

# W physics in ATLAS

Fabrice Balli, CEA Saclay

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## **Outline**

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	- W production at LHC
	- W detection at ATLAS
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	- W+jets at 8 TeV
	- W mass measurement at 7 TeV
	- Wjj and aGC at 7 and 8 TeV
- Conclusive remarks



## Introduction



**F.Balli — W physics at ATLAS — Grodno, 12–24 August , 2018** Le Modele Standard (MS) de la physics at ATLAS — Grodno, 12–24 August , 2018

#### The Standard Model (SM)



- Standard Model : basic model of elementary particle physics constituents à l'antication de la constitution de l a matura de la materiale de la materiale della del<br>Del model of elements
- electromagnetic force photon
- in interaction  $=$   $\angle$ , vv $\cdot$ , vv • weak interaction — Z, W+, W-
- $\log$  interaction  $\sim$  o giuons • strong interaction — 8 gluons
- Higgs boson confers mass to ULLION PULLIOIUU ; UNUUVIULU I<br>2 hy ATI AS and CMS the other particles ; discovered in 2012 by ATLAS and CMS
- interaction distance  $\alpha$  individual must define the matrix  $\alpha$ interaction direction  $\alpha$  interaction  $\alpha$ • Gravitational interaction — not described by the SM



### The W boson

- Discovered in UA1 and UA2 at CERN SPS in 1983
	- 1984 Nobel prize awarded to Carlo Rubbia and Simon van der Meer
- Charge : ±1e
- Width: 2.085±0.042 GeV
- Spin : 1
- Mass: 80.385±0.015 GeV
- Decay channels :
	- $e, \mu, \tau$ : BR ~11 % each
	- hadrons : BR  $\sim$  67 %



- A well-known pillar of SM : why worry about further measurements ?
	- Background to other processes (Higgs, ttbar...)
	- Stringent tests of SM consistency (EW fit), probe of pQCD, of anomalous gauge couplings —> **More precise** measurements of differential cross-sections and mass is **necessary**

![](_page_5_Picture_0.jpeg)

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### The Large Hadron Collider at CERN

![](_page_5_Figure_3.jpeg)

- 27 km circumference
- The only high-energy physics proton-proton collider currently running Le plus ienergietigetigue en fonctionnement :<br>Le plus dans le plus de masseur : relevaire en rendere renadere realisateur en massath : renadire en massiment en 2011 (2011) (14 Tenergie nominale) (14 Tenerale) (14 Tenerale) (14 Tenerale) (14 Tenerale) (14 Tenerale) (1
- Has successfully delivered big amounts of collision data over the last 7 years to the 4 and successions denvered big amodites of comsion data over the fast right.<br>detectors : LHCb, ALICE, CMS and ATLAS
	- Center of mass energy is **7 TeV (2011), 8 TeV (2012)**, 13 TeV (2015 up to now)

![](_page_6_Picture_0.jpeg)

#### The ATLAS experiment at LHC in Run1 (2011-2012) ATLAS

![](_page_6_Figure_3.jpeg)

#### Recorded 4.6/20.2 fb-1 luminosity in 2011/2012

- Inner tracker vertices reconstruction, charged tracks
- Electromagnetic calorimeter - electrons, photons
- Hadronic calorimeter jets
- Muon spectrometer in a toroidal magnetic field

![](_page_6_Figure_9.jpeg)

![](_page_7_Picture_0.jpeg)

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## W events at ATLAS Run1

8

- W cross-section ~ 10<sup>2</sup> nb
	- $\bullet \rightarrow$  7 TeV N<sub>W</sub>  $\sim$  470M
	- $\rightarrow$  8 TeV N<sub>W</sub>  $\sim$  2.02B
	- cleanest signature :  $e, \mu$ 
		- $\sim$  47 (202) M events per channel
			- detector acceptance, event selection and reconstruction efficiencies to be taken into account (conservative factor 10 : still a few M just for 7 TeV)

![](_page_7_Figure_9.jpeg)

![](_page_8_Picture_0.jpeg)

Recommended readings :

- CTEQ Lecture from Jeff Owens (2000)
- QCD and Collider Physics, R.K. Ellis, W.J. Stirling, and B.R. Webber

# Theory context

(an experimentalist's view!)

- Theory motivation to m<sub>w</sub> measurement
- W production at LHC
- W detection at ATLAS

![](_page_9_Picture_0.jpeg)

## mW and the EW fit

![](_page_10_Picture_0.jpeg)

– *v* est la valeur moyenne du Higgs dans le vide (vev).

#### One illustration : W mass (m<sub>W</sub>) and the EW fit ⇡↵(*µ*) = <sup>p</sup> 2*Gµm*<sup>2</sup> W sin2 <del>∪</del> w sin2 ∪

- Electroweak theory (true at all orders) (1)
- Also, one has (2)
- After solving the 2nd order equation in mw<sup>2</sup> one  $\begin{bmatrix} 4 & 4 \\ 4 & 2 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 &$ donnaets, inclusion ou n'inclusion de masse du bosonie de masse du bosonie de masse du bosonie de masse du bos gets (3)
- Le quark top • Where radiative corrections to the W boson contributions) can be expressed as : de Higgs mesur´ee. Les valeurs de *m<sup>W</sup>* et de *m*top et leurs d´eviations standards sont ´egalement propagator (dominated by top and Higgs

$$
\Delta r = \Delta \alpha - \frac{\cos^2 \theta_W}{\sin^2 \theta_W} \Delta \rho + \Delta r_{res},
$$

• Top quark mass dependence la by design during the Higgs (corrections) and determines the Higgs (corrections) and determines of  $\sim$ et *r*<sup>*r*</sup> dominated by : <u>contient des `energies propres des `energies propres de bosons, corrigies propres de bosons, corrig</u>

$$
\Delta \rho^{top} \approx \frac{3\sqrt{2}G_{\mu}m_{\text{top}}^2}{16\pi^2}
$$
 G. Burgers and F. Jegerlehner  
10.5170/CERN 1989,008 V 1.55

• Higgs boson mass dependence<br>dominated by : dominated by : **Firm** top

dominated by :  
\n
$$
\Delta r_{res}^{Higgs} \approx \frac{\sqrt{2}G_{\mu}m_W^2}{16\pi^2} \left[\frac{11}{3}(\ln\frac{m_h^2}{m_W^2} - 5/6)\right].
$$

equation in 
$$
m_W^2
$$
 one  
\n
$$
m_W^2 = \frac{g_W^2 v^2}{4}, \ m_Z^2 = \frac{g_W^2 v^2}{4\rho_0 \cos^2 \theta_W} = \frac{m_W^2}{\rho_0 \cos^2 \theta_W} (2)
$$
\n\no the W boson  
\np and Higgs  
\n
$$
= \frac{m_Z^2}{2} (1 + \sqrt{1 - \frac{4\pi \alpha}{\sqrt{2} G_\mu m_Z^2}})
$$
\n
$$
= \frac{m_Z^2}{2} (1 + \sqrt{1 - \frac{4\pi \alpha_{tree}}{\sqrt{2} G_\mu m_Z^2}}) (3)
$$

 $m_W^2 = \frac{\pi \alpha_{tree}}{\sqrt{2}C \sin^2 \theta_{CC}}$  (1)

(1)

*W*

Corrections `a une boucle au propagateur

Relationship between W mass, top mass and Higgs mass (and EW parameters) !

10.5170/CERN-1989-008-V-1.55

![](_page_11_Picture_0.jpeg)

## The global EW fit

- Idea of electroweak fits
	- Measure many different observables in experiments
	- Calculate the relations between all observables in the Standard Model
	- Probe the consistency of the SM by predicting observables
- Input for the gobal electroweak fit mostly from
	- LEP: Z boson observables (e.g.  $sin^2\theta_W$ )
	- Tevatron: W boson, top quark mass
	- LHC: Higgs boson, top quark mass
- Overall good consistency between indirect determination (i.e. physics parameter left free) and the direct measurements

![](_page_11_Figure_12.jpeg)

![](_page_12_Picture_0.jpeg)

 $\tilde{z}$ 

## The global EW fit

- Test the consistency of the Standard Model
	- e.g. predict mw, provided all other input measurements
- needs 7 MeV precision to compete with indirect determination from theory fit (10-4 relative uncertainty!)
- Electroweak precision measurements also sensitive to several new physics scenarios
	- For this, need a 5 MeV precision on mw
- $\rightarrow$  this measurement needs very accurate prediction for W production and kinematics of decay products :
	- W  $p_T$  and rapidity spectrum
	- polarisation (spin correlations)
	- high order EW (NLO)
- Proton PDFs are an essential ingredient for this
- It also needs detector calibration at the same level of precision! 13

![](_page_12_Figure_14.jpeg)

arXiv:1803.01853

![](_page_13_Picture_0.jpeg)

#### parenthesis : sin2θ<sup>W</sup>

- One of the key inputs to EW and BSM fits
- Observed tension between LEP and SLD measurements (~3 sigmas)
- New preliminary measurement from ATLAS

ATLAS-CONF-2018-037

![](_page_13_Figure_7.jpeg)

![](_page_14_Picture_0.jpeg)

# W production

![](_page_15_Picture_0.jpeg)

#### Drell-Yan production

• factorization theorem :

$$
\sigma_{pp \to X}(\alpha_s, Q^2) = \sum_{a,b} \int_0^1 f_a(x_1, Q^2) f_b(x_2, Q^2) \times \hat{\sigma}_{ab \to X}(\alpha_s, Q^2) dx_1 dx_2
$$

- weight the partonic cross-section by non-perturbative fonctions (parton distribution functions, PDFs) considered at high scale  $Q^2 = \mu F^2$  (separates perturbative and nonperturbative regime)
- partonic cross-section can be calculated perturbatively and is only known up to NNLO and thus depends on a renormalisation scale  $\mu$ <sub>R</sub><sup>2</sup> (=Q<sup>2</sup>):

$$
\hat{\sigma}_{ab\to X}(\hat{s}) = \underbrace{\hat{\sigma}_{0}(\hat{s})}_{LO} + \underbrace{\alpha_{s}(\mu_{R}^{2})\hat{\sigma}_{1}(\hat{s})}_{NLO} + \underbrace{\alpha_{s}^{2}(\mu_{R}^{2})\hat{\sigma}_{2}(\hat{s})}_{NNC}
$$
\n
$$
= \underbrace{\hat{\sigma}_{0}(\hat{s})}_{U}
$$
\n
$$
\underbrace{\mu_{0}}_{U}
$$
\n<math display="</math>

![](_page_16_Picture_0.jpeg)

### W production : PDFs

![](_page_16_Figure_3.jpeg)

- W+(W-) production is dominated by ud (dū)
- $u_v$  is higher at high  $x \rightarrow$  increase of  $W^+$ production at high rapidities (next slide)
- Contribution from 2nd quark generation is not negligible : ~25%
	- induces larger uncertainties than in pp
- Gluon contribution starting at NLO

![](_page_16_Figure_9.jpeg)

![](_page_17_Picture_0.jpeg)

#### W production : rapidity

- Rapidity y is defined as :  $y = 1/2 \ln \left[ \frac{E + p_z}{E p_z} \right]$
- Assuming intrinsic transverse momentum of partons  $k_T=0$ , we have (in center of mass frame) :  $p_a = \sqrt{s/2}$  \* $x_a$  (1,0,0,1) and  $p_b = \sqrt{s/2}$  \* $x_b$  (1,0,0,-1)

• 
$$
\longrightarrow
$$
 y = 1/2 ln (x<sub>a</sub>/x<sub>b</sub>)

• Leading to :  $x_a =$ *M s*  $e^y$  *x<sub>b</sub>* = *M s e*−*<sup>y</sup>*

![](_page_17_Figure_7.jpeg)

- low  $|y|$  : mainly sea quarks  $(x \sim 10^{-2})$
- high |y| : 1 sea quark and 1 valence quark (x~10-4, a few 0.1)
- One has at lowest order :

$$
\frac{d\sigma}{dy} = \frac{1}{s} \sum_{a,b} \sigma_0^{ab \to W}(M) f_a(x_a, M^2) f_b(x_b, M^2)
$$

- Boson rapidity directly sensitive to PDFs
- allows to constrain them from differential cross-section measurements

![](_page_18_Picture_0.jpeg)

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#### W p<sub>T</sub> : where does it come from?

![](_page_18_Figure_3.jpeg)

Fig. 9.2. The lepton pair transverse momentum from the CFS collaboration [4]. The curve corresponds to a Gaussian intrinsic  $k_T$  distribution for the annihilating partons.

- Intrinsic  $k<sub>T</sub>$  of the partons : good agreement with the data (fixedtarget pn collisions) up to  $p_T \sim 2$  GeV  $\rightarrow$  assume Gaussian form with  $\langle k_T(\text{parton}) \rangle \sim 760 \text{ MeV}$
- Not sufficient to describe higher values of W  $p_T$

![](_page_19_Picture_0.jpeg)

#### W p<sub>T</sub> : where does it come from?

- Need additional hard parton emissions to explain higher  $p_T$  region
	- NLO (qg—>Wq, qq'—>Wg), NNLO QCD
	- NNLO diagrams are typically :

![](_page_19_Figure_6.jpeg)

![](_page_20_Picture_0.jpeg)

#### W p<sub>T</sub> : let's diverge a little bit

- spoiled by large logarithms of the type  $\alpha_S^n\ln^m(M^2/p_T^2)$ • Pure fixed-order NNLO predictions : diverge when  $p_T \rightarrow 0$  ( $p_T$ << M ) due to the presence of soft and collinear emissions —
- This can (has to) be resummed at all orders and gives

$$
\frac{d\sigma}{d\tau dy dp_T^2} = \left(\frac{d\sigma}{d\tau dy}\right)_{Born} \frac{4\alpha_s}{3\pi} \frac{\ln s/p_T^2}{p_T^2} \exp\left(-\frac{2\alpha_s}{3\pi} \ln^2 s/p_T^2\right)
$$

#### ATL-PHYS-PUB-2014-015

![](_page_20_Figure_7.jpeg)

- Where the exponential is referred to as the 'Sudakov form factor'
- However, does not include the cases where multiple gluons are emitted with  $k_T \sim$ p<sub>T</sub>, nor the cases where gluon momenta add to 0.
- Several resummation formalisms and calculations to resum the leading, next-toleading and next-to-next-to leading logs
	- e.g., RESBOS, DYRES, Geneva, RADISH…
- Can also use parton showers (typically done in simulations) : Sherpa, Pythia, Herwig…
	- Evolution of fragmentation functions through DGLAP formalism

![](_page_21_Picture_0.jpeg)

![](_page_21_Figure_2.jpeg)

- high p<sub>T</sub>~M : fixed-order V+1 jet (MC : fixed-order matrix elements) ; resummation does not work
	- d $\sigma$ /d $p_T$ <sup>2</sup> goes as  $1/p_T$ <sup>2</sup>
- low  $p_T << M$ : fixed-order breaks down, resummation comes in (MC : Parton showers)
- Transition region : no fixed boundary
	- resummation works but fixed-order gives sensible results as well
		- Best prediction from consistent combination of the two
	- 22 • MC : Matrix element + parton shower merging/matching

![](_page_22_Picture_0.jpeg)

#### NLO EW emissions

• Corrections to W production take into account  $O(\alpha_{EM})$  corrections : photon radiation (ISR, FSR), loop corrections (pure weak), ISR/FSR interference (IFI)

![](_page_22_Figure_4.jpeg)

- ISR factorised in the PDF evolution, FSR is the most important numerically among the rest
- Several tools to handle this. In ATLAS W/Z simulated samples, typically :
	- Pythia8 parton shower handles the ISR
	- Photos handles FSR and electron pair emissions from virtual photon (higher order, uncertainty only)
	- pure weak and IFI corrections not included but impact is estimated with other tools (YFS, Winhac…) (and often simply added as uncertainty)

![](_page_23_Picture_0.jpeg)

## W detection

![](_page_24_Picture_0.jpeg)

#### Event topology, definitions of observables

- Detect single (inclusive) W boson decaying into a lepton and a neutrino
- The ATLAS detector measures :
	- The **lepton** charge and 4-vector (transverse momentum  $\vec{\mathsf{p}}$ <sub>T</sub>e)
	- The activity recoiling against the W **(hadronic recoil** ū̇̃⊤)
		- measures additional jets from signal
		- Sensitive to additional interactions (pile-up) and underlying event
		- Enables to indirectly reconstruct the neutrino transverse momentum  $\vec{p}_{\text{T}}$ <sup>miss</sup> = - ( $\vec{u}_{\text{T}}$  +  $\vec{p}_{\text{T}}\ell$ )
		- Some analyses use a direct reconstruction of p<sub>T</sub>miss, some use a direct reconstruction of the recoil —> different algorithms, impact is significant  $\sim$ only for mw here 5, Impact is

![](_page_24_Figure_11.jpeg)

ruct the neutrino  $\overrightarrow{u}$  : vector sum of calorimeter  $\epsilon = -\vec{u}_{\text{T}} + \vec{p}_{\text{T}}\ell$  deposits excluding lepton deposits

 $\text{m}_\text{T} = \sqrt{[2 \text{ pr} \ell \text{ pr}^{\text{miss}} (1-\text{cos} \Delta \varphi)]}$ 

Impossible to fully reconstruct mw because of the neutrino

![](_page_25_Picture_0.jpeg)

#### The Jacobian peaks :  $m<sub>T</sub>$  and lepton  $p<sub>T</sub>$

![](_page_25_Figure_3.jpeg)

![](_page_26_Picture_0.jpeg)

#### The Jacobian peaks : m<sub>T</sub> and lepton p<sub>T</sub>

- Unlike the Z, not possible to fully reconstruct the W mass due to neutrino
- One can use the transverse mass m<sub>T</sub>(*iv*):

 $m_T(\ell \nu)^2 = (\mid p_T^{\ell} \mid + \mid \overline{p_T^{\nu}} \mid)^2 - (p_T^{\ell} + \overline{p_T^{\nu}})^2 = 2 \mid p_T^{\ell} \mid \mid \overline{p_T^{\nu}} \mid (1 - \cos \Delta \phi_{\ell \nu})$ 

• Can be compared to the invariant mass :

$$
m(\ell \nu) = (|\overrightarrow{p^{\ell}}| + |\overrightarrow{p^{\nu}}|)^2 - (\overrightarrow{p^{\ell}} + \overrightarrow{p^{\nu}})^2
$$

- For small  $p_T(W)$ ,  $m_T$  is invariant to leading order
- $\bullet$  Jacobian peak at m<sub>w</sub>
- In the context of  $m_w$ measurement :
	- $\bullet$  lepton  $p_T$  sensitive to modelling of W  $p_T$  and not so much to detector resolution
	- $m<sub>T</sub>$  is not sensitive to W  $p<sub>T</sub>$ but is very sensitive to MET resolution

![](_page_26_Figure_13.jpeg)

![](_page_27_Picture_0.jpeg)

# Recent ATLAS measurements

![](_page_28_Picture_0.jpeg)

## W cross-sections at 7 TeV

Eur. Phys. J. C (2017) 77:367

![](_page_29_Picture_0.jpeg)

#### W cross-sections at 7 TeV : introduction

- Very precise measurement (<1%) of W and Z cross-sections at 7 TeV
- W cross-section is measured in a fiducial volume, extrapolated to full phasespace and differentially in lepton η absolute value, |η*l*|
- QCD analysis: PDF fits, strange-quark density, determination of CKM matrix  $|V_{cs}|$
- Signal MC sample: Powheg+Pythia6 using CT10 NLO PDF set for the matrix element, CTEQ6L1 for the parton shower
	- Assess uncertainties with alternative samples : MC@NLO+Herwig, Powheg+Herwig : matrix element variation, parton shower/underlying event
	- W p<sub>T</sub> is reweighed to Powheg+Pythia8AZNLO, generator tuned to ATLAS Z data
	- Normalised to NNLO prediction from FEWZ with a 5% uncertainty (PDF, scales)
- Background is also using simulated samples except for multi-jet (data-driven)

![](_page_30_Picture_0.jpeg)

#### Event selection

- at least one primary vertex with 3 tracks of  $p_T > 500$  MeV
- MET > 25 GeV
- $m_T > 40$  GeV
- electron channel :
- $\Delta R = \sqrt{\Delta \phi^2 + \Delta y^2}$
- single electron trigger with medium ID and  $p<sub>T</sub> > 20$  or 22 GeV (depending on the period)
	- Identification (ID) : based on shower shape, track properties and track-to-cluster matching
- Calorimeter isolation : upper limit on sum of energy in a cone of size  $\Delta R = 0.2$ ('topoetcone20')
- Tracking isolation : upper limit on sum of  $p<sub>T</sub>$  of tracks in a 0.4 cone ('ptvarcone40')
- exactly one tight ID electron that matches the trigger, in the acceptance ( $|\eta|$  < 1.37 or  $1.52 < |n| < 2.47$ , with  $p_T > 25$  GeV
- reject events with  $>=1$  medium ID electron with  $p_T > 20$  GeV (cuts the Z background)
- Charge-separated analyses (W+ vs W-) : sensitive to charge misID
	- Evaluated from same-sign Z—>ee events in the data and the MC as a function of η, corrected for in the MC

![](_page_31_Picture_0.jpeg)

#### Event selection

- Muon channel
	- single muon trigger with  $p_T > 18$  GeV
	- reconstructed using a combination of muon spectrometer and inner detector information
	- $|z_0-z_{PV}|$  < 1cm (remove background from cosmic rays) : z extrapolated to the beam line
	- $p_T > 25$  GeV,  $|\eta| < 2.4$
	- isolation : ptvarcone40/ $p_T < 0.1$
	- events with  $>= 1$  muon with  $p<sub>T</sub> > 20$  GeV are rejected (cuts the Z background)
	- charge misID negligible

Fiducial volume : phase-space to where the distributions are unfolded at generation level (i.e. remove detector effects):

- lepton  $p_T > 25$  GeV, lepton  $|\eta| < 2.5$  Born level for the leptons
- $p_T v > 25$  GeV
- $m_T > 40$  GeV
- 11 lepton |η bins (common to electron and muon channels) : [0.00, 0.21, 0.42, 0.63, 0.84, 1.05, 1.37, 1.52, 1.74, 1.95, 2.18, 2.50]

![](_page_32_Figure_0.jpeg)

#### Muon calibration among the correlation uncertainties can appear. These correlations were investigated and found to be correlated

- Use a combination of ID and MS, calibrate transverse momentum as a function of η
- Momentum resolution : obtained by fitting the Z invariant mass, as well as 1/p<sub>T</sub> <sup>*ID*</sup> − 1/p<sub>T</sub> <sup>*MS*</sup> for both  $\mu$ <sup>+</sup> and  $\mu^{\scriptscriptstyle +}$  in Z and W  $\,$
- Momentum scale : compare Z peak in data and MC As described in Section 5.1, the kinematic parameters of selected muons are determined from the as $s_{\text{S}}$  information of accuracy of the momentum measurement is limited by imperfective momentum measurement is limited by imperfective momentum measurement is limited by imperfective momentum measurement in  $\mathcal{L}$
- **longitudinal biases (sagitta biases**, from systematic misalignment modes) **comparise the magnetic field, and re**
- muon sagitta bias correction uses W events (E/p) and Z events material in the detector.
- Momentum corrections at the level of 0.1-0.4% and uncertainty of  $\sim$ 2.10<sup>-4</sup>
- Use tag-and-probe methods  $(Z \rightarrow \mu\mu)$  for the scale factors (reconstruction, trigger, isolation) and uncertainties charge-independent momentum-scale correction. The latter typically original  $\alpha$  is a correction of the latter typical corrections or  $\alpha$
- Same level of correction, a bit higher for trigger, 5-10% (still known with a relative uncertainty of 0.1-0.8%) factors proportional to **q**  $\theta$  *p*  $(0.8\%)$  and  $(0.8\%)$

![](_page_32_Figure_11.jpeg)

![](_page_32_Figure_12.jpeg)

![](_page_33_Picture_0.jpeg)

#### Electron calibration

- Electron measurement : energy from the EM calorimeter; eta and phi from the ID
- Calibration sequence :
	- Calorimeter longitudinal intercalibration using muon energy deposits ( $Z \rightarrow \mu \mu$  events)
	- Passive material and presampler response corrections derived using longitudinal shower profiles of electrons and photons
	- Overall energy scale and resolution from Z—>ee decays
- Selection efficiencies for reconstruction, identification, trigger, isolation
	- use tag-and-probe methods (Z—>ee) for the scale factors and uncertainties

![](_page_33_Figure_10.jpeg)

![](_page_34_Picture_0.jpeg)

#### MET calibration

- MET built from a soft term (tracks) + hard term that comprises leptons and jets
- Uncertainties from each hard object is propagated to the hard term
- Soft term uncertainty is obtained by looking at Z in MC and data (response and resolution)

![](_page_34_Figure_6.jpeg)

![](_page_35_Picture_0.jpeg)

## Multijet background (QCD)

- non-isolated electrons, converted photons or hadrons misidentified as signal electrons, or heavy quarks or hadron decays into muons + MET cut is passed thanks to neutrinos from hadron decays/resolution effect
- multijet is poorly modelled in most ATLAS analyses (huge cross-section, tiny selection efficiency, bad modelling of non-prompt muons…)
	- Need to use the data to estimate its contribution
- Use a control region (CR) enriched in multijet to build shape templates
	- loosen lepton ID and invert isolation requirement, subtract EW/top contributions
- Fit the fraction in a 'normalization (or fit) region'  $==$  signal region with relaxed  $m<sub>T</sub>$  and MET cuts (enriched in QCD)
- Extrapolate the fitted fraction to the signal region by taking into account the selection efficiency
- Scan in  $m<sub>T</sub>/MET$  relaxing cut values  $\rightarrow$  dependence  $\rightarrow$  additional linear extrapolation to signal region (10% correction, added as uncertainty here)
- Additional uncertainties for the template shapes : different requirements for CR, detector calibration uncertainties, alternative signal MC


## Analysis method

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- Measured fiducial cross-section where :
	- $\bullet$  L is the integrated luminosity N<sub>W</sub> is the total number of events, B<sub>w</sub> is the estimated background, C<sub>w</sub> is defined as :
- Then extrapolated to common fiducial volume, where E<sub>W</sub> accounts for the different eta acceptances
- Total cross-sections can be then inferred from :
- where  $A_W = \frac{N}{N}$   $\frac{1}{N}$   $\frac{1}{N}$ *NMC*,*generated*,*fiducial W NMC*,*generated*,*total*
- Luminosity uncertainty is 1.8%
- Theory uncertainties relate to : PDF, NLO ME/PS matching, hadronisation, underlying event

*σ fiducial*,*e*/*<sup>μ</sup> W* =  $N_W - B_W$  $C_W \times L$ 

 $C_W =$ *NMC*,*reconstructed W NMC*,*generated*,*fiducial W*

$$
\sigma_W^{fiducial,\ell\nu} = \frac{\sigma_W^{fiducial, e/\mu}}{E_W^{e/\mu}}
$$

$$
\sigma_W^{total, \ell\nu} = \frac{\sigma_W^{fiducial, \ell\nu}}{A_W}
$$

- C<sub>w</sub> uncertainties amount to ~0.2(0.6)% in the  $\mu$  (e) channel
- Ew has small uncertainties w.r.t. experimental
- Aw has larger uncertainties : 1.5-2%



## **Results**



- Combine using technique introduced at HERA
- Properly taking into account the correlations between the measurements
- $\chi^2$  minimisation allowing to have contributions of the correlated uncertainty sources to shift





## **Results**

$$
\sigma_{W\to\ell\nu}^{\rm tot}~({\rm pb})
$$

 $W^+ \rightarrow \ell^+ \nu$  $6350 \pm 2$  (stat)  $\pm 30$  (syst)  $\pm 110$  (lumi)  $\pm 100$  (acc)

 $W^- \rightarrow \ell^- \bar{\nu}$  $4376 \pm 2$  (stat)  $\pm 25$  (syst)  $\pm 79$  (lumi)  $\pm 90$  (acc)

 $W \rightarrow \ell \nu$  $10720 \pm 3$  (stat)  $\pm 60$  (syst)  $\pm 190$  (lumi)  $\pm 130$  (acc)

$$
R_W = \frac{\sigma_{W \to e\nu}^{\text{fid,e}}/E_W^{\text{e}}}{\sigma_{W \to \mu\nu}^{\text{fid},\mu}/E_W^{\mu}} = \frac{\sigma_{W \to e\nu}^{\text{fid}}}{\sigma_{W \to \mu\nu}^{\text{fid}}} = \frac{BR(W \to e\nu)}{BR(W \to \mu\nu)}
$$
  
= 0.9967 ± 0.0004 (stat) ± 0.0101 (syst)  
= 0.997 ± 0.010.

- Integrated cross-section : dominated by luminosity and acceptance factor uncertainties
- Test of lepton universality
	- Ratio result is more precise than LEP result of  $1.007 \pm 0.019$



#### Systematic uncertainties : electron channel



• Dominated by multijet background and signal modelling (MC@NLO vs Powheg)

- Luminosity dominates
- Total experimental uncertainty (excluding luminosity) is  $~1\%$



#### Systematic uncertainties : muon channel



- Dominated by multijet background, followed by various uncertainties ~at the same level
- Total uncertainty excluding luminosity is ~0.6%



Results : unfolded differential distributions

- Bayesian unfolding, purity > 90%
	- Unfolding is almost an efficiency correction





## Comparison with theory

- Fixed-order NNLO QCD predictions use DYNNLO 1.5 (baseline) and FEWZ 3.1.b2 (used for NNLO uncertainty evaluation)
	- NLO EW corrections provided by MCSANC
- In DY cross-section calculations the value of  $\alpha_{EW}$  can be fixed in different input-parameter schemes
	- Here 'G $\mu$ ' scheme (primary parameters are particle masses and Fermi constant with values taken from PDG) —>see I Dittmaier, S. & Huber, M. J. High Energ. Phys. (2010) 2010: 60

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• Uncertainties in these plots is the dominating PDF uncertainty only









#### Comparison with theory : lepton η and asymmetry

Lepton charge asymmetry defined as :

 $d\sigma_{W+}/d|\eta_{\ell}| - d\sigma_{W-}/d|\eta_{\ell}|$  $A_{\ell} =$  $d\sigma_{W+}/d|\eta_\ell| + d\sigma_{W-}/d|\eta_\ell|$ 

- Significant constrain on  $u/d$  PDFs between  $x \sim 10^{-1}$  and  $10^{-3}$
- Data overall well described (luminosity uncertainty of 1.8% is excluded in the plots)
- HERAPDF2.0, NNPDF3.0, MMHT14 and CT14 more or less agree with the data within uncertainties
	- ABM12 remarkably good, but does a poorer job describing Z distributions





## PDF profiling : methodology

• PDF profiling and full PDF fits heavily rely on methods used in previous experiments (Tevatron, HERA) and global PDF fit groups, see e.g

Eur. Phys. J. C **75**, 458 (2015), JHEP **12**, 100 (2014), JHEP **09**, 061 (2012)

- Use of APPLGRID for theory predictions, together with k-factors from the accurate theory tools described before (NNLO QCD from DYNNLO 1.5, NLO EW from PHOTOS and MCSANC)
- Use of xFitter software
	- start from existing PDF sets
	- Use a  $\chi^2(b_{\text{exp}}, b_{\text{th}})$  that minimises difference between observed and predicted cross-section allowing nuisance parameters ( $b_{\text{exp}}$ ,  $b_{\text{th}}$ ) to shift
	- Allows for quantitative estimate of the agreement between the data and the PDF sets from global fits, and study further constraining power from the new measurement
		- Best  $\chi^2$  is obtained with CT14nnlo, CT10nnlo, and reasonable with MMHT14nnlo



### PDF profiling results (exemple of MMHT14)







## Full PDF fit

- More complex than PDF profiling
	- Requires parametrisation of the PDFs at a starting scale  $Q_0^2 = 1.9 GeV^2$

$$
xu_{v}(x) = A_{u_{v}}x^{B_{u_{v}}}(1-x)^{C_{u_{v}}}(1+E_{u_{v}}x^{2}),
$$
  
\n
$$
xd_{v}(x) = A_{d_{v}}x^{B_{d_{v}}}(1-x)^{C_{d_{v}}},
$$
  
\n
$$
x\bar{u}(x) = A_{\bar{u}}x^{B_{\bar{u}}}(1-x)^{C_{\bar{u}}},
$$
  
\n
$$
x\bar{d}(x) = A_{\bar{d}}x^{B_{\bar{d}}}(1-x)^{C_{\bar{d}}},
$$
  
\n
$$
xg(x) = A_{g}x^{B_{g}}(1-x)^{C_{g}} - A'_{g}x^{B'_{g}}(1-x)^{C'_{g}},
$$
  
\n
$$
x\bar{s}(x) = A_{\bar{s}}x^{B_{\bar{s}}}(1-x)^{C_{\bar{s}}},
$$

- PDFs are evolved to the scale of the measurements and convolved with hardscattering coefficients to obtain the theoretical cross-section predictions
- Then, fit using similar (but different)  $\chi^2$  as in profiling, with parameters left free
- Fit done with HERA and new ATLAS data
	- New set termed ATLAS-epWZ16
	- Includes experimental and theory uncertainties



## Full PDF fit : result

- Theory uncertainties include : variation on the heavy quark masses, on  $Q_0$ , on  $Q_{min}$  (smallest scale for HERA data), different parametrisation,  $\alpha_S(m_Z)$  (= 0.118  $\pm$  0.002), NLO EW, FEWZ vs DYNNLO
- Experimental uncertainties on the new PDF set greatly reduced by a factor 3 w.r.t. the previous one (ATLAS-epWZ12)
- Data well described by the theory





## strange quark density

- ATLAS 2010 W,Z : unsuppressed strangeness at  $x \approx 0.023$  and  $Q^2 = 1.9$  GeV<sup>2</sup>, —> strange, down and up sea quarks of similar strength in that kinematic range
	- Supported by ATLAS W+c measurement
	- Not expected from neutrino scattering experiments, which have big weight in global PDF fits

$$
r_s = \frac{s + \bar{s}}{2\bar{d}}
$$

$$
r_s = 1.19 \pm 0.07
$$
 (exp)  $\frac{+0.13}{-0.14}$  (mod + par + thy)

- Many checks were performed :
	- remove the constraint  $\bar{u} = \bar{d}$  for  $x \rightarrow 0$
	- Inclusion of E866 data for which there are tensions  $(\bar{u} \bar{d})$
	- remove low/high mass Drell-Yan
	- How much the  $\chi^2$  increases when imposing suppressed strangeness ( $r_s = 0.5$  and  $C_{\bar{s}}$  $= C_{\bar{d}}$
- Everything points to strangeness values consistent with this measurement



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## strange quark density

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- Uncertainties include :
	- experimental
	- model variations (heavy quark masses, starting scale, minimum scale of HERA data)
	- **• PDF parametrisation**
	- $\overline{\alpha_S}$
	- EW corrections



- **• factorisation and renormalisation scales**
- **• FEWZ vs DYNNLO**





## |Vcs| measurement

- W production mainly from ud and cs quarks
	- $|V_{ud}|$  already measured with high precision, but not  $|V_{cs}|$
	- W production rate and lepton  $\eta$  distributions are sensitive to  $|V_{cs}|$
	- ->PDF fit with  $|V_{cs}|$  allowed to vary (other CKM matrix elements fixed to the 2012 PDG value)
	- parametrisation variations ->Dominant uncertainty
	- Competitive with other measurements

$$
|V_{cs}| = 0.969 \pm 0.013 \text{ (exp)} \, {}^{+0.006}_{-0.003} \text{ (mod)} \, {}^{+0.003}_{-0.027} \text{ (par)} \, \frac{+0.011}{-0.005} \text{ (thy)}.
$$







## W cross-section at 13 TeV

Physics Letters B 759 (2016) 601–621

- Shorter analysis was performed using 81pb-1 of early Run2 2015 data: W,Z fiducial and inclusive cross-sections
- Peak delivered instantaneous luminosity was L =  $1.7 \times 10^{33}$  cm<sup>-1</sup> s<sup>-1</sup>,  $\langle \mu \rangle = 19$ .
- measurement compared with NNLO fixed-order predictions from DYNNLO using different PDF sets





#### W cross-section : conclusion and summary

- Remarkable precision on Drell-Yan cross-section, below 1% (when excluding luminosity uncertainty) even for W
- Allows for stringent tests of Standard Model
- High number of experimental points is a good input to constrain parton distribution functions and thus reduce their related uncertainty
	- Key ingredient for most physics analyses
- Competitive measurement of  $|V_{cs}|$  CKM matrix element



## W+jets at 8 TeV

JHEP05(2018)077



## W+jets at 8 TeV : introduction

- Several differential cross-sections for  $W + \geq 1$  jet in the electron channel
	- H<sub>T</sub> (scalar sum of the transverse momenta of electron, neutrino and jets),  $p_T(W)$ ,  $p_T(j1)$ , y(j1) (transverse momentum and rapidity of the leading jet, i.e. with highest  $p_T$ )
- And for  $W + \geq 2$  jets
	- $p_T(j2)$ ,  $y(j2)$ ,  $\Delta R(j1,j2)$ ,  $M_{ij}$
- W+/W- cross-section ratio
	- Cancelling of dominant systematic uncertainties
- Motivation : stringent tests of pQCD, constraints to PDFs fits, sensitivity to ME/PS matching schemes
- detector calibration for reconstructed objects (leptons, jets, missing  $E_T$ ) are all from standard Combined Performance groups
- Main signal MC sample is Alpgen+Pythia6 with up to 5 partons in the ME
- All backgrounds estimated with MC except for multi jet (data driven)



## W+jets at 8 TeV : event selection

- Trigger: electron, isolated (p<sub>T</sub>>24 GeV) or not (p<sub>T</sub>>60 GeV)
- exactly 1 electron with  $p_T > 25$  GeV, within detector acceptance and matching the trigger
	- 'tight' identification
	- cut on impact parameters of associated track
	- isolation : ptvarcone30/p $_T$  < 0.07, topoetcone30/p $_T$  < 0.14
	- no other 'medium' electron with p<sub>T</sub>>20 GeV
- Anti-k<sub>T</sub> jets with  $R=0.4$ ,  $p_T>30$  GeV,  $|y|<4.4$ , separated from the electron
	- pile-up rejection for jets within tracking acceptance ('JVF' cut)
	- Veto on events with b-tagged jets ( $p_T > 20$  GeV,  $|\eta| < 2.5$ ) —>rejects tt
- jets-electron overlap removal
- MET  $> 25$  GeV,  $m_T > 40$  GeV



## W+jets at 8 TeV : event yields



• Dominant background is multijet for lower jet multiplicity, tt becomes more important for events with higher multiplicity



#### W+jets at 8 TeV : comparison to various predictions

- All observables use Bayesian unfolding to account for detector effects (bin migrations)
- Comparison to various prediction (first 3 have non-perturbative corrections using Sherpa 2.2.1 to account for hadronisation and underlying event)
- NLO EW corrections investigated with Sherpa 2.2.1, PDF sensitivity investigated with MCFM
- Theory uncertainties include : renormalisation/factorisation scale variations,  $\alpha_{\rm S}$ , PDF uncertainties (only statistical uncertainty is shown for LO generators)





#### Theoretical predictions and the set of the space of the Fiducial phase space



## W+jets cross-section and W+/W- ratio

- Overall good agreement of measured values with predictions
- NLO vs LO improves agreement (Sherpa 2.2.1)
- Better agreement for the ratio
	- $\rightarrow$  probable cancelling of theoretical mismodeling (jet emission)







- High gain in precision from the ratio
- multijet uncertainty dominant at large jet multiplicity (ratio)
- Significant impact of jet energy scale at high multiplicity (dominant for W cross-section)

W cross-section



### W+jets at 8 TeV : differential cross-sections

- Just a few exemples of distributions : W  $p_T$  for W  $+ \ge 1$  jet (W and ratio),  $ΔR, M<sub>ij</sub>$  for W +  $≥$  2 jets
- W  $p_T$ : sensitive to PDFs, interesting for fits
	- best described by Alpgen+Pythia and LO Sherpa 1.4
	- Ratio : most predictions are off
- ΔR, M<sub>ii</sub>: test hard parton radiation at large angles and matrix-element/ parton-shower matching
	- Good description from BlackHat+Sherpa
	- Much better description of large M<sub>ii</sub>/ΔR values from Sherpa 2.2.1 as compared to Sherpa 1.4







## W+jets at 8 TeV : a few conclusions

- High precision reached especially for the W+/W- ratio
- Overall good agreement between the data and the theory predictions
- Degradation at large jet rapidity, angular separation and energy
- Sensitivity of W+/W- ratio to PDFs
- Multi-leg generators (Alpgen, Sherpa) do best in many places
	- High multiplicities in the ME
- No single prediction describes each and every measured observable



 $H_r$  [GeV]



#### W cross-section : conclusion and summary

#### **Standard Model Production Cross Section Measurements** Status: July 2018  $\sigma$  [pb]  $\Delta$  O total (2x) **ATLAS** Preliminary  $10^{11}$ ╶╟┉┙╢╟  $\mathbf{\mathsf{\Gamma}}_{\boldsymbol{\Delta}}\mathbf{\mathsf{\mathsf{O}}}$  inelastic Theory Run 1,2  $\sqrt{s}$  = 7,8,13 TeV o  $10<sup>6</sup>$ incl LHC pp  $\sqrt{s}$  = 7 TeV  $\overline{\Delta}_{\mathbf{O}}$ ю Data  $4.5 - 4.9$  fb<sup>-1</sup> dijets  $10<sup>5</sup>$  $\overline{O}$  $\sum_{25}$  GeV LHC pp  $\sqrt{s} = 8$  TeV  $10<sup>4</sup>$ Data  $20.2 - 20.3$  fb<sup>-1</sup>  $\blacksquare$  $\Delta$ -O LHC pp  $\sqrt{s}$  = 13 TeV  $10<sup>3</sup>$ Ð  $p_T > 125$  GeV  $\Delta$ -O ᅭი Data  $3.2 - 79.8$  fb<sup>-1</sup> О ጟ<sub>Ⴍ</sub>⊡  $\frac{p_T}{\Delta}$  $10<sup>2</sup>$  $\Delta$  o  $\Box$ œ  $\blacktriangle$ A $n_j \geq i$  $10^1$ п  $\blacktriangle$ Δ -∩  $\begin{array}{c} n_j \geq 4 \\ \bigcirc \\ n_j \geq 5 \\ \bigcirc \end{array}$  $n_i \geq 3$  $W<sub>\gamma</sub>$ s-chan  $\mathbf 1$  $\overline{O}$   $\overline{O}$   $tZj$ О  $\blacktriangle$  $\overline{Q}$ **AO** O  $10^{-1}$  $n_j \geq 8$  $\Delta$  $\mathbf{O}$  $\bullet$ o  $\overline{\mathbf{o}}$  $H\rightarrow \gamma\gamma$  $10^{-2}$  $\blacktriangle$ A A A  $\prod_{i=1}^{H}$ ◻  $10^{-3}$ H WV V $\gamma$  ttw ttz tth tty v Y Wjj Zjj WWZ y W y y W z y jjVVjj pp  $\gamma$ W Z  $t\bar{t}$ **VV**  $\gamma\gamma$ **Jets**  $\mathbf t$

tot.

tot.

tot. tot. tot.

EWK EWK Excl.

tot.

**EWK** 



# W mass at 7 TeV

Eur.Phys.J. C78 (2018) no.2, 110



## Analysis strategy

- Measurement's methodology :
	- obtain predictions with simulated events for signal and background (except data-driven multijet background)
	- $\bullet$  to extract the result, compare data and predictions for distributions sensitive to m<sub>w</sub> by performing a template  $\chi$ 2 fit
- Very simple in principle, but extremely challenging in practice as it requires at the 1/10,000 level :
	- Accurate theoretical description of W production and decay kinematics in the simulation
	- Precise calibration of the detector
- Fully reconstructed mass in Z-boson sample to validate the analysis and to provide significant experimental and theoretical constraints (ancillary measurements)





## Measurement's categories



- Measurement performed in 2 channels, using 2 observables, 2 charge categories, 3 (4) |η(lepton)| bins in the electron (muon) channel
	- In total, 28 different values of mw are extracted
	- Allows to :
		- Thoroughly validate the physics modelling
		- benefit from different sensitivities to systematic uncertainties



## **Motivation**

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- See Introduction
- Current world average (Tevatron + LEP) provides the most precise value :
	- $m_W = 80.385 \pm 0.015$  GeV
- The natural goal for the measurement's precision is set by the EW fit prediction's uncertainty (7 MeV)









#### W mass at LHC : more data, larger challenges M mosco at LHC : more data larger chall WUITCOO CULTUS COMPARED V moso ot LLC : more dota lerger challenge  $H_{\text{H}}$  and  $H_{\text{H}}$  and  $H_{\text{H}}$  and  $H_{\text{H}}$





- In proton-proton, W+/W- boson production is asymmetric  $\mathsf{p}$  produce the  $\mathsf{p}$  collisions they are equally  $\mathsf{p}$
- Different contributions from sea/valence quarks e intervals in the same helicity states.<br>• Different contributions from sea/valence quark
- Charge dependence of  $p_T$  spectrum and thus on the  $p_T\ell$  and  $m_T$ die eerste bestervables in die verskap van die<br>Heavy-flavour-initiated processes in die verskap van die verskap van die verskap van die verskap van die versk Further Charge dependence  $\bullet$  Charge dependence of  $p_T$  sp Heavy-flavour-initiated processes
	- More heavy flavour initiated production (25% of the W production is a political M-political M-politi induced by at least one second generation quark s or c) wore neavy navour miniated productivity to the produced by at laser one second nane  $\ddot{a}$ uced by at least one second generation q
	- W+, W- and Z are produced by different light flavour fractions
	- Larger gluon-induced W production Larger gluon-induced W production ger gluon-induced W production
	- Large PDF-induced W-polarisation uncertainty (valence vs sea quarks) **0.2**
	- Strange quark pdf uncertainty —> uncertainty on the relative fraction of charm-initiated W boson  $\longrightarrow$  alter the balance between valence quark  $\begin{array}{ccc} \hline \end{array}$ and sea quark  $k$   $\alpha$   $\alpha$   $\alpha$   $\beta$  most of the most optical is the measurement is the transfer from  $\alpha$  to  $\alpha$   $\beta$ Stefano Camarda 3 (1991) 1995 - Stefano Camarda 3 (1995) 1996 - Stefano Camarda 3 (1996) 1997 - Stefano Camard<br>Stefano Camarda 3 (1996) 1997 - Stefano Camarda 3 (1997) 1997 - Stefano Camarda 3 (1997) 1997 - Stefano Camard ange quark pdf uncertainty  $\longrightarrow$  uncertainty on the relative traction of  $\frac{1}{10^{4}}$   $\frac{1}{10^{3}}$   $\frac{10^{2}}{10^{1}}$   $\frac{1}{10^{1}}$





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### Spectra differences between 'purely sea' and 'standard' quark induced W production

#### Eur. Phys. J. C (2010) 69: 379–397



• Uncertainty on sea and valence PDFs —>on the measured spectra

$$
\sigma_{W^+}(y) \propto u(x_1) \cdot \bar{d}(x_2) + \bar{d}(x_1) \cdot u(x_2)
$$
  

$$
\sigma_{W^-}(y) \propto d(x_1) \cdot \bar{u}(x_2) + \bar{u}(x_1) \cdot d(x_2)
$$



# MODELING ASPECTS



#### **and introduction to the modeling** the data. The e↵ect of virtual photon production and *Z*/⇤ interference is included in both the predictions

• Factorisation of cross-section under 4 terms and the Powheg+Pythia 8 simulated *Z*-boson samples. The reweighting procedure used to include the

spherical harmonics

 $\mathcal{L}$ 

• Approximation checked and valid at 2 MeV level for mw

$$
\frac{d\sigma}{dp_1 dp_2} = \left[\frac{d\sigma(m)}{dm}\right] \left[\frac{d\sigma(y)}{dy}\right] \left[\frac{d\sigma(p_T, y)}{dp_T dy} \left(\frac{d\sigma(y)}{dy}\right)^{-1}\right] \left[(1 + \cos^2\theta) + \sum_{i=0}^{7} A_i(p_T, y) P_i(\cos\theta, \phi)\right]
$$

- dσ(*m*)/d*m* modeled with Breit Wigner and and azimuth of the polar angle and azimuth of the po
- lepton<sup>1</sup> in any given rest frame of the dilepton system; *Ai* are numerical coecients, and *Pi* are spherical harmonics of order zero, one and two. • Other terms : reweight MC according to various predictions
- The div(y//dy . invadiated invariant prediction (D invariant), and in a Breit–Chief-1. dσ(y)/dy : fixed-order NNLO prediction (DYNNLO)
- **2.** p<sub>T</sub> at a given y : Pythia8 with 'AZ' tune **propagator propagator propagator** propagator propag
- $\Omega$  is clear of our  $\Lambda$ : further running energy  $\Lambda$ cussed in Section 6.1. The divided in the dividity of boson rapidity, and the boson rapidity, depending on the 3. polarisation Ai : fixed-order NNLO prediction (DYNNLO)
- (NB : baseline MC is Powheg+Pythia) (NB : baseline MC is Powheg+Pythia)

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## Validity check of the reweighting

cea



Use NNPDF3 prediction as pseudo-data, perform the various reweightings (y,  $p_{T}$ , polarisation) to CT10 sample : strongly validates the modeling procedure

 $\Delta m_W = 1.5 \pm 2.0$  MeV



## **Rapidity**

- Use of DYNNLO (Fixed-order pQCD, NNLO)
- Validate against 7 TeV ATLAS W, Z cross-section measurements
	- Distributions sensitive to PDF effects, critical for validation

Eur. Phys. J. C 77 (2017) 367

• PDF set : CT10nnlo (best agreement), MMHT14nnlo and CT14nnlo used for uncertainties (other global sets disfavoured by the data)






## W/Z polarisation

- Kinematic of the decay leptons depend on polarisation of the (W or  $\overline{Z)}$  boson
	- NLO QCD brings  $p<sub>T</sub>$  to boson, in turn affecting polarisation
- Full 5-dimensional cross-section can be written as :



- $\theta$  and  $\phi$  : angles of the charged lepton (W $\cdot$ , Z) or neutrino (W $\cdot$ ) in the rest frame of the boson
- $A_i(m, y, p)$ : dimensionless angular coefficients (m dependence is small)
	- A $\rightarrow$  0 when p<sub>T</sub>~0 except for A<sub>4</sub> (responsible of forward-backward lepton asymmetry, sensitive to sin<sup>2</sup> $\theta$ w)
	- A<sub>5</sub>-A<sub>7</sub> small, only appear at NNLO in  $\alpha$ s



## **Polarisation**

- Crucial to get right in *pp* collisions, otherwise miss some effects
- ATLAS measurement of Z angular coefficients validates fixedorder pQCD NNLO prediction
	- except for  $A_2$ : additional uncertainty
		- data/prediction difference is added to the uncertainty ; pseudo-experiments show no correlation with other coefficients
	- Uncertainties on the Z measurement are propagated to the W









#### W boson transverse momentum an NLO plus parton shower generator shower generator setup such as Pythia 8, and of resummed predictions at Py<br>Pythia 8, and of resummed predictions at Pythia 8, and of resummed predictions at Pythia 8, and of resummed pr next-to-leading logarithmic order  $\mathsf{V}\mathsf{V}$

#### • Pythia8 tuned on Z pT ATLAS data (AZ tune) verse modern distribution measured with the A<sub>TLAS</sub> detector at a centre-of-mass energy of products energy of p<br>ATLAS detector at a centre-of-mass energy of psychology of psychology of psychology of psychology of psycholog  $\frac{1}{2}$   $\frac{1}{2}$   $\frac{1}{2}$  Data uncertainty



• Good agreement is obtained for the ratio of **For all and the experiments** differential cross-sections using this tune:

$$
R_{W/Z}(p_{\rm T}) = \left(\frac{1}{\sigma_W} \cdot \frac{\mathrm{d}\sigma_W(p_{\rm T})}{\mathrm{d}p_{\rm T}}\right) \left(\frac{1}{\sigma_Z} \cdot \frac{\mathrm{d}\sigma_Z(p_{\rm T})}{\mathrm{d}p_{\rm T}}\right)^{-1}
$$

- $p_T(W)$  is obtained via the product of the predicted  $I_{\text{F}} = \frac{1}{\sqrt{2\pi}}$  $\frac{1}{2}$  and the experimental Z p<sub>T</sub> spectrum in  $\frac{0.95}{0.85}$  and the experimental Z p<sub>T</sub> spectrum
- The total uncertainty being the sum in **participal was measured using the** *p*  $\frac{0.65}{0.65}$ quadrature of these two components,  $\sim$  1-2%  $\frac{1}{10}$   $\frac{1}{20}$   $\frac{1}{30}$   $\frac{1}{30}$







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# Uncertainties to pT(W)

- Only modelling uncertainties which are uncorrelated between Z and W give sizeable uncertainties on the measurement
	- Induced by heavy flavour initiated production : 6/3% of cc/bb for Z, 20% of cs for W production
- Missing higher orders in QCD ISR : factorisation scale  $(\mu_F)$ variations taken as correlated between W and Z for light quark, independently for heavy quarks
- other sources : uncertainty on mc, choice of parton shower LO PDF
- Central prediction and uncertainty well validated with the recoil distribution in the data







# Uncertainty from the PDFs

- Each PDF uncertainty eigenvector from CT10nnlo propagated simultaneously to Ai, rapidity and pT reweighing distributions
	- Overall uncertainty evaluated with the Hessian method
	- Consider MMHT14nnlo and CT14nnlo as alternative
- For  $p_T(W)$ , only relative variations of  $p_T(W)$  and  $p_T(Z)$  are considered, reducing the impact of this uncertainty







## QCD modelling uncertainties

- Biggest uncertainties arise from the PDFs, modelling of  $p_T(W)$  and polarisation
	- Strong anti-correlated effect of PDF on W+ vs W- !
- May be further reduced in the future thanks to:
	- More ancillary measurements : polarisation @13 TeV (higher statistics, uncorrelated systematics), direct W  $p_T$  measurement (requires a low pile-up run at LHC)
	- Simultaneous profiling of the PDFs in the fit





## Electroweak uncertainties

- QED FSR photons implemented with Photos, negligible uncertainty on it
- FSR pair production impact checked with Photos and Sanc, taken as uncertainty
	- can be implemented in the baseline MC for future measurements (Photospp 3.61)
- Combined NLO EW⊗QCD corrections are assessed with Winhac and taken as an uncertainty - they include IFI + pure weak corrections but need to be checked using a realistic  $p_T$  (NLO in QCD)
	- same remark (available recently in Powheg-EW)





# EXPERIMENTAL ASPECTS



## Event selection

- Lepton selection
	- muon :  $p_T > 30$  GeV,  $|\eta| < 2.4$ , track-based isolation
	- electron :  $p_T > 30$  GeV,  $|\eta| < 1.2$  or  $1.8 < |\eta| < 2.4$ , track and calorimeter-based isolation
- Kinematic requirements :
	- Recoil:  $u_T < 30$  GeV
	- $m_T > 60$  GeV,  $p_T$ <sup>miss</sup>  $> 30$  GeV



7.8M events

5.9M events



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## Background fractions





### Electron calibration







• φ modulation due to mechanical deformation under gravity of the calorimeter ('pear-shape') corrected with W and Z events





### Muon calibration



As expected, uncertainties are smaller than for electron



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## Hadronic recoil calibration

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- Several steps of the correction :
	- Correct pile-up activity
	- Correct  $\Sigma E_T$  distribution in  $p_T$  bins
	- residual response and resolution corrections in Z events, extrapolated to W
		- Includes a correction of recoil phi









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### Hadronic recoil calibration



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### Hadronic recoil calibration

• Validation check using Powheg+Herwig6 as pseudo-data







## Multijet background

- Use of a data-driven technique :
	- 2 different fitting regions to extract multijet fraction (FR1, FR2) mainly for uncertainty purposes
		- Templates obtained in an MJ enriched region by reverting isolation
			- EW and top contamination subtracted with MC estimation
		- Scan in lepton isolation variable for MJ shape template building
	- 3 different observables :  $m<sub>T</sub>$ , p<sub>T</sub> $\ell/m<sub>T</sub>$ , p<sub>T</sub>miss
	- Fitted fraction corrected for signal region selection efficiency
	- linear extrapolation to signal region





## Multijet background

- Differences in the extrapolation to signal region are taken as (main) fraction uncertainty
- Shapes of observables obtained by an linear extrapolation from CR to SR using the ratio of different anti-isolated regions
	- Uncertainty dominated by statistics, evaluated by fluctuating the bin contents within stat. uncertainty before extrapolation



- Background fraction (η-dependent)
	- 0.6 1.7 % (e channel)
	- 0.5 0.7 % (mu channel)





## Extraction of mW



## Preliminary cross-check with Z

#### • Several cross-checks :

- $\chi$ <sup>2</sup> template fit to the reconstructed m<sub>II</sub> distribution :  $\Delta m_Z = 1(3) \pm 3(5)$  MeV for the electron (muon) channel
- Treat alternatively the negative or positive lepton as 'invisible' to mimic a neutrino,
	- fit to  $p_T\ell$  and  $m_T$  distributions as in W analysis
- Results consistent with combined LEP value of mz within experimental uncertainties





Difference between measured Z boson mass in ATLAS and the combined LEP result



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### Z ee plots after all corrections

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### Z mumu plots after all corrections

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# Fitting range stability

- Fitting range of the distributions optimised in the simulation
- Check the stability of the full combination when varying either  $m<sub>T</sub>$  or  $p<sub>T</sub>$ *l* fitting range
	- Check that the result on the difference with respect to the central value is within  $\sim$  1-2 standard deviations (fully uncorrelated uncertainty)





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## Distributions - hadronic recoil

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## Distributions - lepton eta

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## mW extraction

- $\chi$ 2 template fit to the data in each category (distribution, charge, lepton channel, η*l* bin)
- All categories give consistent result —> strength of detector calibration and physics modelling
- Also cross-check in more categories (bins of recoil, mu…)
- Several combinations performed, using BLUE method











# CONCLUSION AND SUMMARY



## ΔmW result

- Measurement of mw<sup>+</sup> mw<sup>-</sup>
	- Not blinded
	- Many uncertainties cancel as they are not charge-dependent
	- Dominant uncertainty from the PDFs
	- Result compatible with 0 within  $\sim$  1 sigma









urement of the mass di↵erence between the *<sup>W</sup>*<sup>+</sup> and *<sup>W</sup>* bosons yields *mW*<sup>+</sup> *mW* <sup>=</sup> 29 ± 29 MeV.<br>29 **29 PM** • Dominant uncertainty comes from the physics modelling

• Largest contribution from QCD



## Comparison to other measurements

- Same precision as single best measurement from CDF
- Pulling the mw towards a value close to Standard Model prediction from EW fit







## What's next?

- What can be done to improve the precision in the coming  $($  <10  $)$ years ?
- **• More progress on theory side : resummation, incorporation of NLO EW**⊗**QCD effects in simulation**
- Experimental innovations : e.g. pile-up mitigation techniques, more and more **ancillary measurements like W p**<sub>T</sub>, polarisation...
	- Run at low pile-up (mu~1) very interesting in this respect
- Combinations with existing/future measurements (*e.g* Tevatron)



### Low mu runs : W pt

- Strong motivation to measure W  $p_T$  with an uncertainty below 1% in the Sudakov region ( $p_T \sim$  20 GeV), with a resolution of ~5 GeV to catch the variation of the Sudakov peak
	- crucial input to  $m_W$  : with comparable precision, does not need to rely on Z anymore (—>no extrapolation uncertainty)
	- stringent test of resummed predictions and parton showers
- Pile-up degrades the reconstructed  $p_T(W)$  resolution (== hadronic recoil) as  $\sqrt{(} < \mu >)$ 
	- $\rightarrow$  need very low pile-up data taking
		- ~300 pb<sup>-1</sup> at 13 TeV needed for enough statistical precision
			- $\cdot$  ~150 pb $^{-1}$  taken in 2017, ~200 taken last month (mu~2) : stay tune!
			- 250 pb<sup>-1</sup> at 5 TeV (2017) : allows to check energy dependence of the models for  $p_T(W)$





# And now for something completely different…

Eur. Phys. J. C (2017) 77:474



### Electroweak *W j j* production and constraints on anomalous gauge couplings (aGC) with 7 and 8 TeV data

- Opportunity to probe anomalous (triple) gauge couplings through VBF topology
	- probe for new fundamental interactions
	- First VBF W observation
- Largest background from strong Wij production (~O(10) times larger than EW)
- Signal :
	- Large interval in jet rapidities (rapidity gap) —>large invariant dijet mass M<sub>ii</sub>
	- *W* boson in rapidity gap
	- Little hadronic activity in the rapidity gap due to the absence of colour flow between the interacting partons 105





Leading order diagrams for Wjj strong production

Eur. Phys. J. C (2017) 77:474



## Modeling

- Wij MC : Powheg+Pythia8, using MiNLO for strong production (sets the QCD emission scales), PDF set is CT10
	- strong production
- Sherpa 1.4 at leading order is used to estimate interference between strong and EW (using QCD+EW vs separate QCD and EW Wjj samples)
	- Also used in cross-section comparisons
- Other backgrounds are modelled with MC except for multi jet (data-driven)

### Cross-section measurement

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- Binned likelihood fit to the M<sub>ii</sub> distribution with Gaussian constraints to every background
- Determination of  $\mu_i$ 
	- $\bullet$  i = QCD or EW
	- N<sub>i</sub>: measured number of (background subtracted) events
	- $C_i$ : N(reco)/N(truth) passing the selection
	- A<sub>i</sub>: acceptance of fiducial volume

$$
(\sigma_i^{\ell \nu j j} \times A_i)^{\text{meas}} = \mu_i (\sigma_i^{\ell \nu j j} \times A_i)^{\text{theo}} = \frac{N_i}{C_i \mathcal{L}},
$$

fiducial cross-section

normalisation factor  $(=1$  in SM)



## Event selection

- Using single lepton trigger
- Leptons and jets reconstruction/calibration/preselection similar as the W+jets analysis presented before
- Use of lepton and jet centrality defined as

 $C_{\ell(jet)} = |$  $y_{\ell(jet)} - \frac{y_1 + y_2}{2}$ *y*<sup>1</sup> − *y*<sup>2</sup>  $\vert$ 

• Such that  $C < C$  max = 0.4 implies y to be in the range :

$$
\left[\frac{y_1 + y_2}{2} - C_{max} \times |y_1 - y_2|, \frac{y_1 + y_2}{2} + C_{max} \times |y_1 - y_2|\right]
$$



- Forward lepton CR used to constrain strong Wjj production
- Validation region (for multi jet and QCD Wjj) has >=1 central jet



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## Event yields

- EW signal represents  $\sim$  15 % of the total number of events
- Dominant background is strong Wjj production (~50% of total number of events)




### SM Cross-section result







### Unfolded differential cross-sections : Wjj EW vs QCD

- Many unfolded differential cross-sections obtained in the fiducial region
	- Direct probes of new physics
	- Here, just one exemple : for M<sub>ii</sub>> 500GeV (left), dominated by QCD; EW Wjj appears for  $M_{ii} > 2TeV$  (right) !





### Constraints on aGC

- Effective Lagrangian for WWV coupling with operators up to dimension 6 (V=Z, gamma)
- Deviation from SM :  $g_1^V \neq 1$ ,  $\kappa_V \neq 1$ ,  $\lambda_V \neq 0$ ,  $\tilde{\kappa}_V \neq 0$

$$
i\mathcal{L}_{\text{eff}}^{WWV} = g_{WWV} \left\{ \left[ g_1^V V^\mu (W^-_{\mu\nu} W^{+\nu} - W^+_{\mu\nu} W^{-\nu}) \right. \right. \\ + \kappa_V W^+_{\mu} W^-_{\nu} V^{\mu\nu} + \frac{\lambda_V}{m_W^2} V^{\mu\nu} W^+_{\nu} \rho W^-_{\rho\mu} \right\} \left. \left. \left( -\left[ \frac{\tilde{\kappa}_V}{2} W^-_{\mu} W^+_{\nu} \epsilon^{\mu\nu\rho\sigma} V_{\rho\sigma} \right. \right. \right. \\ \left. \left. \left( -\frac{\tilde{\kappa}_V}{2} W^-_{\mu} W^+_{\nu} \epsilon^{\mu\nu\rho\sigma} V_{\rho\sigma} \right) \right\} \right\} \right\}
$$
  
\n
$$
W^{\pm}_{\mu\nu} = \partial_{\mu} W^{\pm}_{\nu} - \partial_{\nu} W^{\pm}_{\mu}.
$$

Gauge invariance requires :

$$
\Delta g_1^Z = \Delta \kappa_Z + \Delta \kappa_\gamma \tan^2 \theta_W, \quad \lambda_\gamma = \lambda_Z \equiv \lambda_V, \quad g_1^\gamma = 1,
$$
  

$$
\tilde{\kappa}_\gamma = -\tilde{\kappa}_Z \cot^2 \theta_W, \quad \text{and} \quad \tilde{\lambda}_\gamma = \tilde{\lambda}_Z \equiv \tilde{\lambda}_V.
$$

To preserve unitarity one can introduce :  $\alpha(q^2) = \frac{\alpha}{(1 + q^2/\Lambda^2)^2}$ 

- With  $\Lambda = 4$  TeV (preserves unitarity) for all parameters in sensitivity range
- Or in EFT : assume perturbative coupling coefficients ci and scale of new interaction  $\Lambda$ ,  $O_i$ are dimension 6 field operators :
- One can relate the c<sub>i</sub> to the aTGC parameters

$$
\mathcal{L}_{\text{EFT}} = \sum_{i} \frac{c_i}{\Lambda^2} O_i
$$



### Constraints on aGC

- Effect of anomalous couplings modelled within Sherpa
- Fit using the yield in signal region with  $M_{ii} > 1$  TeV,  $p_T$ of leading jet > 600 GeV
	- last requirement maximises sensitivity to aTGC
	- Fit each aTGC parameter, fixing the other parameters to SM value



 $\stackrel{>}{\sim}$ 

 $0.2$ 

**ATLAS** 







## EW Wjj : summary

- First observation (>5 sigma) of EW Wij
	- systematic uncertainty dominates the 8 TeV measurement
- Many unfolded differential cross-sections provided
	- Allowing to check theoretical models (high order calculations)
- aTGC constraint
	- $\bullet$   $\lambda$ <sub>V</sub> intervals competitive with those from WW
	- No deviation from SM is observed



### Concluding remarks

- Wide range of physics 'touched' using Ws at LHC
	- I just talked about a few selected items (linked with the most recent publications, mainly) …
		- Precise tests of SM
		- Key input to PDFs
		- validation of predictions that are essential for signal/background modelling in other analyses
		- W mass measurement and global EW fit
		- Wij electroweak direct probe of new physics
	- Many other topics
		- W+charm
		- new physics searches at high  $m<sub>T</sub>$
		- $\bullet$  . . .



# Thank you for your attention!!





## BACKUP



### PDF profiling : methodology

- PDF profiling and full PDF fits heavily rely on methods used in DIS experiments (esp. HERA)
- Software used is xFitter, that makes use of APPLGRID for theory predictions, together with k-factors from the accurate theory tools described before (NNLO QCD from DYNNLO 1.5, NLO EW from PHOTOS and MCSANC)
- Profiling : use of xFitter software
	- start from existing PDF set
	- Use a  $\chi^2(b_{\text{exp}}, b_{\text{th}})$  that minimises difference between observed and predicted crosssection allowing nuisance parameters ( $b_{exp}$ ,  $b_{th}$ ) to shift
	- New PDF  $f_0$ ' is :

$$
f_0' = f_0 + \sum_{k} \left[ b_{k,th}^{\min} \left( \frac{f_k^+ - f_k^-}{2} \right) + (b_{k,th}^{\min})^2 \left( \frac{f_k^+ + f_k^- - 2f_0}{2} \right)^2 \right],
$$

- $f_0$  is the central PDF set and  $f_k$ + and  $f_k$  the corresponding up/down eigenvectors
- Allows for quantitative estimate of the agreement between the data and the PDF sets from global fits, and study further constraining power from the new measurement
	- Best  $\chi^2$  is obtained with CT14nnlo, CT10nnlo, and reasonable with MMHT14nnlo 117



### PDF profiling : results



- 'correlated chi2' is the contribution from the penalty term
- left of '|' is including PDF set uncertainties, right is excluding PDF set uncertainties



## Full PDF fit

- More complex than PDF profiling
	- Requires parametrisation of the PDFs at a starting scale  $Q_0^2 = 1.9$ GeV<sup>2</sup>

$$
x u_{v}(x) = A_{u_{v}} x^{B_{u_{v}}}(1 - x)^{C_{u_{v}}}(1 + E_{u_{v}} x^{2}),
$$
  
\n
$$
x d_{v}(x) = A_{d_{v}} x^{B_{d_{v}}}(1 - x)^{C_{d_{v}}},
$$
  
\n
$$
x \bar{u}(x) = A_{\bar{u}} x^{B_{\bar{u}}}(1 - x)^{C_{\bar{u}}},
$$
  
\n
$$
x \bar{d}(x) = A_{\bar{d}} x^{B_{\bar{d}}}(1 - x)^{C_{\bar{d}}},
$$
  
\n
$$
x g(x) = A_{g} x^{B_{g}}(1 - x)^{C_{g}} - A'_{g} x^{B'_{g}}(1 - x)^{C'_{g}},
$$
  
\n
$$
x \bar{s}(x) = A_{\bar{s}} x^{B_{\bar{s}}}(1 - x)^{C_{\bar{s}}},
$$

- Parameter scans requiring  $\chi^2$  saturation (no  $\chi^2$  decrease when adding free parameters)
- Some constraints by sum rules
	- $\rightarrow$  15 parameters
- $A_{\bar{s}}$  and  $B_{\bar{s}}$  appear as free parameters
- assume  $\bar{s} = s$
- PDFs are evolved to the scale of the measurements and convolved with hard-scattering coefficients to obtain the theoretical cross- section predictions
	- Done with a variable flavour number scheme that switches on c- and b-quark PDFs at  $m_c$  and  $m_b$
- Then, fit using similar (but different)  $\chi^2$  as in profiling, with parameters left free
- Fit done with HERA and new ATLAS data
	- New set termed ATLAS-epWZ16
	- Includes experimental and theory uncertainties



# Comparison CDF vs ATLAS

#### **Similar PDF uncertainties**

#### $p_{\tau}$  W uncertainties are larger for  $p_{\tau}$  lepton

#### than  $m<sub>z</sub>$  at CDF, but similar in ATLAS



#### Includes also Ai uncertainties





## Comparison D0 vs ATLAS



#### Includes also Ai uncertainties





## MSSM fit



#### **MSSM** band:

scan over **SUSY masses** 

#### overlap:

SM is MSSM-like MSSM is SM-like

#### SM band:

variation of  $M_H^{\text{SM}}$ 



## Comparison CDF vs ATLAS

#### **Similar PDF uncertainties**

courtesy of M. Boonekamp

#### $p_{\tau}$  W uncertainties are larger for  $p_{\tau}$  lepton

#### than m<sub>-</sub> at CDF, but similar in ATLAS









# Comparison D0 vs ATLAS



#### Includes also Ai uncertainties



## Difference Tevatron vs LHC

cea



## Difference Tevatron vs LHC

cea





### Spectra differences between 'purely sea' and 'standard' quark induced W production

#### Eur. Phys. J. C (2010) 69: 379–397



• Uncertainty on sea and valence PDFs —>uncertainty on W helicity —> on the measured spectra

$$
\sigma_{W^+}(y) \propto u(x_1) \cdot \bar{d}(x_2) + \bar{d}(x_1) \cdot u(x_2)
$$
  

$$
\sigma_{W^-}(y) \propto d(x_1) \cdot \bar{u}(x_2) + \bar{u}(x_1) \cdot d(x_2)
$$



# Full uncertainty table





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### Weights of all categories



# p<sub>T</sub> modeling strategy

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### observables



### various predictions of the 2<br>Recoil distributions without modelling corrections





## p<sub>T</sub> modeling strategy







- Very different prediction of  $p_T(W)/p_T(Z)$  ratio from resummed technique or Powheg MiNLO with respect to Pythia 8 AZ
	- This is a high topic of interest and subject to many discussions in the LHC EW working group
- Pythia8 AZ is validated by the data  $(u_{//})$  contrary to other predictions
- Negligible impact of the parton shower model (Herwig 7)





### Low mu runs

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• Dependence of uncertainty (statistical + recoil calibration systematic) in first pT(W) bin (0-5 GeV) vs mu for 300 pb-1 integrated luminaosity

Dependence of SET with sqrts

