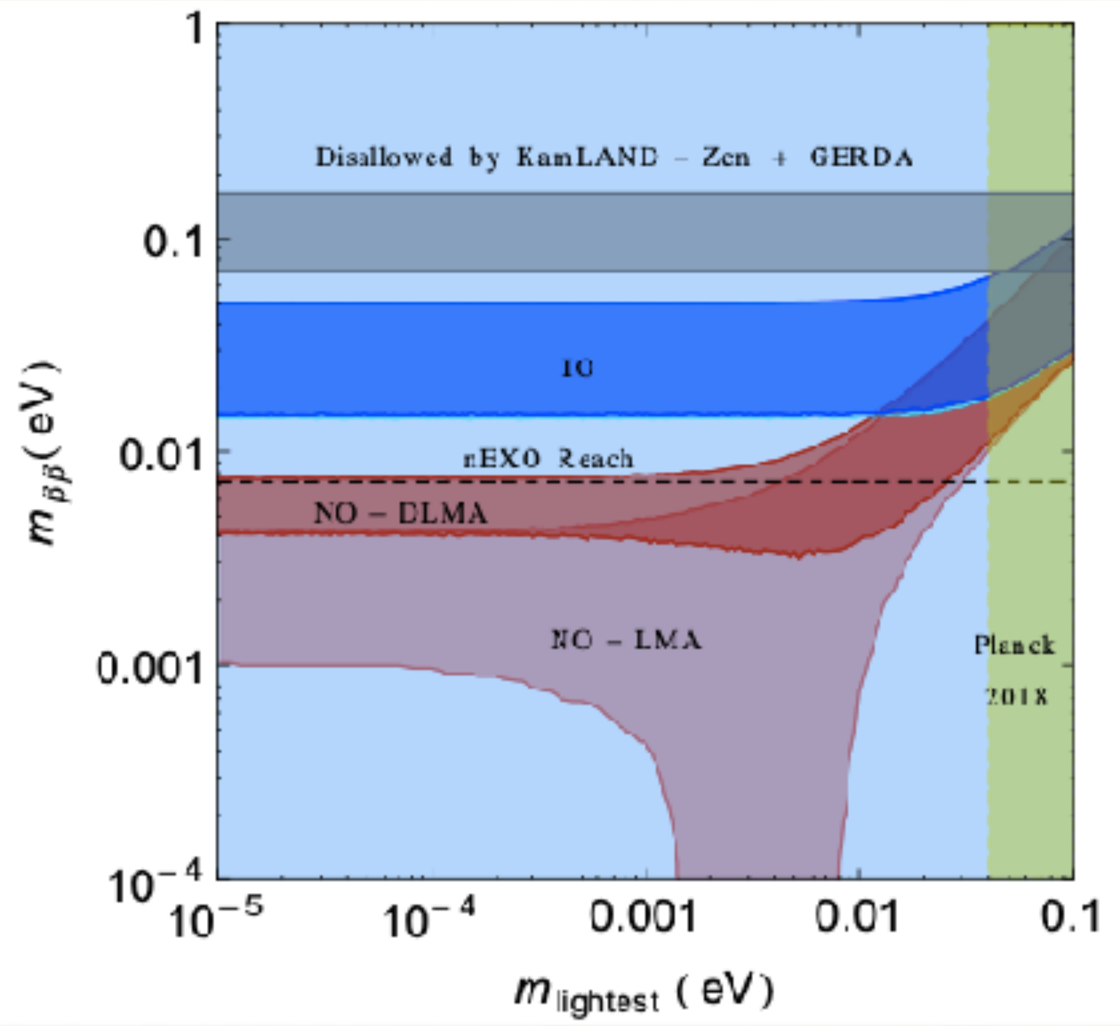
Future sensitivity

0.008 - 0.3 eV : IH can be confirmed
0.003 - .008 eV : 1-10 ton detector
0.001 - .003 eV : 10-100 ton detector
ultimate sensitivity

Barabash, 1901.11342

NSI and Neutrinoless double beta decay

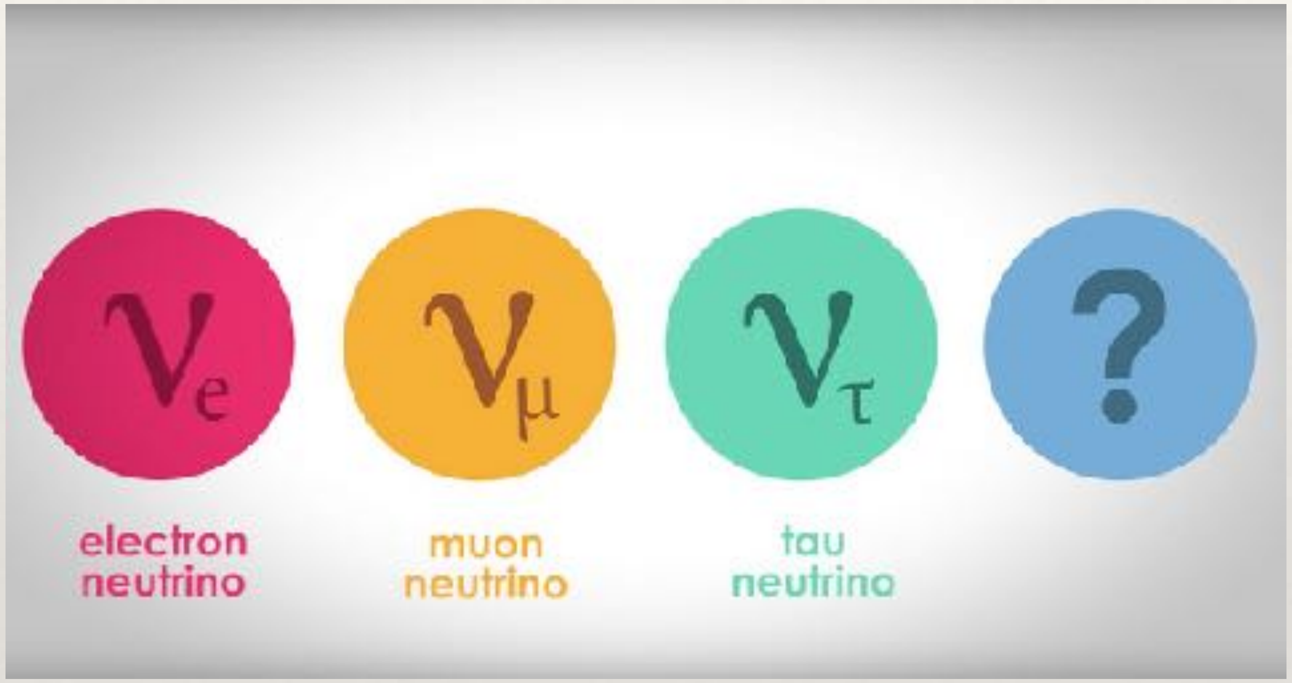
$$m_{\beta\beta} = |m_1 c_{12}^2 c_{13}^2 + m_2 s_{12}^2 c_{13}^2 e^{2i\alpha_2} + m_3 s_{13}^2 e^{2i\alpha_3}|$$



- New predictions in presence of NSI for NH
- Within reach of 10 kt detectors
- New sensitivity goal
- For NH degeneracy between LMA and DLMA can be broken for lower values of lightest neutrino mass
- Model independent

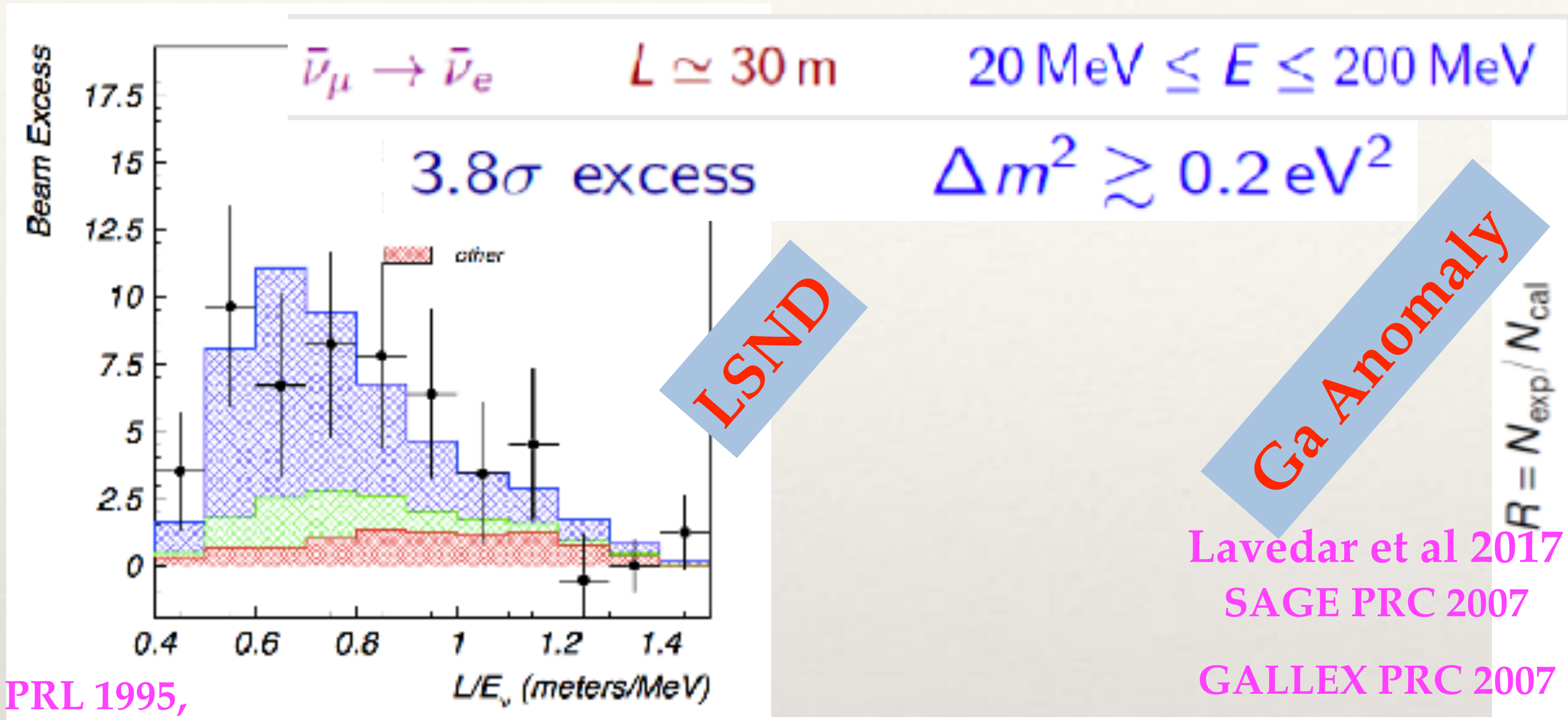
Vishnudath, Choubey, Goswami, Phys, Rev. D. (2019)

Are there more than 3 neutrinos ?

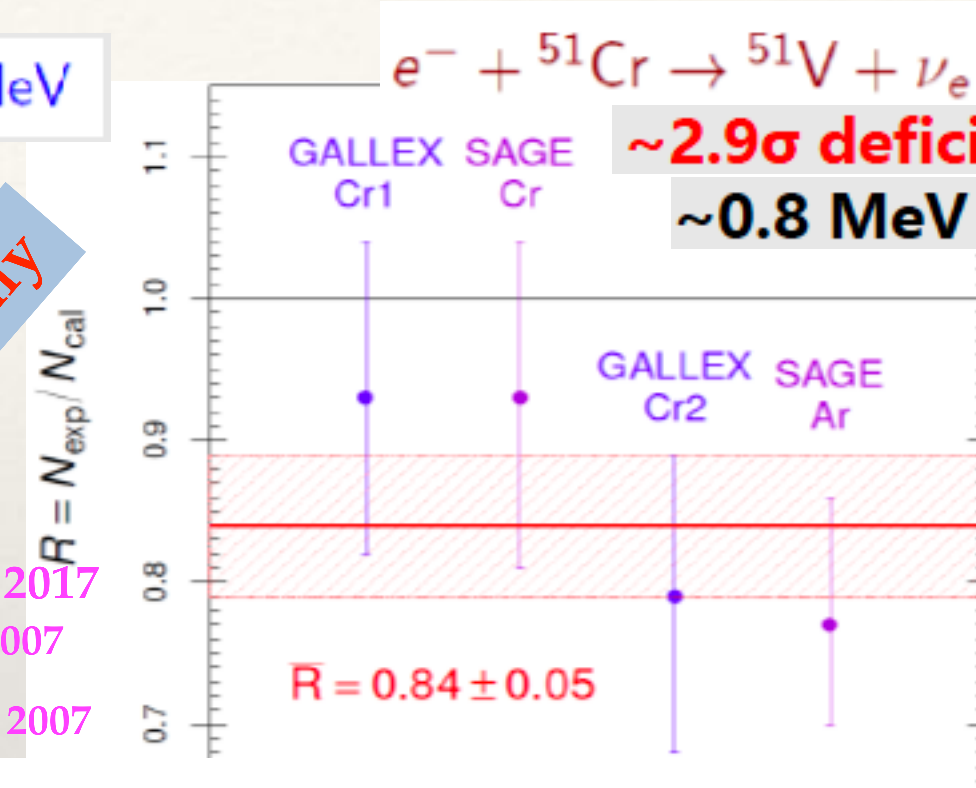


Extra sterile neutrino ?
Light or heavy or both ?

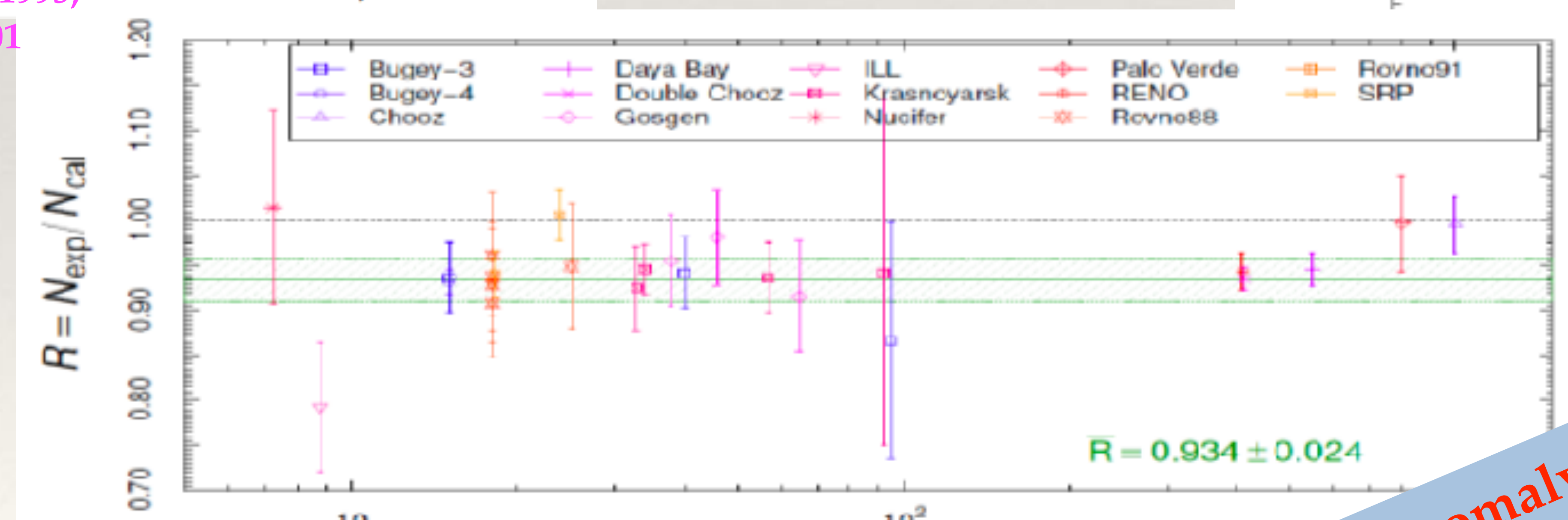
Sterile Neutrinos : indications



LSND PRL 1995, PRD 2001



Lavedar et al 2017
SAGE PRC 2007
GALLEX PRC 2007



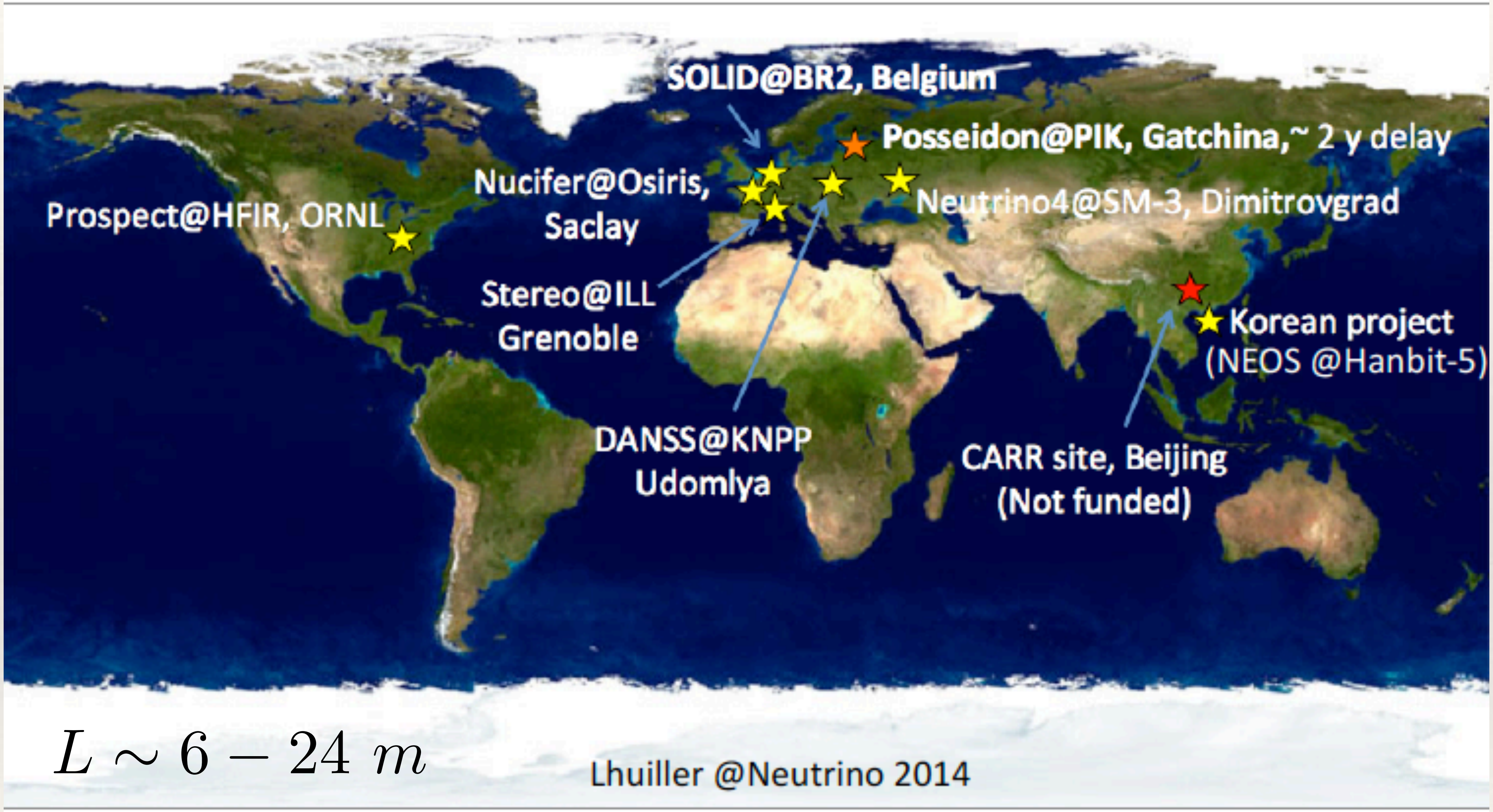
Mueller et al PRC 2011
Huber PRC 2011
Mention et al PRD 2011

$\Delta m^2_{41} = 2.4\ \mathrm{eV}^2$
$\sin^2(2\theta_{14}) = 0.14$

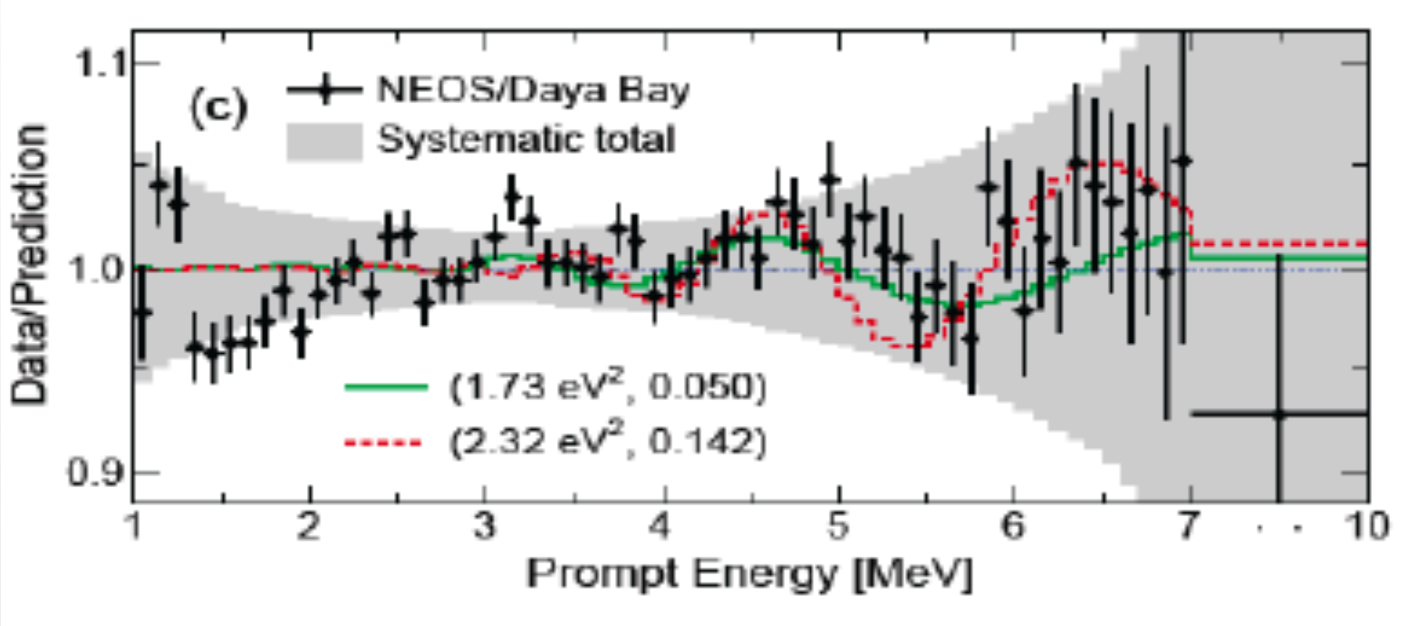
Reactor Anomaly

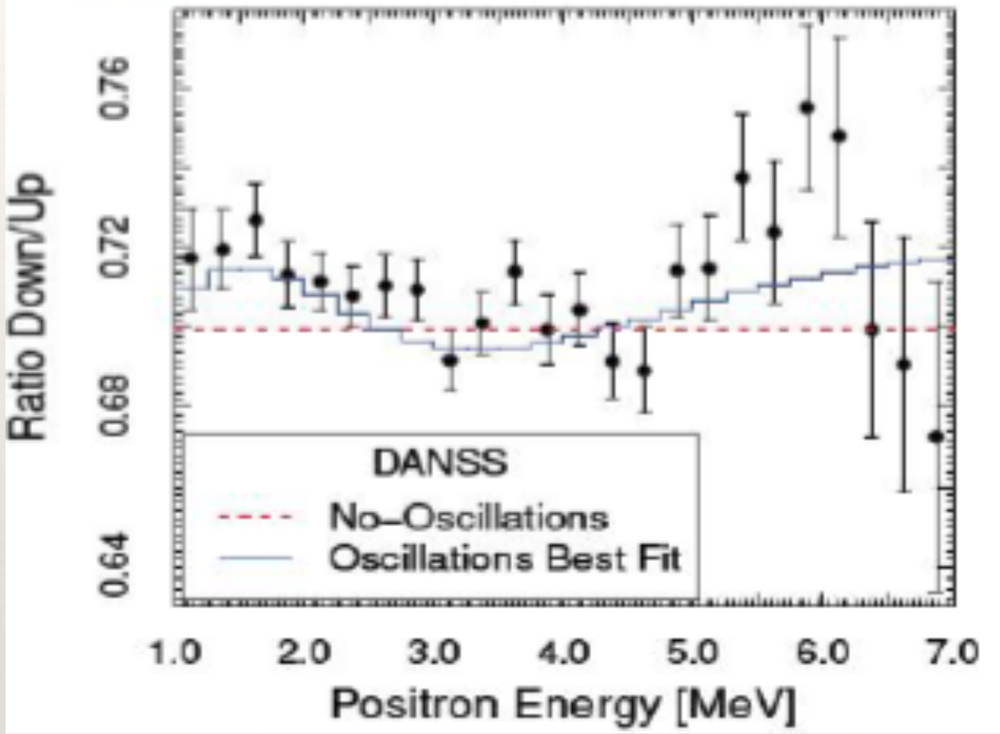
$$P \simeq 1 - \sin^2 2\theta_{14} \sin^2 \left[1.27 \frac{\Delta m^2_{41} L}{E_\nu} \left(\frac{\mathrm{eV}^2 \cdot \mathrm{m}}{\mathrm{MeV}} \right) \right]$$

Very short baseline reactor experiments



NEOS and DANSS





Neos Collaboration PRL 2016

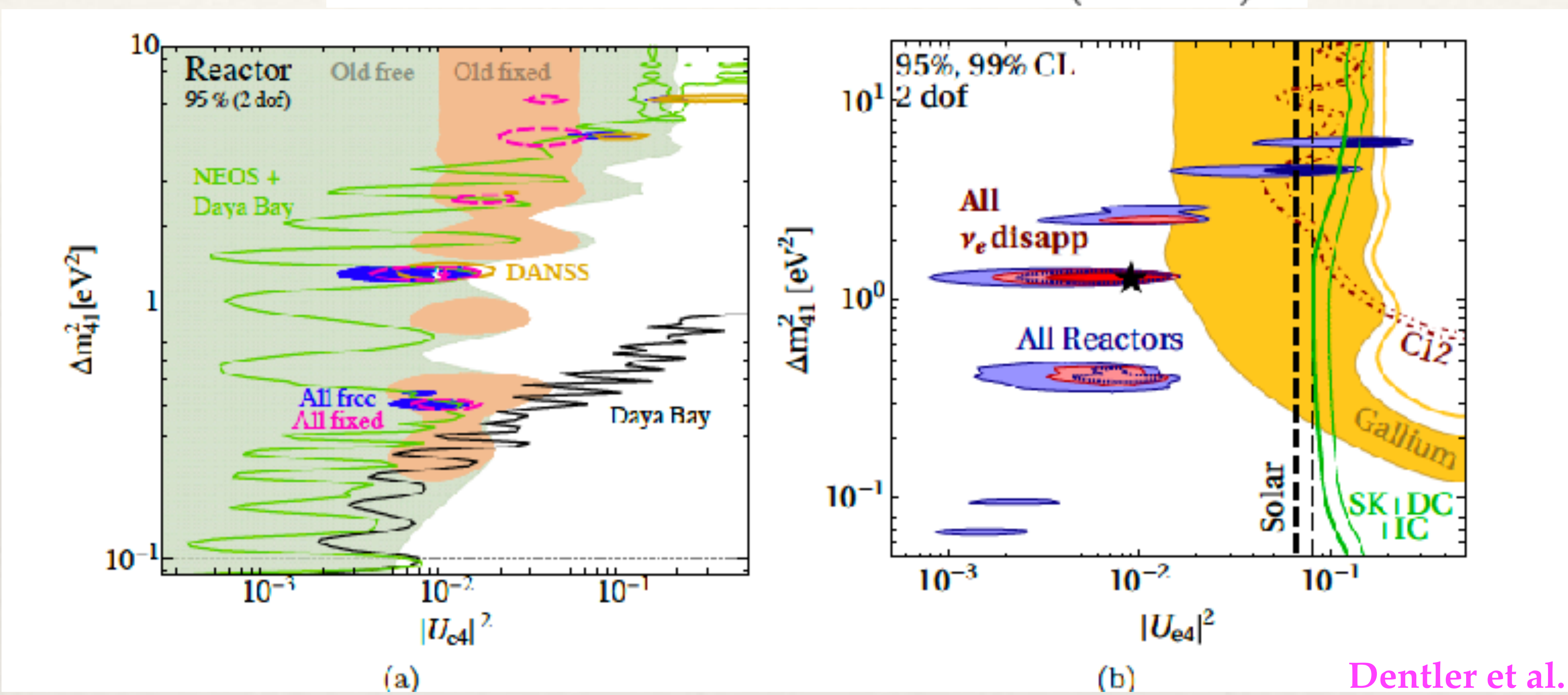
Danss collaboration, PLB 2018

Comparison of measured spectra at different baselines

Insensitive to flux calculation uncertainty

Bounds from reactor searches

$$P_{\alpha\alpha}^{\rm SBL} = 1 - 4|U_{\alpha 4}|^2(1-|U_{\alpha 4}|^2)\,\sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$



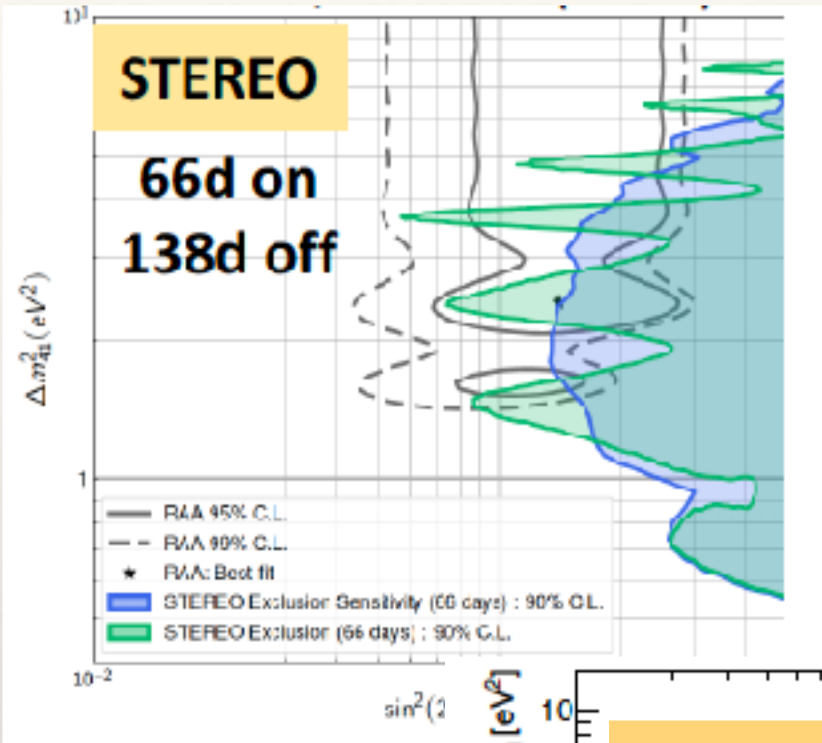
Dentler et al. JHEP 2018

Blue shaded regions allowed by fitting all reactor data with free fluxes

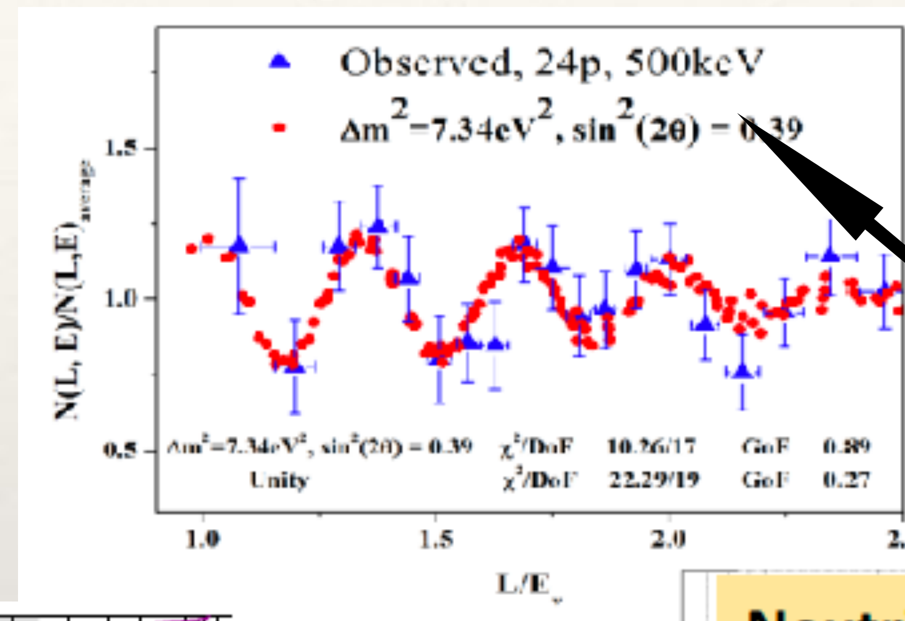
DANSS 2019 results give a lower Δm^2_{41}

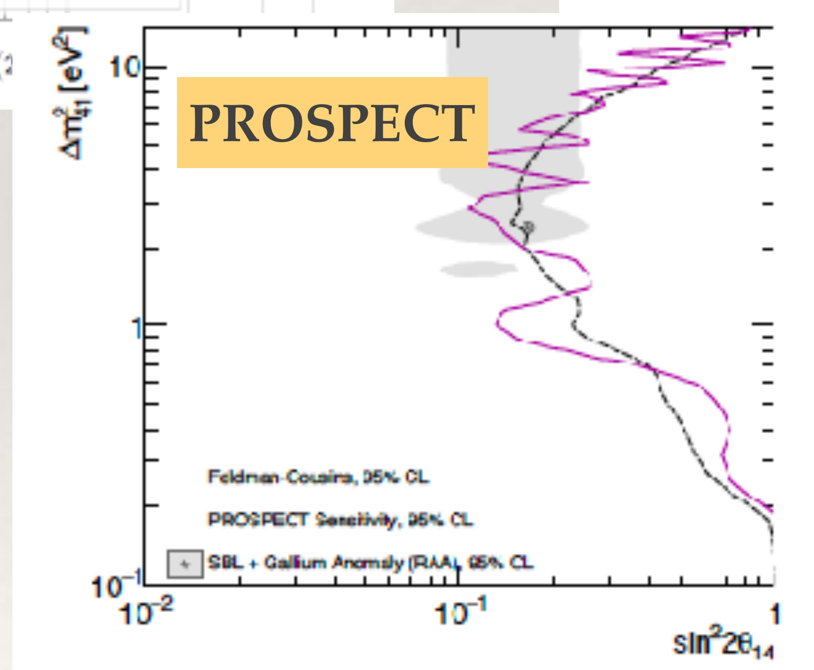
Danilov, talk at EPSHEP 2019
Ternes talk at CERN 2019

New results from reactor experiments

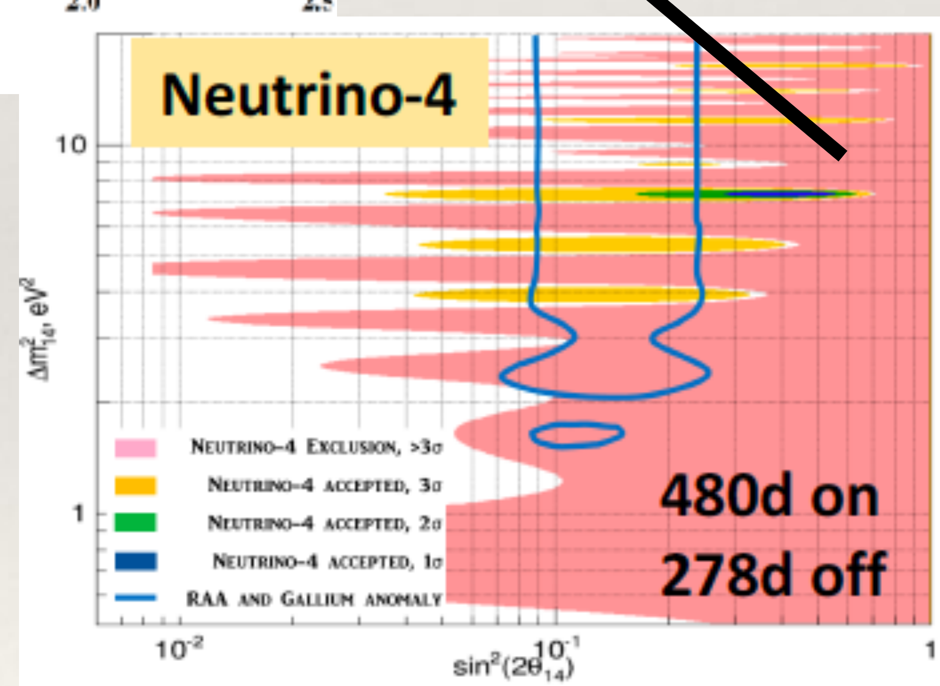


PRL 121 160821 2018





PRL 12 251802, 2018

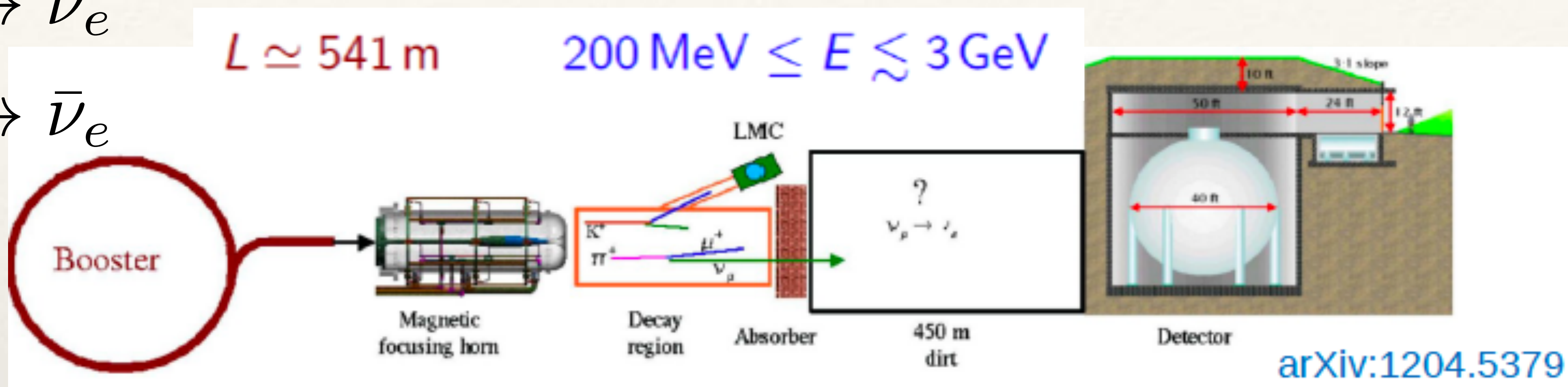


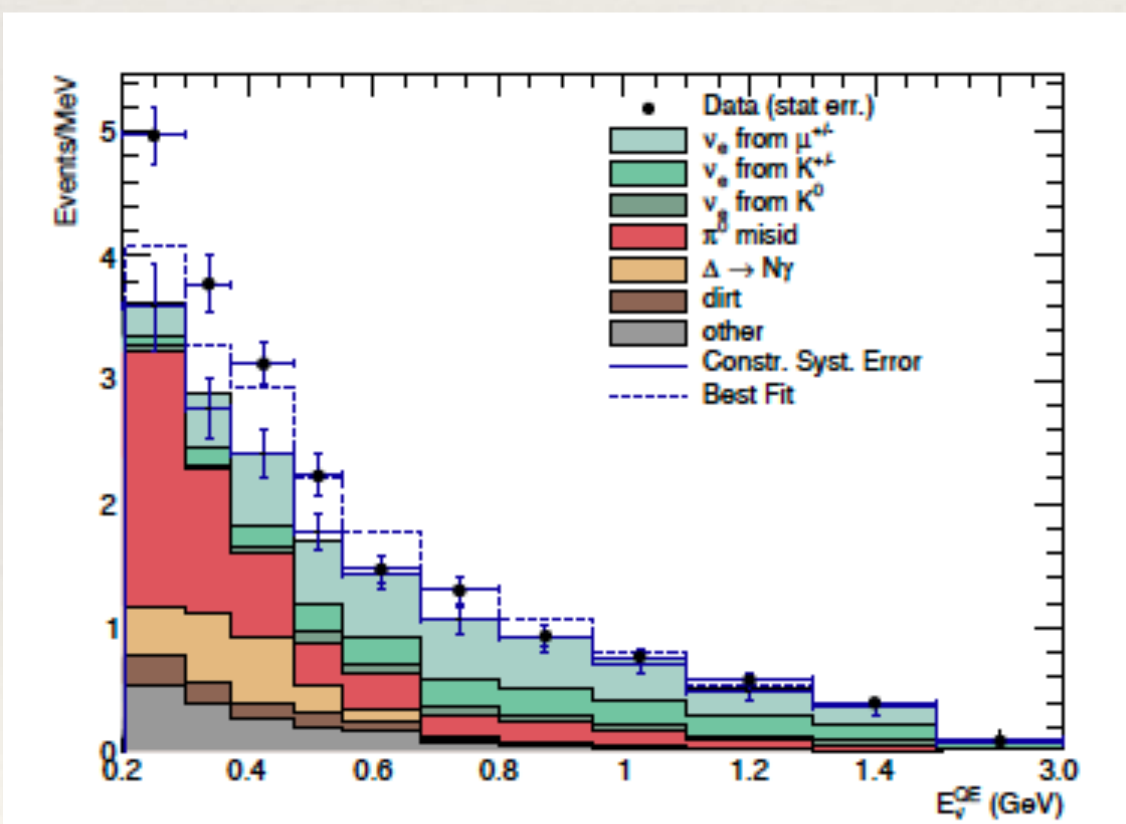
JETP letters 109, 213 2019

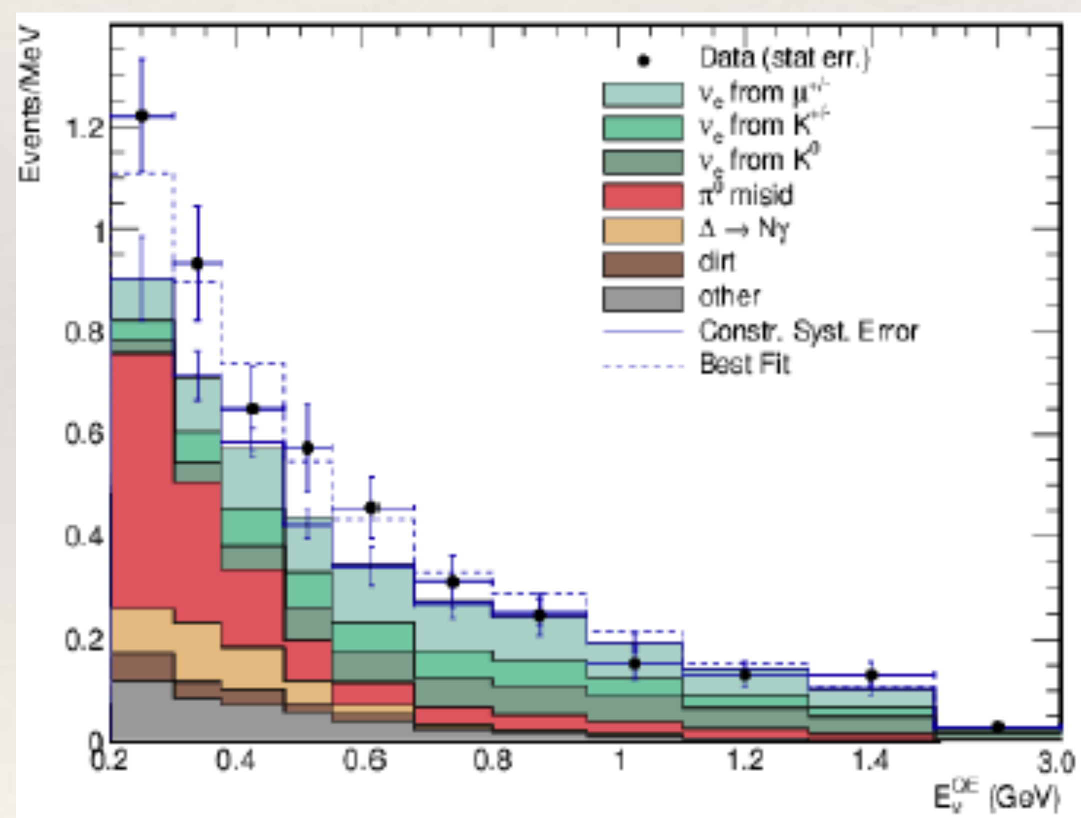
MiniBoone

$\nu_\mu \to \nu_e$

$\bar{\nu}_\mu \to \bar{\nu}_e$



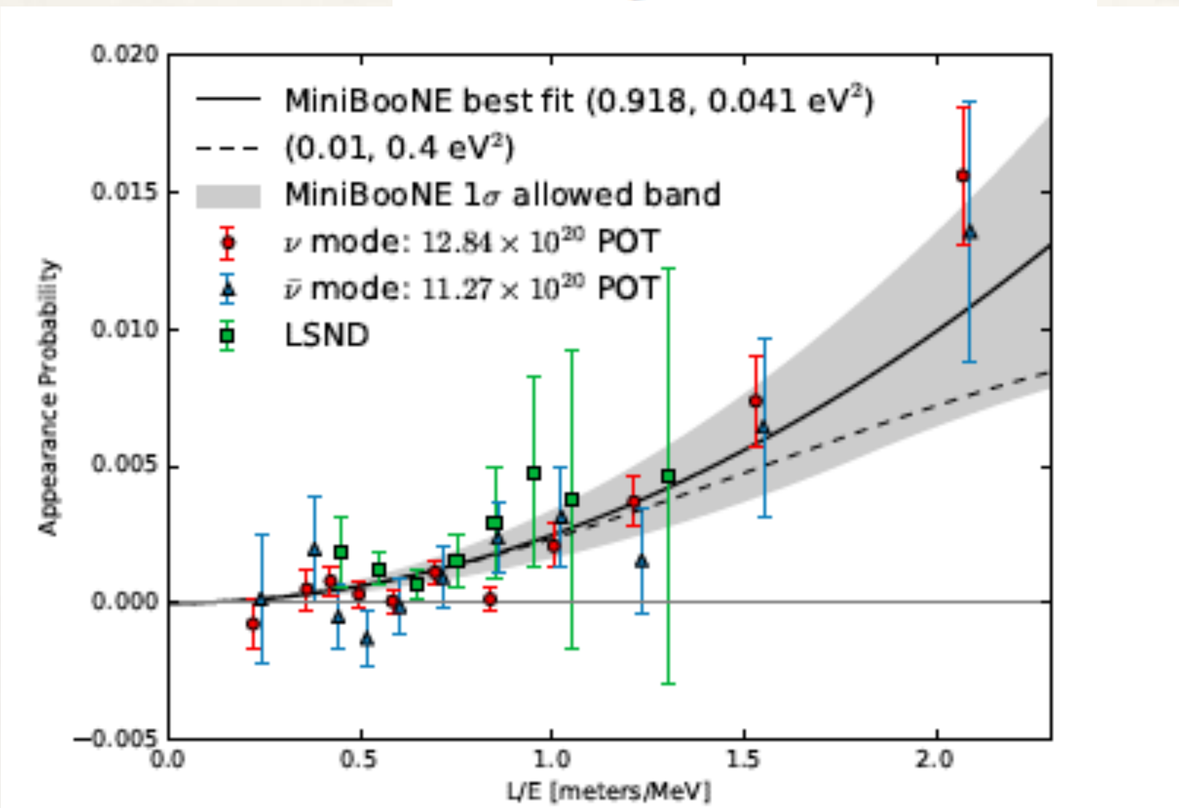




A.A. Aguilar Arevalo, PRL 121, 221801, 2018.

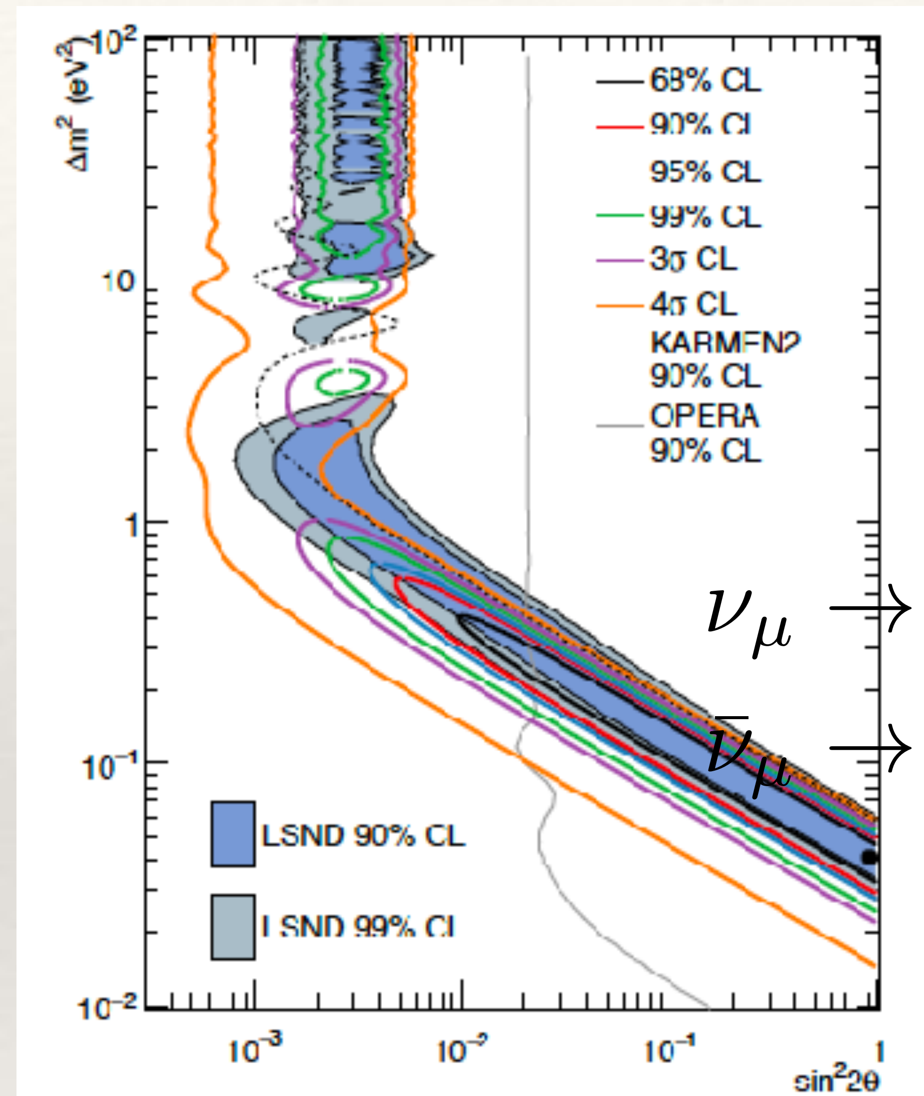
LSND and MiniBoone

$$200 < E_\nu^{QE} < 3000 \text{ MeV}$$



Combined significance $\sim 6\sigma$

A.A. Aguilar Arevalo, PRL 121, 221801, 2018.



$\nu_\mu \to \nu_e$

$\bar{\nu}_\mu \to \bar{\nu}_e$

Two neutrino fit

MiniBoone : neutrino + antineutrino

Disappearance and appearance tension

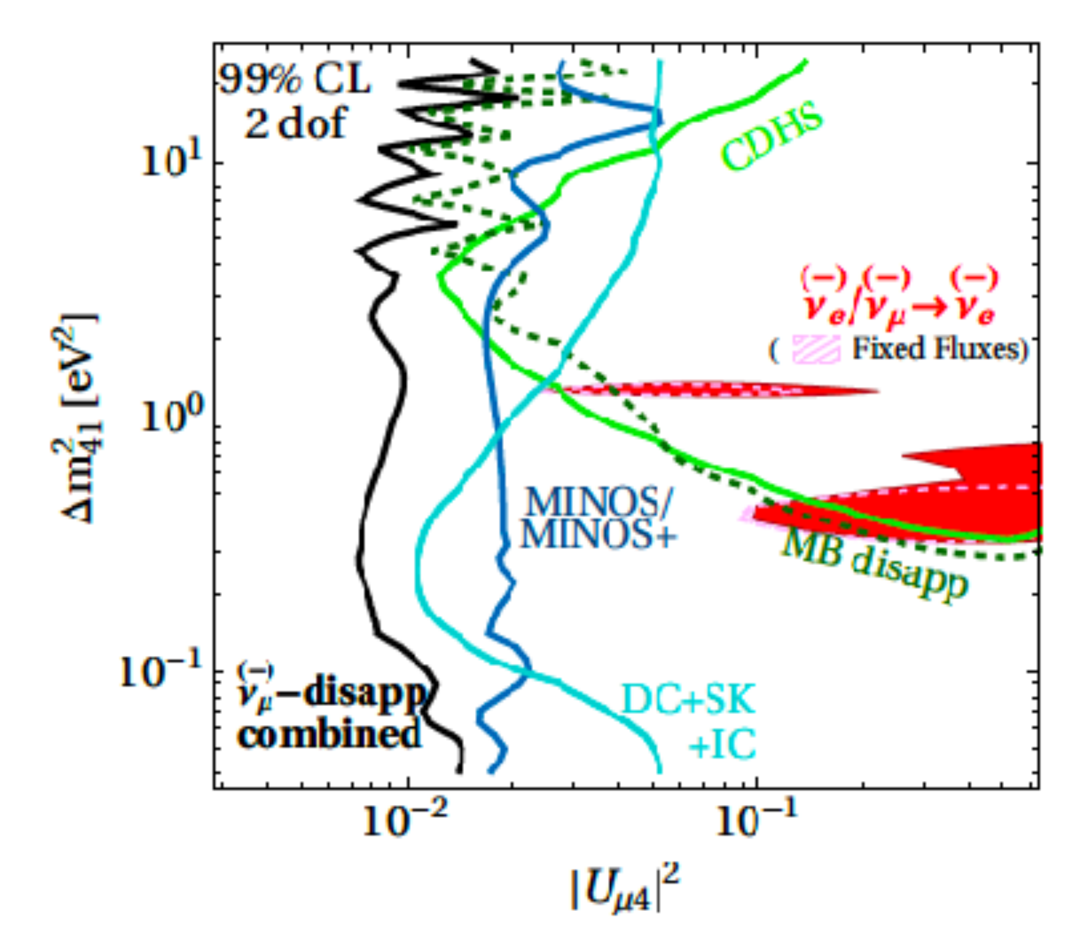
$$P_{\alpha\alpha}^{\rm SBL} = 1 - 4|U_{\alpha 4}|^2(1 - |U_{\alpha 4}|^2)\,\sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

$$P_{\alpha\beta}^{\rm SBL} = 4|U_{\alpha 4}|^2|U_{\beta 4}|^2\,\sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right).$$

P_{ee} depends on $|U_{e4}|^2$

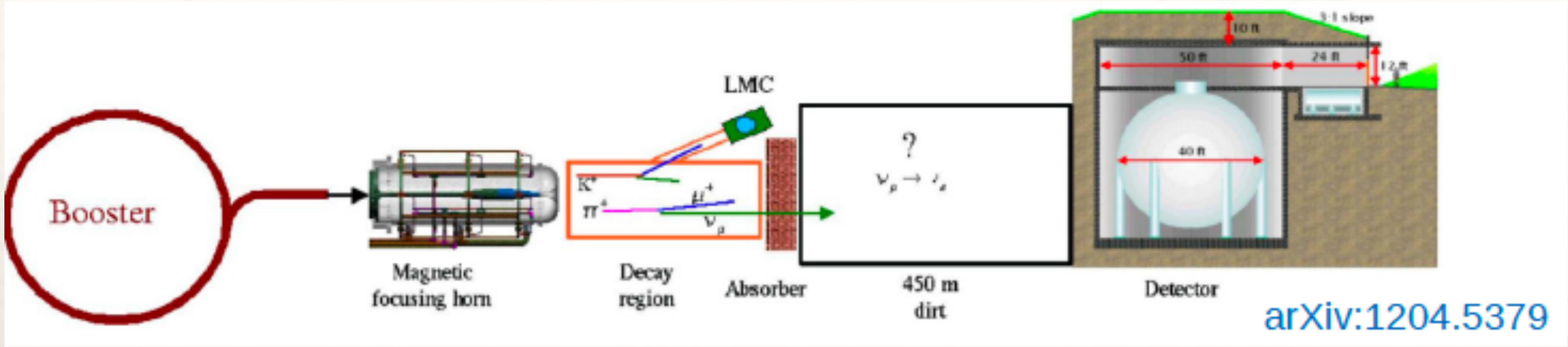
$P_{\mu\mu}$ depends on $|U_{\mu 4}|^2$

$P_{\mu e}$ depends on $|U_{\mu 4}|^2|U_{e4}|^2$



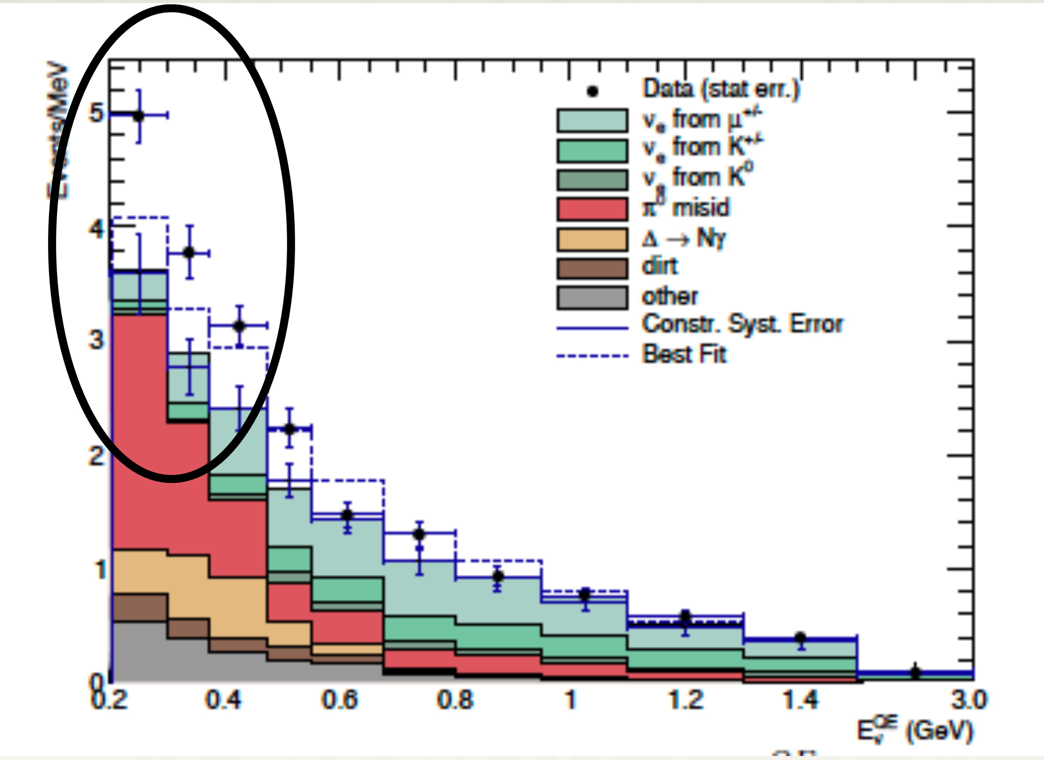
Dentler, Hernandez-Cabezudo,Kopp, Maltoni, Schwetz,JHEP 2017

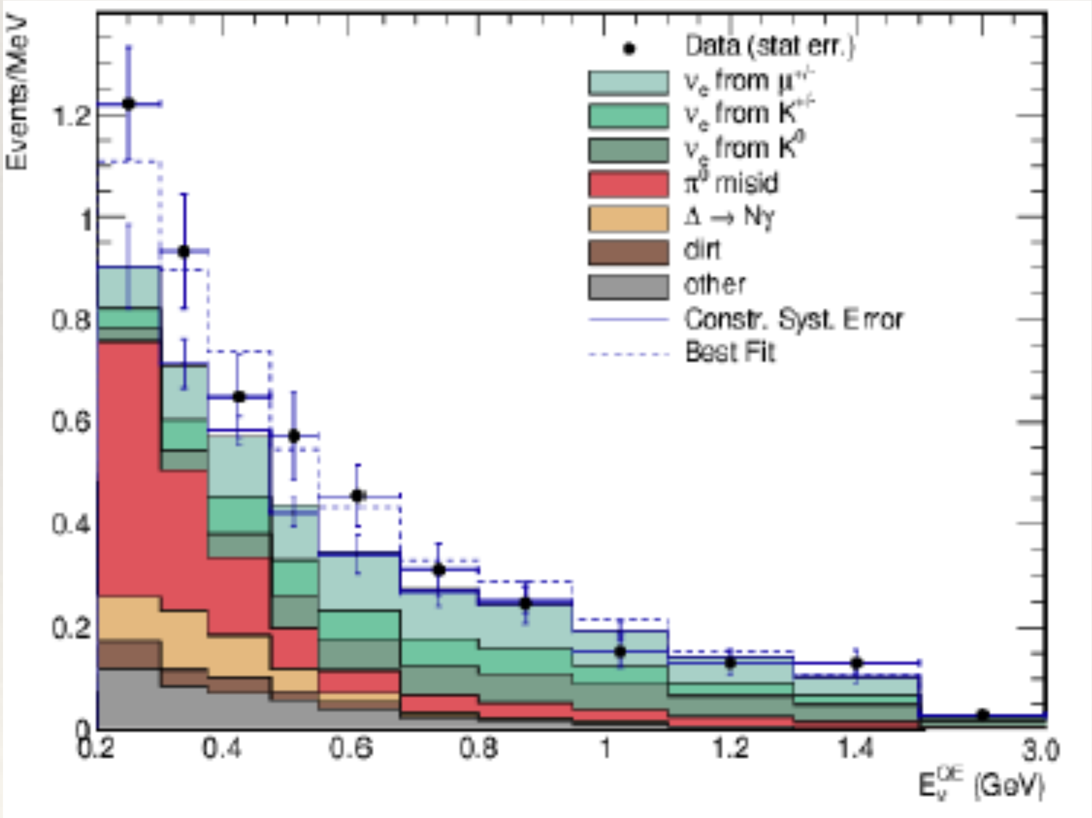
Low Energy Excess in MiniBoone



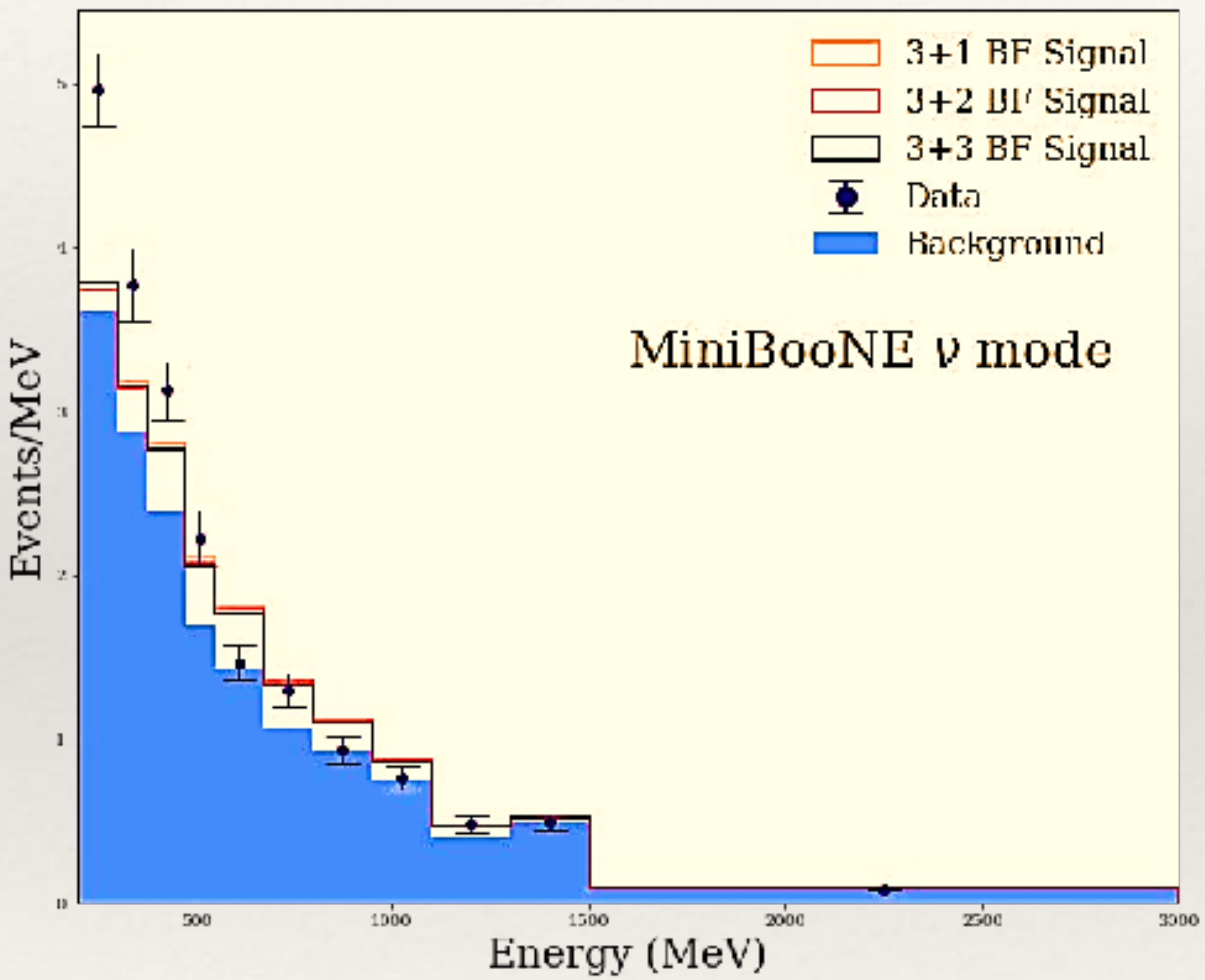
4.5σ

Excess





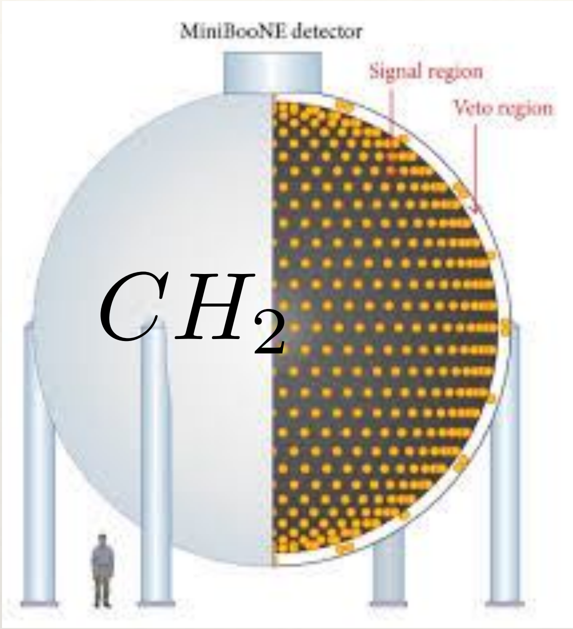
Can sterile neutrinos explain this ?

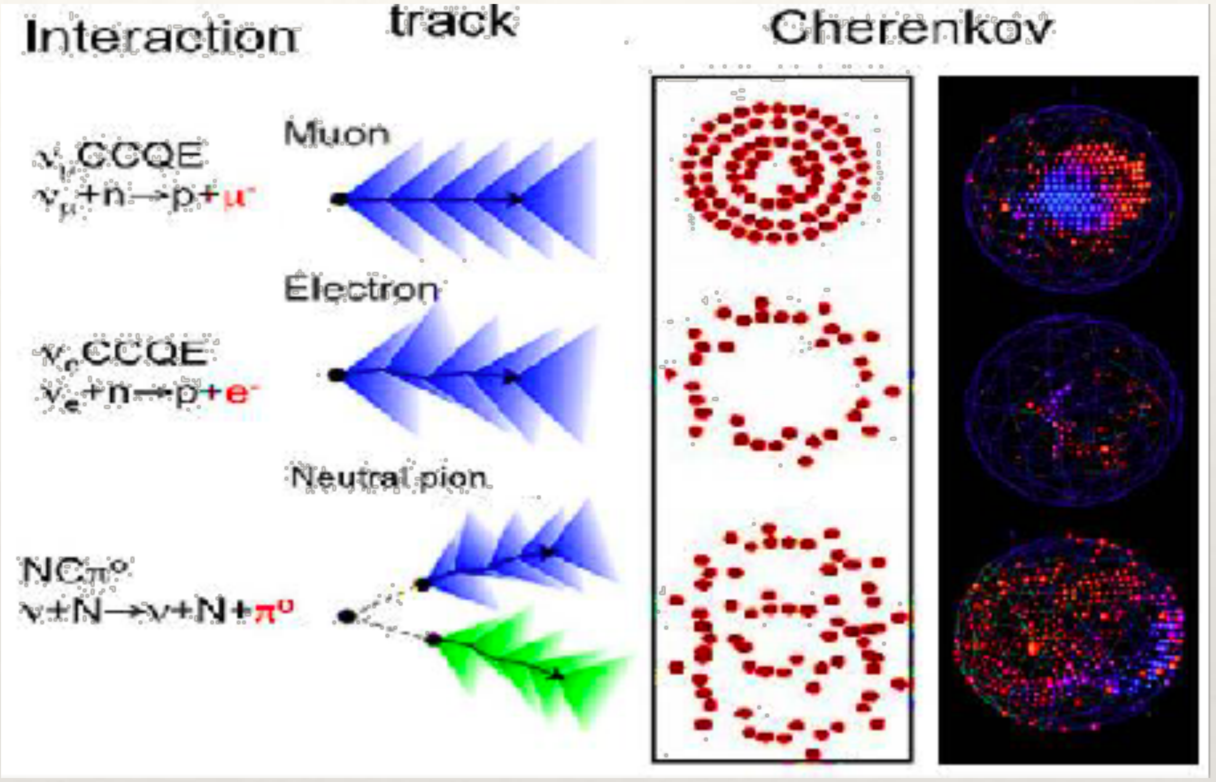


3+N sterile neutrino scenario cannot explain the MiniBoone low energy excess

D. Cianci , talk at Applied Antineutrino Physics Workshop 2018

Is it due to background effect ?





From : S. Jana , Pheno 2019

Cannot distinguish between Cherenkov cone of electrons and single photon

The single photons coming from NC background cannot explain the excess

Alternative explanations

Dark Neutrino Portal to Explain MiniBooNE excess
Enrico Bertuzzo (Sao Paulo U.), Sudip Jana (Oklahoma Ctr. High Energy Phys. & Oklahoma State
Published in **Phys.Rev.Lett. 121 (2018) no.24, 241801**

Explaining the MiniBooNE excess by a decaying sterile neutrino with mass in the 250 MeV range
Oliver Fischer, Álvaro Hernández-Cabezudo, Thomas Schwetz (KIT, Karlsruhe, IKP). Sep 20, 2019. 26 pp.
e-Print: arXiv:1909.09561 [hep-ph] | PDF
References | BibTeX | LaTeX(US) | LaTeX(EU) | Harvmac | EndNote
ADS Abstract Service

U(1)' mediated decays of heavy sterile neutrinos in MiniBooNE
Peter Ballett, Silvia Pascoli (Durham U., IPPP), Mark Ross-Lonergan (Nevis Labs, Columbia U.). Aug 8, 2018. 8 pp.
Published in **Phys.Rev. D99 (2019) 071701**
IPPP/18/70

Testing New Physics Explanations of MiniBooNE Anomaly at Neutrino Scattering Experiments
Carlos A. Argüelles (MIT, Cambridge, Dept. Phys.), Matheus Hostert (Durham U., IPPP), Yu-Dai Tsai (Fermilab). Dec 20, 2018. 7 pp.
IPPP/18/113, FERMILAB-PUB-18-686-A-ND-PPD-T
e-Print: arXiv:1812.08768 [hep-ph] | PDF

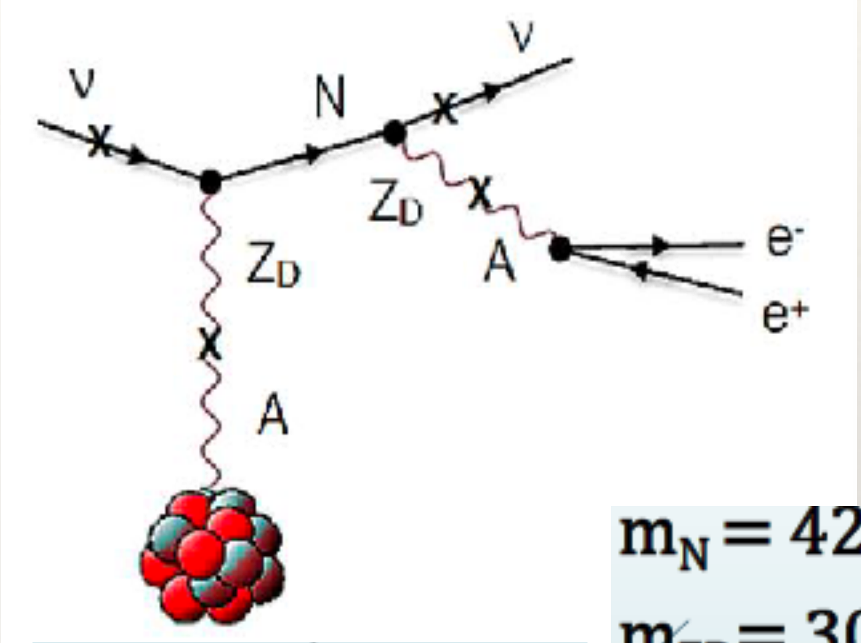
Severe Constraints on New Physics Explanations of the MiniBooNE Excess
Johnathon R. Jordan (Michigan U.), Yonatan Kahn (Princeton U. & Chicago U., KICP & Illinois U., Urbana (main))
2018. 7 pp.
Published in **Phys.Rev.Lett. 122 (2019) no.8, 081801**

Many more, apologies if your paper is not listed

Slide: D. Pramanik, Whepp 2019

Dark neutrino portal

$$\mathcal{L}_D \supset \frac{m_{Z_D}^2}{2} Z_{D\mu} Z_D^\mu + g_D Z_D^\mu J_{D\mu} + e\epsilon\, Z_D^\mu J_\mu^{\rm em} + \frac{g}{c_W}\epsilon'\, Z_D^\mu J_\mu^Z$$



$\alpha_D = 0.25$

$\alpha\varepsilon^2 = 2 \times 10^{-10}$

$\chi^2/\text{dof} = 33.2/36$

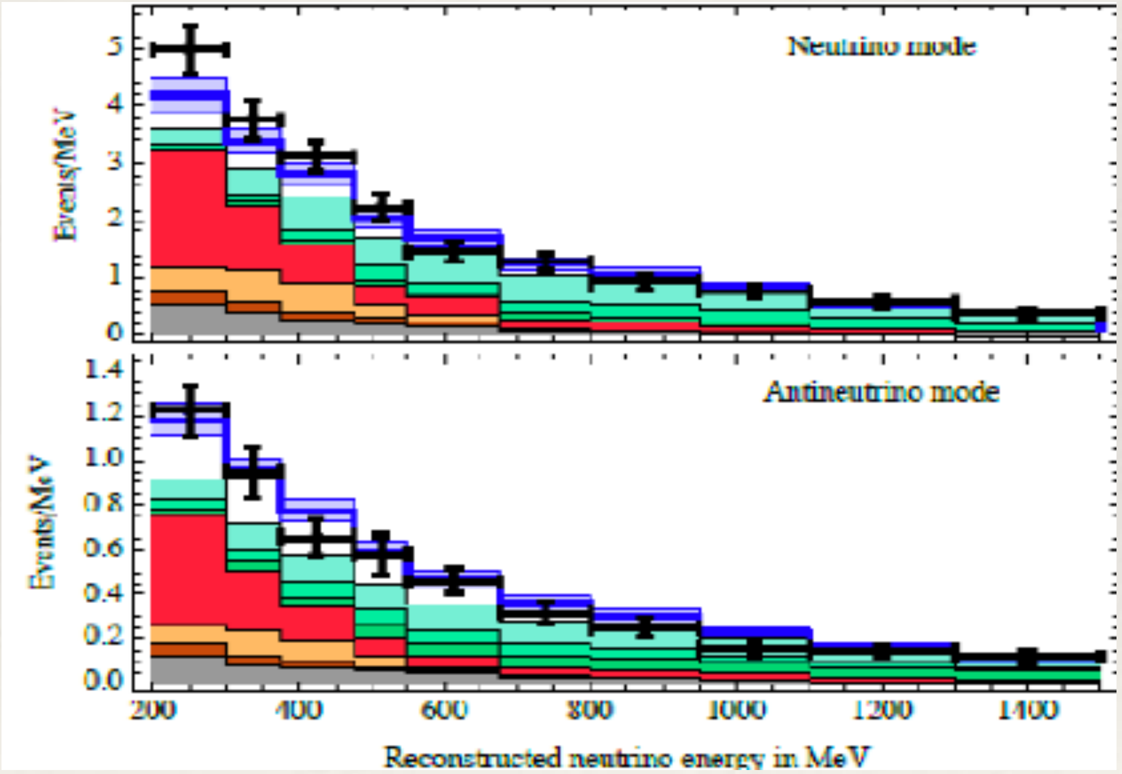
$m_N = 420$ MeV

$m_{ZD} = 30$ MeV

$|U_{\mu 4}|^2 = 9\text{x}10^{-7}$

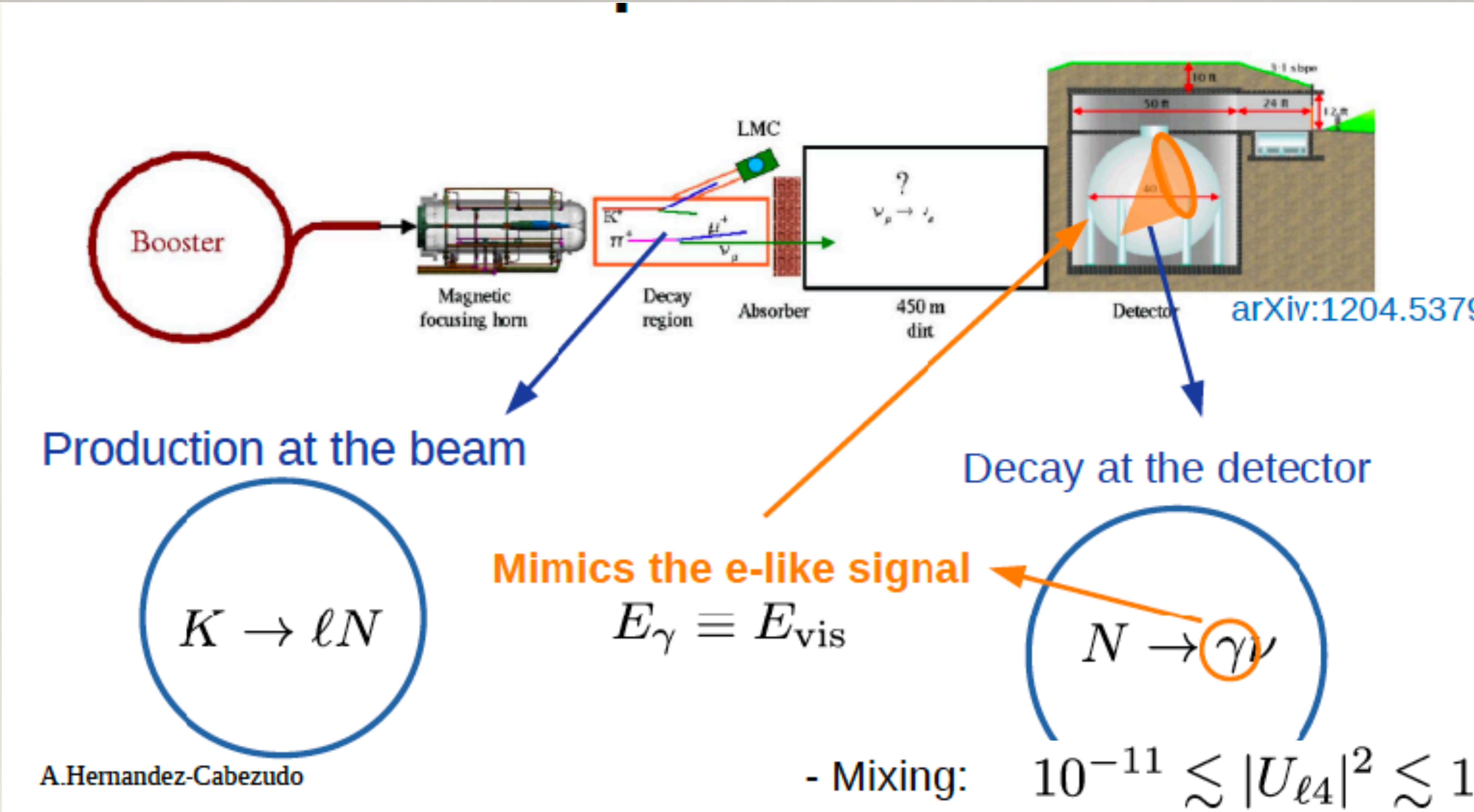
Explains the observed distribution

Right handed neutrinos part of dark sector



Bartuzzo, Jana, Machado, Zukanovich-Funchal, PRL 2018

Sterile neutrino decay

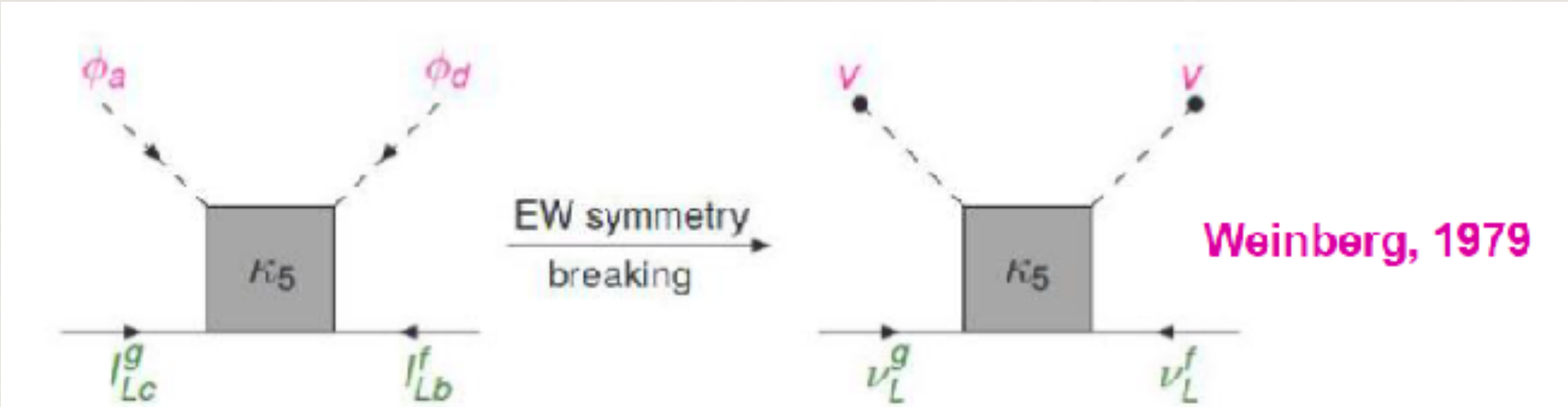


Fischer, Hernandez-Cabezudo, Schwetz, .1909 09501

- Mixing: $10^{-11} \lesssim |U_{\ell 4}|^2 \lesssim 10^{-7}$
- Mass: $\sim 250\,\text{MeV}$
- New physics scale: $10^4\,\text{TeV} \lesssim \Lambda \lesssim 10^7\,\text{TeV}$

Why neutrinos masses are small ?

- A natural way to explain small neutrino masses is via seesaw mechanism
- Relates smallness of neutrino masses with new physics at a high scale
- Tree level exchange of some heavy particle gives rise than effective dimension 5 operator at the low scale
- $\mathcal{L} = \kappa_5 l_L l_L \phi\phi, \quad \kappa_5 = y_\kappa/\Lambda$ $\qquad m_\nu \sim \kappa_5 v^2/\Lambda$



- Violation of lepton number → Majorana nature of neutrinos
- Radiative mass generation, Models with higher dimensional operators

Talk by M. Schimdt

Implementing Seesaw

- Models with extended particle content

Type -I seesaw : mediated by right handed neutrinos

Type-II seesaw : mediated by triplet Higgs

Type-III seesaw : mediated by triplet fermions

- Models with extended gauge groups

Grand Unified Theories , Left-Right symmetric Models , Models with extra U(1)

Implications of seesaw

- Neutrinos are Majorana particles $\Rightarrow$ neutrinoless double beta decay
- LNV decay of heavy mediators $\Rightarrow$ leptogenesis
- GUT scale seesaw no testability at colliders $\longrightarrow$ TeV Scale seesaw
- Extra singlets — inverse/linear seesaw
- Large light-heavy mixing $\Rightarrow$ testability at colliders
- Large light-heavy mixing $\Rightarrow$ Lepton flavour violation
- Stability of electroweak vacuum
- Combined constraints on model parameters
- Can one obtain thermal leptogenesis ? **Talk by T. Dutka**

What explains the mixing pattern ?

- Why two large and one small mixing unlike quark sector where all mixing angles are small ?
- Flavour symmetry — discrete non-abelian symmetries

Group	d	Irr. Repr.'s	Presentation
$D_3 \sim S_3$	6	$1,\ 1',\ 2$	$A^3 = B^2 = (AB)^2 = 1$
D_4	8	$1_1, \ldots 1_4, 2$	$A^4 = B^2 = (AB)^2 = 1$
D_7	14	$1,\ 1',\ 2,\ 2',\ 2''$	$A^7 = B^2 = (AB)^2 = 1$
A_4	12	$1,\ 1',\ 1'',\ 3$	$A^3 = B^2 = (AB)^3 = 1$
$A_5 \sim PSL_2(5)$	60	$1,\ 3,\ 3',\ 4,\ 5$	$A^3 = B^2 = (BA)^5 = 1$
T'	24	$1,\ 1',\ 1'',\ 2,\ 2',\ 2'',\ 3$	$A^3 = (AB)^3 = R^2 = 1,\ B^2 = R$
S_4	24	$1,\ 1',\ 2,\ 3,\ 3'$	$BM: A^4 = B^2 = (AB)^3 = 1$ $TB: A^3 = B^4 = (BA^2)^2 = 1$
$\Delta(27) \sim Z_3 \rtimes Z_3$	27	$1_1, \ldots 1_9, 3, \bar{3}$	
$PSL_2(7)$	168	$1, 3, \bar{3}, 6, 7, 8$	$A^3 = B^2 = (BA)^7 = (B^{-1}A^{-1}BA)^4 = 1$
$T_7 \sim Z_7 \rtimes Z_3$	21	$1, 1', \bar{1'}, 3, \bar{3}$	$A^7 = B^3 = 1,\ AB = BA^4$

Many options

Many scalars (flavons)

Modular invariance

Feruglio. 1706.08749

Altarelli, Feruglio, Rev. Mod. Phys. 2010

Generation of non-zero θ_{13}

Concluding Remarks

- Three neutrino oscillation paradigm well established
- Future experiments expected to determine hierarchy, octant and CP
- Can new physics be probed in these experiments — extra parameters giving rise to additional degeneracies
- Complementary information
- Sterile neutrino — oscillation explanation is trouble
- MiniBoone low energy excess — many ideas
- Future neutrinoless double beta decay experiments can test IO — new physics can give different predictions

Concluding Remarks

- Origin of neutrino masses and mixing — still under mist.


Thank
you!!


BACKUP

TeV scale seesaw

- Interesting from the point of view of testability in colliders
- Add extra gauge singlets with opposite lepton number

$$-L_{mass} = \overline{\nu}_L M_D N_R + \overline{\nu}_L M_s \nu_s + \overline{N_R^c} M_R \nu_s + \frac{1}{2} \overline{\nu_s^c} M_\mu \nu_s + \frac{1}{2} \overline{N_R^c} M_N N_R + \text{h.c.}$$

-

$$M_\nu = \begin{pmatrix} 0 & m_D^T & m_S^T \\ m_D & M_N & M_R^T \\ m_S & M_R & \mu \end{pmatrix}$$

$$m_D = Y_\nu v/\sqrt{2} \text{ and } m_S = Y_s v/\sqrt{2}.$$

$$m_D = Y_\nu v/\sqrt{2} \text{ and } m_S = Y_s v/\sqrt{2}.$$

Neutrinoless double beta decay experiments

Experiment	Isotope	Techinique	Total mass [kg]	Exposure [kg yr]	FWHM @ $Q_{\beta\beta}$ [keV]	Background [counts/keV/kg/yr]	$S^{0\nu}$ (90% C. L.) [10^{25} yr]
Past							
Cuoricino, [177]	^{130}Te	bolometers	40.7 (TeO_2)	19.75	5.8 ± 2.1	0.153 ± 0.006	0.24
CUORE-0, [178]	^{130}Te	bolometers	39 (TeO_2)	9.8	5.1 ± 0.3	0.058 ± 0.006	0.29
Heidelberg-Moscow, [179]	^{76}Ge	Ge diodes	11 (^{enr}Ge)	35.5	4.23 ± 0.14	0.06 ± 0.01	1.9
IGEX, [180, 181]	^{76}Ge	Ge diodes	8.1 (^{enr}Ge)	8.9	~ 4	$\lesssim 0.06$	1.57
GERDA-I, [165, 182]	^{76}Ge	Ge diodes	17.7 (^{enr}Ge)	21.64	3.2 ± 0.2	~ 0.01	2.1
NEMO-3, [183]	^{100}Mo	tracker + calorimeter	6.9 (^{100}Mo)	34.7	350	0.013	0.11
Present							
EXO-200, [184]	^{136}Xe	LXe TPC	175 (^{enr}Xe)	100	89 ± 3	$(1.7 \pm 0.2) \cdot 10^{-3}$	1.1
KamLAND-Zen, [185, 186]	^{136}Xe	loaded liquid scintillator	348 (^{enr}Xe)	89.5	244 ± 11	~ 0.01	1.9
Future							
CUORE, [187]	^{130}Te	bolometers	741 (TeO_2)	1030	5	0.01	9.5
GERDA-II, [172]	^{76}Ge	Ge diodes	37.8 (^{enr}Ge)	100	3	0.001	15
LUCIFER, [188]	^{82}Se	bolometers	17 (Zn^{82}Se)	18	10	0.001	1.8
MAJORANA D., [189]	^{76}Ge	Ge diodes	44.8 ($^{enr/nat}$Ge)	100 [a]	4	0.003	12
NEXT, [190, 191]	^{136}Xe	Xe TPC	100 (^{enr}Xe)	300	$12.3 - 17.2$	$5 \cdot 10^{-4}$	5
AMoRE, [192]	^{100}Mo	bolometers	200 (Ca^{enr}MoO_4)	295	9	$1 \cdot 10^{-4}$	5
nEXO, [193]	^{136}Xe	LXe TPC	4780 (^{enr}Xe)	12150 [b]	58	$1.7 \cdot 10^{-5}$ [b]	66
PandaX-III, [194]	^{136}Xe	Xe TPC	1000 (^{enr}Xe)	3000 [c]	$12 - 76$	0.001	11 [c]
SNO+, [195]	^{130}Te	loaded liquid scintillator	2340 (^{nat}Te)	3980	270	$2 \cdot 10^{-4}$	9
SuperNEMO, [196, 197]	^{82}Se	tracker + calorimeter	100 (^{82}Se)	500	120	0.01	10

[a] our assumption (corresponding sensitivity from Fig. 14 of Ref. [189]).
[b] we assume 3 tons fiducial volume.
[c] our assumption by rescaling NEXT.

Dell'Oro , Marcocci, Viel, Vissani, 2016