

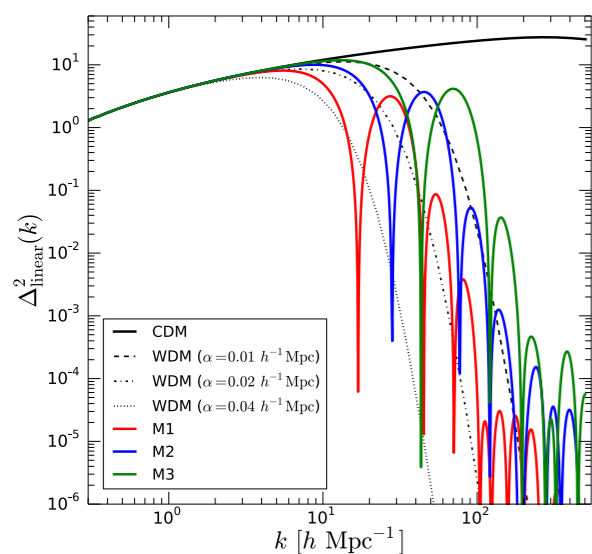
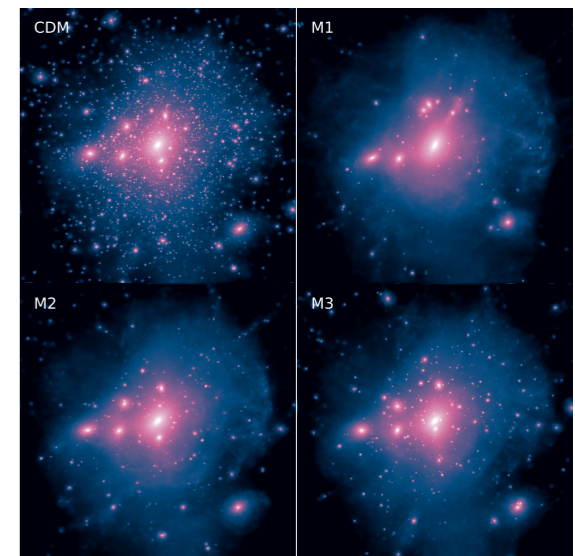
An *E*ffective *T*heory *O*f *S*tructure *F*ormation

or

Structure formation with non-gravitational dark matter interactions

Torsten Bringmann

With Francis-Yan Cyr-Racine, Christoph Pfrommer,
Mark Vogelsberger, Jesús Zavala, ...



‘ETHOS 1’: theory set-up

Cyr-Racine+ [1512.05344]

‘ETHOS 2’: first N-body simulations

Vogelsberger+ [1512.05349]

hydrodynamical
simulations

‘ETHOS 3’: Galaxy abundance & re-ionization

Lovell+ [1711.10497]

‘ETHOS 4’: Ly- α signatures

Bose+ [1811.10630]

‘ETHOS 5’: First star formation

Lovell+ [1812.04627]

Λ CDM cosmology

- Dark matter (DM) is a **crucial ingredient**

- constant** co-moving energy **density**
- only gravitational** interactions
- cold + dissipation-less

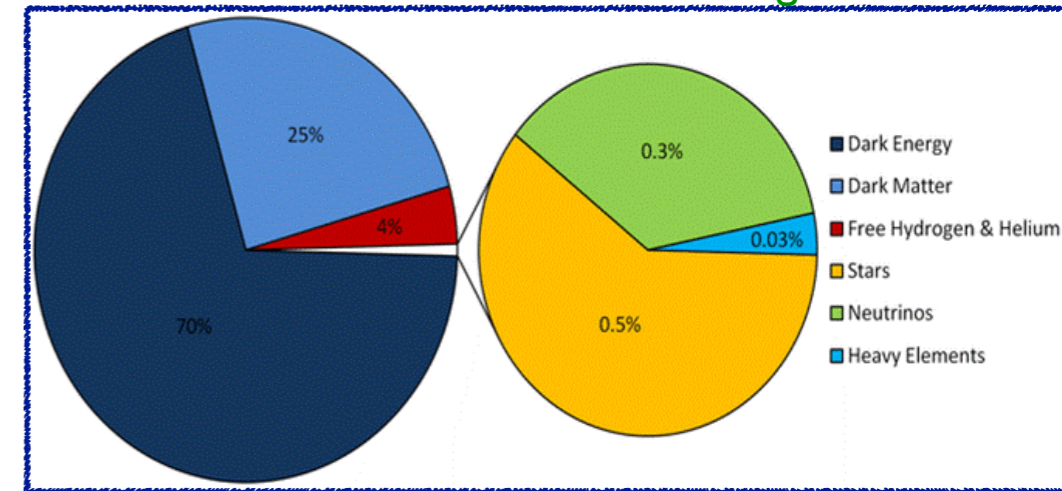


$$\Omega_{\text{CDM}} h^2 = 0.1188 \pm 0.0010$$

Ade+ [Planck Coll.], A&A '16

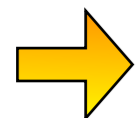
Percent-level measurements of a single parameter!

Image credit: KIAS



- DM **conversion** into (in)visible energy?

- E.g. decays, late-time annihilation, coalescing PBHs, ...



Ω_{CDM} decrease of **up to 10%** possible during matter domination!
(model-independent; much more allowed during RD)

TB, Kahlhoefer, Schmidt-Hoberg & Walia, PRD '18

- Non-gravitational** interactions?

- DM — SM: strong constraints from standard DM searches
- DM self-interactions
- DM — dark radiation interactions

———— **cf. yesterday's talks**

———— **many possible models** \rightsquigarrow **ETHOS**

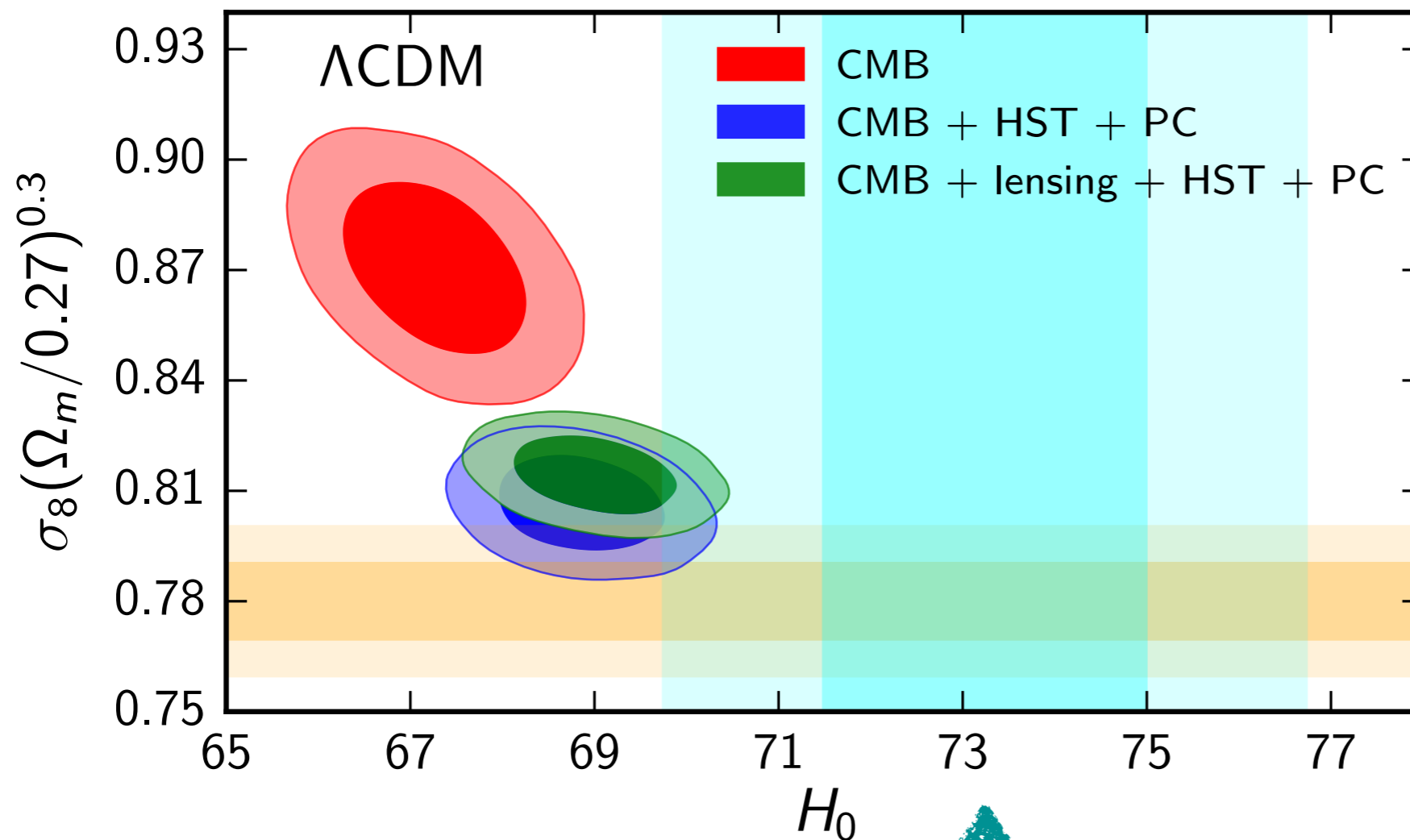
Λ CDM = the perfect success story?

Every party has
a pooper...



- Measurements at high and low redshifts do not seem to *quite* agree

Ade+ [Planck coll.], A&A '14



(Fig from 1803.03644)

normalisation of
power spectrum
(from Planck
clusters)

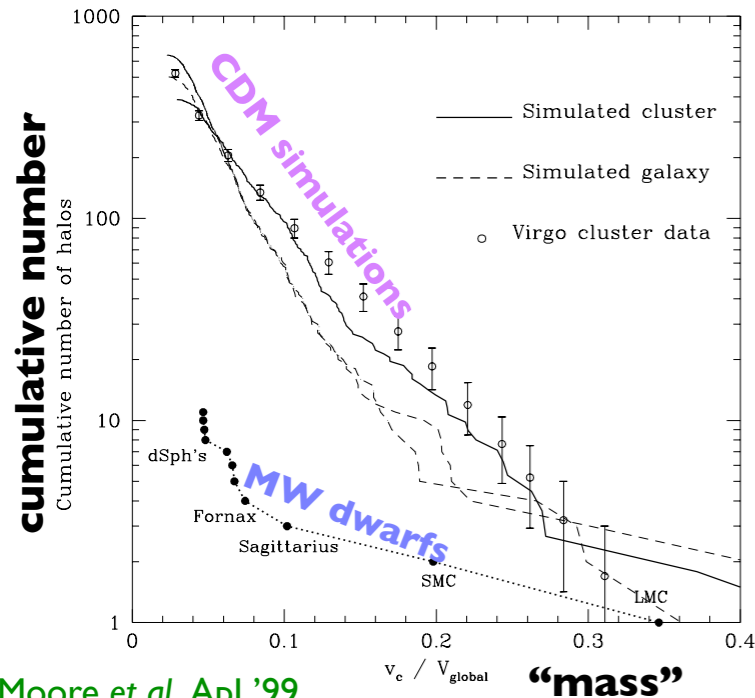
Direct measurements of
Hubble constant (HST)

Health warning:
DES *is* consistent
with CMB ... !

Macaulay+, MNRAS '19
see also Rameez & Sarkar,
1911.06456

Small-scale problems ?

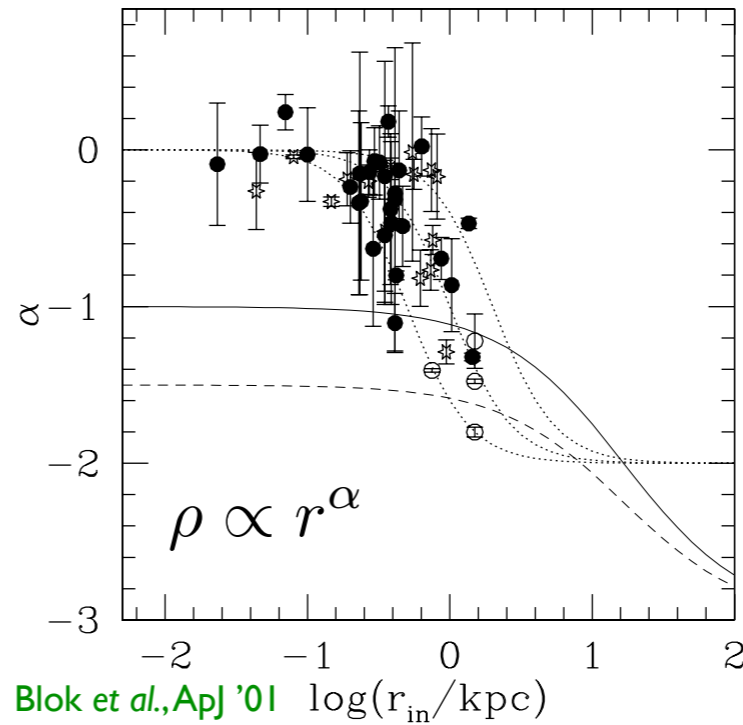
1. Missing satellites?



Moore et al., ApJ '99

More satellites in simulations of MW-like galaxies than observed

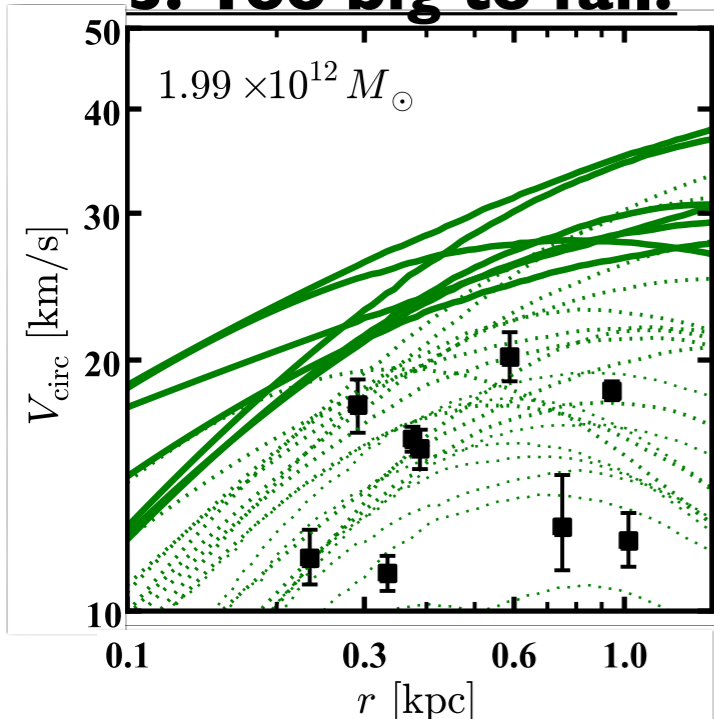
2. Cusps or cores?



Blok et al., ApJ '01

Cuspy inner density profiles predicted by simulations not found in (all) observations

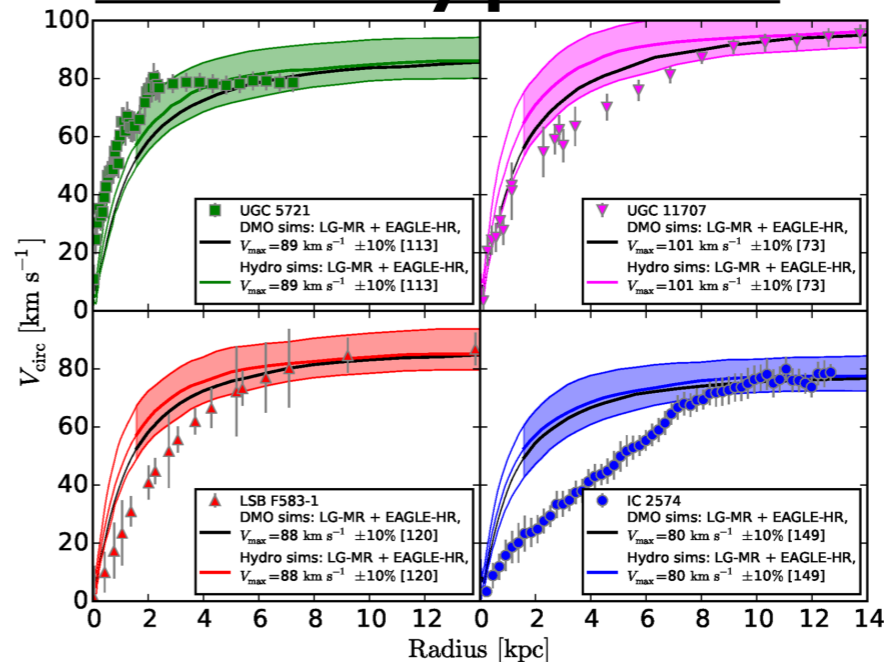
3. Too big to fail?



Most massive simulated sub-halos too dense to form observed brightest dwarf galaxies

Boylan-Kolchin, Bullock & Kaplinghat, '11

4. Diversity problem?



Oman+, MNRAS '15

Real rotation curves vary more than in simulations with(!) baryons



Disclaimer

***ETHOS** does **not**
primarily attempt to
address these issues!*

(Nor claims that this would be necessary)

*But being able to do so serves as possible
proof-of-principle that relevant **observables**
connected to non-gravitational,
'dark' interactions can be identified...*

Generic dark sector models

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

Standard Model

- SM particles

- A 'portal' typically still ensures thermalisation at high temperatures

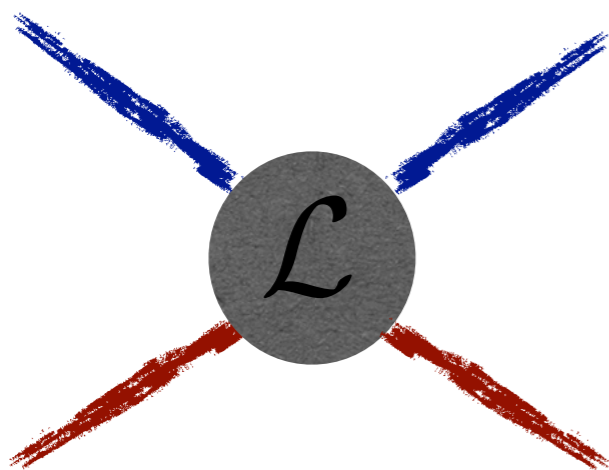
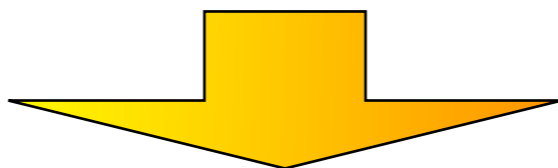
- Separate entropy conservation after decoupling $\rightsquigarrow T_{\text{photon}} \neq T_{\text{dark}}$

$$\text{e.g. } U(1)_X \times \dots$$

Dark Sector

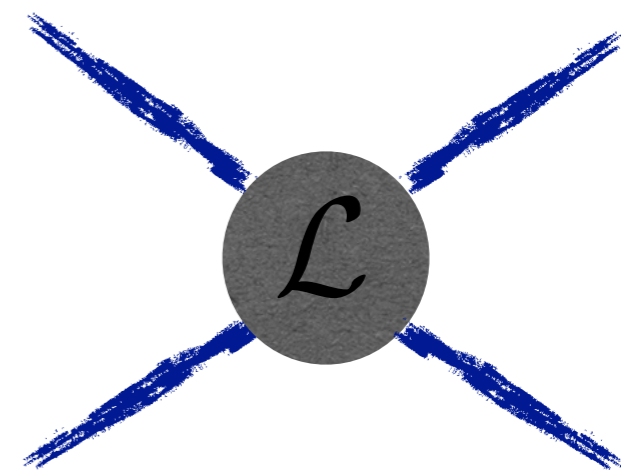
- Dark matter
- Dark radiation ('sterile neutrinos', 'dark photons', ...)

.....
e.g. $\mathcal{L}_{\text{Higgs}} \supset \kappa |\phi|^2 |\Theta|^2$



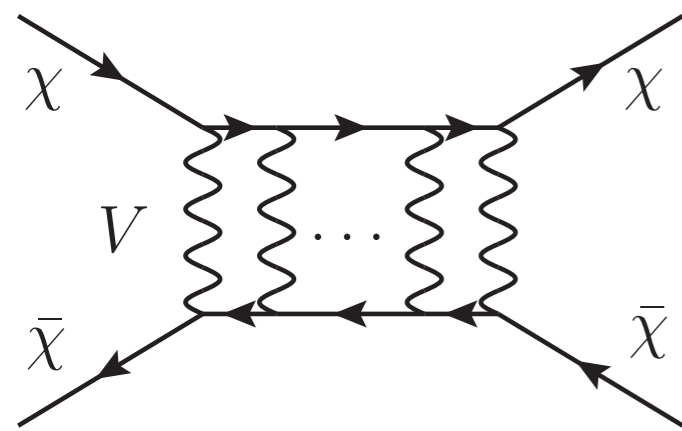
- imprints on linear $\mathcal{P}(k)$

need to treat consistently!



- imprints on inner (sub-)halo structure

From theory to observations



particle model

input:
masses, spins,
coupling constants



cosmological
simulations

input:
consistent initial
conditions, non-
gravitational forces
between "particles"



astrophysical
observables

input
(for interpretation of data):
**output from
simulations**

- The first task can be **demanding**, the second in addition computationally very expensive
- But expect large degeneracies, so **very inefficient**...
- **Idea of ETHOS**: identify **effective parameters** and provide **maps** for each of those steps (\rightsquigarrow no need to re-compute each model!)

Cyr-Racine+, PRD'16; Vogelsberger+, MNRAS '16

Linear perturbations - setup

- Fundamentally, have to solve coupled Boltzmann equations:

$$\frac{df_\chi}{d\lambda} = C_{\chi\tilde{\gamma}\leftrightarrow\chi\tilde{\gamma}}[f_\chi, f_{\text{DR}}], \quad \frac{df_{\text{DR}}}{d\lambda} = C_{\chi\tilde{\gamma}\leftrightarrow\chi\tilde{\gamma}}[f_{\text{DR}}, f_\chi] + C_{\tilde{\gamma}\tilde{\gamma}\leftrightarrow\tilde{\gamma}\tilde{\gamma}}[f_{\text{DR}}]$$

- rewrite as differential equations for DM **density**, **velocity** and **'temperature'**:

$$n_\chi \equiv \eta_\chi \int \frac{d^3p}{(2\pi)^3} f_\chi(\mathbf{p}) \quad \vec{v}_\chi \equiv \frac{\eta_\chi}{n_\chi} \int \frac{d^3p}{(2\pi)^3} f_\chi(\mathbf{p}) \frac{p \hat{\mathbf{p}}}{E} \quad T_\chi \equiv \frac{\eta_\chi}{3n_\chi^{(0)}} \int \frac{d^3p}{(2\pi)^3} \frac{\mathbf{p}^2}{m_\chi} f_\chi^{(0)}(p)$$

- keep terms up to first order in perturbations

- Take advantage of various simplifications

- Neglect (subdominant) DR-DR iterations

- Assume DR close to EQ: $f_{\text{DR}}(\mathbf{x}, \mathbf{q}, \tau) = f_{\text{DR}}^{(0)}(q, \tau)[1 + \Theta_{\text{DR}}(\mathbf{x}, \mathbf{q}, \tau)]$

- Momentum transfer in DM-DR scatterings must be small!

- Derive hierarchy of Boltzmann moments

- Expand in Legendre polynomials: $\Theta_{\text{DR}}(k, \hat{q}, q, \tau) = \sum_{l=0}^{\infty} (-i)^l (2l+1) F_l(k, q, \tau) P_l(\mu)$

$$\left(\frac{1}{\eta_\chi \eta_{\text{DR}}} \sum_{\text{states}} |\mathcal{M}|^2 \right) \Big|_{\substack{t=2p_1^2(\tilde{\mu}-1) \\ s=m_\chi^2+2p_1 m_\chi}} = \sum_{n=0}^{\infty} (2n+1) A_n(p_1) P_n(\tilde{\mu})$$

$$\mu \equiv \hat{q} \cdot \hat{k}$$

- Integrate BEs on both sides with $\frac{1}{2(-i)^l} \int_{-1}^1 d\mu P_l(\mu)$

Linear perturbations - summary

Summary (1st order equations):

DM:

$$\begin{aligned} \dot{\delta}_\chi + \theta_\chi - 3\dot{\phi} &= 0, \\ \dot{\theta}_\chi - c_\chi^2 k^2 \delta_\chi + \mathcal{H}\theta_\chi - k^2\psi &= \dot{\kappa}_\chi [\theta_\chi - \theta_{\text{DR}}] \end{aligned}$$

Only way that
(non-gravitational)
DM physics enters! **'ETHOS'**

DR:

$$\dot{\Pi}_{\text{DR},l} + \frac{k}{2l+1} ((l+1)\Pi_{\text{DR},l+1} - l\Pi_{\text{DR},l-1}) = (\alpha_l \dot{\kappa}_{\text{DR-DM}} + \beta_l \dot{\kappa}_{\text{DR-DR}}) \Pi_{\text{DR},l}$$

$\sim \kappa_\chi$ ~ 0

$$\delta \equiv (\rho - \bar{\rho}) / \bar{\rho}$$

$$\theta_j \equiv i \mathbf{k} \cdot \mathbf{v}_j$$

ϕ, ψ : metric perturb.
in conformal
Newtonian gauge

$$\Pi_{\text{DR},l}(k, \tau) = \frac{\int dq q^3 f_{\text{DR}}^{(0)}(q, \tau) F_l(k, q, \tau)}{\int dq q^3 f_{\text{DR}}^{(0)}(q, \tau)}$$

$$\delta \sim \Pi_0 \quad \theta \sim \Pi_1$$

Step-by-step procedure:

calculate κ, α_l from $A_l(p) = \frac{1}{2} \int_{-1}^1 d\tilde{\mu} P_l(\tilde{\mu}) \left(\frac{1}{\eta_\chi \eta_{\text{DR}}} \sum_{\text{states}} |\mathcal{M}|^2 \right) \Big|_{\substack{t=2p^2(\tilde{\mu}-1) \\ s=m_\chi^2+2pm_\chi}}$

Details: Cyr-Racine+, PRD '16

calculate $c_\chi^2 = \frac{T_\chi}{m_\chi} \left(1 - \frac{\dot{T}_\chi}{3\mathcal{H}T_\chi} \right)$ from $\frac{dT_\chi}{d\tau} = -2\mathcal{H}T_\chi + \Gamma_{\text{heat}}(T_{\text{DR}}) (T_{\text{DR}} - T_\chi)$

Details: TB, NJP '09,
TB+, PRD '16

aka **'momentum exchange rate'** Υ



Linear perturbations - results

- ETHOS comes with a dedicated **Boltzmann solver**:

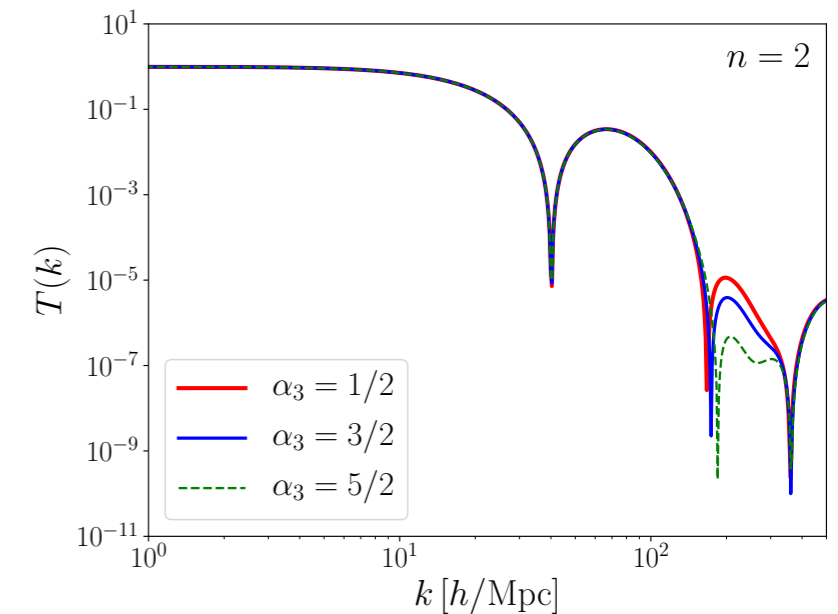
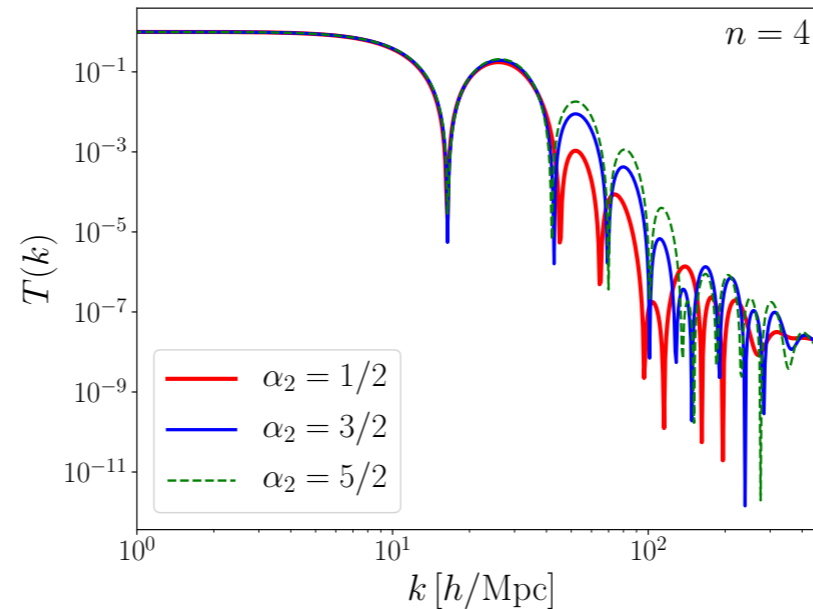
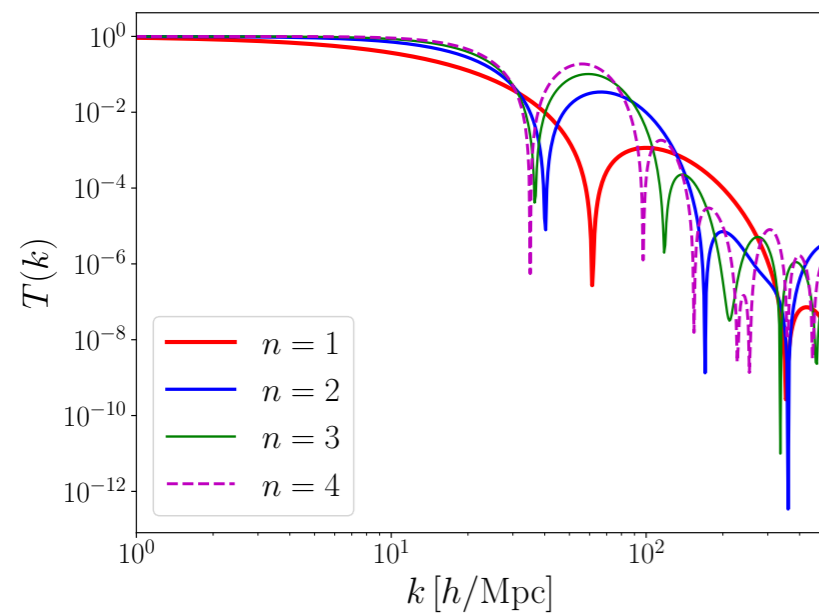
https://bitbucket.org/franyancr/ethos_camb

- modified version of CAMB Lewis & Bridle, PRD '02
- Actual implementation based on phenomenological power-law ansatzes:

$$\mathcal{K} \sim \sum_n a_n \left(\frac{1+z}{1+z_D} \right)^n \quad \Gamma_{\text{heat}} \sim \sum_n d_n \frac{(1+z)^{n+1}}{(1+z_D)^n} \quad \text{etc}$$

- detailed examples for calculating a_n, d_n from given model ($|\mathcal{M}|^2$) Cyr-Racine+, PRD '16
TB+, PRD '16

Example spectra:



➔ **‘Dark acoustic oscillations’**

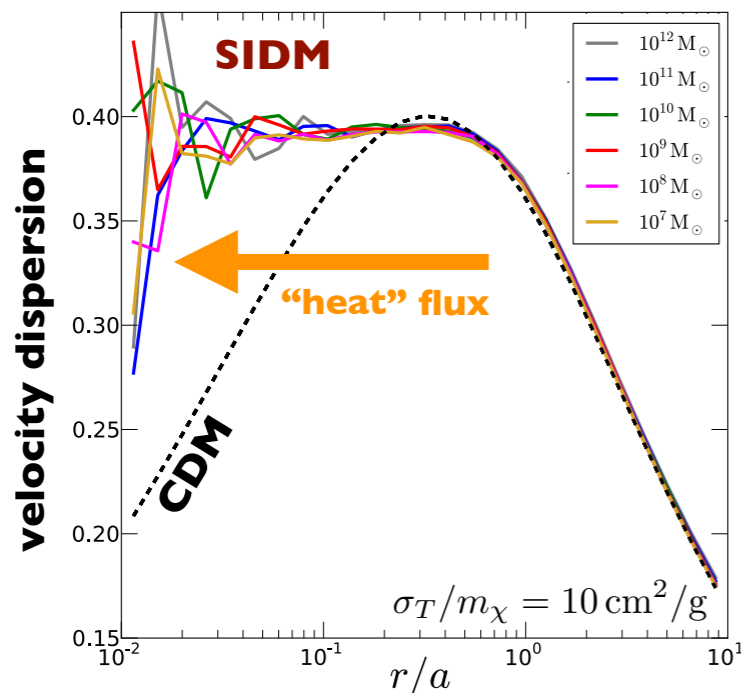
(Physics very similar to CMB photons scattering on electrons around decoupling!)

Self-interacting DM (SIDM)

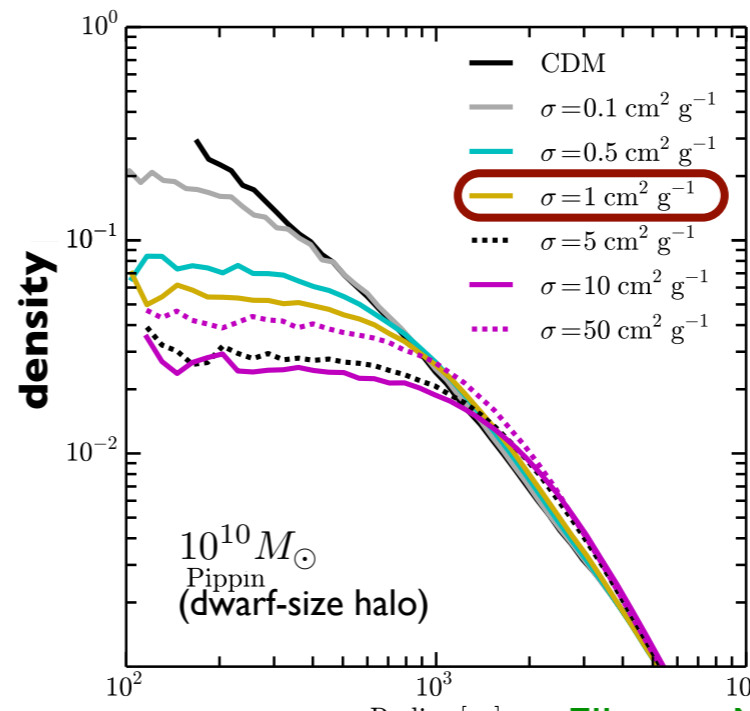
- DM-DM scatterings [Spergel & Steinhardt, PRL '99](#)

- do not affect linear perturbations (number densities!)
- but isotropise DM distribution in inner parts of halo:

→ core formation once $\mathcal{O}(1)$ scatters per dynamical time



[Vogelsberger, Zavala & Loeb, MNRAS '12](#)




roughly needed for cusp/core

[Elbert+, MNRAS '15](#)

- Simple analytic models to predict core radius from σ_{SIDM}

- reproduce CDM simulation results for $\rho_\chi(r)$ remarkably well [Kaplinghat, Tulin & Yu, PRL '15](#)
- but underlying (microphysics) assumptions not really satisfied [Sokolenko+, JCAP '18](#)

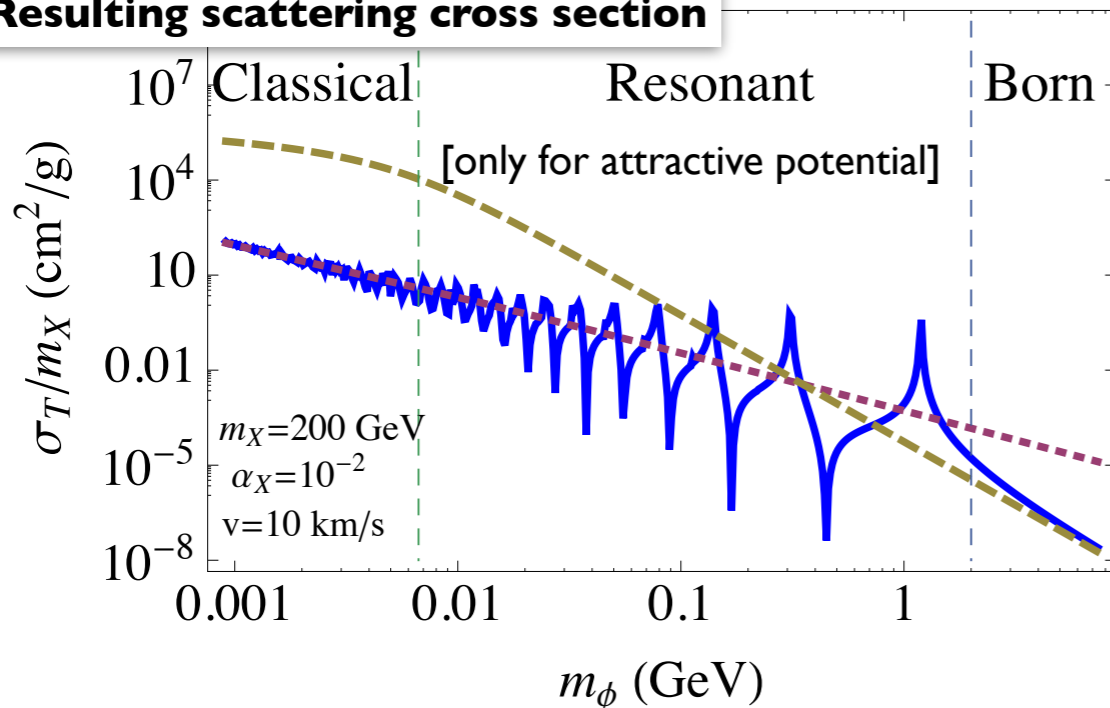
→ Use caution when applied to systems including baryons! 

Velocity dependence

- Massive mediators induce a **Yukawa potential** between DM particles.

$$\left(-\frac{\nabla^2}{m_\chi} + V \right) \psi(r) = m_\chi v^2 \psi(r)$$

Resulting scattering cross section



$$\sigma_T \equiv \int d\Omega (1 - \cos \theta) \frac{d\sigma}{d\Omega}$$

see e.g. Tulin, Yu & Zurek, PRD '13

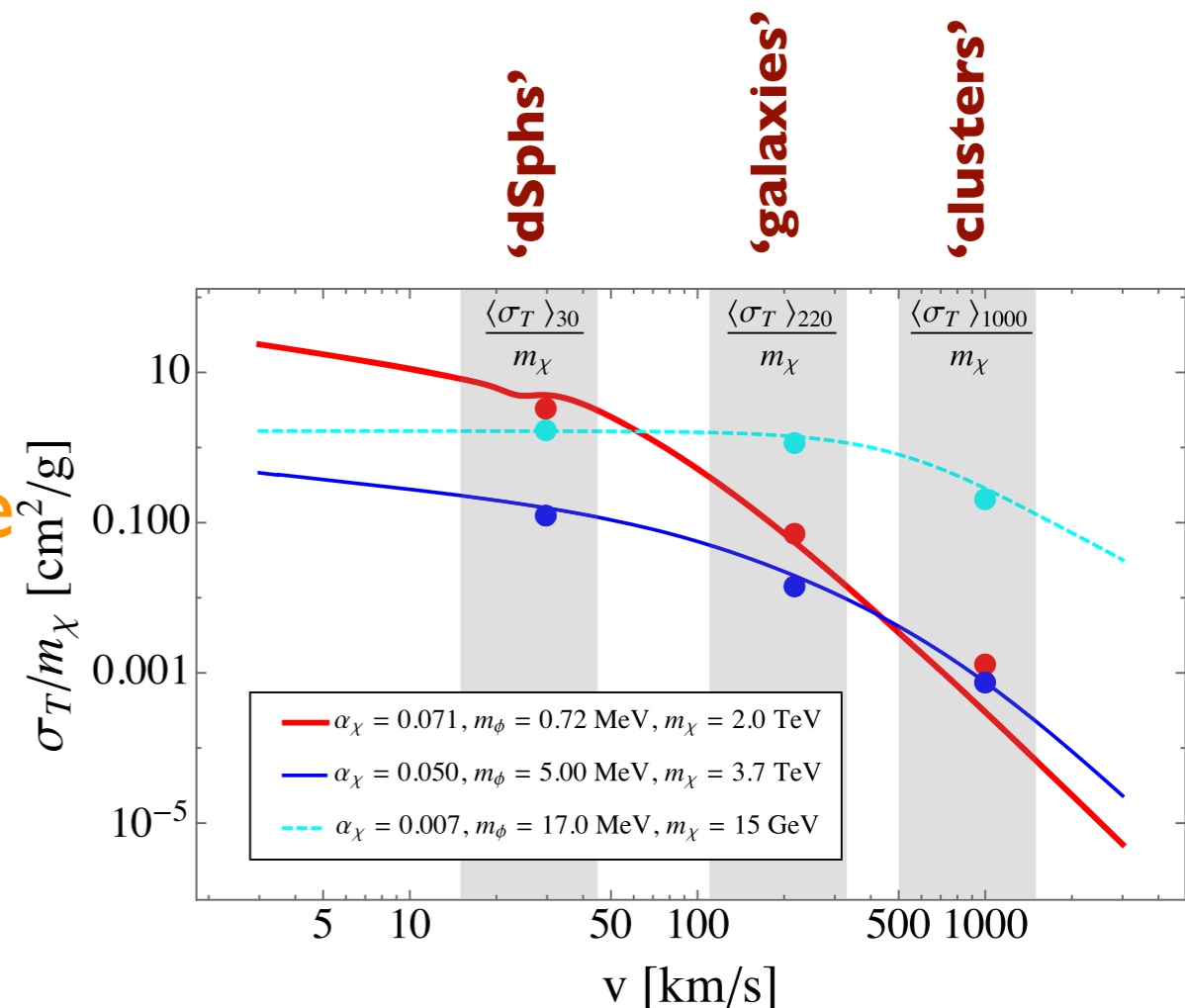
- Phenomenologically important: **characteristic velocity dependence**

[not only for Yukawa potentials!]

Chu, Garcia-Cely & Murayama, PRL '19]

- ETHOS** implementation:

- sufficient to take **average values** for 3 characteristic velocities



Implementation

- Translate power spectrum to **initial particle distribution**

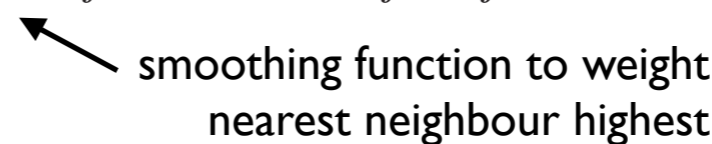
- use MUSIC code [Hahn & Abel, MNRAS '11](#) [see also Dolag+, '08]

- Probabilistic method** to account for elastic scattering

[Vogelsberger, Zavala & Loeb, MNRAS '12](#)

- isotropic scattering of macroscopic 'particles' with mass m_i

$$P_{ij} = \frac{m_i}{m_\chi} W(r_{ij}, h_i) \sigma_T(v_{ij}) v_{ij} \Delta t_i$$

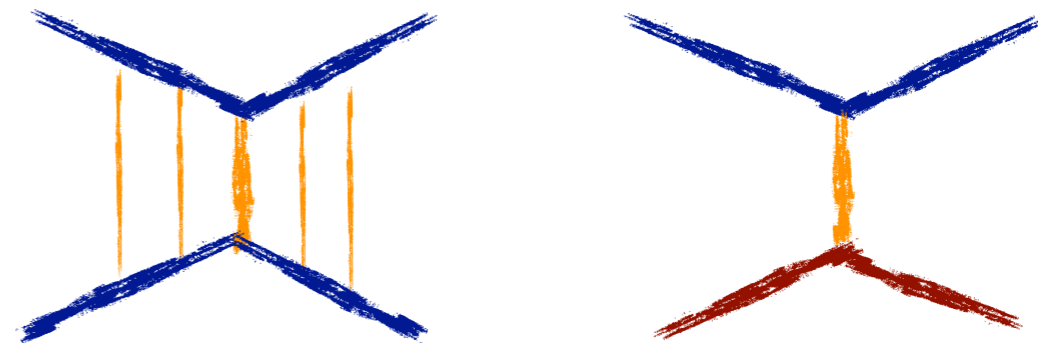

 smoothing function to weight nearest neighbour highest

- Cosmological simulation with $m_i \sim 10^8 M_\odot$ ($\epsilon \sim 3$ kpc), **zoom-in** of MW-like halos down to $m_i \sim 3 \times 10^4 M_\odot$ ($\epsilon \sim 70$ pc)

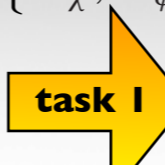
- First ETHOS example:**

- TeV-scale **DM** particle
 - MeV-scale vector **mediator**
 - massless** (sterile) neutrino-like fermion

[van den Aarssen, TB & Pfrommer, PRL '12](#)
[TB, Hasenkamp & Kersten, JCAP '14](#)

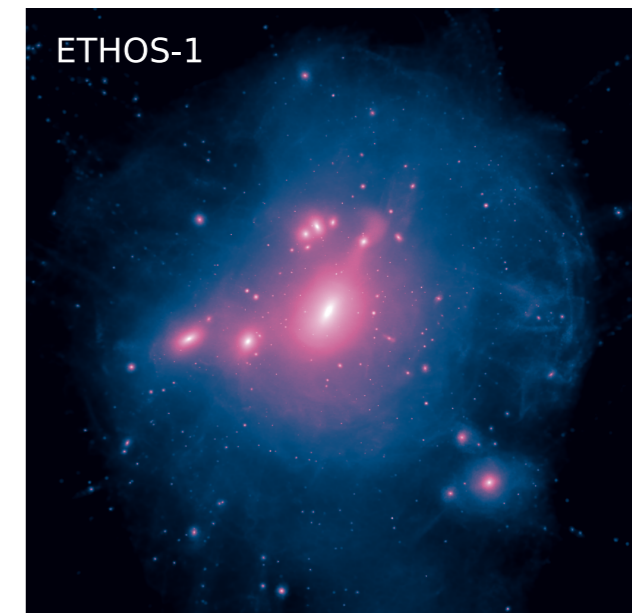
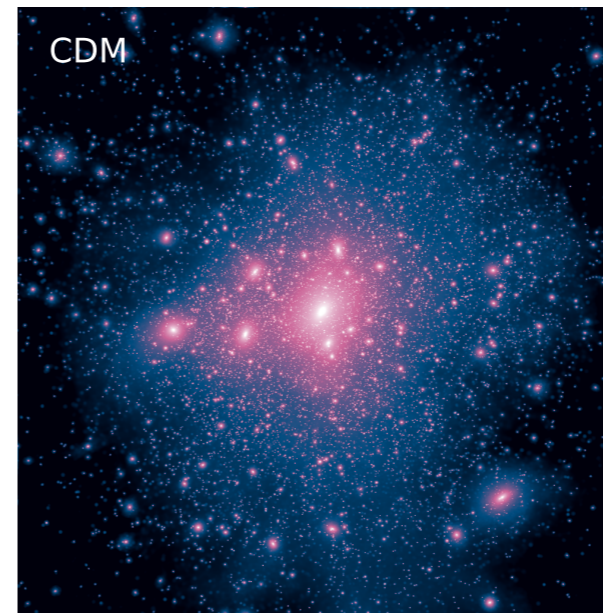
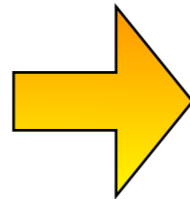
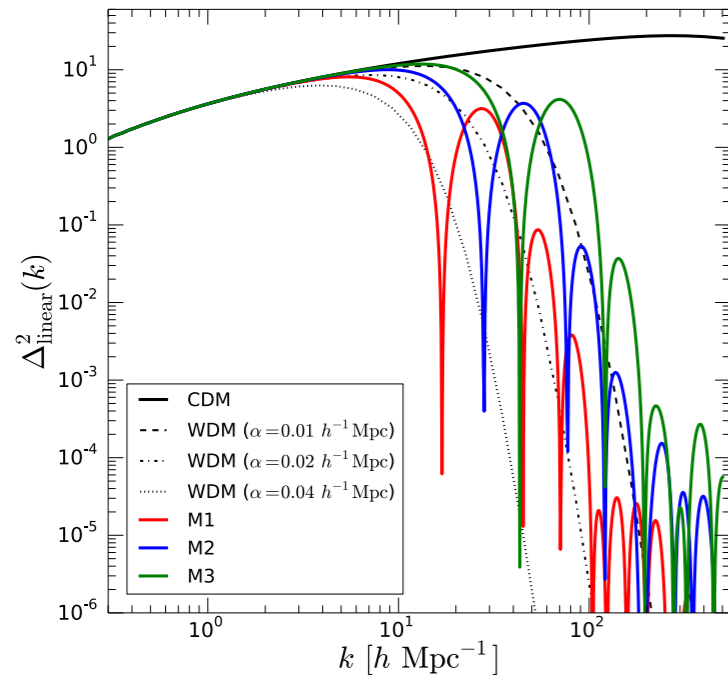


$\{m_\chi, m_\phi, g_\chi, g_\nu, \eta_\chi, \eta_\nu, \xi\}$


 $\left\{ \omega_{\text{DR}}, a_4, \alpha_{l \geq 2} = \frac{3}{2}, \frac{\langle \sigma_T \rangle_{30}}{m_\chi}, \frac{\langle \sigma_T \rangle_{220}}{m_\chi}, \frac{\langle \sigma_T \rangle_{1000}}{m_\chi} \right\}$

Late kinetic decoupling

- Select four benchmarks: Vogelsberger+, MNRAS'16



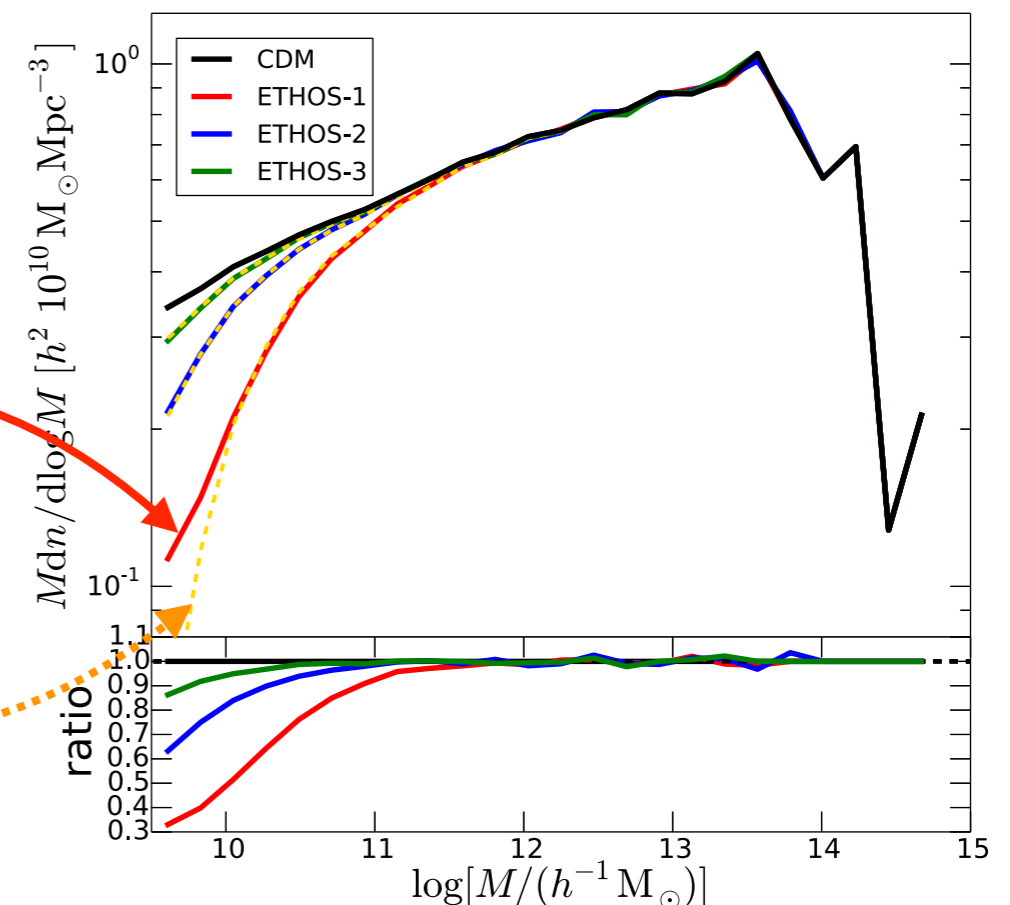
- Almost identical suppression of halo mass function as for WDM cosmology:

$$M_{\text{cut,kd}} = 5 \cdot 10^{10} \left(\frac{T_{\text{kd}}}{100 \text{ eV}} \right)^{-3} h^{-1} M_{\odot}$$

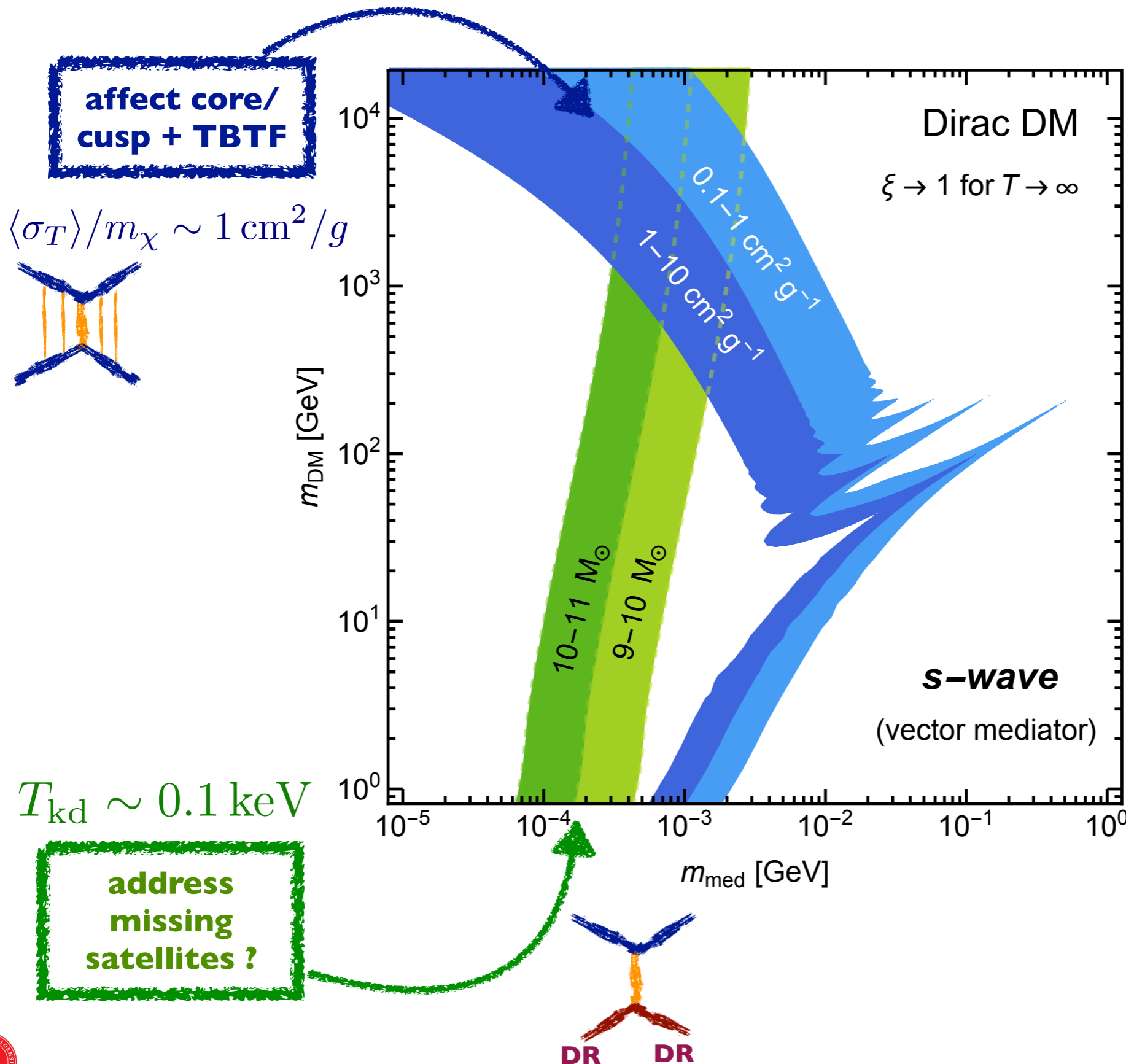
[solid lines; NB: up to factor ~2 same as analytic estimate!]

$$M_{\text{cut,WDM}} = 10^{11} \left(\frac{m_{\text{WDM}}}{\text{keV}} \right)^{-4} h^{-1} M_{\odot}$$

[dashed lines; would-be result from WDM free-streaming]



Full parameter scan



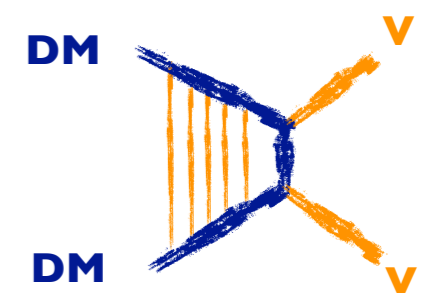
TB, Edsjö,
 Gondolo, Ullio
 & Bergström,
 JCAP '18

● **NEW** since v6.1:

- SIDM
- Sommerfeld
- handle varying

$$\xi \equiv T_{\text{dark}} / T_{\text{photon}}$$

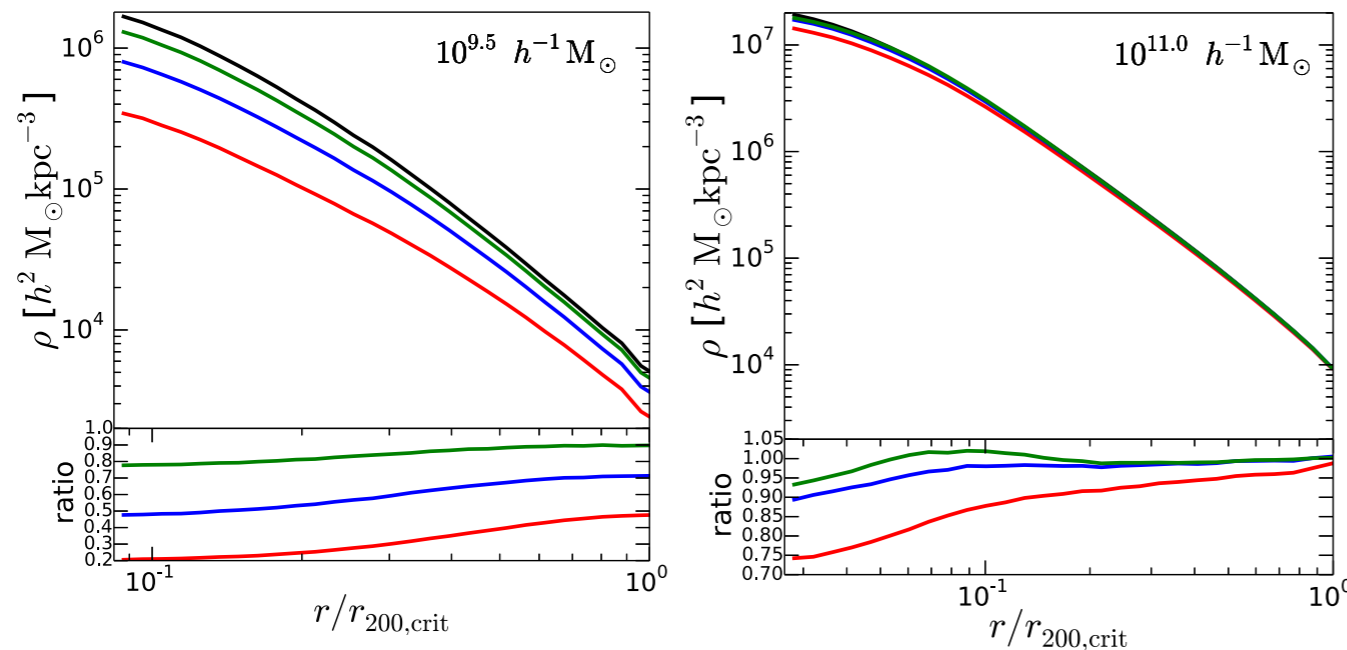
➔ coupling fixed by thermal relic density



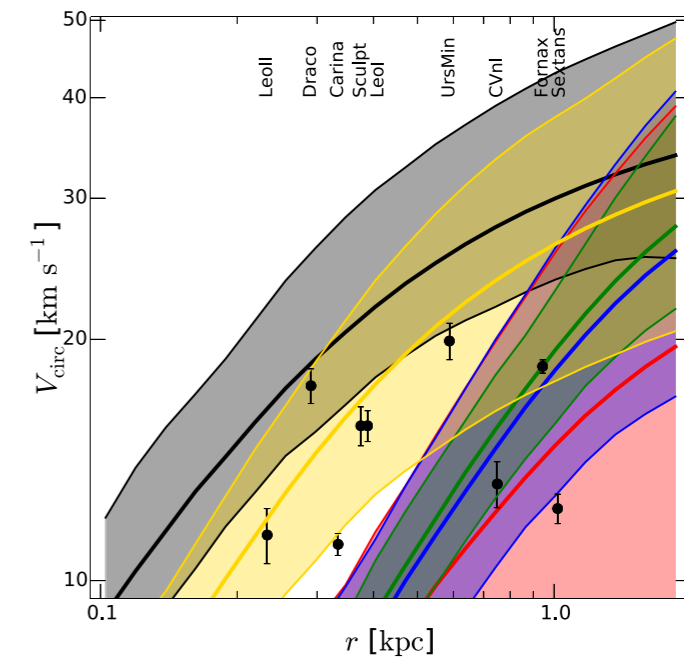
Inner halo structure

Vogelsberger+, MNRAS'16

- Closer look: can indeed address CDM abundance and structural 'problems' **simultaneously**, in a **consistent** particle framework:



central (sub)halo densities reduced
(→ **core/cusp**)



most massive subhalos less dense
(→ **too-big-to-fail**)

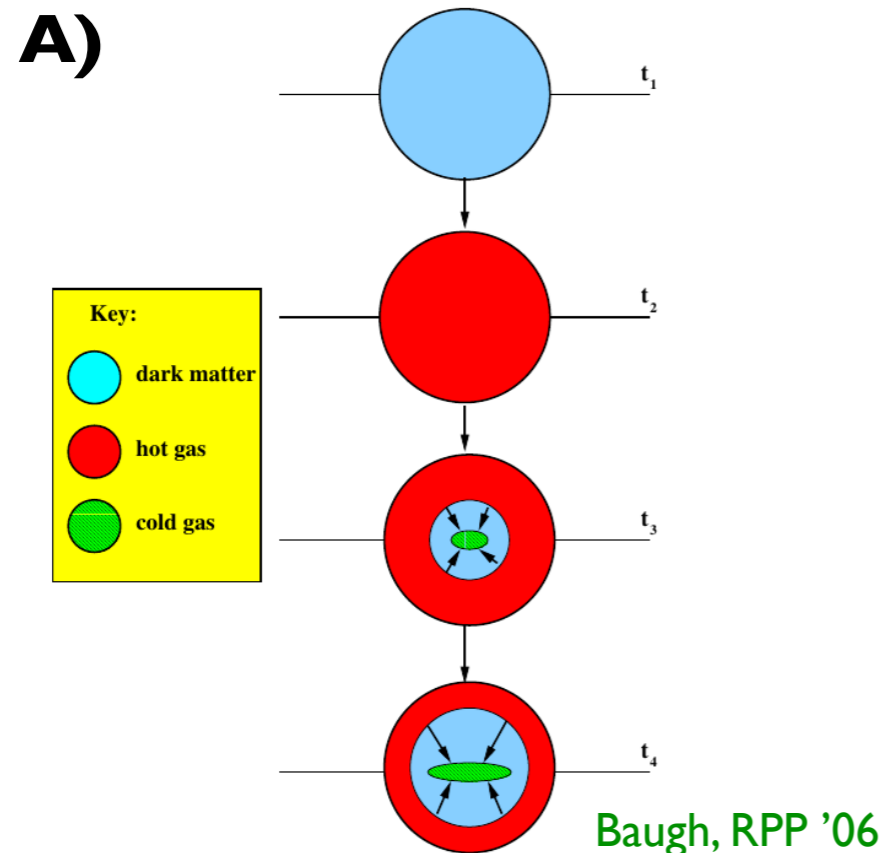
- NB: **Non-trivial interplay** between modified power spectrum and self-interactions

➔ *Details more complicated than the usual 'need $\sim 1 \text{ cm}^2/\text{g}$ '!*

- Also, this is still without **baryonic physics**...
[though dSphs highly DM dominated]

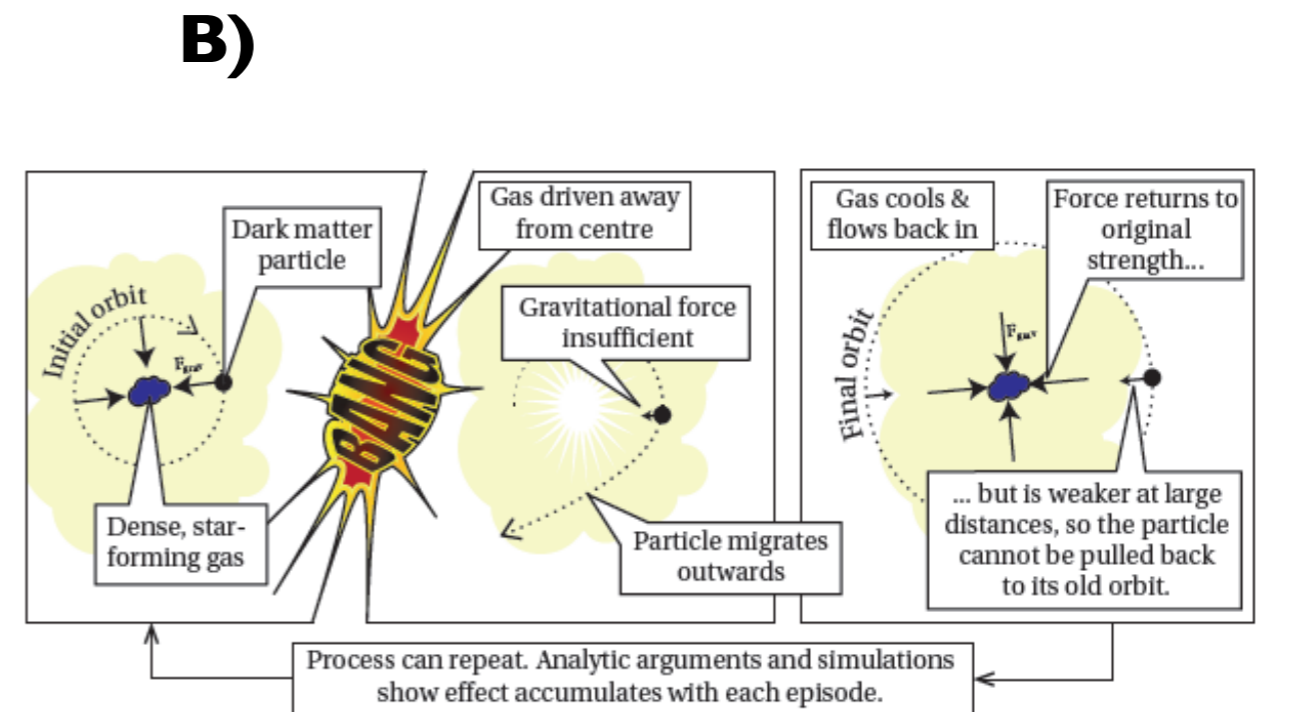
Adding Baryons

- Simplest picture: two competing effects



Adiabatic contraction due to disk assembly

↑ *increase*
of inner DM density



Pontzen & Governato, Nature '14

gas and DM heating due to supernova feedback

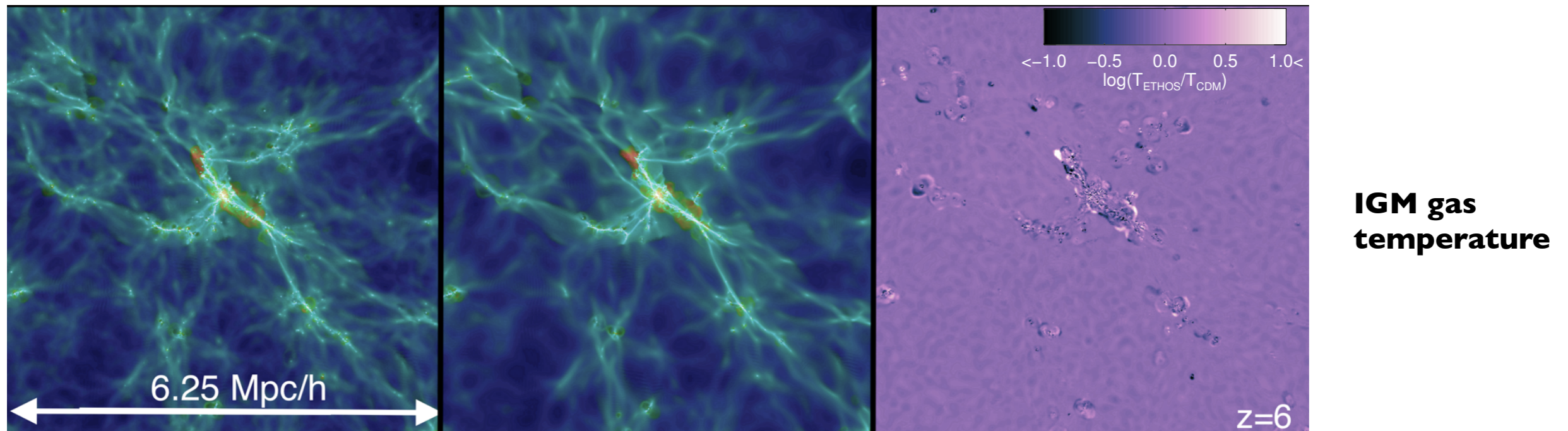
↓ *decrease*
of inner DM density

- SIDM + A)** may lead to core collapse Elbert+, ApJ '18

⇒ A way to address the *diversity* problem? Creasey+, MNRAS '17
Kamada+, PRL '17

Re-ionization history

- Lack of small halos should delay onset of structure formation



Lovell+, MNRAS '18

- Hydrodynamical simulations: Indeed — but effect on **reionisation history** is surprisingly small [similar to **WDM!**]
 - Suppression of high- z , low-mass galaxies: maybe visible with JWST
 - **Brighter starbursts** in these galaxies compensate effect on optical depth
- Follow-up: **halo collapse comparison** on individual basis
 - Virial masses of ETHOS halos are suppressed, but not stellar mass
 - Promising way to test/constrain ETHOS: large populations of very old stars ($z > 17$)

Lovell, Vogelsberger & Zavala, MNRAS '19

Imprint on Lyman alpha spectra ?

Bose+, MNRAS '19

- Need strong features in linear $\Delta^2(k)$ to survive in non-linear regime

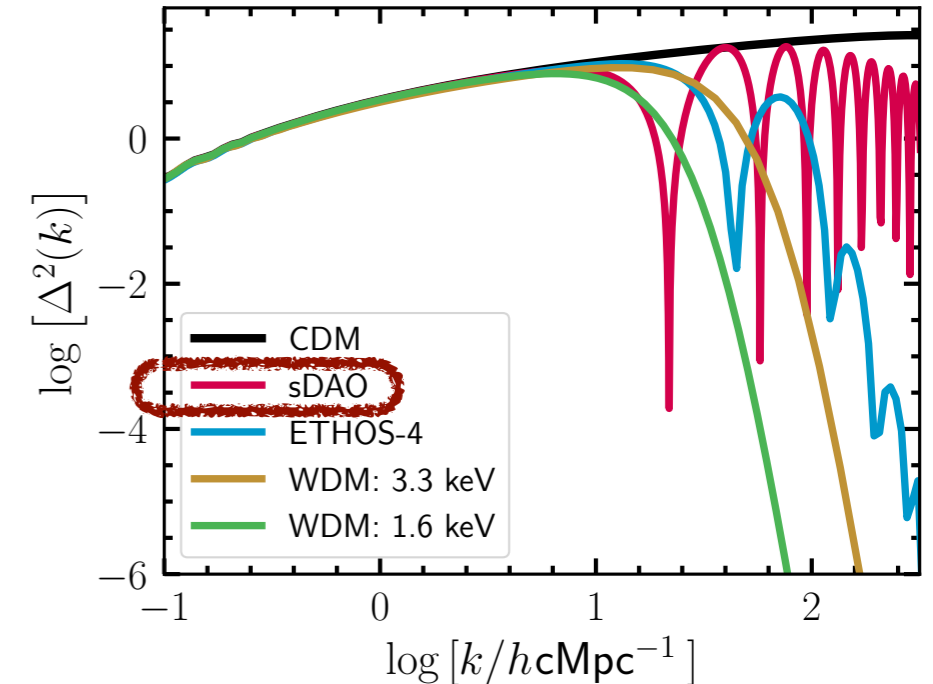
- use **atomic DM** benchmark (**sDAO**)

Kaplan+, JCAP '10

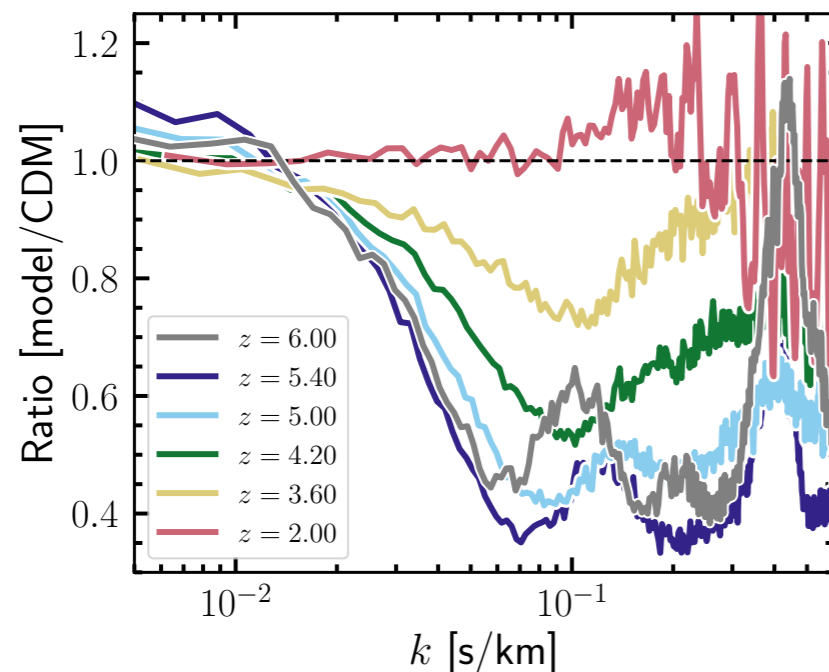
Cyr-Racine & Sigurdson, PRD '13

- galaxy formation model as in IllustrisTNG

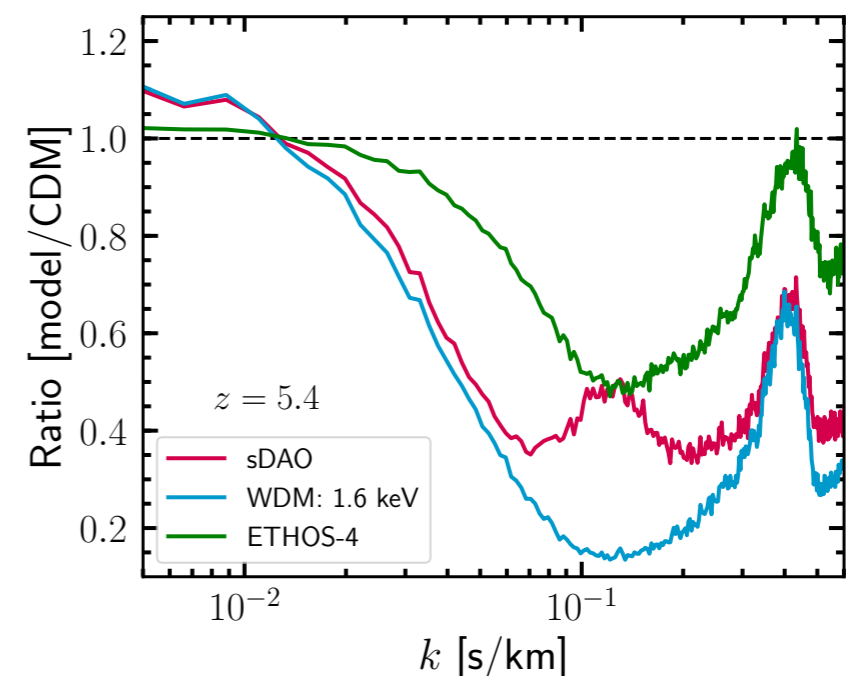
Marinacci+, MNRAS '18; ...



- DAO bump** visible in 1D Ly- α flux spectra only for $z \gtrsim 5$



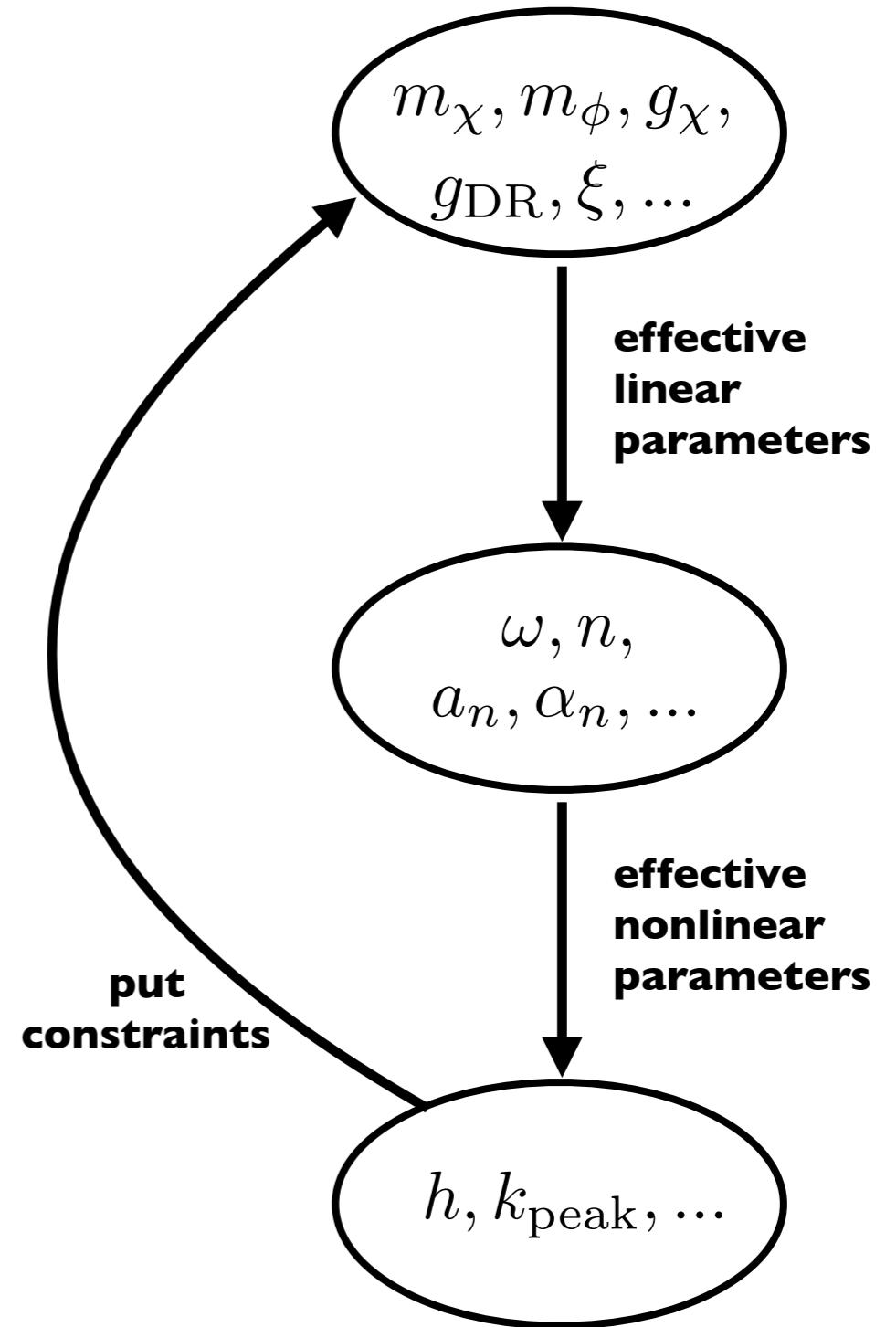
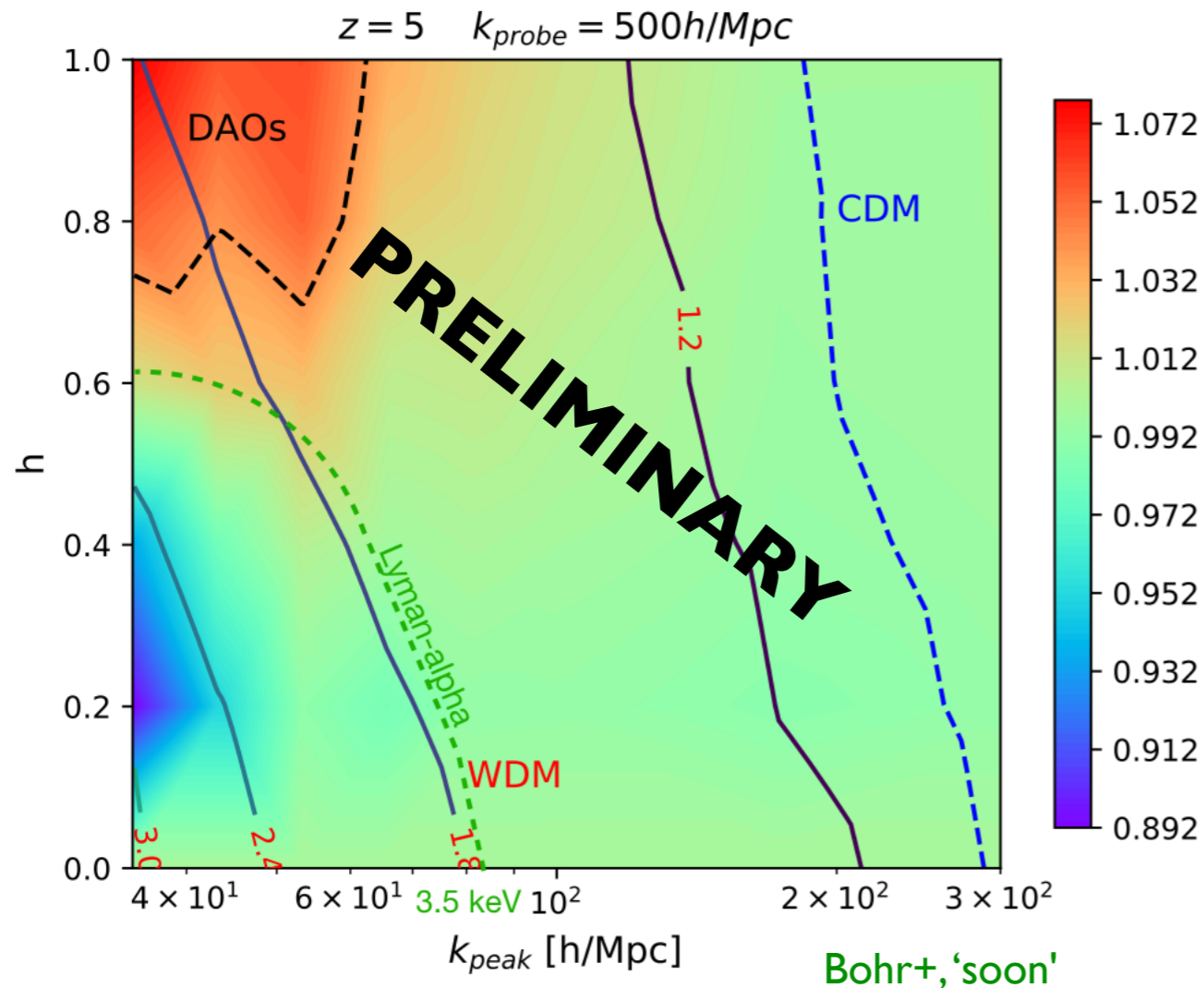
NB: bump at ~ 0.4 s/km set by numerical resolution!



➔ *In principle, this allows to disentangle WDM from sDAO!*

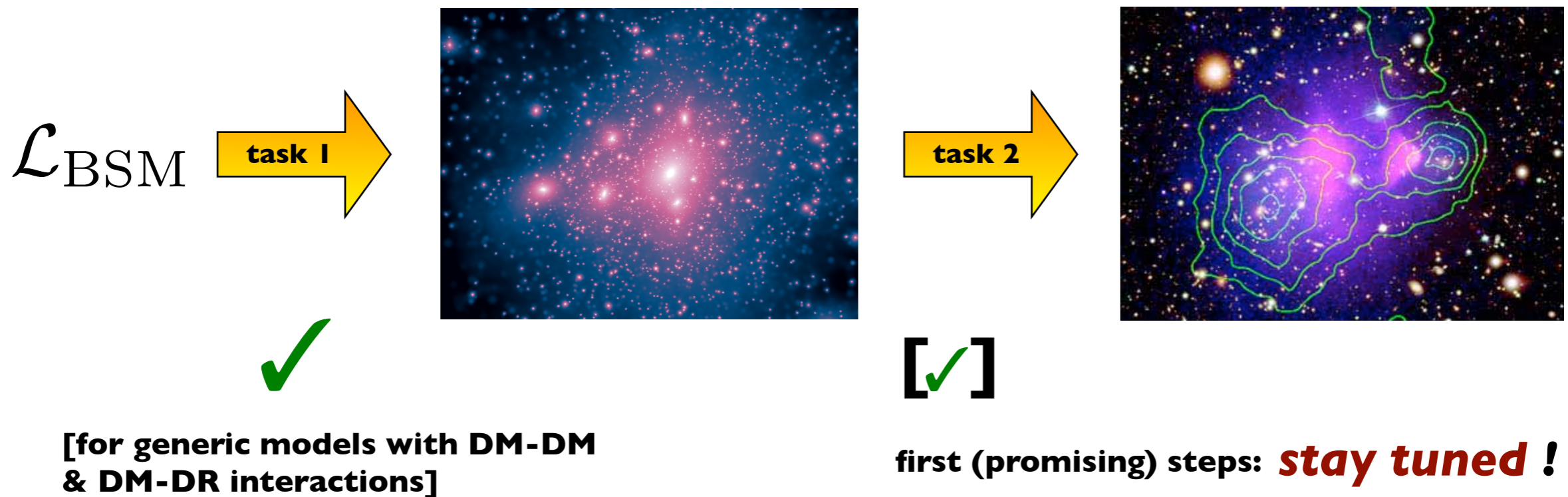
Outlook ?

- Goal: a fast and **automated map** instead of running expensive simulations!



Conclusions

- Cosmological observations are a fascinating, and unique, tool to **test** 'invisible' **dark sector** interactions
- Goal of **ETHOS**: provide a **consistent framework** for this



Thanks for your attention!