

The Migdal effect and bremsstrahlung in effective field theories of dark matter scattering



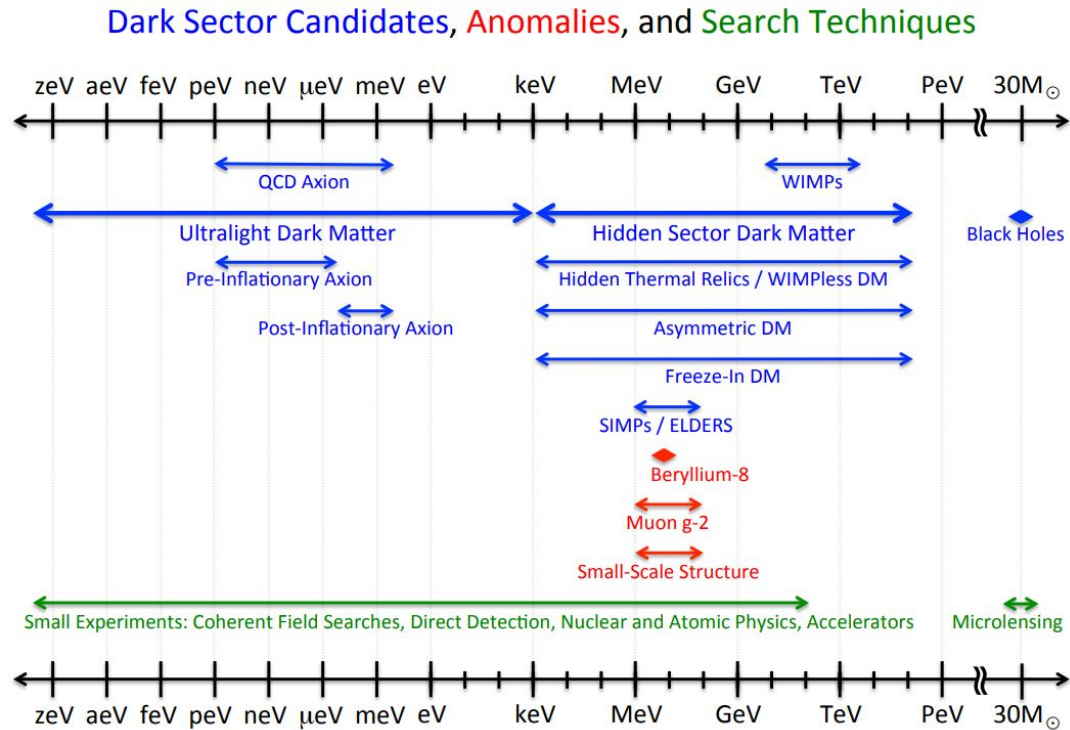
TeVPA 2019
Sydney

Jayden Newstead
University of Melbourne
Purdue University
Arizona State University

arXiv:1905.00046
N.F.Bell, J.B.Dent,
S.Sabharwal, T.Weiler

Overview

There has been a concerted effort to extend our WIMP horizons, notably in the direction of lighter masses.

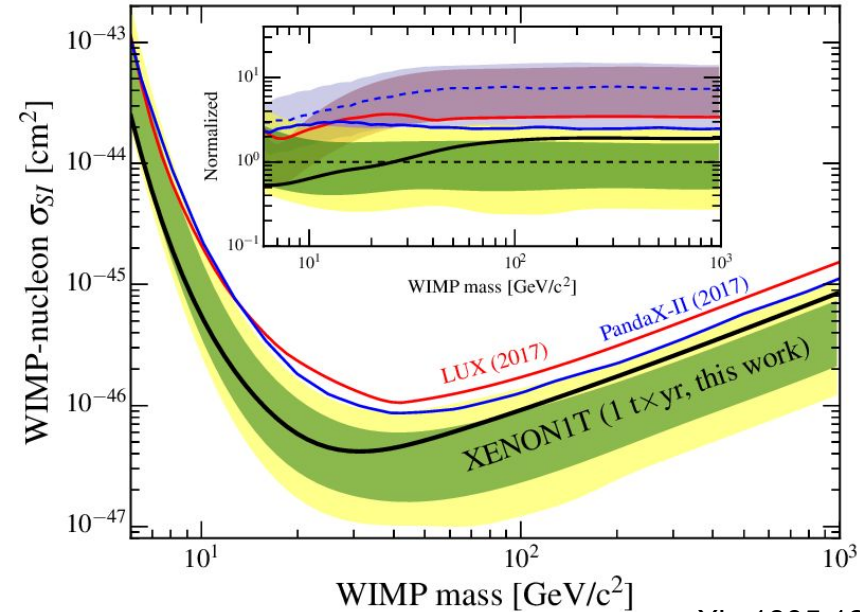
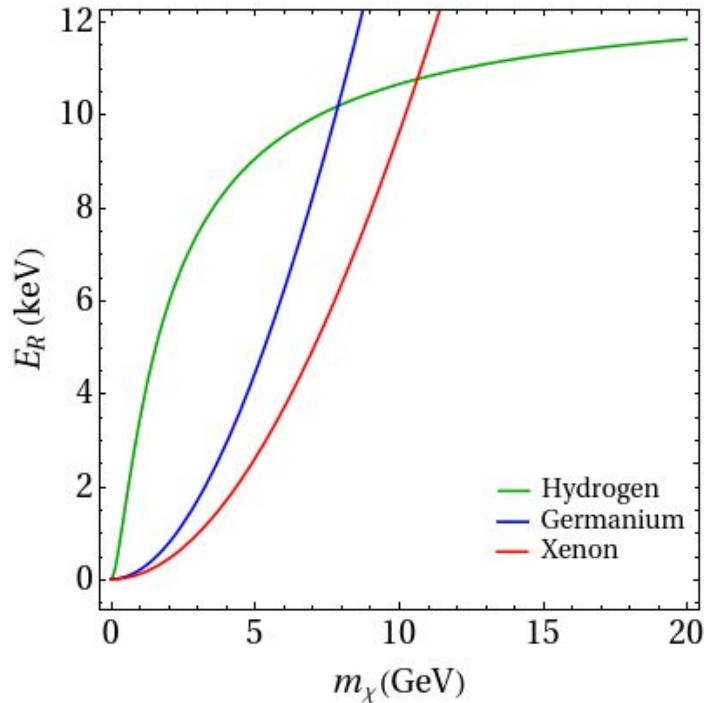


From US Cosmic Visions
arXiv:1707.04591

The kinematic problem

- Light dark matter does not pack much of a punch:

$$E_{R_{\max}} = \frac{2\mu_T^2}{m_T} v_{\max}^2$$



arXiv:1805.12562

Methods for probing lighter dark matter

(an incomplete list)

Existing direct detection:

- Cosmic ray dark matter (arXiv:1810.10543)
- The Migdal effect (arXiv:1707.07258)
- Bremsstrahlung (arXiv:1607.01789)
- Electron scattering (arXiv:1703.00910)

Cosmology/Astrophysics/other

- DM-proton CMB (arXiv:1712.07133)
- Reverse direct detection (arXiv:1810.07705)
- Neutron star heating (arXiv:1704.01577)
- Boosted dark matter (arXiv:1405.7370)
- Colliders (arXiv:1903.01400)

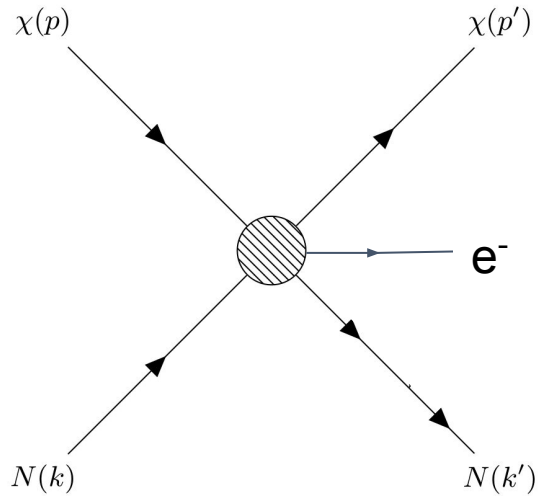
Improved (traditional) direct detection:

- Diamond detectors (arXiv:1901.07569)
- Single charge germanium (arXiv:1804.10697)
- Single charge xenon (e.g. ALBECA)
- Helium detectors (arXiv:1302.0534)

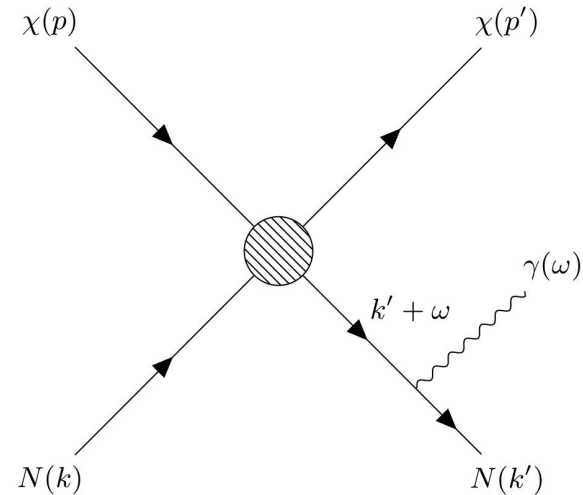
New Ideas

- Multi-exciton/rotons in LHe (arXiv:1611.06228)
- Absorption (arXiv:1608.01994)
- Chemical bond breaking (arXiv:1608.02940)
- Spin-flip avalanche (arXiv:1701.06566)

The Migdal effect and bremsstrahlung



The Migdal effect



Bremsstrahlung

Why are these detection modes useful?

1. Kinematic advantage: light particles carry away more energy per unit momentum

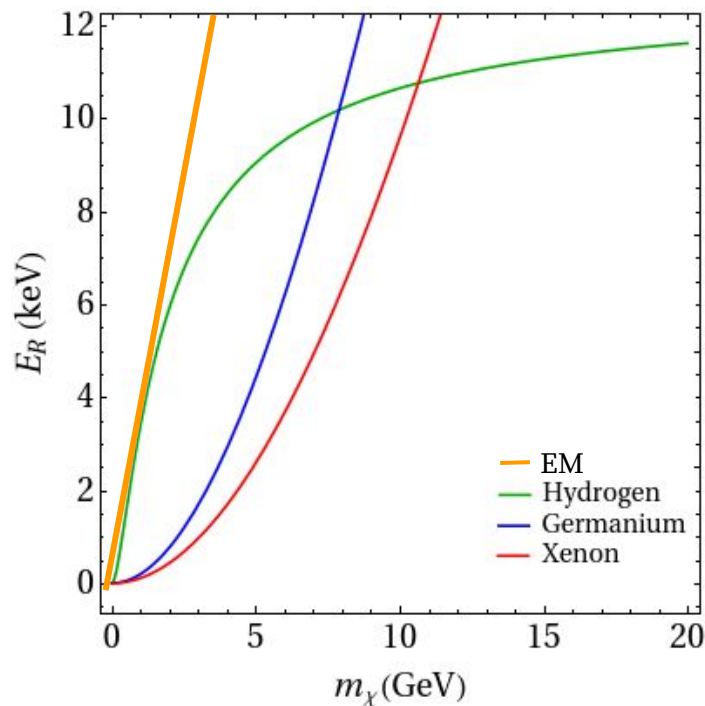
Nuclear:

$$E_{R_{\max}} = \frac{2\mu_T^2}{m_T} v_{\max}^2$$

Electronic
(Migdal/Brem):

$$E_{EM_{\max}} = \frac{\mu_T}{2} v_{\max}^2$$

2. Electronic energy isn't quenched



A brief history of the Migdal effect

1941: A.B. Migdal, J. Phys. USSR 4 449

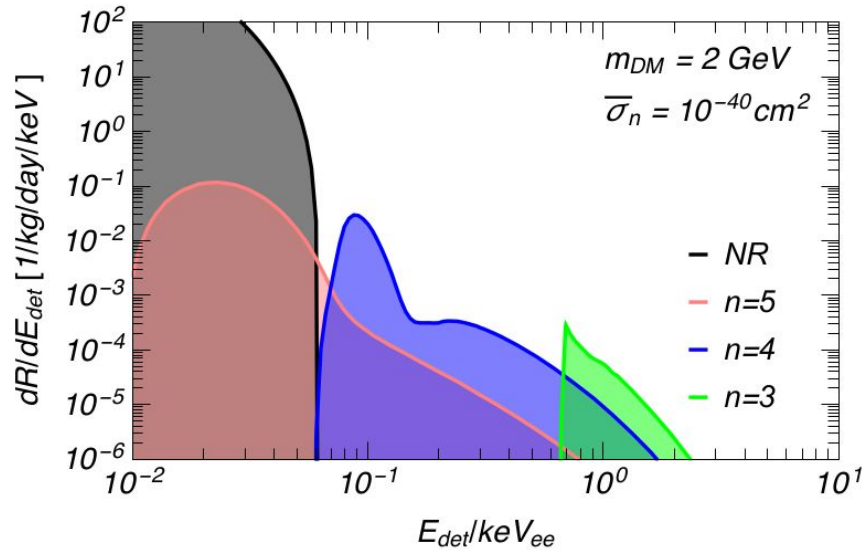
1958: Landau and Lifshitz Vol. 3: Quantum Mechanics, sec. 41:

PROBLEM 2. The nucleus of an atom in the normal state receives an impulse which gives it a velocity v ; the duration τ of the impulse is assumed short in comparison both with the electron periods and with a/v , where a is the dimension of the atom. Determine the probability of excitation of the atom under the influence of such a “jolt” (A. B. MIGDAL 1939).

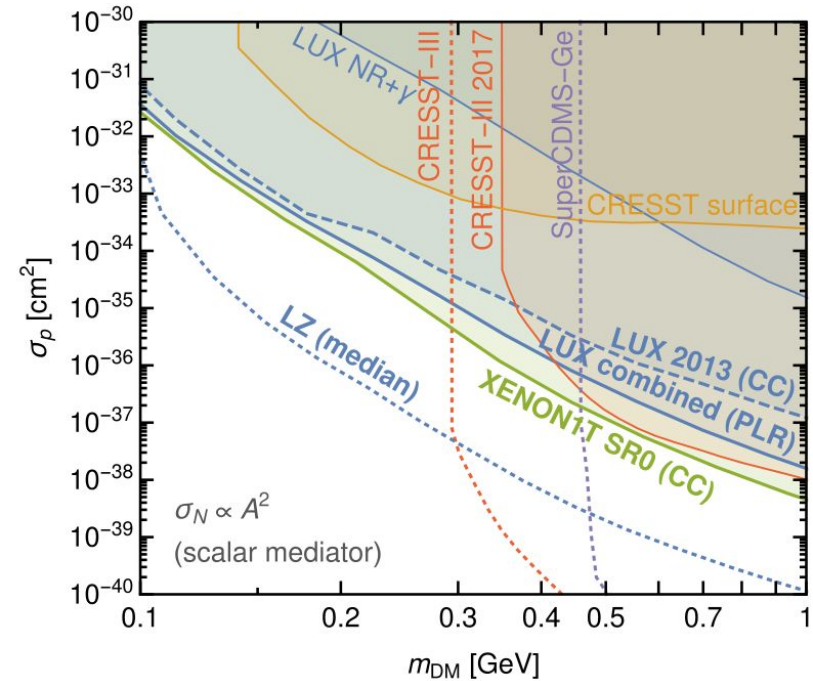
2005: J.D. Vergados and H. Ejiri, Phys. Lett. B 606, 313, [[hep-ph/0401151](https://arxiv.org/abs/hep-ph/0401151)]

2018: M. Ibe, W. Nakano, Y. Shoji and K. Suzuki, JHEP 1803 (2018) 194 [arXiv:1707.07258]

Bounds from the Migdal effect

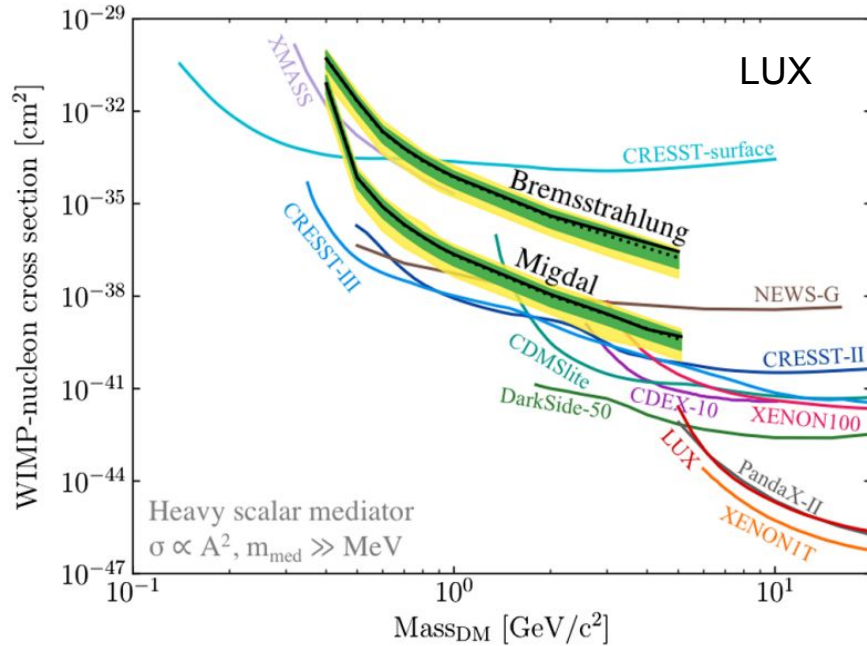


M. Ibe, W. Nakano, Y. Shoji, and K. Suzuki,
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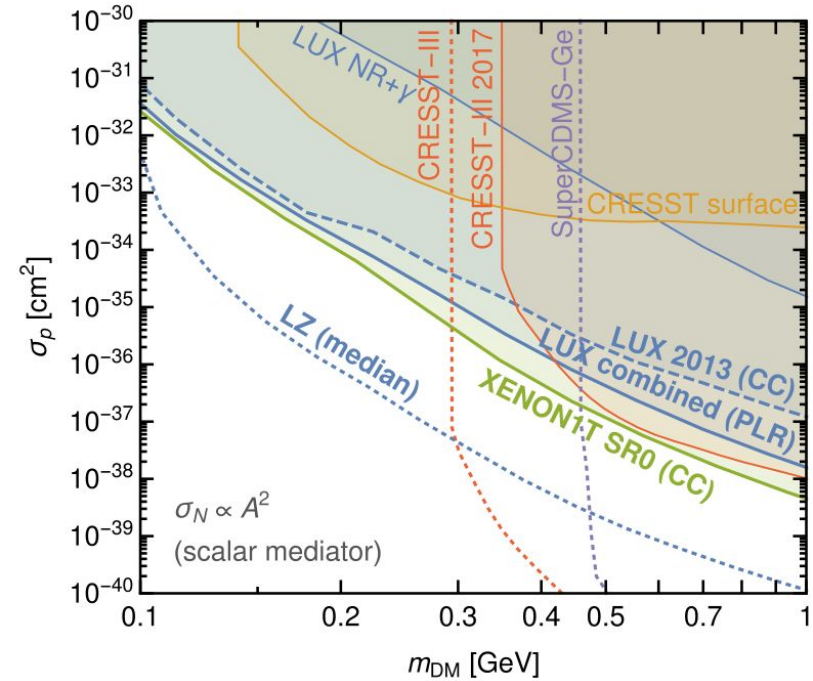


M. J. Dolan, F. Kahlhoefer, and C. McCabe, (PRL)
arXiv:1711.09906

Bounds from the Migdal effect

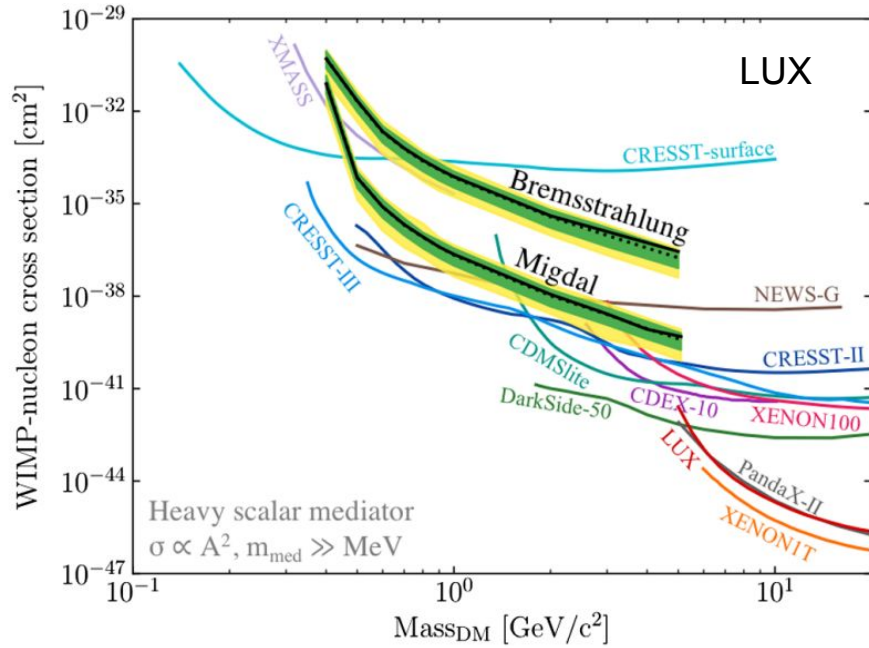


Akerib et al. Phys. Rev. Lett. **122**, arXiv:1811.11241



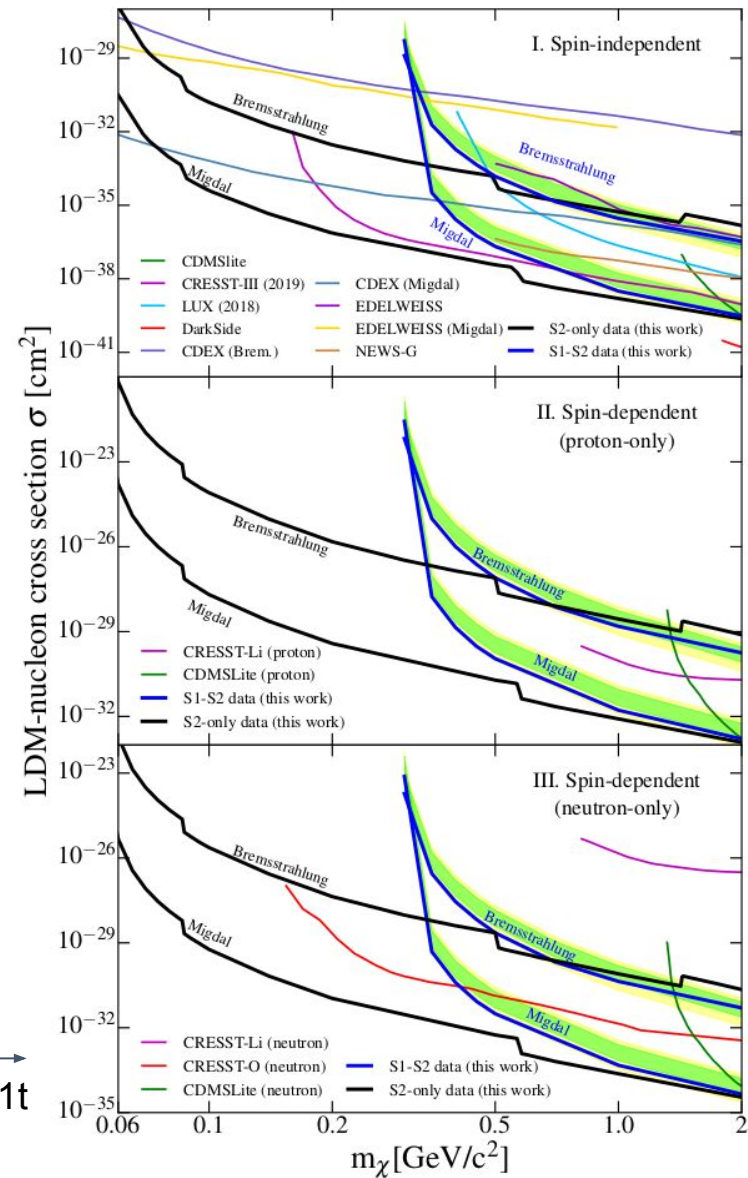
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Bounds from the Migdal effect



Akerib et al. Phys. Rev. Lett. **122**, arXiv:1811.11241

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XENON1t



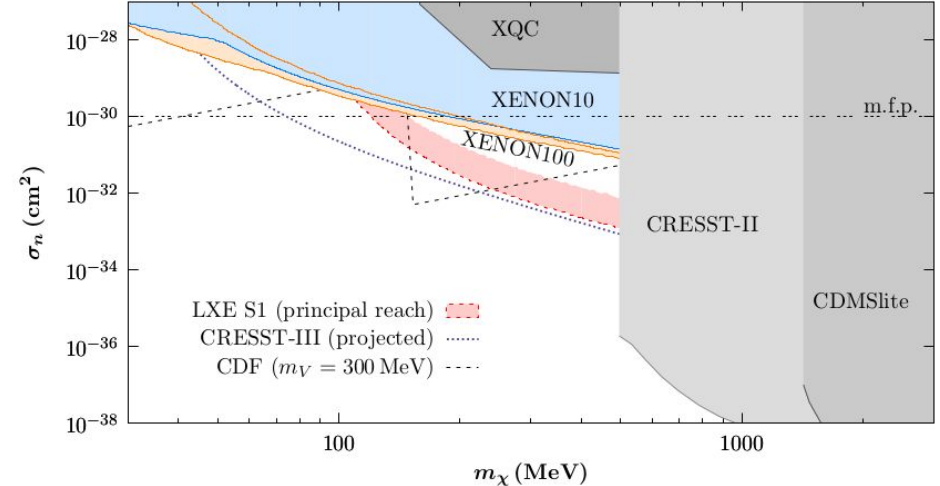
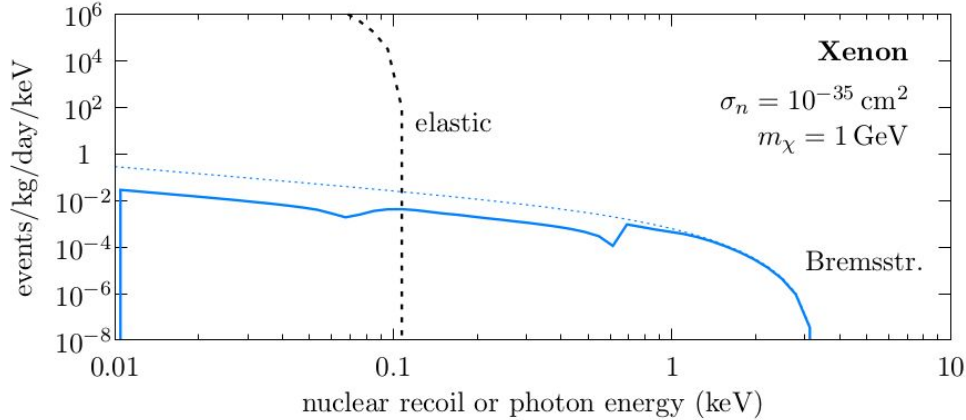
Bounds from bremsstrahlung

Kouvaris and Pradler calculate the rate of nuclear bremsstrahlung due to DM nuclear recoil from
 The factorized cross section:

$$\left. \frac{d^2\sigma}{dE_R d\omega} \right|_{\text{naive}} = \frac{4Z^2\alpha}{3\pi} \frac{1}{\omega} \frac{E_R}{m_N} \times \frac{d\sigma}{dE_R} \Theta(\omega_{\text{max}} - \omega)$$

arXiv:1607.01789

Phys. Rev. Lett. 118, 031803 (2017)



Calculating Migdal and Brem. rates

Bremsstrahlung

$$\frac{d^3 R}{dE_R d\omega dv} = \frac{d^2 R_{\chi T}}{dE_R dv} \frac{4\alpha Z^2}{3\pi} \frac{E_R}{m_T \omega}$$

The Migdal effect

$$\frac{dR}{dE_R dE_e dv_{DM}} \simeq \frac{dR_0}{dE_R dv_{DM}} \times \frac{1}{2\pi} \sum_{n,\ell} \frac{d}{dE_e} p_{q_e}^c(n\ell \rightarrow E_e)$$

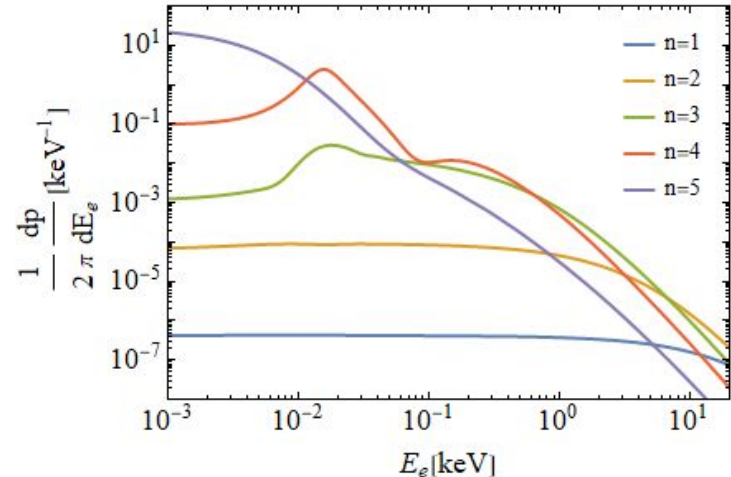
Calculating Migdal and Brem. rates

Bremsstrahlung

$$\frac{d^3 R}{dE_R d\omega dv} = \frac{d^2 R_{\chi T}}{dE_R dv} \frac{4\alpha Z^2}{3\pi} \frac{E_R}{m_T \omega}$$

The Migdal effect

$$\frac{dR}{dE_R dE_e dv_{DM}} \simeq \frac{dR_0}{dE_R dv_{DM}} \times$$



What does a Migdal event look like?

There are three components:

- i. Nuclear recoil
- ii. Electron energy
- iii. Atomic de-excitation

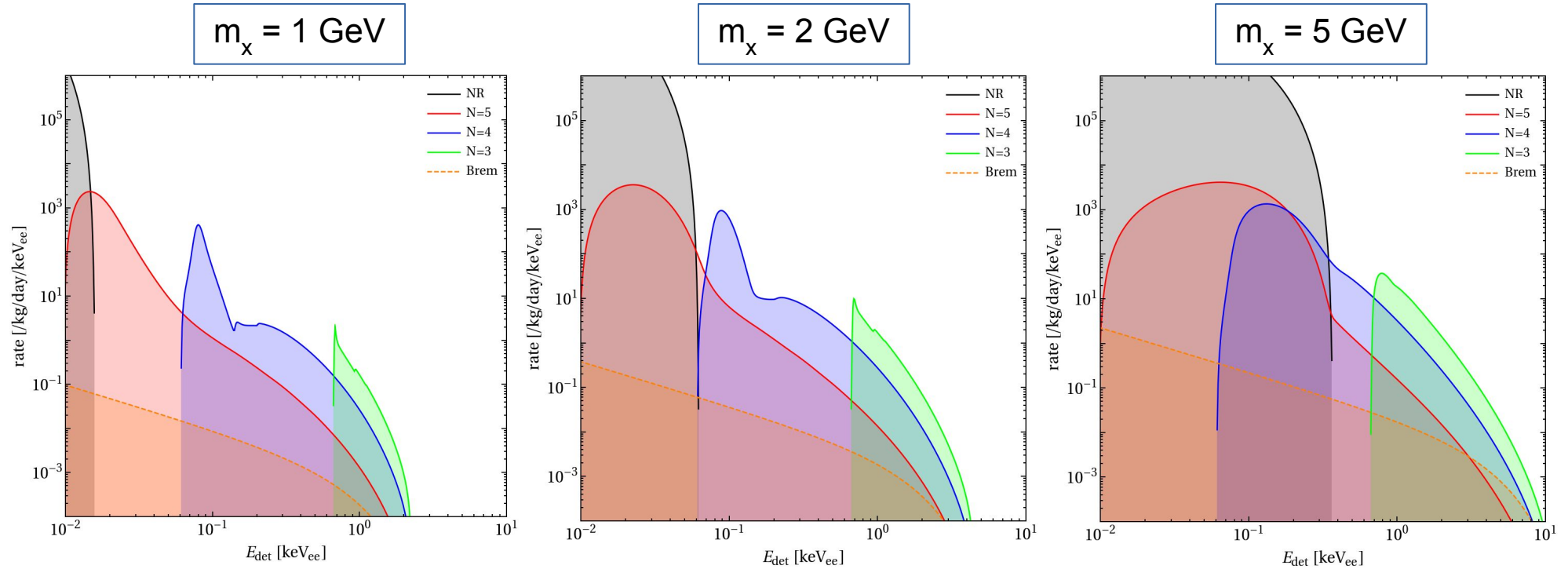
De-excitation proceeds via radiative, Auger or Coster-Kronig transition, dependent on which shell the vacancy is in.

Total detected energy:

$$E_{\text{det}} = q_{\text{nr}} E_{\text{R}} + E_{\text{EM}}$$

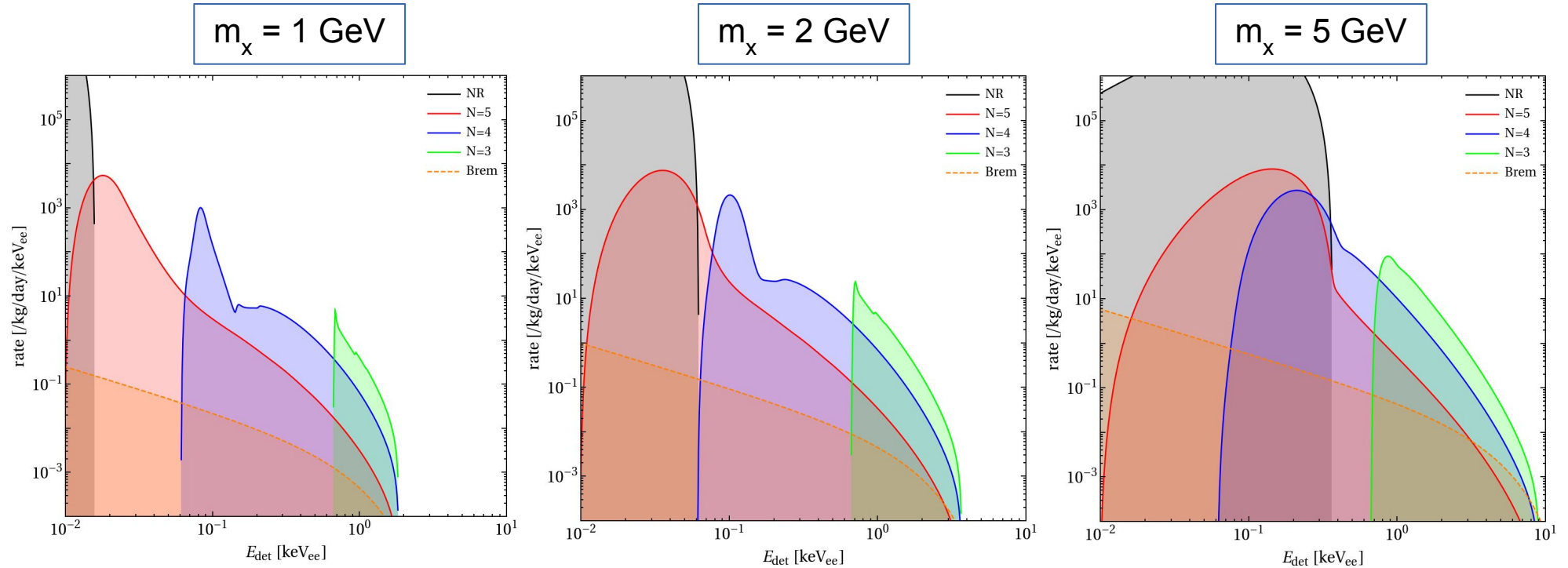
$$E_{\text{EM}} = E_{\text{e}} + E_{\text{nl}}$$

Migdal and Brem. rates for WIMPs



Spin-independent

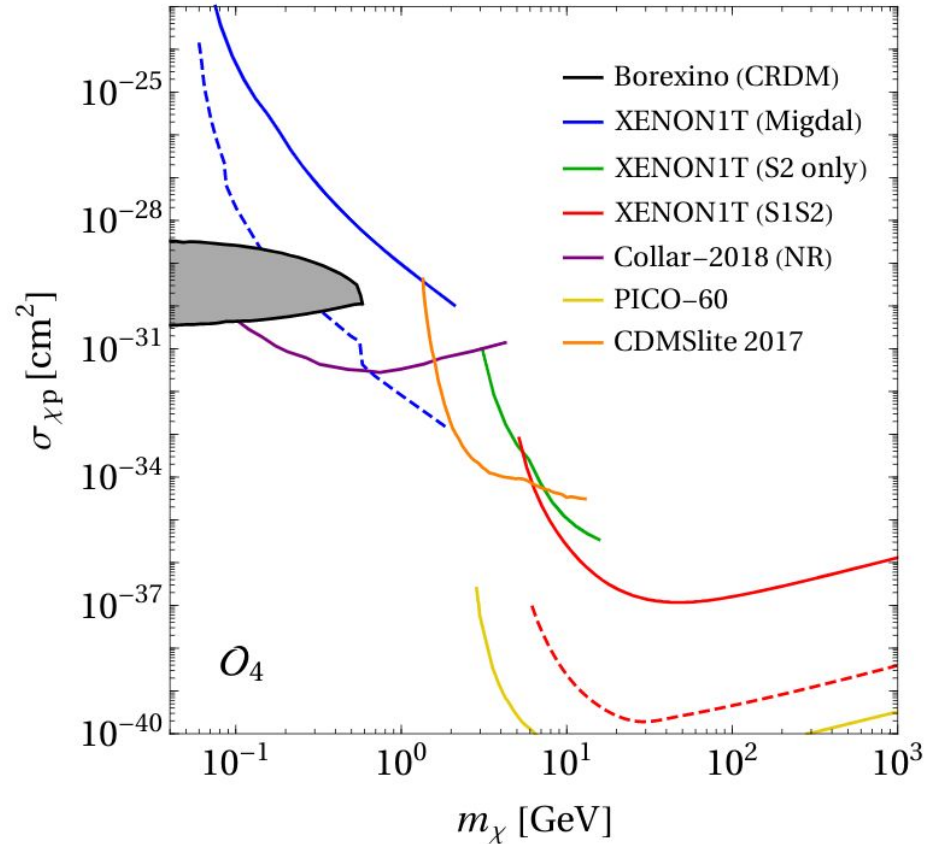
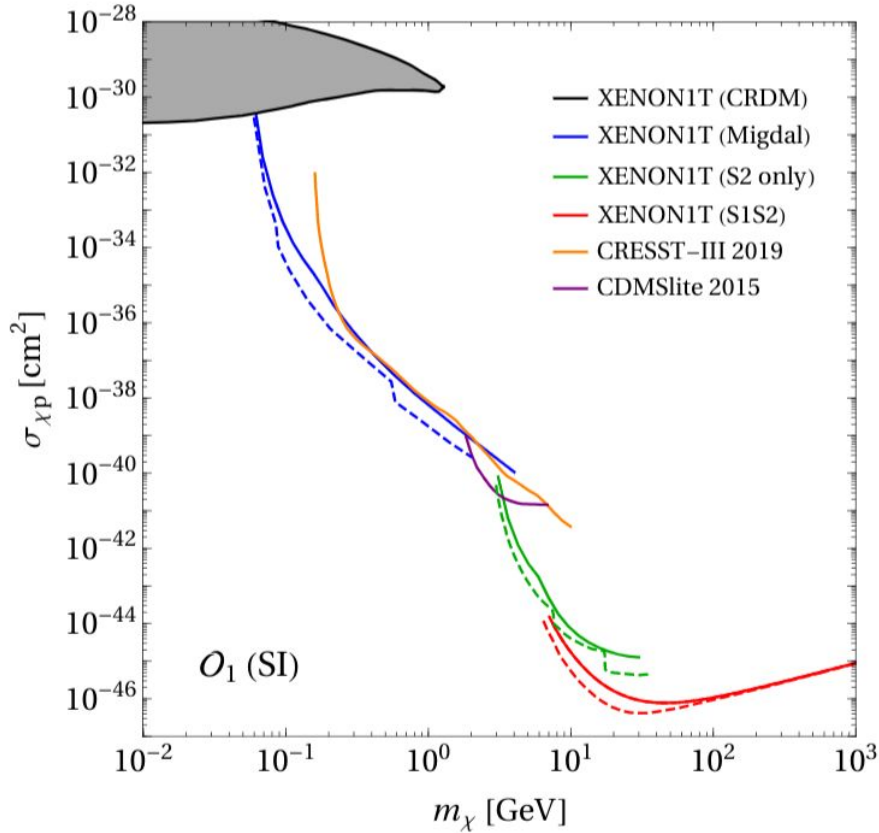
Migdal and Brem. rates for WIMPs



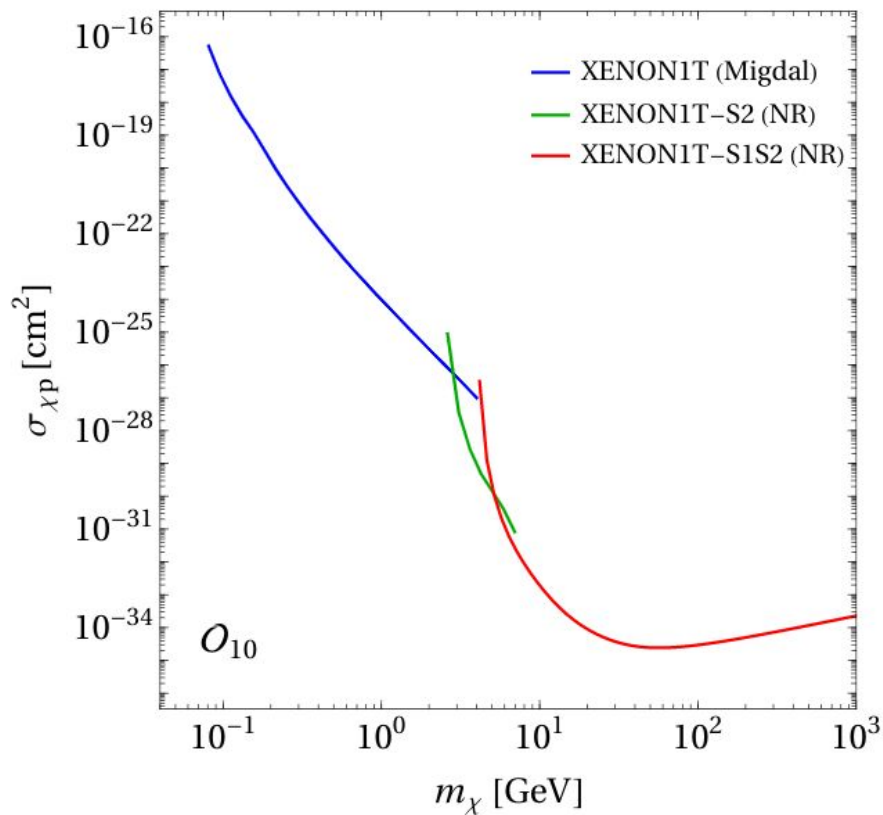
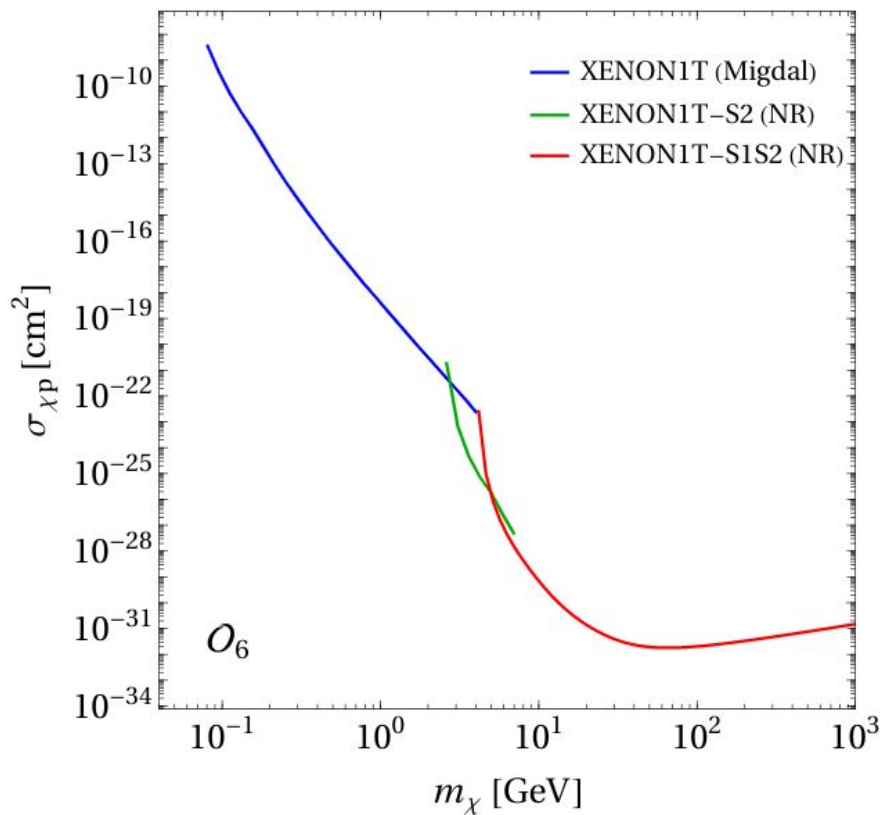
$$O_6 = (S_x \cdot q/m_N)(S_N \cdot q/m_N)$$

- spin and momentum dependent

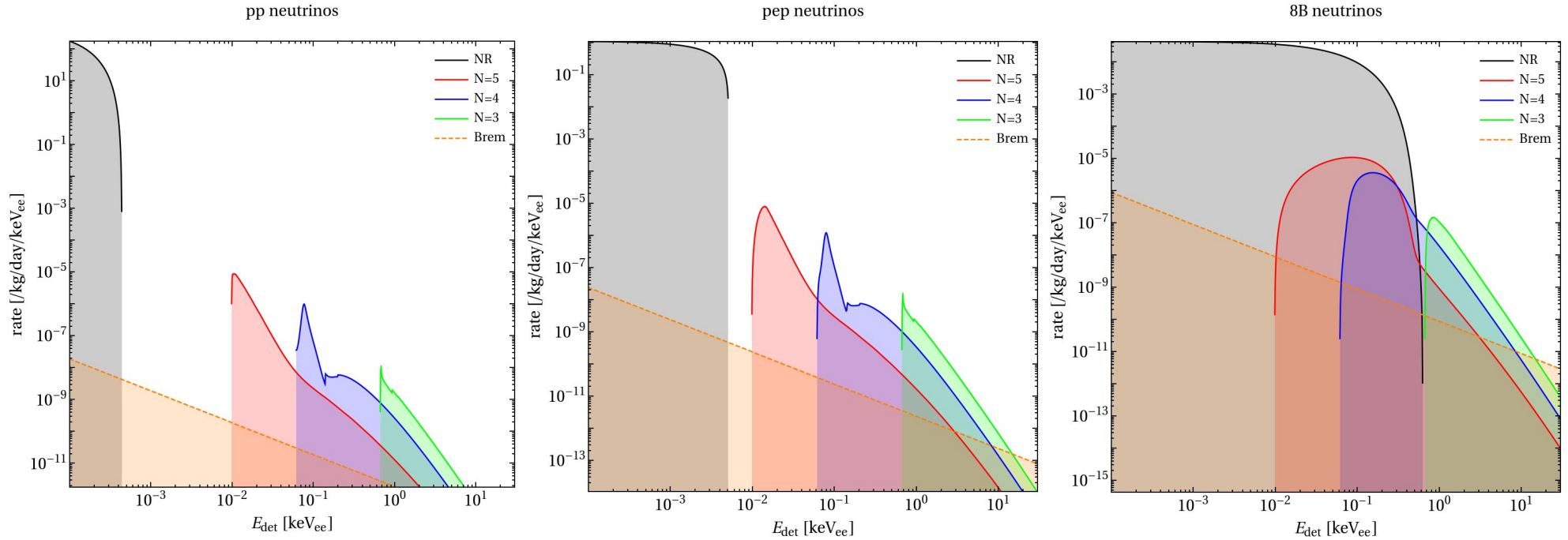
Our bounds from the Migdal effect



Our bounds from the Migdal effect

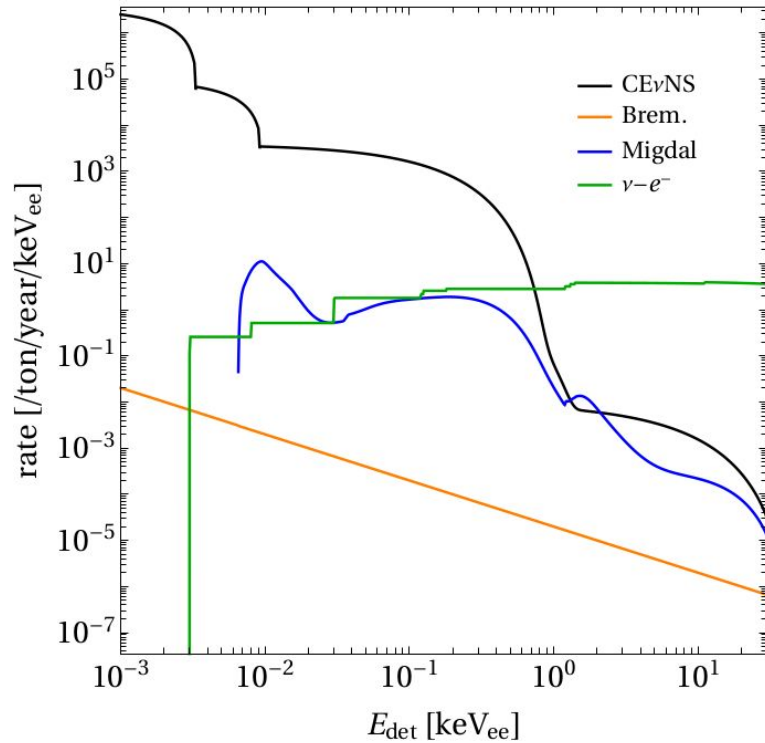


Migdal and Brem. rates for neutrinos

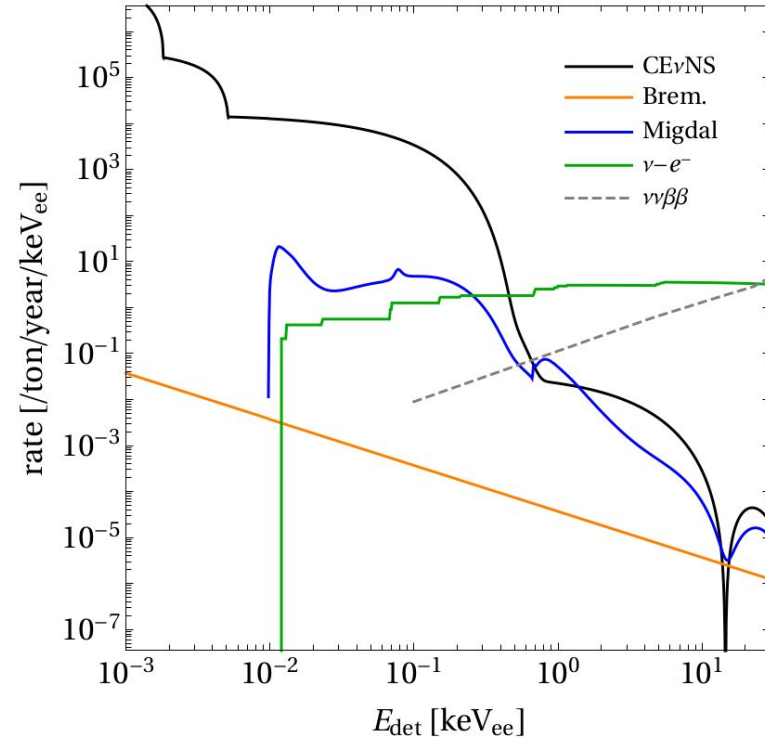


Relevance of the Migdal effect for neutrino rates

Germanium

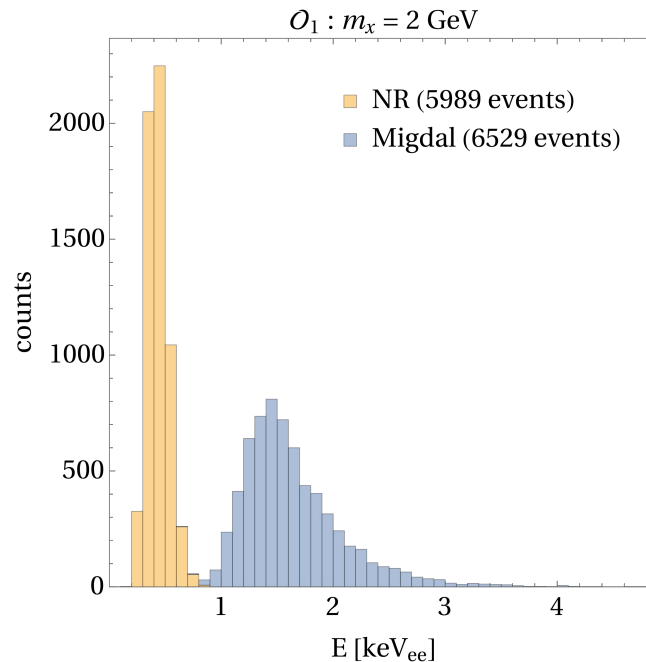
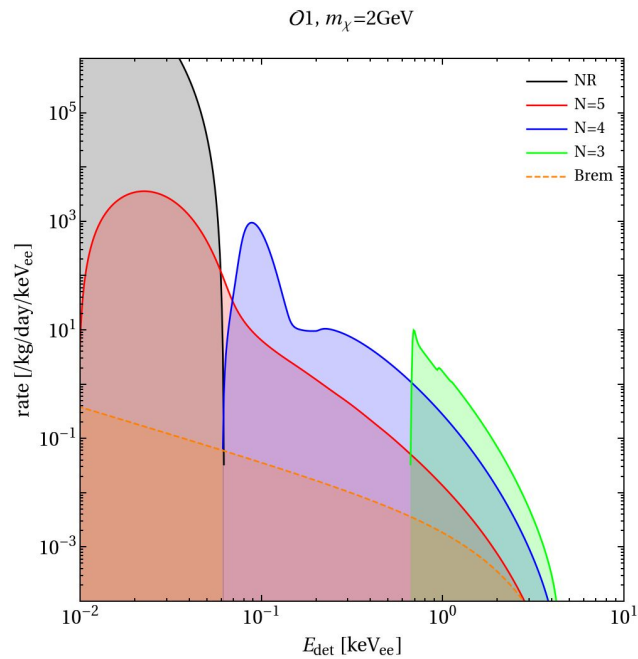


Xenon



Using NEST to simulate Migdal events

1. Sample the NR and Migdal spectrum directly, treating all events individually



- Ignores that part of Migdal event's energy comes from NR

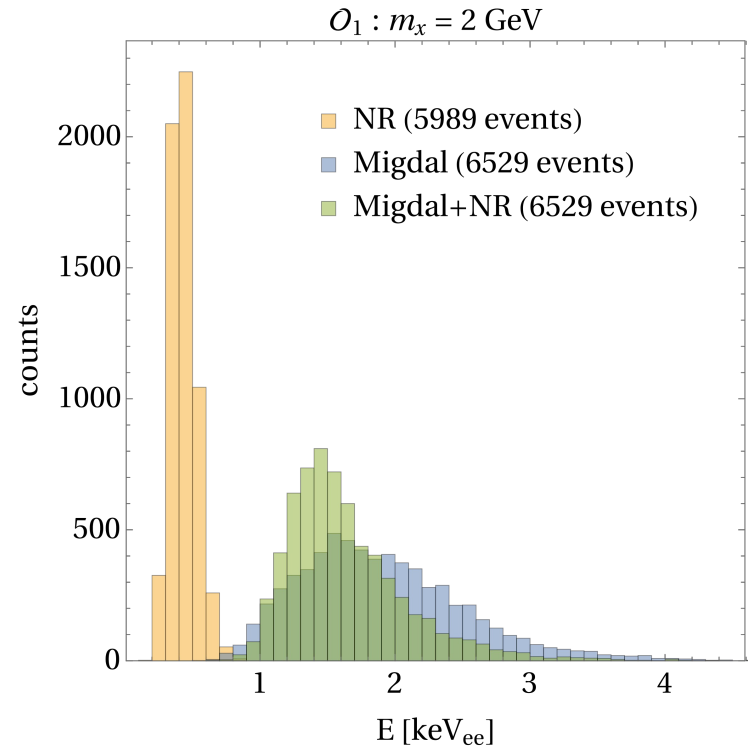
$$E_{\text{det}} = q_{\text{nr}} E_{\text{R}} + E_{\text{EM}}$$

- Only suitable if NR is *well* below threshold

Using NEST to simulate Migdal events

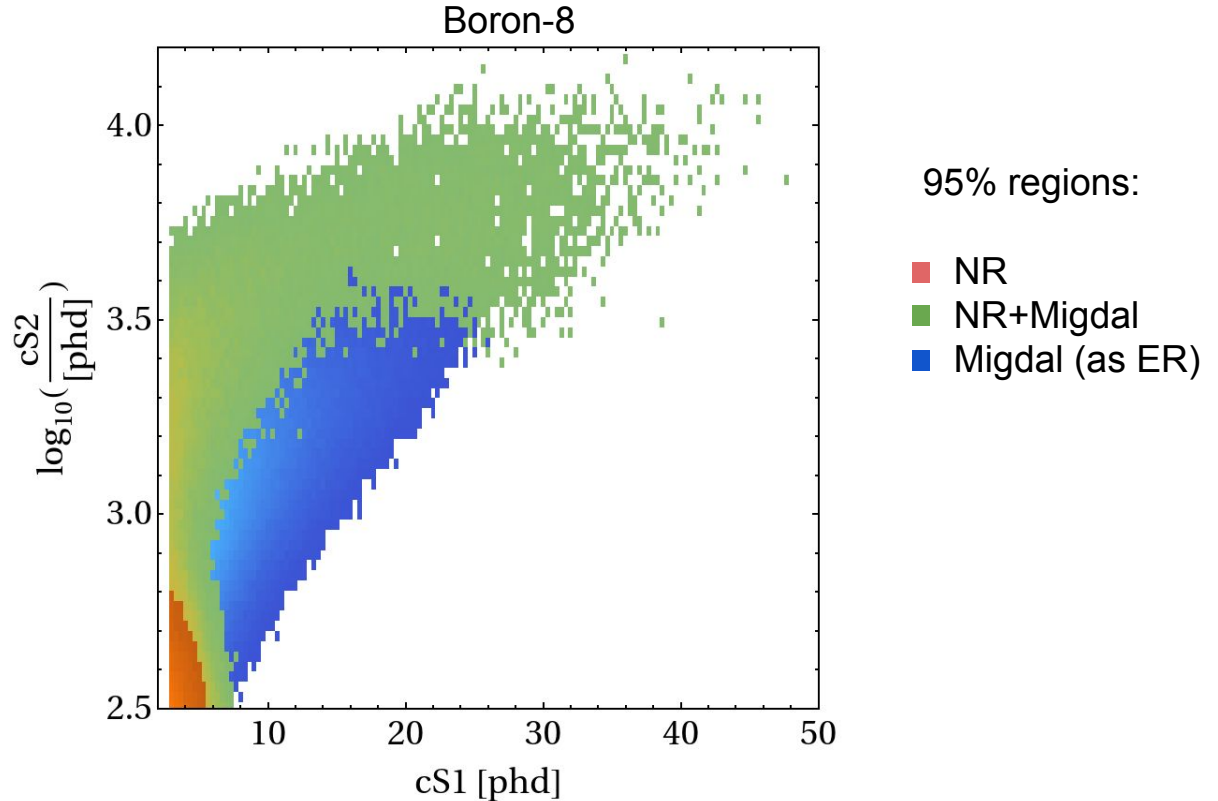
2. Event-by-event simulation: use NRs to produce Migdal events

- i. MC sample E_{NR} and ν
- ii. Randomly select an electron (or none) with scaled probability
- iii. MC sample E_e for selected electron shell
- iv. De-excite atom, producing gamma or Auger electrons



Using NEST to simulate Migdal events

- In S1/S2 space we see that the Migdal events occupy a different region when simulated event-by-event



Three take-aways:

- The Migdal effect is a useful probe of light dark matter, and dominates over bremsstrahlung for all cases explored
- The Migdal effect from solar neutrinos has a rate comparable to the atmospheric rate
- Need to model the Migdal effect event-by-event with the NR, not as a separate signal channel in ER