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# SOME PHYSICS RESULTS & BOSE-EINSTEIN CORRELATIONS STUDY AT 13 TEV WITH ATLAS

BEC@13: ATL-COM-PHYS-2016-1621; ATL-COM-PHYS-2018-044

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# A TOROIDAL LHC APPARATUS (ATLAS)

## Muon spectrometer

( $\mu$  Trigger/tracking and Toroid Magnets)

**Precision Tracking:**

- **MDT** (Monitored Drift Tubes)
- **CSC** (Cathode Strip Chambers)  $|\eta| > 2.4$

**Trigger:**

- **RPC** (Resistive Plate Chamber) barrel
- **TGC** (Thin Gas Chamber) endcap

## Inner Detector (ID)

Tracking; 2T Solenoid Magnet

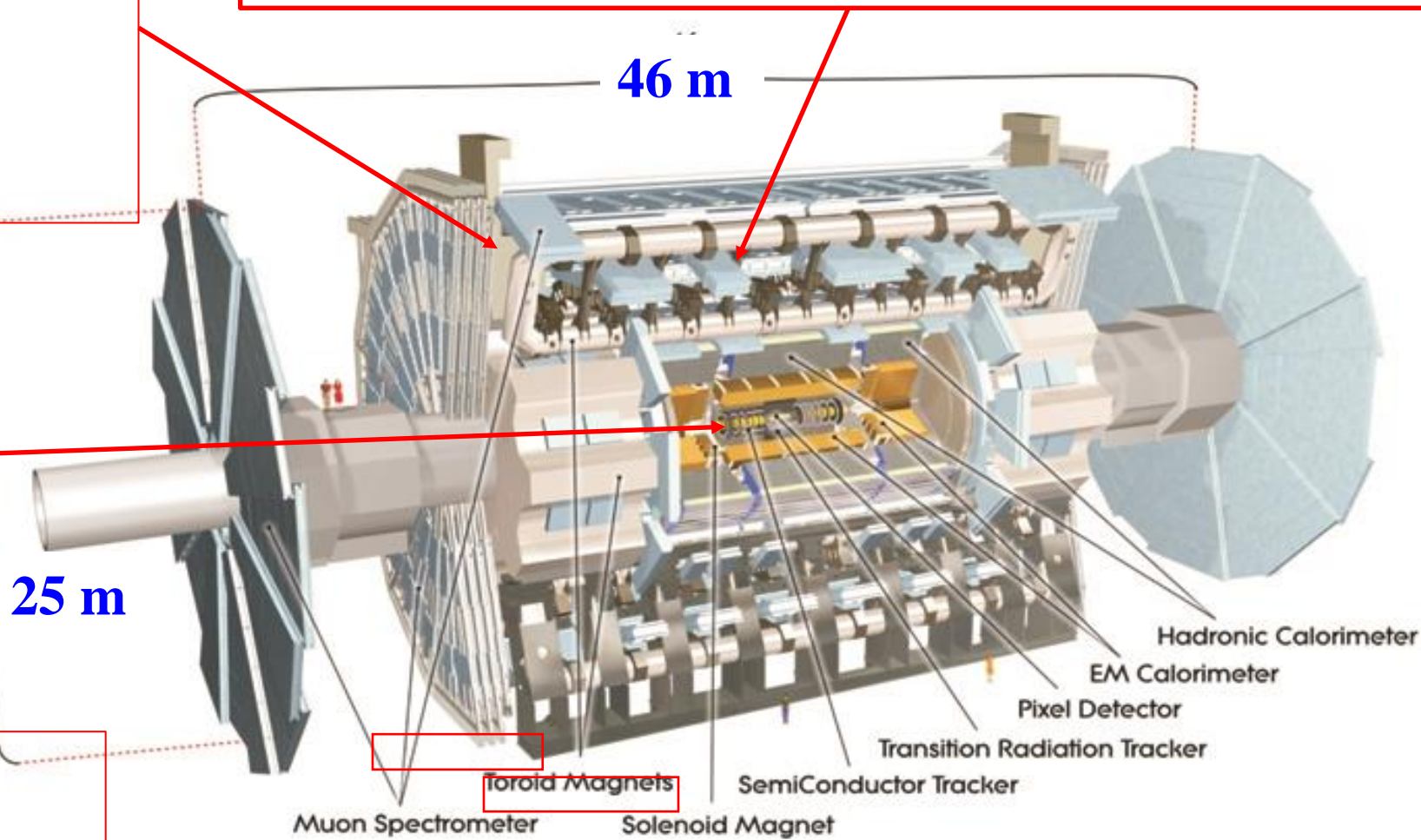
- **Silicon Pixels**  $50 \times 400 \mu\text{m}^2$
- **Silicon Strips** (SCT)  $40 \mu\text{m}$  rad stereo strips
- **Transition Radiation Tracker (TRT)** up to 36 points/track

## Two Level Trigger system

- **L1** – hardware: **100 kHz**, **2.5  $\mu\text{s}$**  latency
- **HLT** – farm: merge the former **L2** and **Event Filter** **1.5 kHz**, **0.2 s** latency

## Calorimeter: EM and Hadronic energy

- **Liquid Ar (LAr)** EM barrel and End-cap & Hadronic End-cap
- **Tile** calorimeter (Fe-scintillator) Hadronic barrel

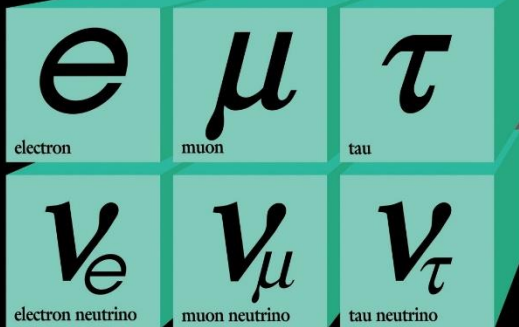
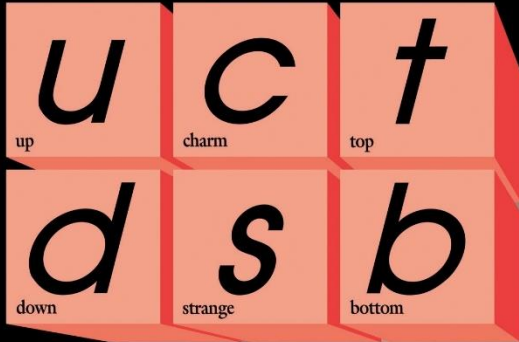


**7 000 tons**

# STANDARD MODEL

Fermions: spin = 1/2 particles

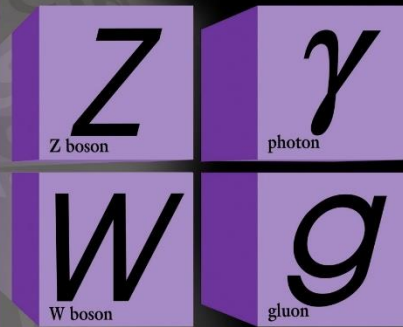
## Quarks



## Leptons

Vector Bosons: spin = 1 particles

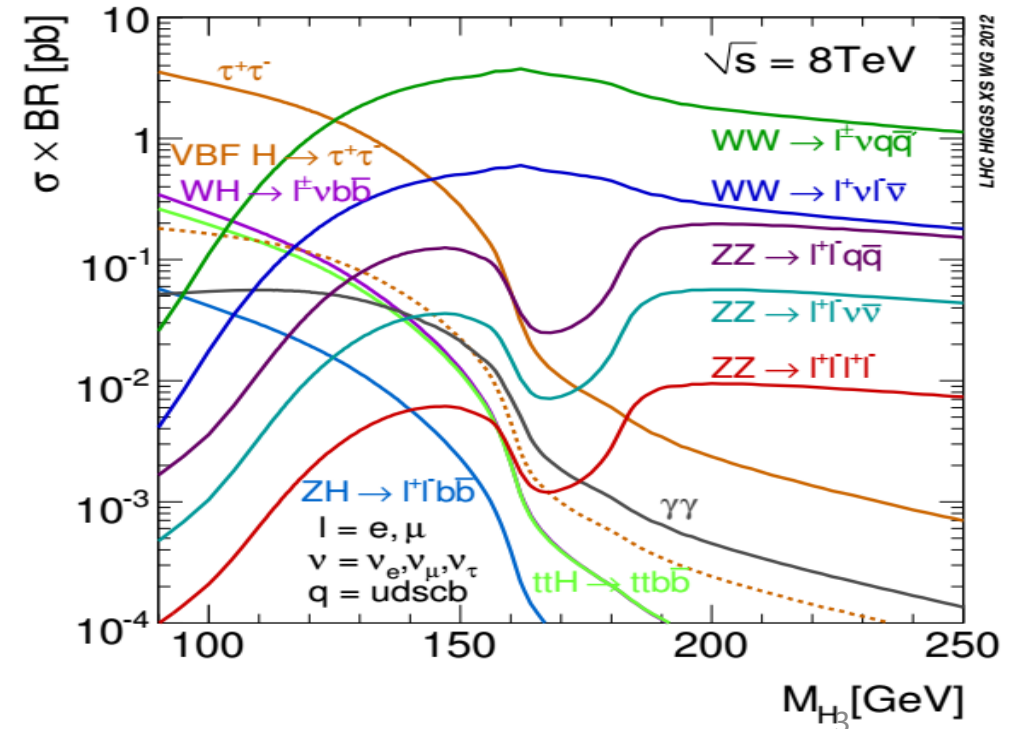
## Forces



Higgs Boson:  
spin = 0  
fundamental  
scalar particle

**Mass of the Higgs boson  
was not predicted!**

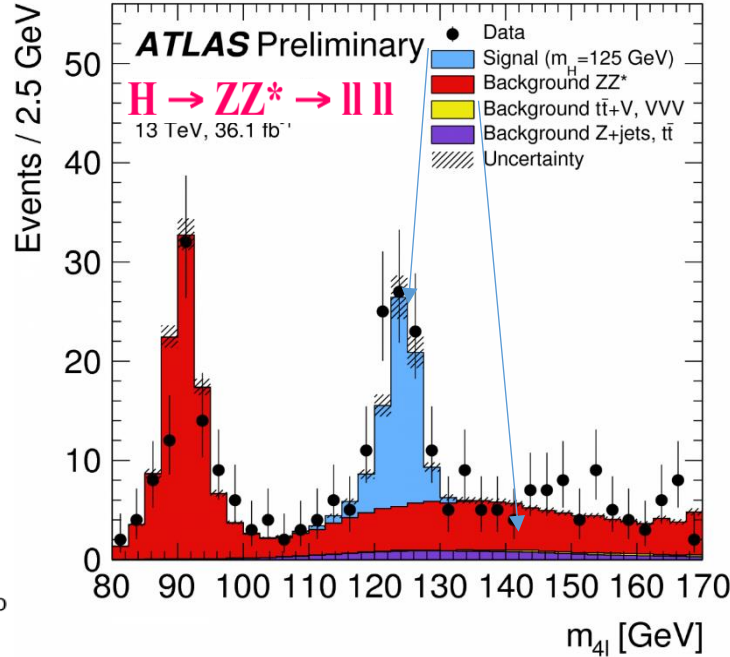
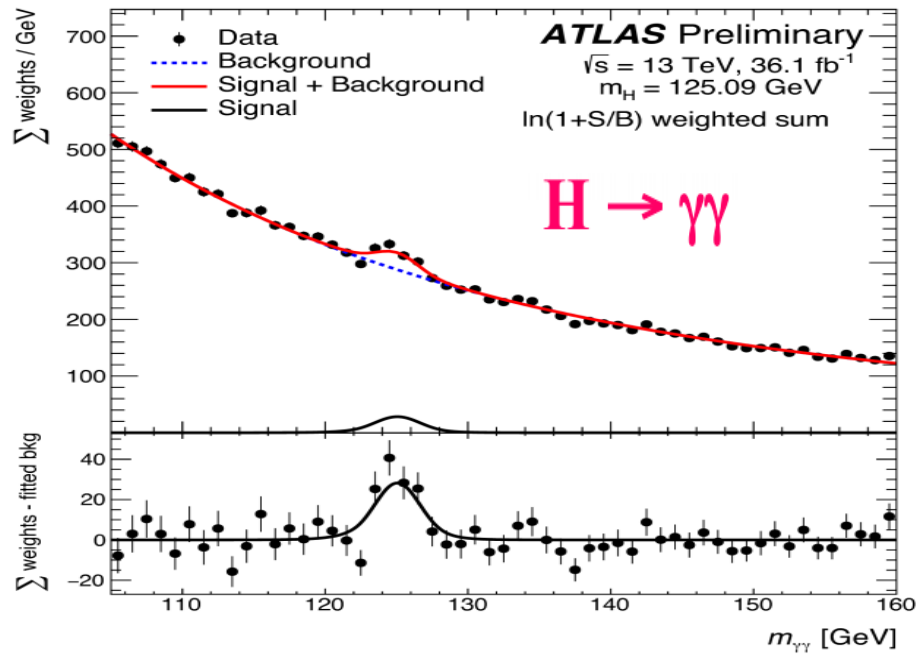
Higgs Decay Channels prediction for  
Higgs boson mass from 90 to 250 GeV





# DISCOVERY OF HIGGS BOSON

arxiv:1708.02810  
 ATLAS-CONF-2017-045  
 ATLAS-CONF-2017-047

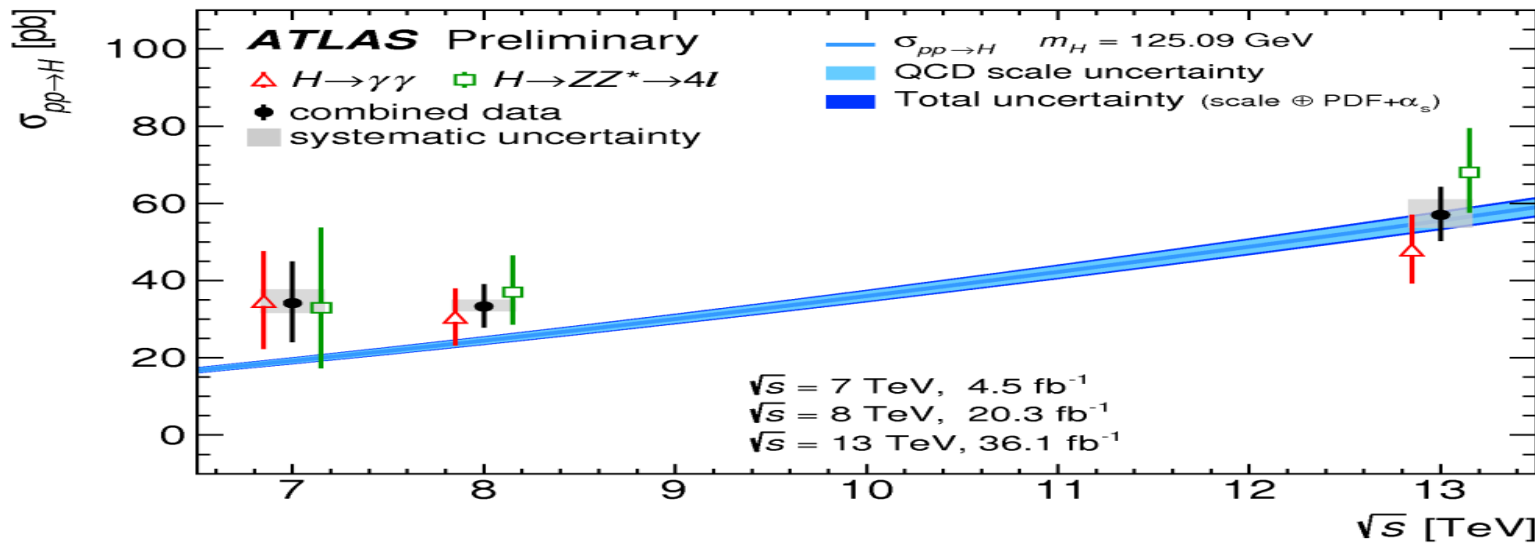


New results in the

- $H \rightarrow ZZ^* \rightarrow 4l$  &
- $H \rightarrow \gamma\gamma$  channels.
- Combined measurements of fiducial and total production cross sections (assuming SM branching ratios)

$m_H = 125.09 \pm 0.24 \text{ GeV}$

Combined global signal strength compatible with the Standard Model:

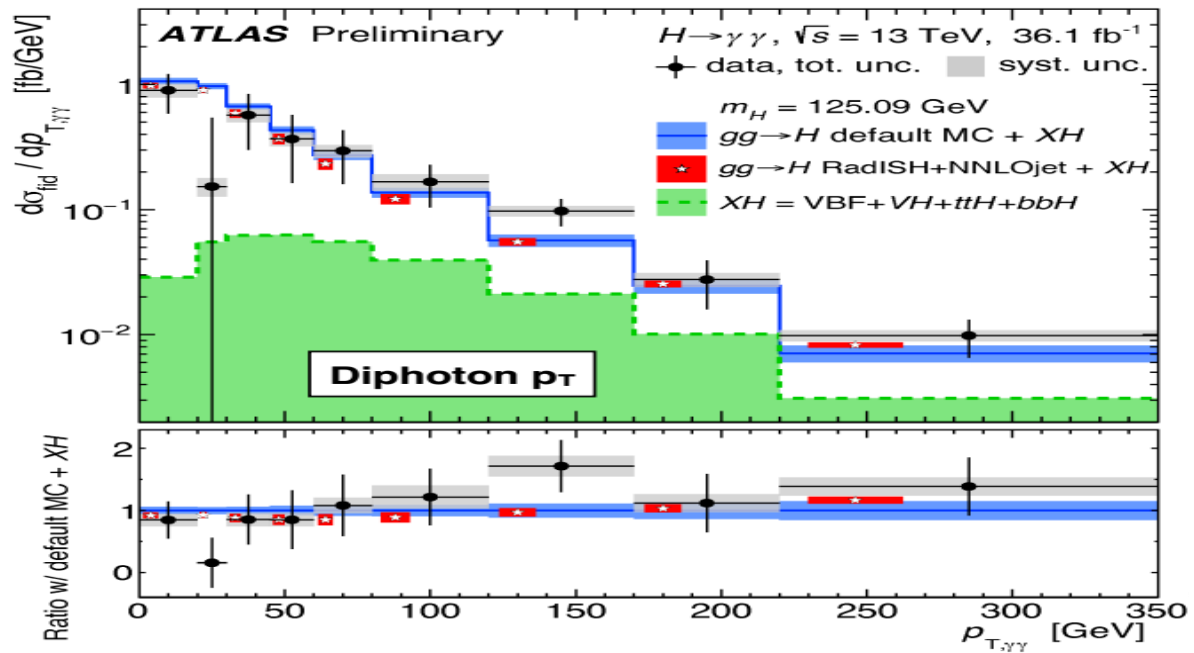


$\mu = 1.09 \pm 0.12$   
 $= 1.09 \pm 0.09 \text{ (stat.) } {}^{+0.06}_{-0.05} \text{ (exp.) } {}^{+0.06}_{-0.05} \text{ (th.)}$

Total  $pp \rightarrow H+X$  cross sections measured at  $\sqrt{s} = 7, 8, \text{ and } 13 \text{ TeV}$ , compared to SM predictions

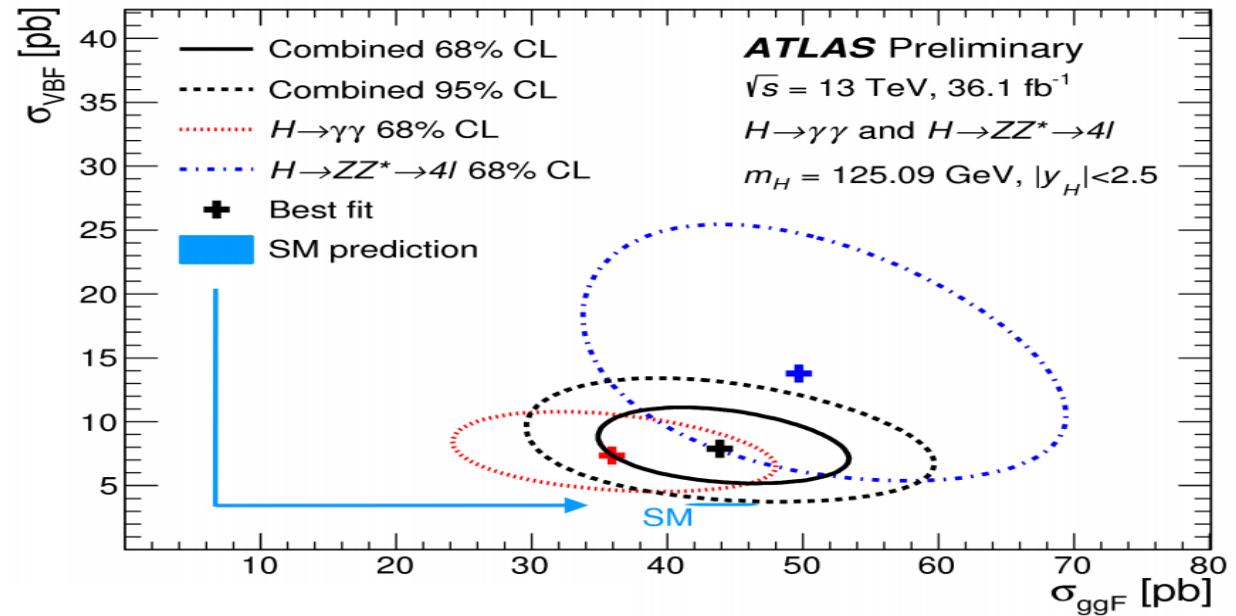


- The increased collision energy and large integrated luminosity provides the opportunity to measure Higgs differential cross sections
- Combinations of measurements in different kinematical regions can be used to probe different production modes, notably *gluon-gluon fusion (ggF)*, *vector boson fusion (VBF)*
- Results also interpreted using new simplified template cross section framework, reducing theoretical systematic uncertainties



The differential cross section for  $pp \rightarrow H \rightarrow \gamma\gamma$  as a function of  $p_T^{\gamma\gamma}$  and compared to SM expectation

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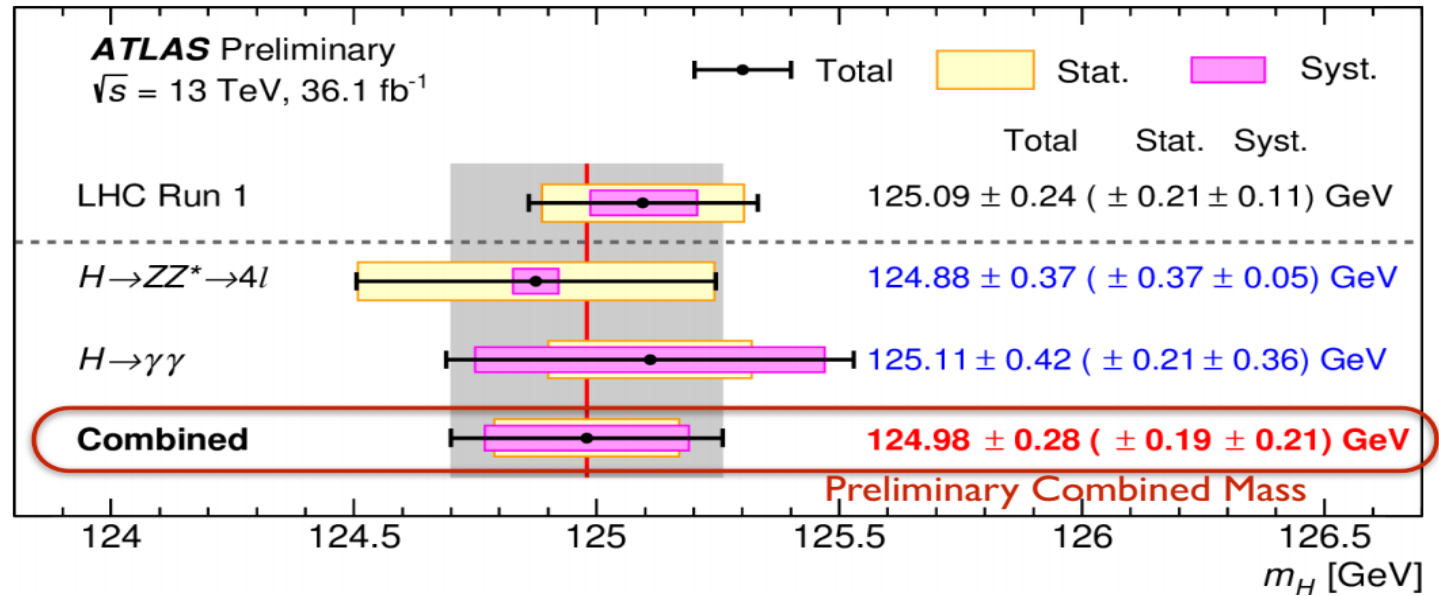
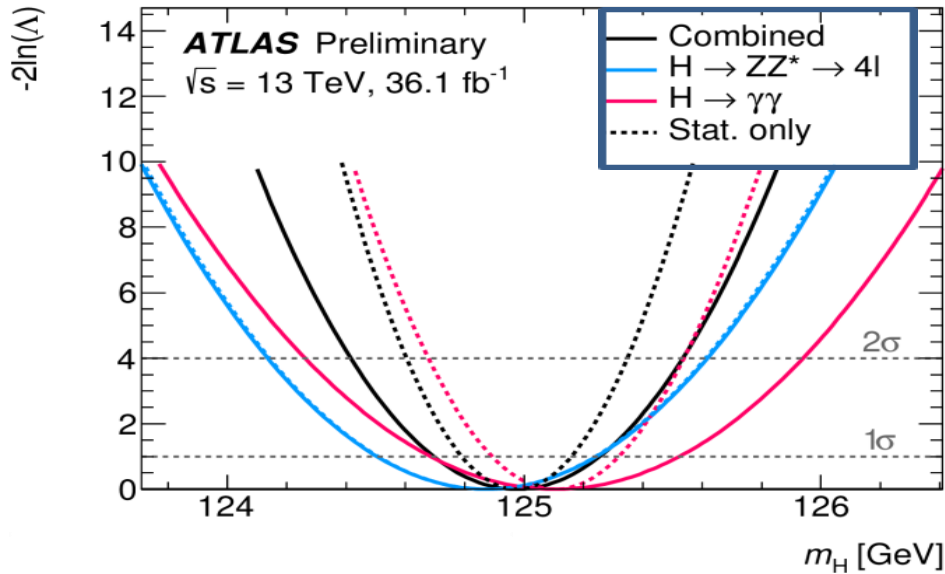
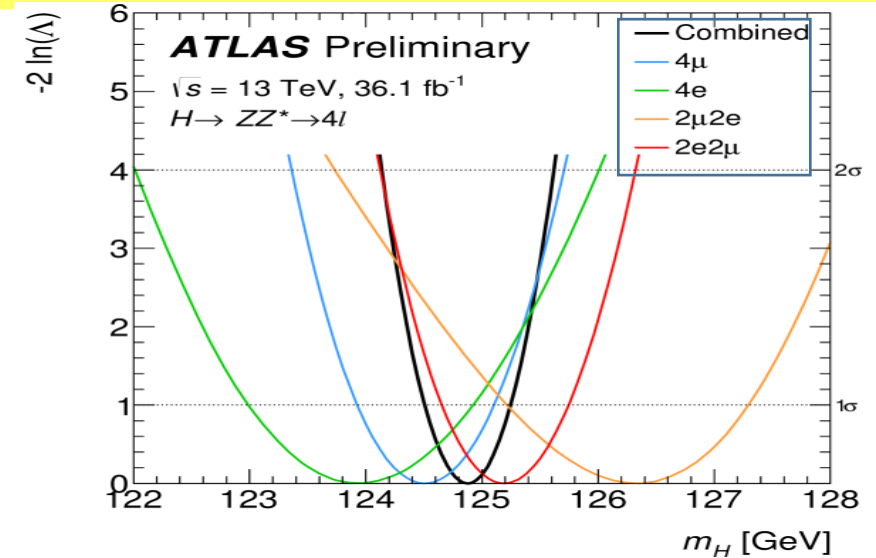


Likelihood contours in  $\sigma_{\text{VBF}}$  vs.  $\sigma_{\text{ggF}}$  plane in  $H \rightarrow \gamma\gamma$  (red) and  $H \rightarrow ZZ^* \rightarrow 4l$  (blue), their combination (black), SM prediction (light blue)

□ Higgs boson mass measurements in the  $H \rightarrow ZZ^* \rightarrow 4\ell$  &  $H \rightarrow \gamma\gamma$  channels complementary

- $4\ell$  channel the stat uncertainty dominates
- $\gamma\gamma$  channel dominated by systematic uncertainties (most notably the  $\gamma$  energy scale calibration)

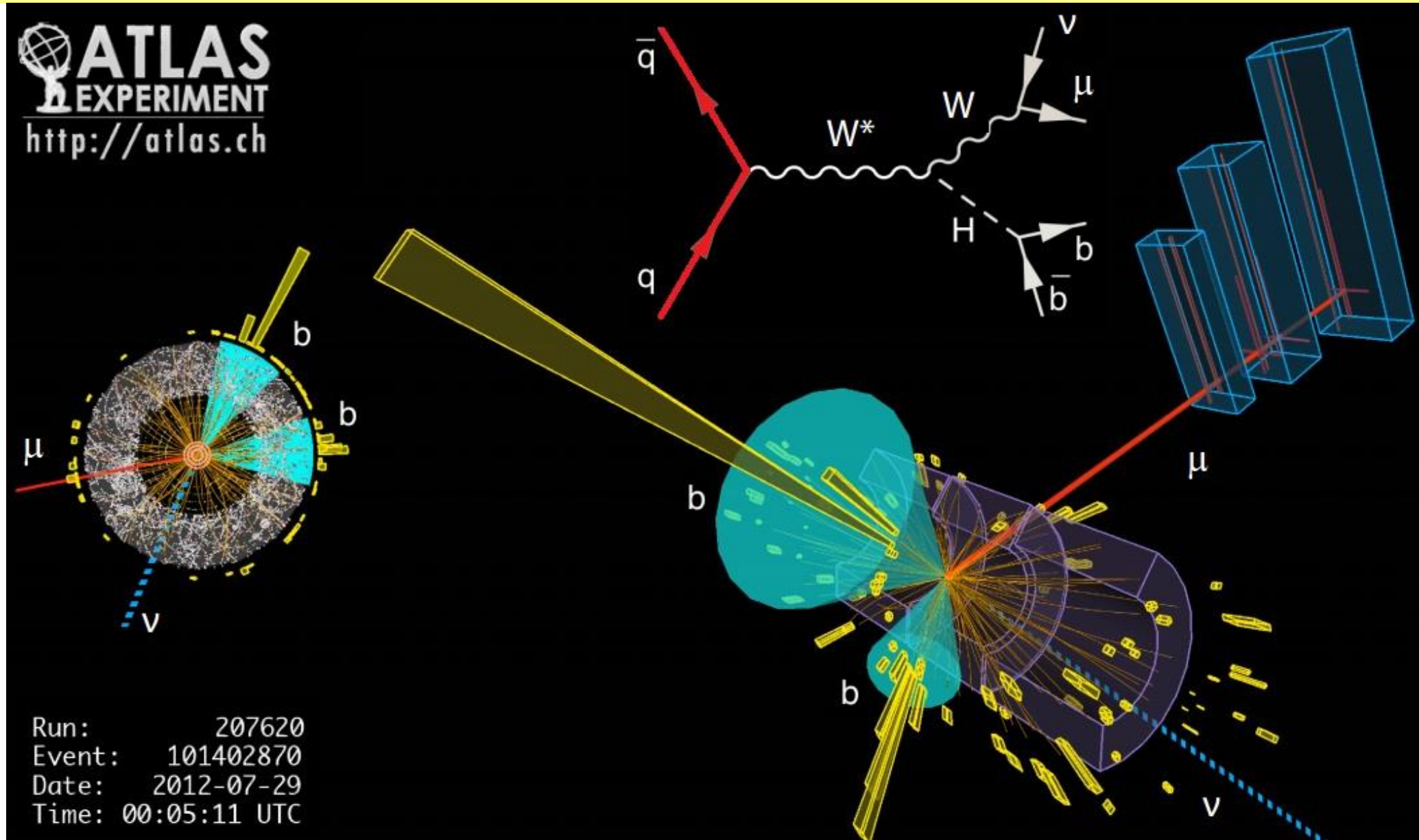
□ Measurements consistent between sub-channels, and consistent with the Run 1 combined result



The value of  $-2\ln\Lambda$  as a function of  $m_H$  for the individual channels & their combination

The Higgs boson mass measurements from the individual and combined analyses, compared to the combined *Run 1* measurement by ATLAS & CMS

# HIGGS BOSON PRODUCTION: $QQ \rightarrow W^* \rightarrow WH, W \rightarrow \nu\mu, H \rightarrow b\bar{b}$

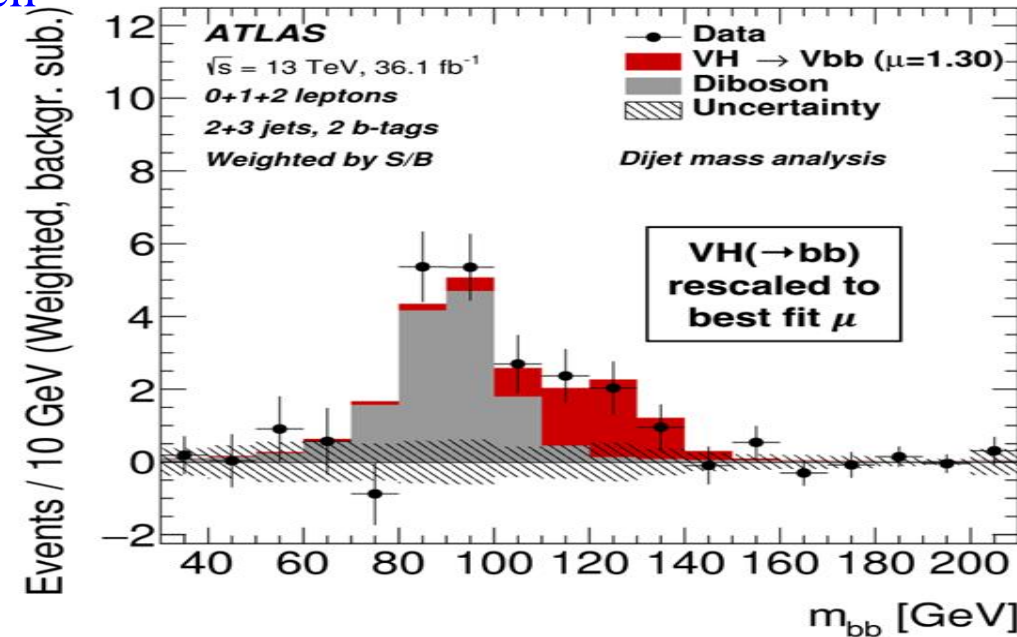
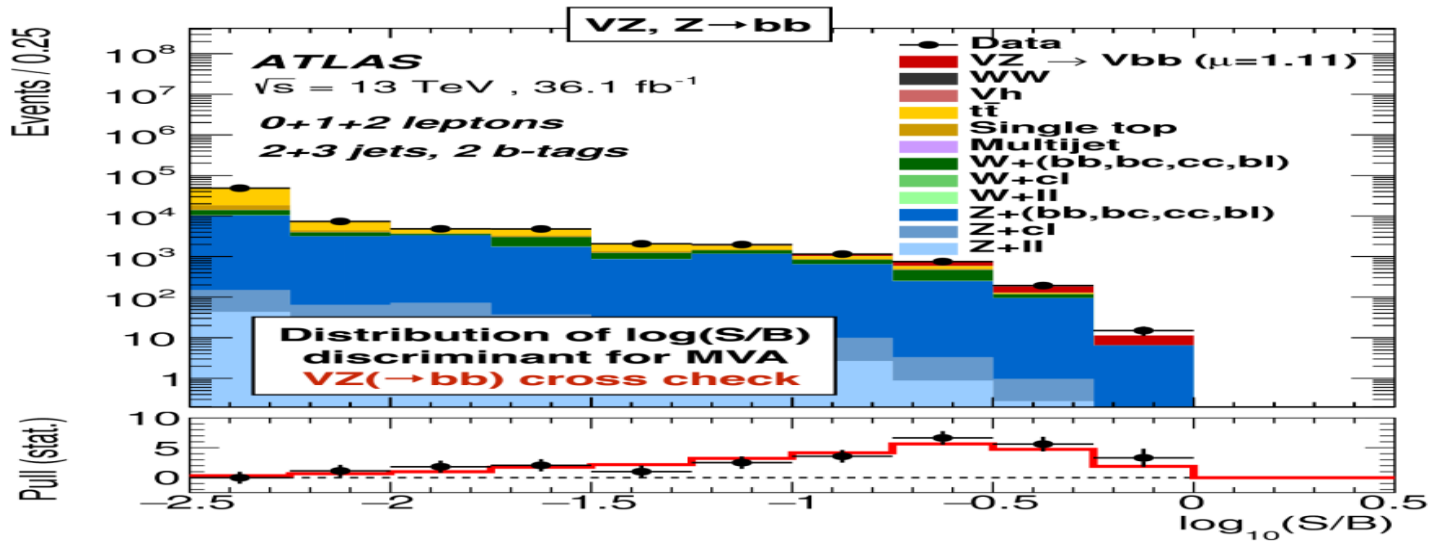
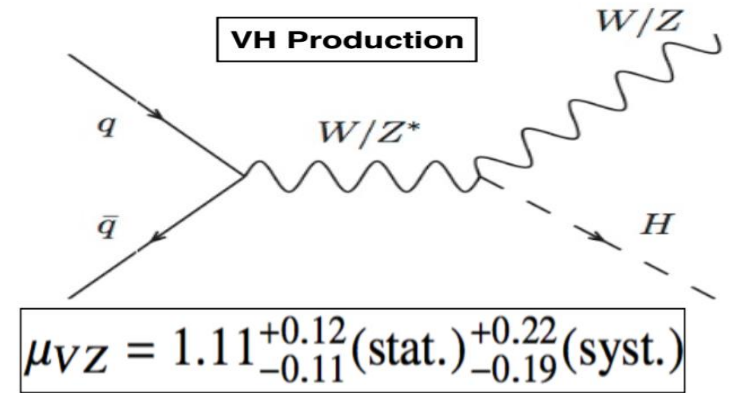




# SEARCHES FOR HIGGS BOSON DECAYS TO $\bar{b}b$

arxiv:1708.03299

- ❑ Most sensitive channel to look for the decay  $H \rightarrow \bar{b}b$  is associated production,  $VH$  ( $V=W/Z$ ) with  $H \rightarrow \bar{b}b$
- ❑ Largest Higgs branching ratio  $BR(H \rightarrow \bar{b}b) \approx 58\%$
- ❑ ATLAS analysis combines **Z & W** final states, channels characterized by lepton multiplicity:
  - 2 lepton ( $Z \rightarrow \ell\ell$ ), 1-lepton ( $W \rightarrow \ell\nu$ ), 0-lepton ( $Z \rightarrow \nu\nu$ )
- ❑ Validation of performance and systematics understanding of the Boosted Decision Trees (BDTs) analysis from an independent search for  $VZ$ ,  $Z \rightarrow \bar{b}b$
- ❑ Observed (expected) significance:  **$5.8\sigma$  ( $5.3\sigma$ )**



Event yields as a function of  $\log_{10}(S/B)$  for data, background and a Higgs boson signal with  $m_H=125 \text{ GeV}$

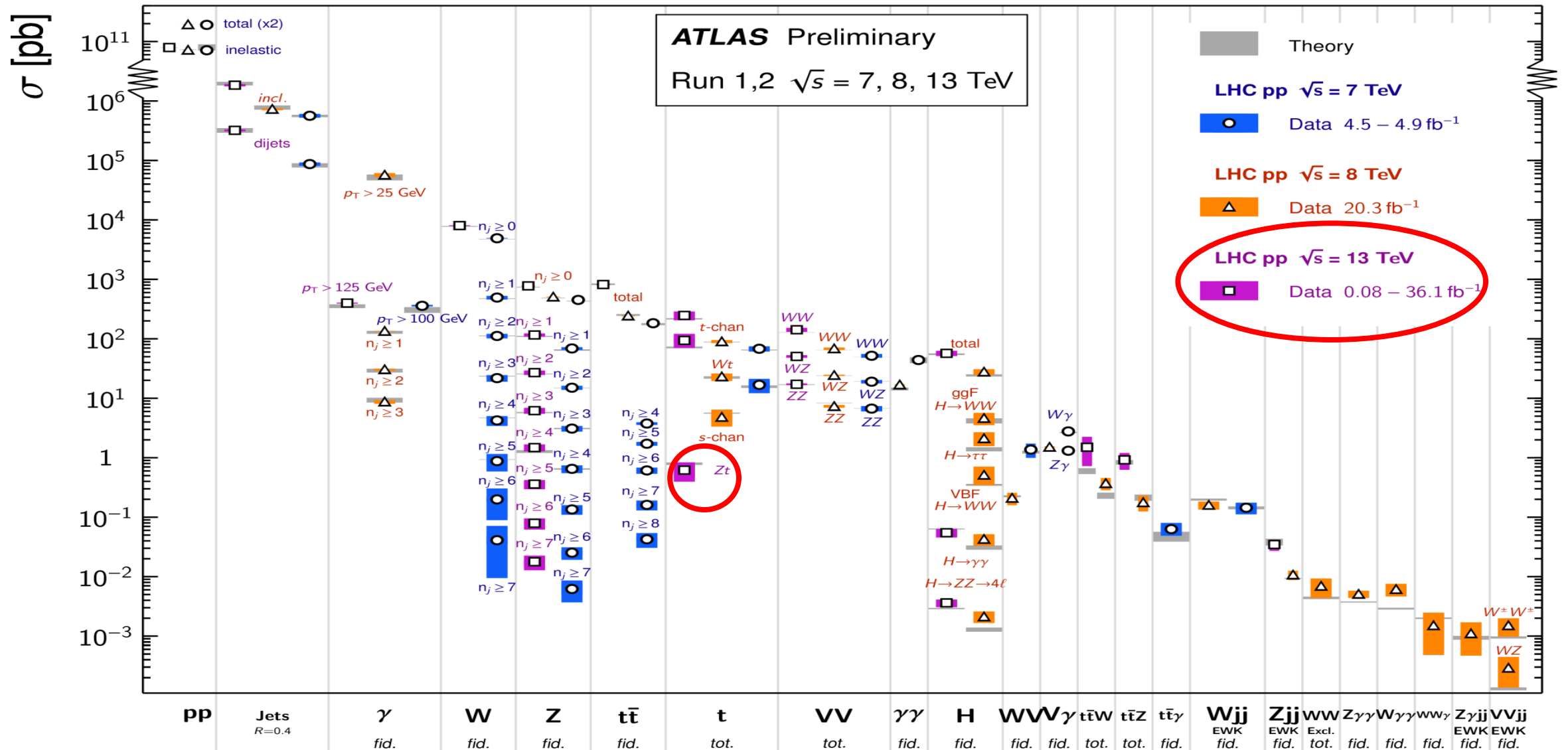
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The distribution of  $m_{bb}$  after subtraction of backgrounds, except for the WZ and ZZ dibosons proc.

# STANDARD MODEL MEASUREMENTS

## Standard Model Production Cross Section Measurements

Status: July 2017



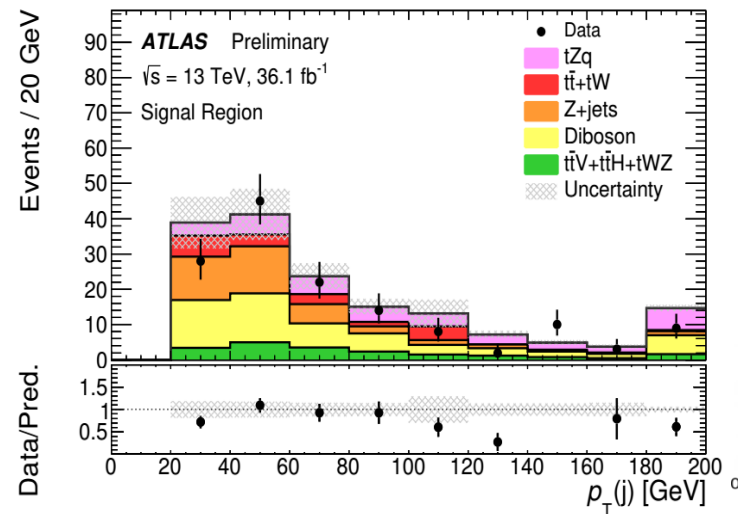
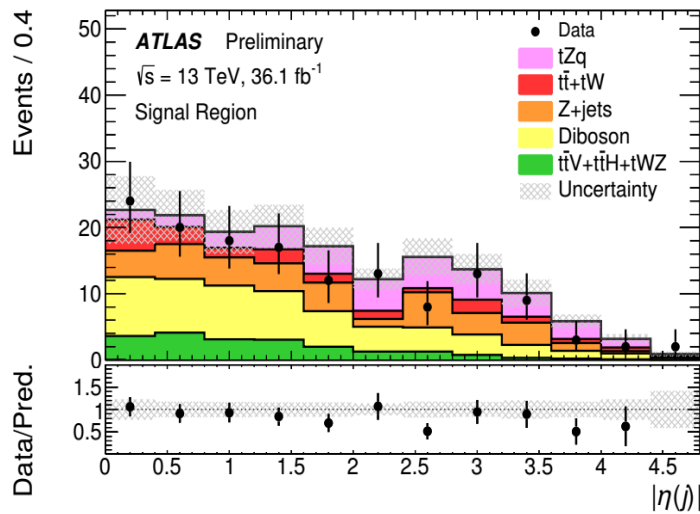
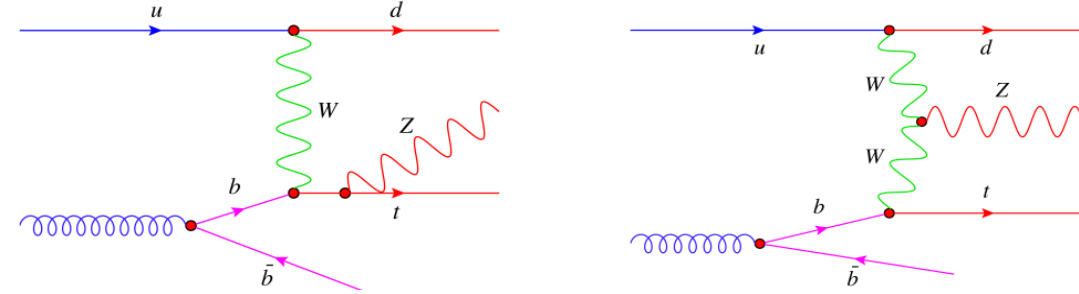
Summary of several SM total and fiducial production cross section measurements, corrected for leptonic branching fractions, compared to the theoretical expectations

# EVIDENCE FOR $Zt$ PRODUCTION

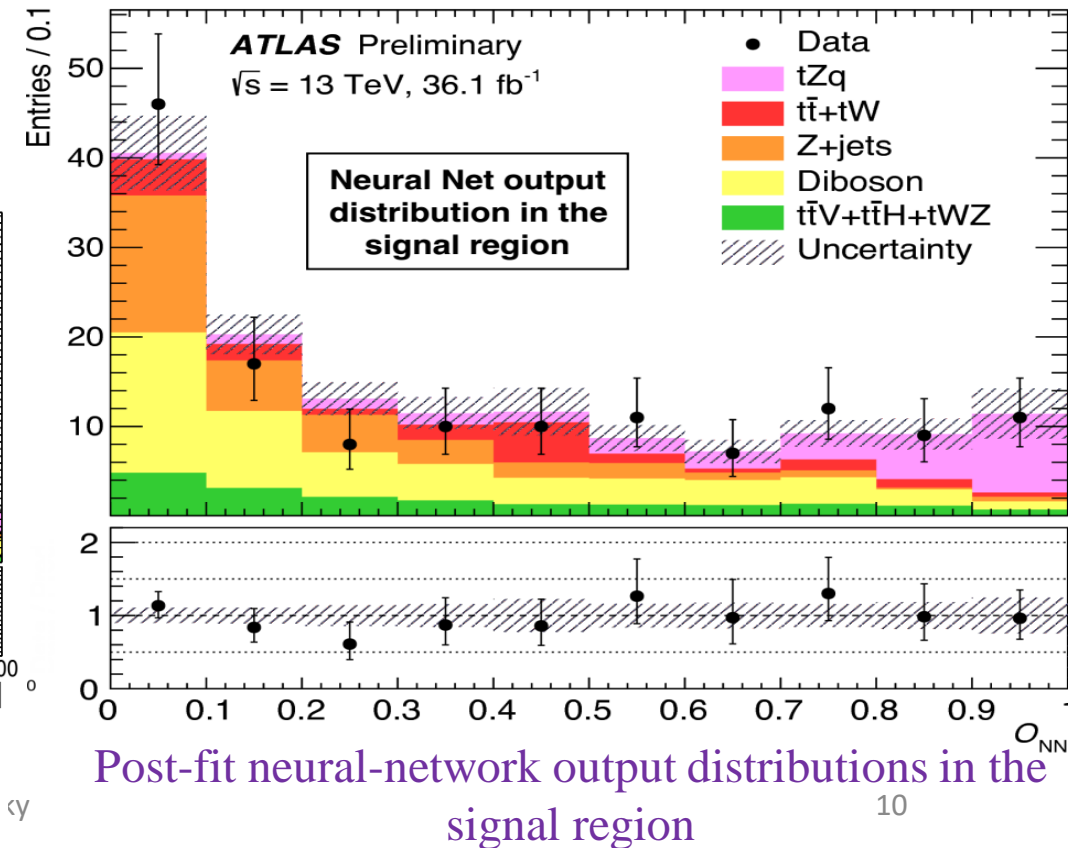
ATLAS-CONF-2017-052

- Events selected by requiring:
  - 3 identified leptons
  - 2 jets, out of which one should be b-tagged
- A *Neural Net* is trained to separate the signal events from the background
- Evidence for  $Zt$  production:
  - Significance  $4.2\sigma$  ( $5.4\sigma$  expected)
  - Cross section  $620 \pm 170 \pm 140 \text{ fb}$
  - ✓ Consistent with the SM expectation

Evidence of the production of a single t-quark in association with a Z boson in the  $pp \rightarrow tZq$



Comparison of Data and Signal + Background model of the neural-network training variables



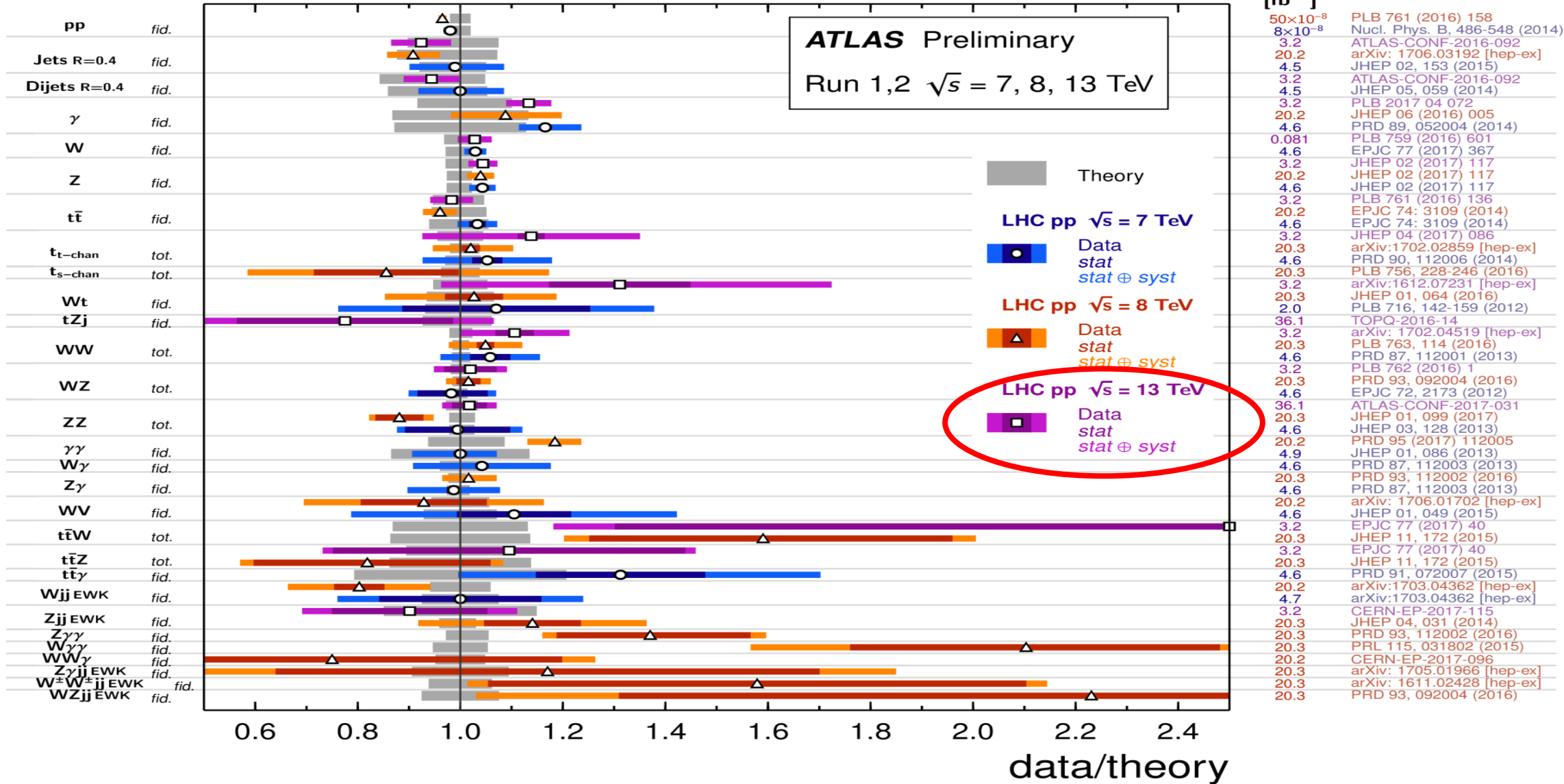
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10



# STANDARD MODEL MEASUREMENTS

## Standard Model Production Cross Section Measurements Status: July 2017



**Ratio of SM cross section measurements to theory**

# STUDY OF MINIMUM-BIAS EVENTS WITH ATLAS

Understanding of soft-QCD interactions has direct impact on

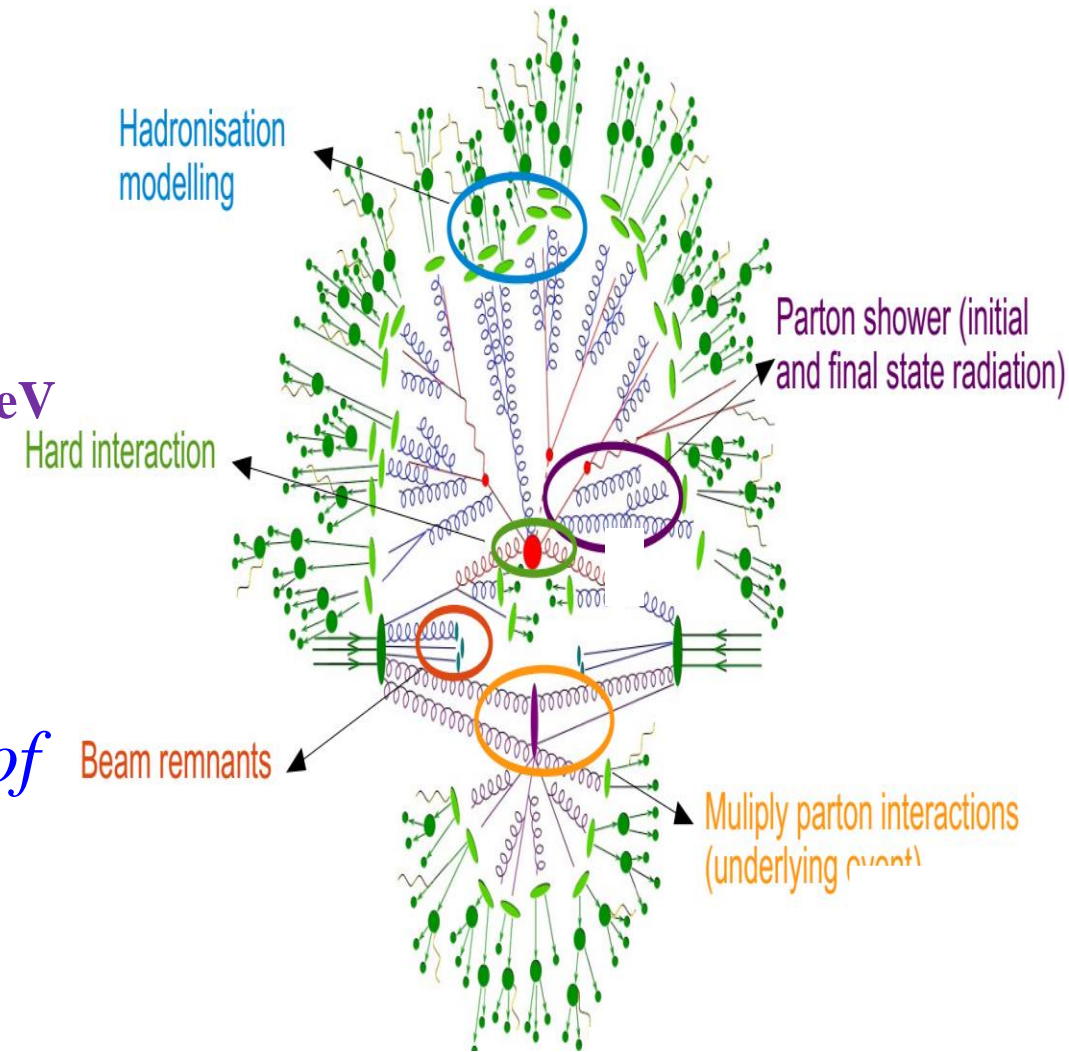
1. precision measurements
2. searches for new physics

## Studies include:

- Charged-particle distributions in pp interactions at 0.9 – 13 TeV
- **Bose-Einstein correlations (BEC)**
  - represent a unique probe of the *space-time geometry* of the *hadronization region*
  - allow the *determination the size and shape of the source* from which particles are *emitted*
- Underlying events distributions in pp interactions

**Provides insight into strong interactions in non-perturbative QCD regime:**

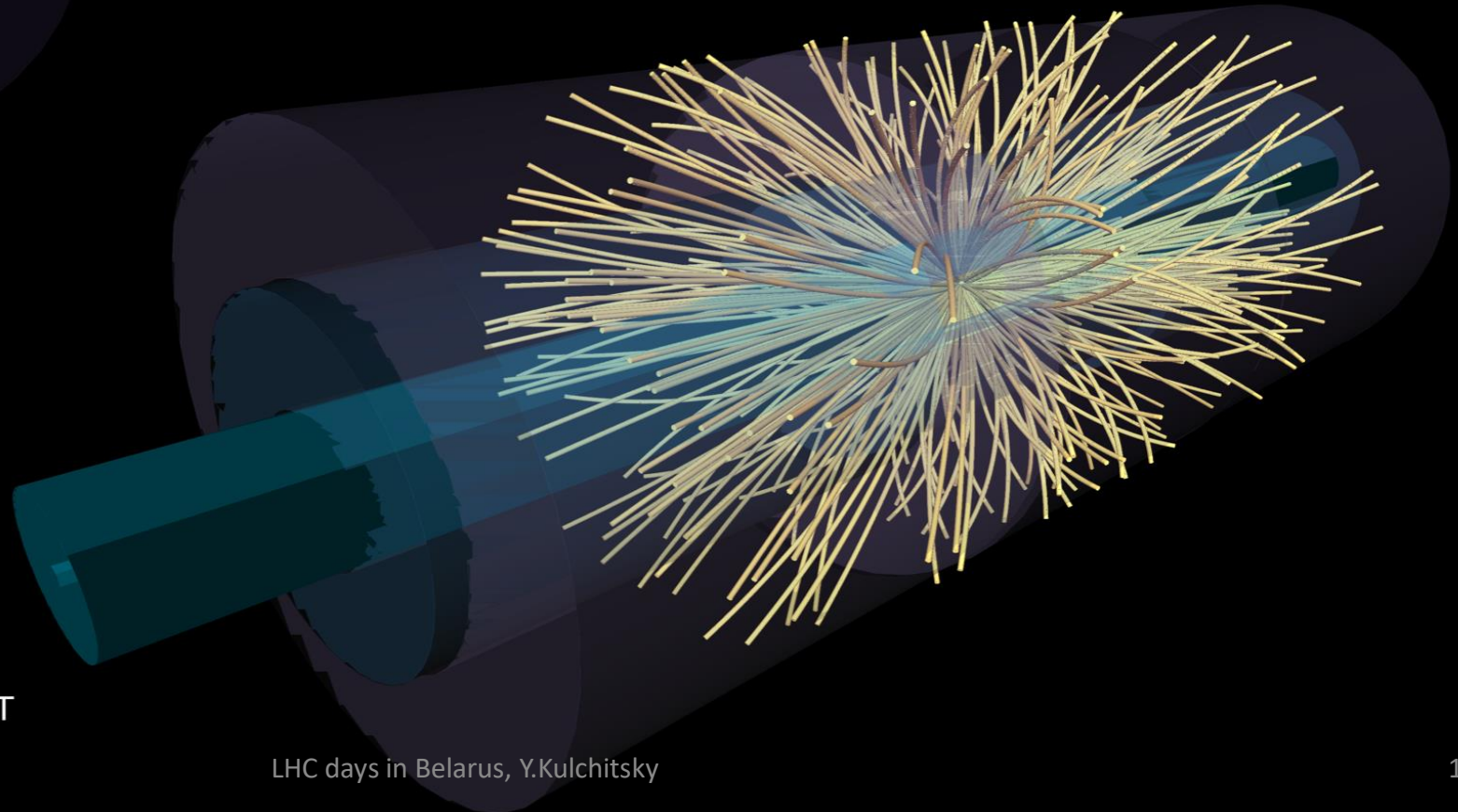
- Soft QCD results used in Monte-Carlo generators tuning,
- Low energy QCD description essential for simulating multiple pp interactions



# EXAMPLE OF VERY-HIGH-MULTIPLICITY EVENT

High-multiplicity event with 319 reconstructed tracks.  
The shown tracks are from a single vertex and have  $p_T > 0.4$  GeV

**319 reconstructed charged-particles!**

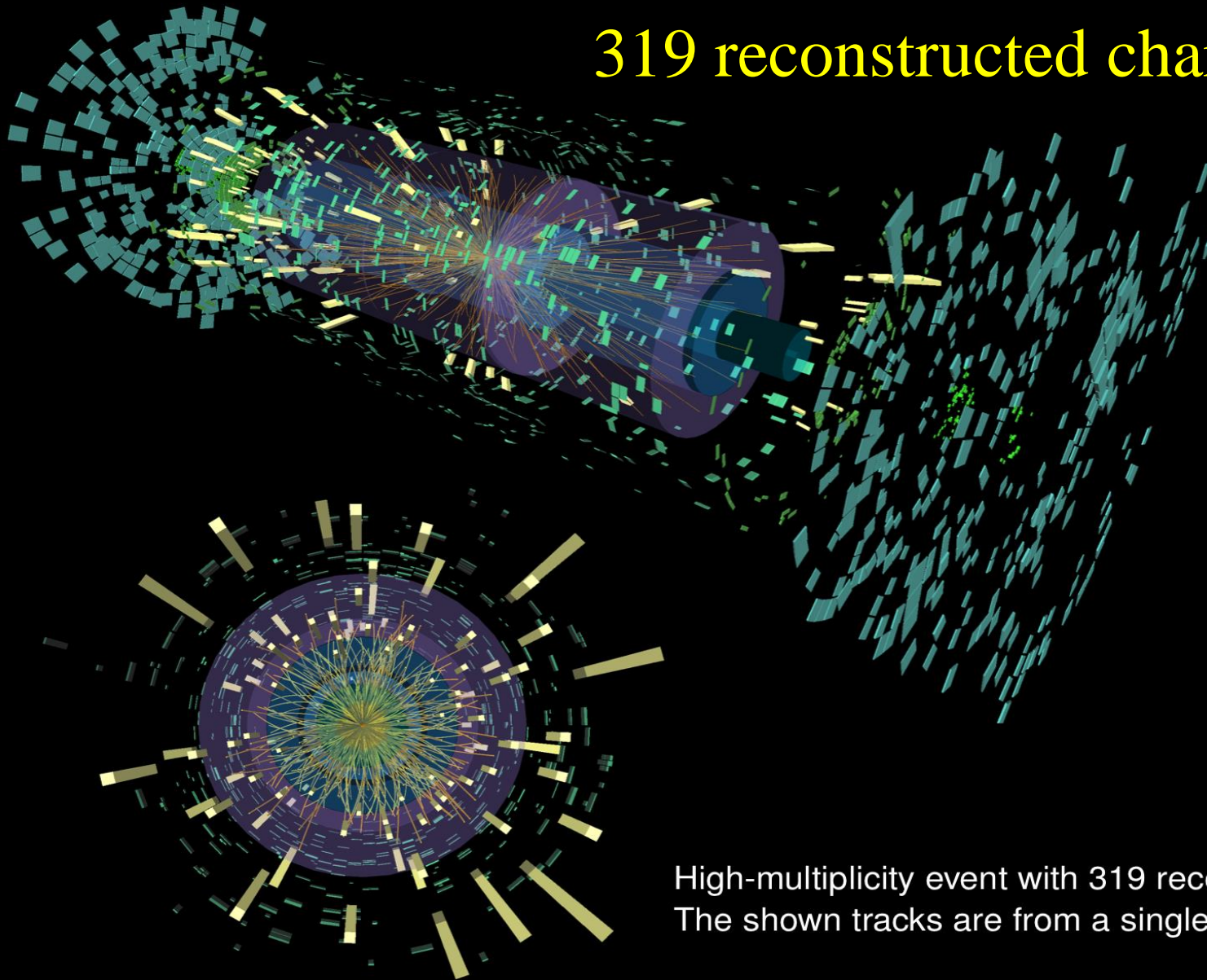


Run: 312837  
Event: 135456971  
2016-11-14 07:42:28 CEST



# EXAMPLE OF VERY-HIGH-MULTIPLICITY EVENT

319 reconstructed charged-particles



Run: 312837

Event: 135456971

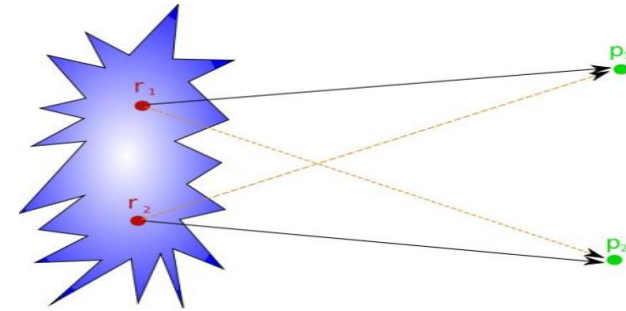
2016-11-14 07:42:28 CEST

High-multiplicity event with 319 reconstructed tracks.  
The shown tracks are from a single vertex and have  $p_T > 0.4$  GeV

# BOSE-EINSTEIN CORRELATIONS

Correlations in phase space between two identical bosons from symmetry of wave functions.

- ▶ Enhances likelihood of two particles close in phase space
- ▶ Allows one to ‘probe’ the source of the bosons in size and shape
- ▶ Dependence on particle multiplicity and transverse momentum probes the production mechanism



Correlation function  $C_2(Q)$  a ratio of probabilities:

$$C_2(Q) = \frac{\rho(p_1, p_2)}{\rho_0(p_1, p_2)} = C_0 (1 + \Omega(\lambda, RQ)) \cdot (1 + Q\varepsilon), \quad Q^2 = -(p_1 - p_2)^2$$

$$\Omega^G(\lambda, RQ) = \lambda e^{-R^2 Q^2}$$

$$\Omega^E(\lambda, RQ) = \lambda e^{-RQ}$$

$C_0$  is a normalisation,  $\varepsilon$  accounts for long range effects,  $\mathbf{R}$  is the effective radius parameter of the source,  $\lambda$  is the strength of the effect parameter, 0/1 for coherent/chaotic source. Two possible parameterisation: Gaussian and Exponential.

$$C_2(Q) = \frac{N^{++,--}(Q)}{N^{ref}(Q)}$$

$N_{ref}$  without BEC effect from: unlike-charge particles (UCP), opposite hemispheres, event mixing. **Basic Reference:** distribution of UCP pairs of non-identical particle taken from the same event.

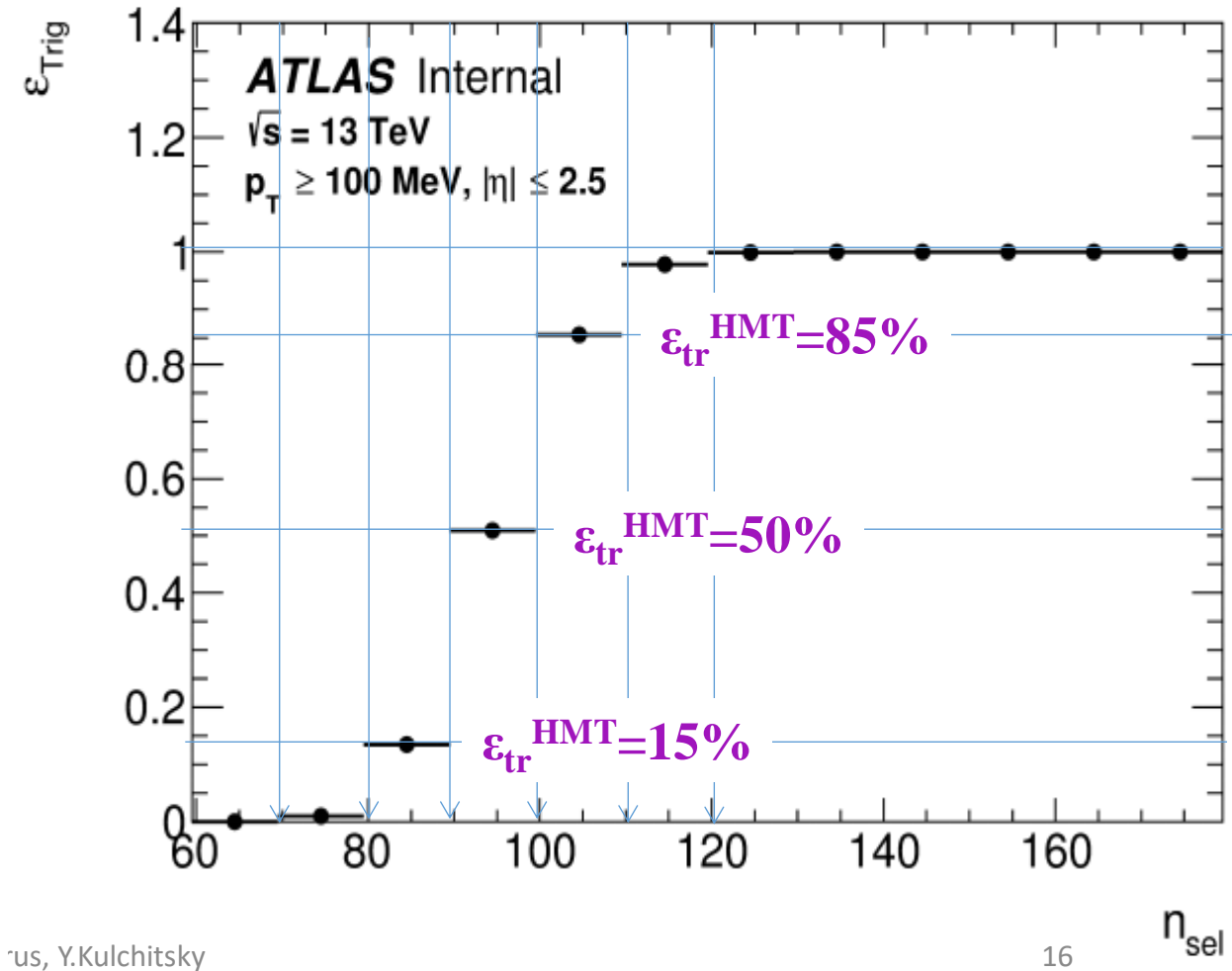
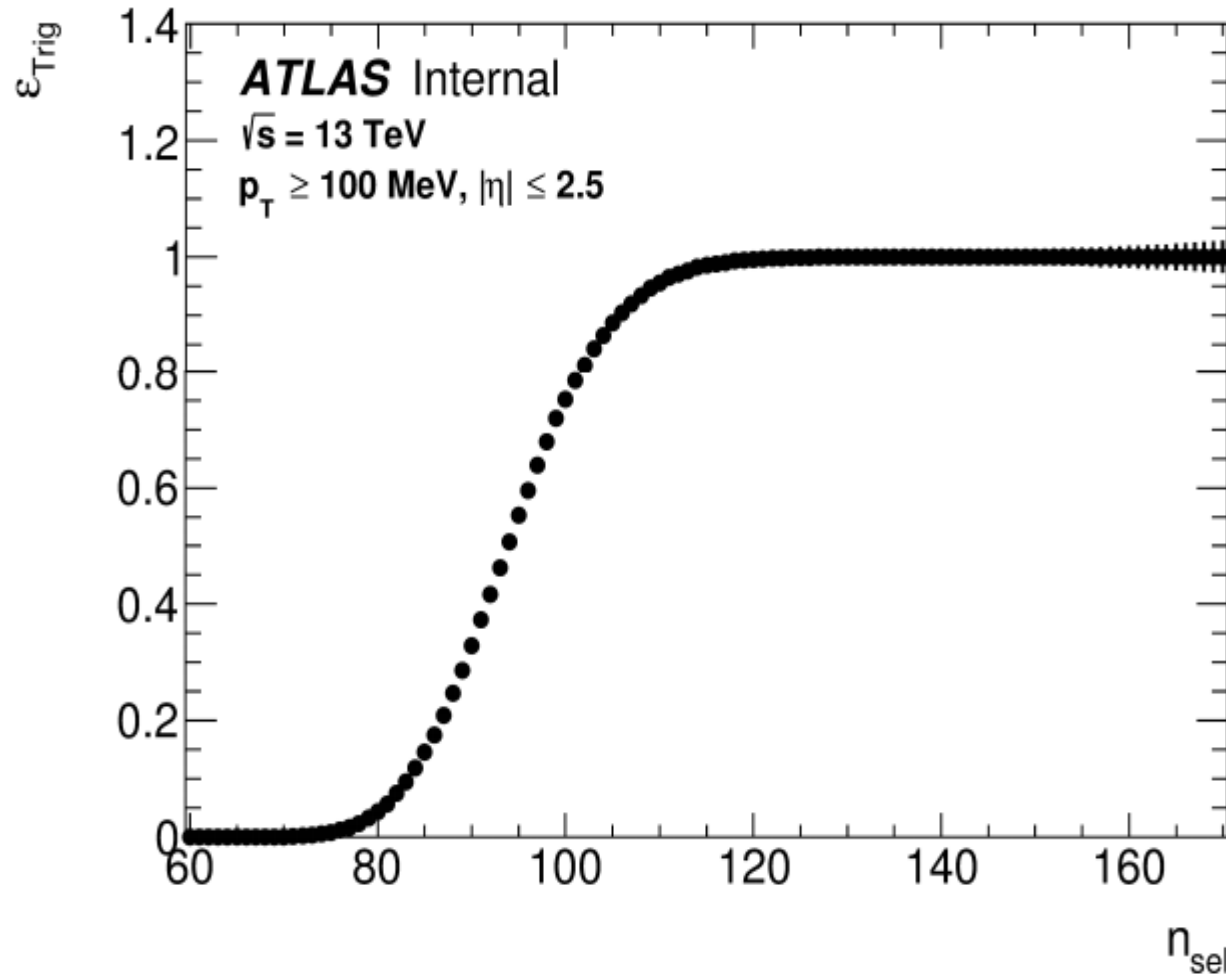
$$R_2(Q) = \frac{C_2^{Data}(Q)}{C_2^{MC}(Q)} = \frac{\rho^{(++,-)} / \rho^{(+-)}}{\rho^{MC}^{(++,-)} / \rho^{MC}^{(+-)}}$$

The studies are carried out using the **double ratio correlation function**. The  $R_2(Q)$  eliminates problems with energy-momentum conservation, topology, resonances etc.

**MC without BEC.**

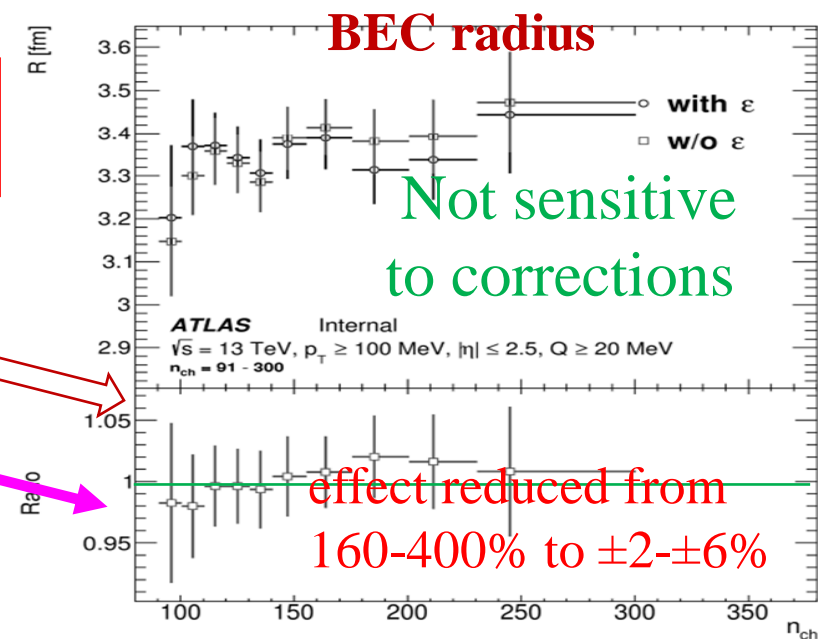
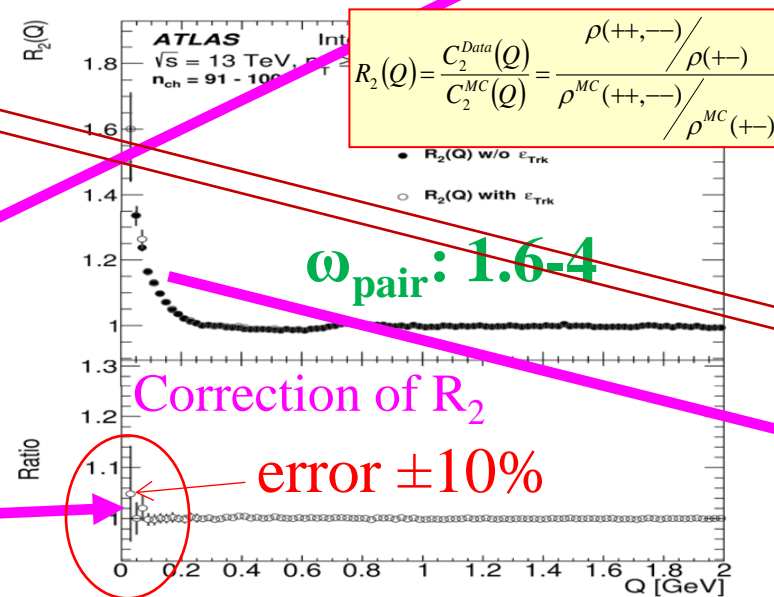
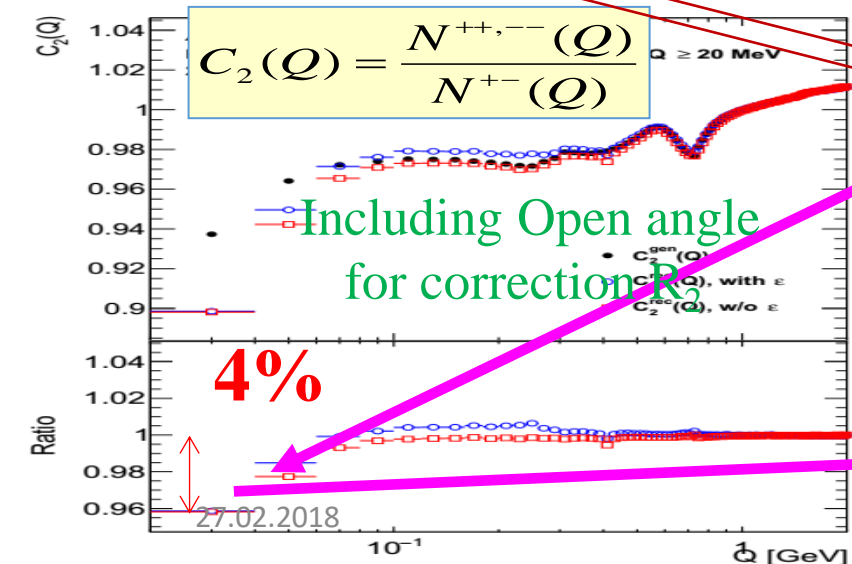
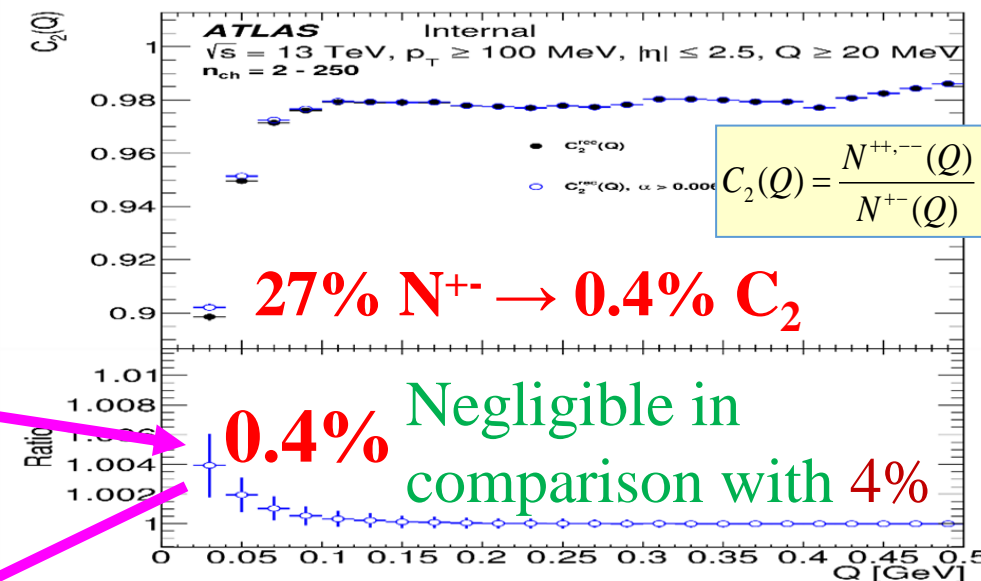
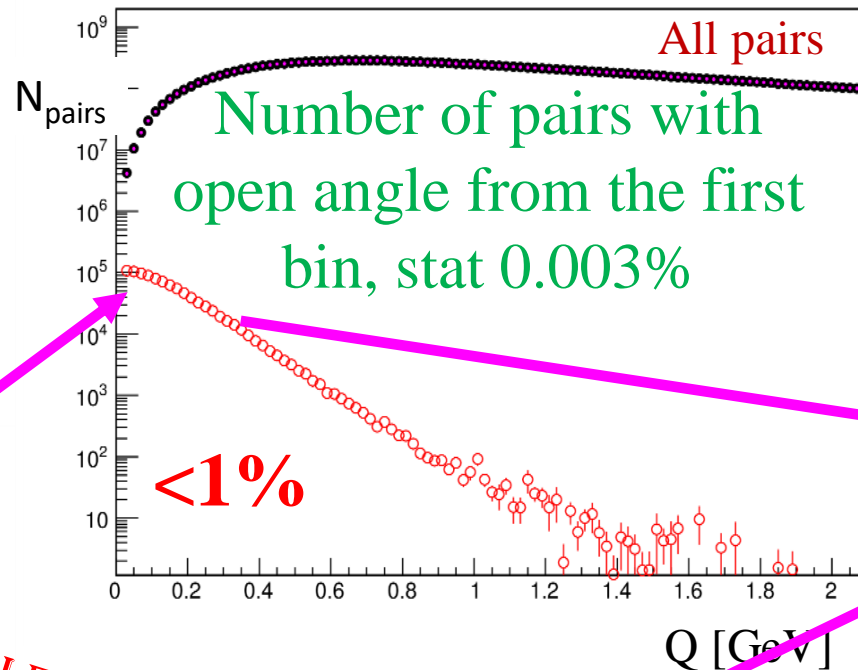
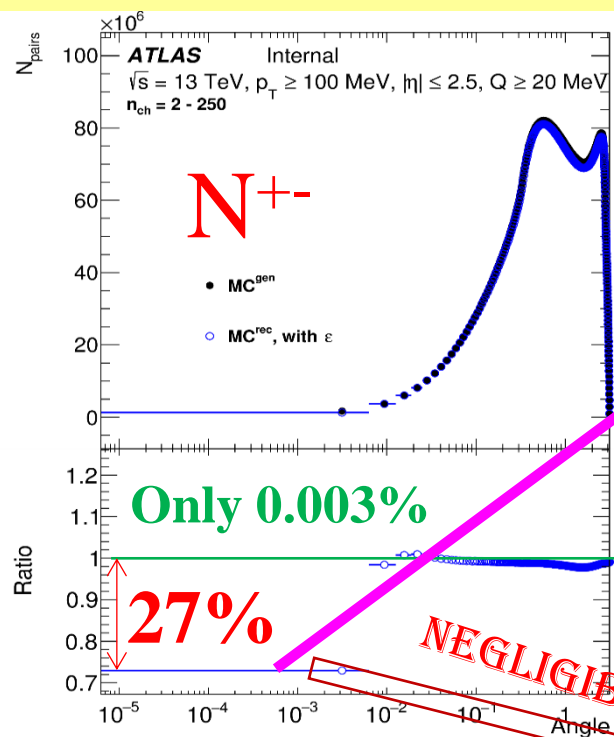
# HIGH MULTIPLICITY TRACK TRIGGER EFFICIENCY

- High multiplicity track trigger (**HLT\_mb\_sp900\_trk60\_hmt\_L1MBTS\_1\_1**) efficiency calculation as difference the HMT trigger to minimum-bias trigger (**L1MBTS\_1\_1**)





# SUMMARY: CORRECTION OF $R_2$ CORRELATION FUNCTION



# COULOMB CORRECTION

The measured  $N(Q)$  distribution for like or unlike signed particle (track) pairs in presence of the Coulomb interaction is given by:

where  $N_{meas}(Q)$  is the measured distribution,  $N(Q)$  is the distribution free of Coulomb correlations.

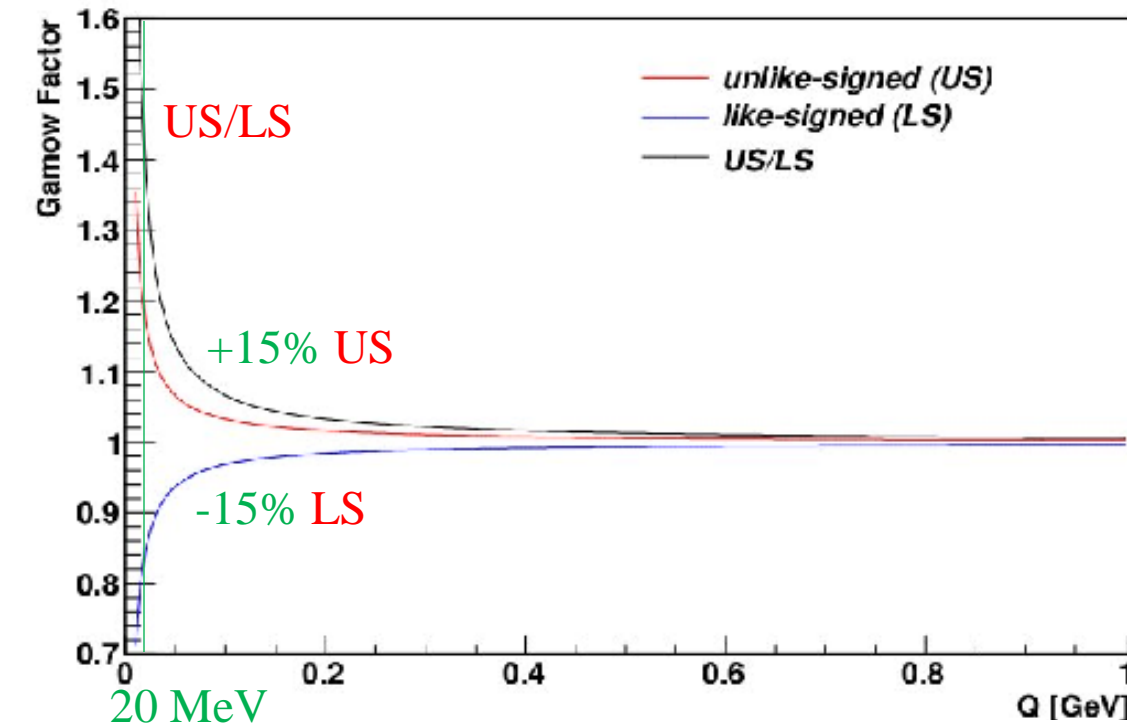
$$N_{meas}(Q) = G(Q)N(Q)$$

Gamow penetration  $G(Q)$  factor

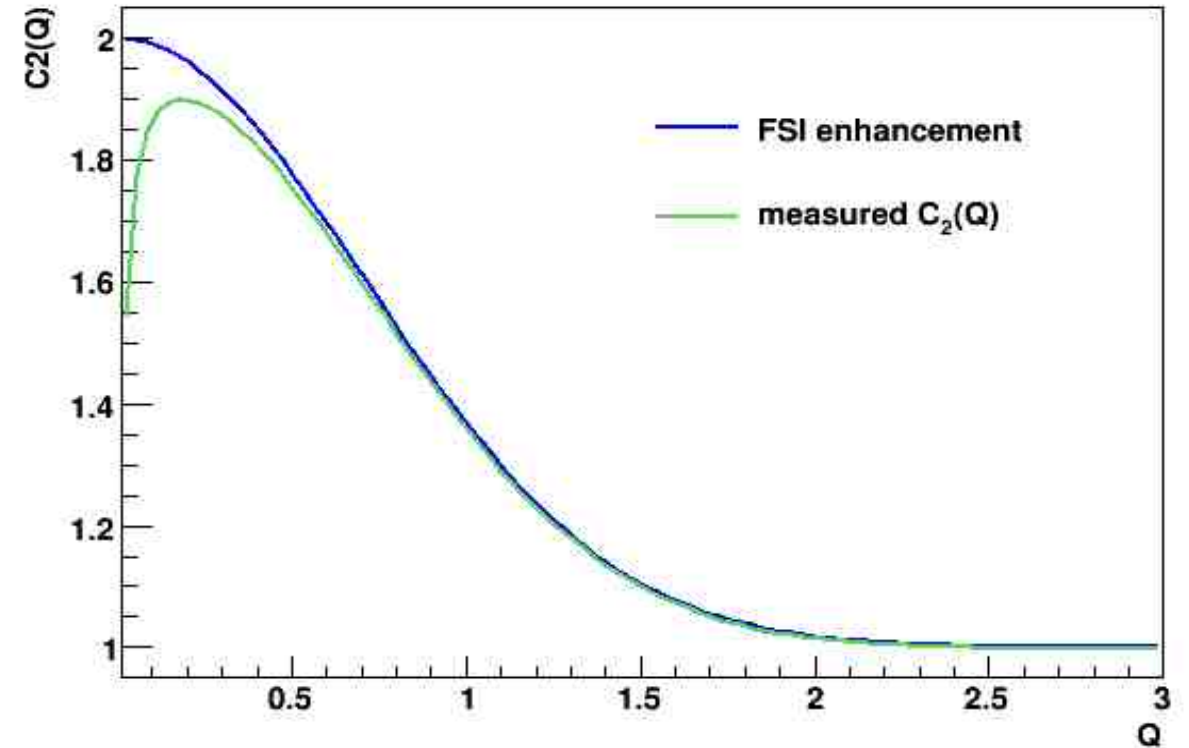
$$G(Q) = \frac{2\pi\eta}{e^{2\pi\eta} - 1}$$

Sommerfeld parameter  $\eta$

$$\eta = \frac{\pm \alpha_m}{|Q|}$$



Gamow factor: **blue curve** – like-signed (LS) particle pairs,  
**red curve** – unlike-signed (US) particle pairs,  
**black curve** – US/LS



Ratio of the same  $C_2(Q)$  function – **with** and **without** correction on Coulomb correction. FSI – final state interactions.

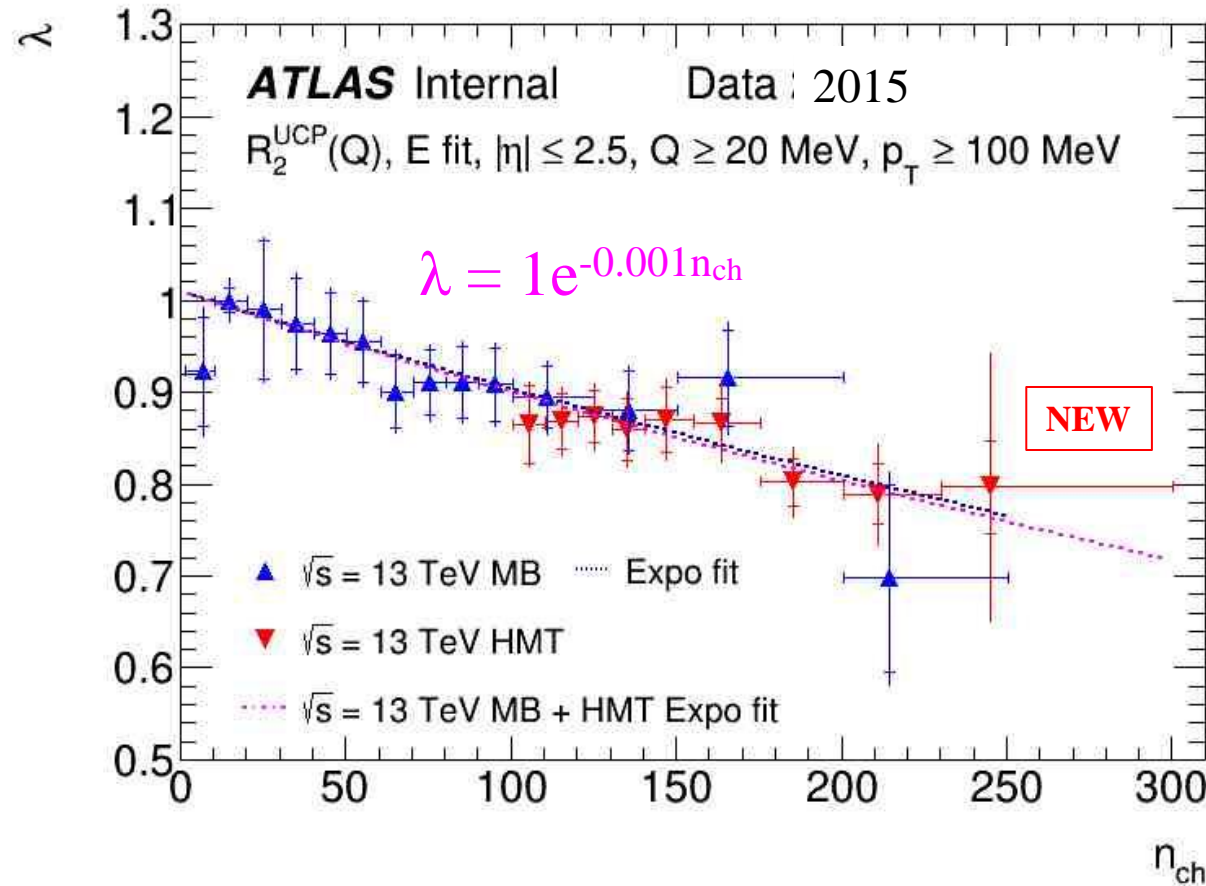
# SYSTEMATIC UNCERTAINTIES FOR BEC AT 13 TeV

The systematic uncertainties of the spread for  $n_{\text{ch}}$  distribution and inclusive fit parameters,  $R$  and  $\lambda$ , of the exponential model are summarized in the Table. The systematic uncertainties are combined by adding them in quadrature and the resulting values are given in the bottom row of the Table. The same sources of uncertainty are considered for the differential measurements in  $n_{\text{ch}}$ ; the average transverse momentum  $k_{\text{T}}$  of a pair; the two-differential measurements in intervals of  $(n_{\text{ch}}; k_{\text{T}})$ , and their impact on the fit parameters is found to be similar in size.

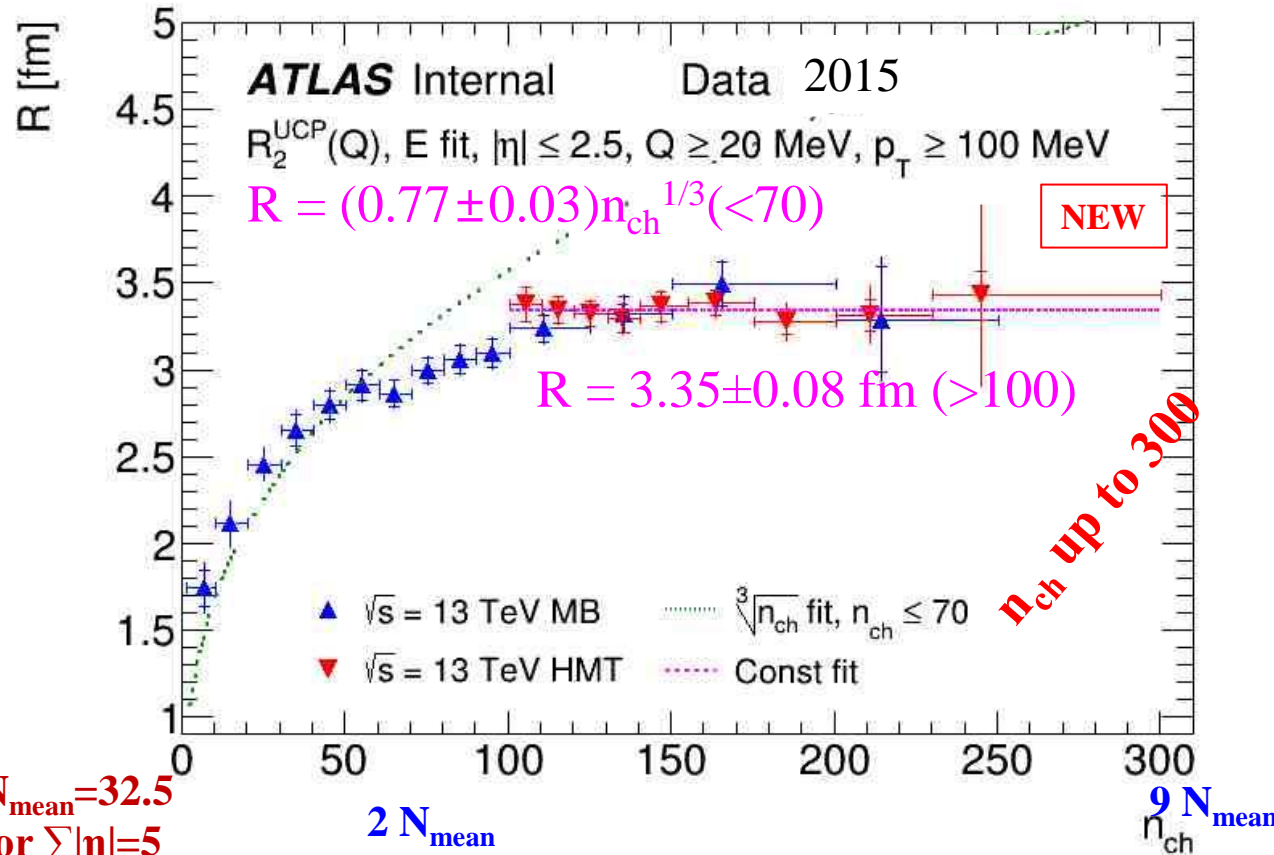
Sources	13 TeV				13 TeV (HMT)			
	$\lambda$	$R$	$\lambda$	$R$	$\lambda$	$R$	$\lambda$	$R$
	(%)	(%)	(%)	(%)	(%)	(%)	(%)	(%)
	$n_{\text{ch}} - \text{spread}$		Inclusive		$n_{\text{ch}} - \text{spread}$		Inclusive	
1. Track reconstruction efficiency: $\omega \pm \delta\omega$	0.0–0.4	0.1–0.4	0.3	0.1	0.1–0.2	0.01–0.1	0.2	0.01
2. Monte Carlo: EPOS, Pythia8 Monash	Track splitting and merging				negligible			
3. Coulomb correction: $\pm 15\%$	0.0–6.9	0.5–7.0	1.1	1.4	1.0–16.	1.0–14.	1.4	0.7
4. Fitted range of $Q$ : $2 \text{ GeV} \pm 3\sigma_Q$ (100 MeV)	1.3–2.0	0.01–0.6	1.8	0.1	1.7–1.9	0.2–0.4	1.8	0.3
5. Starting value, $Q_{\text{min}}$ : 10, 20, 30 MeV	0.0–0.5	0.02–0.9	0.2	0.3	0.0–0.2	0.0–0.2	0.02	0.03
6. Bin size: 10, 20, 30 MeV	0.0–1.9	0.01–1.1	0.3	0.2	0.5–1.4	0.3–0.7	0.7	0.4
7. Excluded intervals: $\pm 20 \text{ MeV}$	0.0–2.4	0.1–1.5	0.8	0.4	1.0–1.7	0.4–0.9	1.3	0.6
	0.0–1.1	0.3–0.8	0.1	0.3	0.4–0.6	0.3–0.6	0.5	0.5
Total	1.3–7.9	0.9–7.2	2.4	1.5	3.0–17.	1.0–15.	2.8	1.2

# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS AT 13 TEV

ATL-COM-PHYS-2016-1621



$N_{\text{mean}}=32.5$   
for  $\sum|\eta|=5$



$2 N_{\text{mean}}$

$9 N_{\text{mean}}$

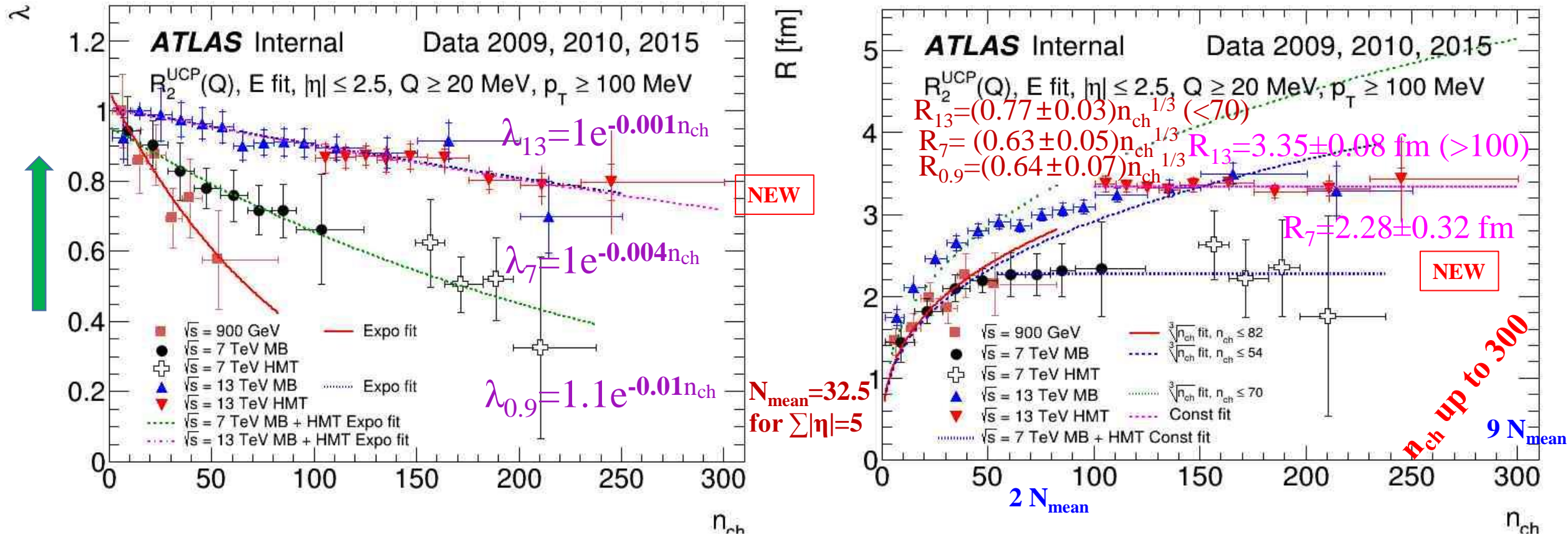
The results for Bose-Einstein Correlations for pairs of like-sign charged particles measured in the kinematic range  $p_T > 100$  MeV and  $|\eta| < 2.5$  in pp collisions at energy 13 TeV

- The multiplicity dependence of the BEC parameters characterizing the *correlation strength* and the *correlation source size* are investigated for multiplicities of up to *very-high number of charged-particles*,  $n_{\text{ch}} \approx 300$
- A *saturation effect* in the multiplicity dependence of the correlation source-size parameter is observed at 13 TeV



# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS AT 0.9 – 13 TEV

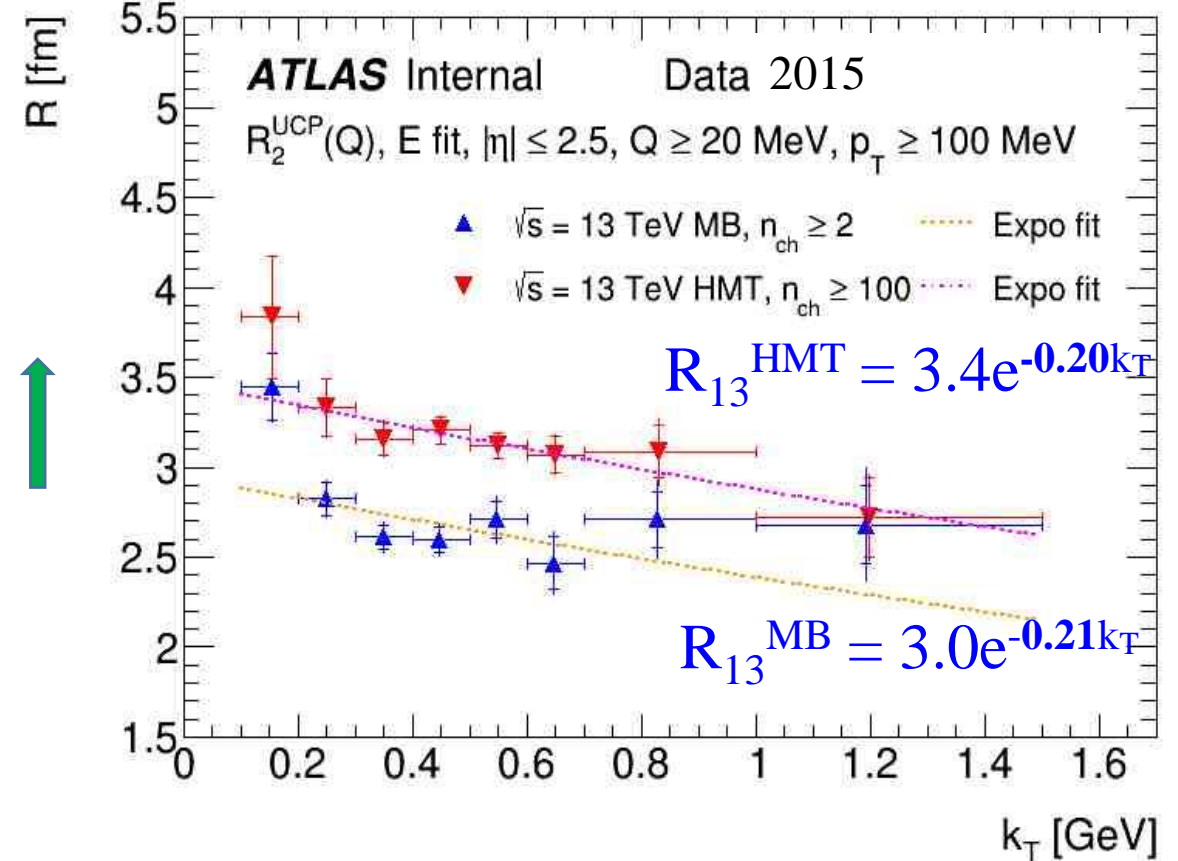
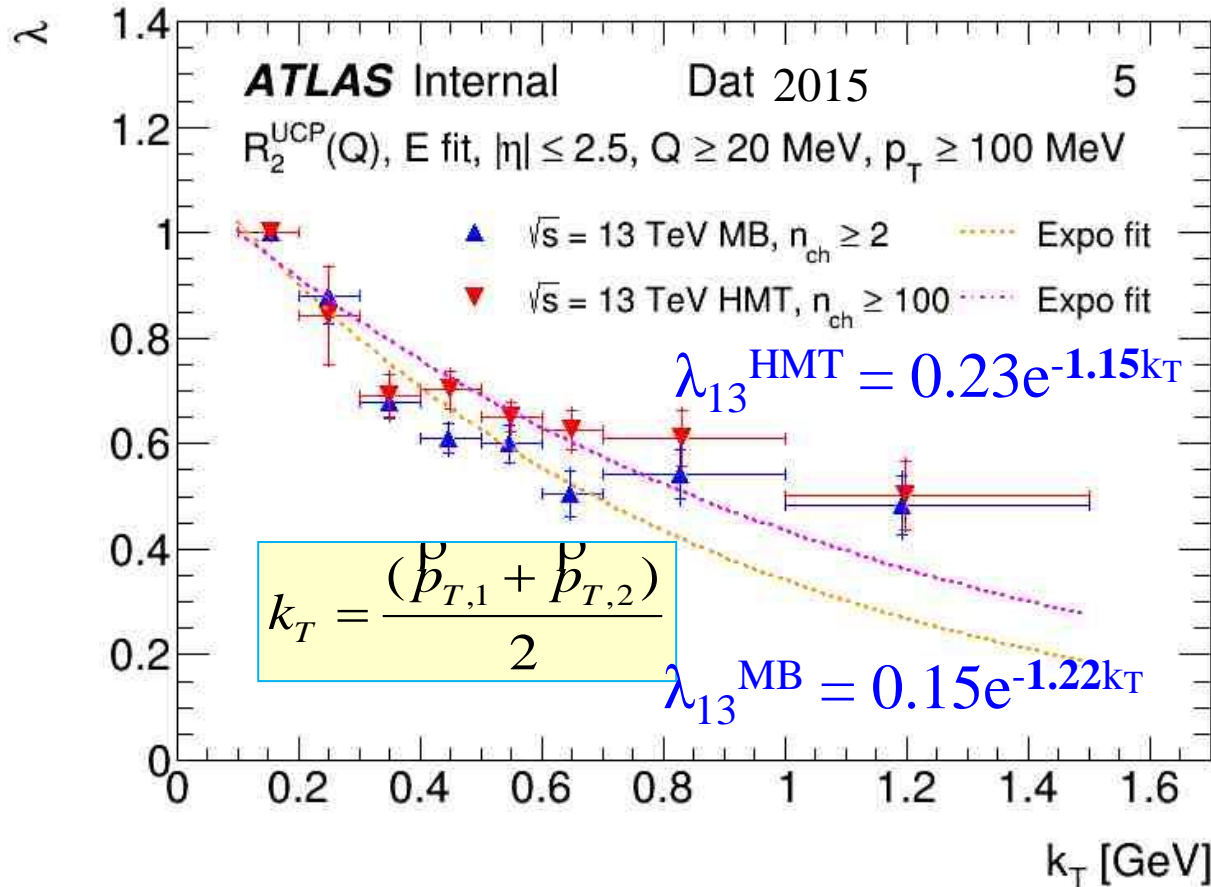
EPJC 75 (2015) 10, 466; ATL-COM-PHYS-2016-1621



- The slope of an exponential fit of the  $\lambda$  vs  $n_{ch}$  distributions decrease with increasing of energy.
- The parameters  $\alpha$  of the  $\alpha \cdot n_{ch}^{1/3}$  fit of  $R$  vs  $n_{ch}$  for  $n_{ch} \leq 55$  at 0.9 TeV is  $\alpha = 0.64 \pm 0.07$  fm, 7 TeV is  $\alpha = 0.63 \pm 0.05$  fm and for  $n_{ch} \leq 70$  at 13 TeV is  $\alpha = 0.77 \pm 0.03$  fm.  
 For multiplicity region  $n_{ch} \leq 70$ , the  $R$  values are systematically higher at 13 TeV than at 7 TeV.
- The  $R$  is a constant for  $n_{ch} > 55$  at 7 TeV  $R = 2.28 \pm 0.32$  fm and for  $n_{ch} > 100$  at 13 TeV  $R = 3.35 \pm 0.08$  fm. The  $R$  is systematically higher at 13 TeV than at 7 TeV but in the error bars is in agreement.

# $k_T$ DEPENDENCE OF BEC PARAMETERS

ATL-COM-PHYS-2016-1621

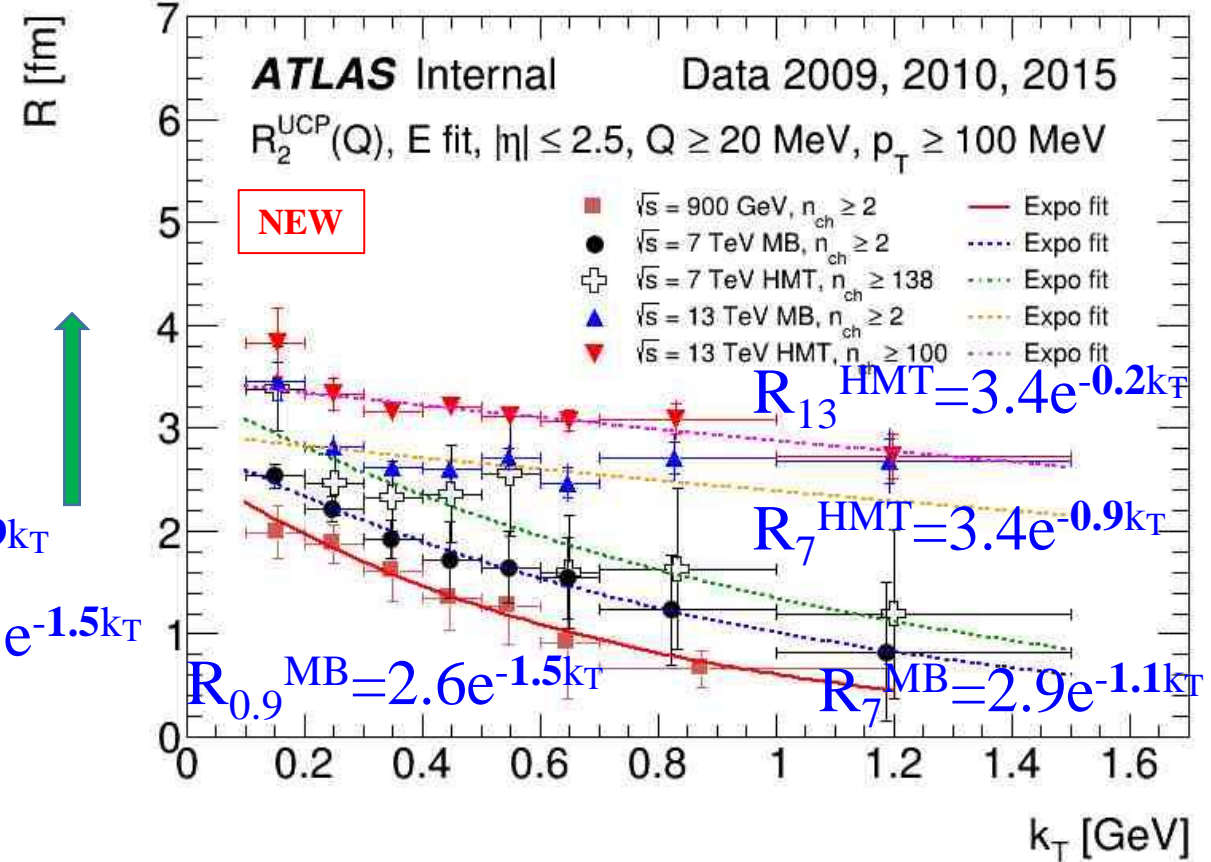
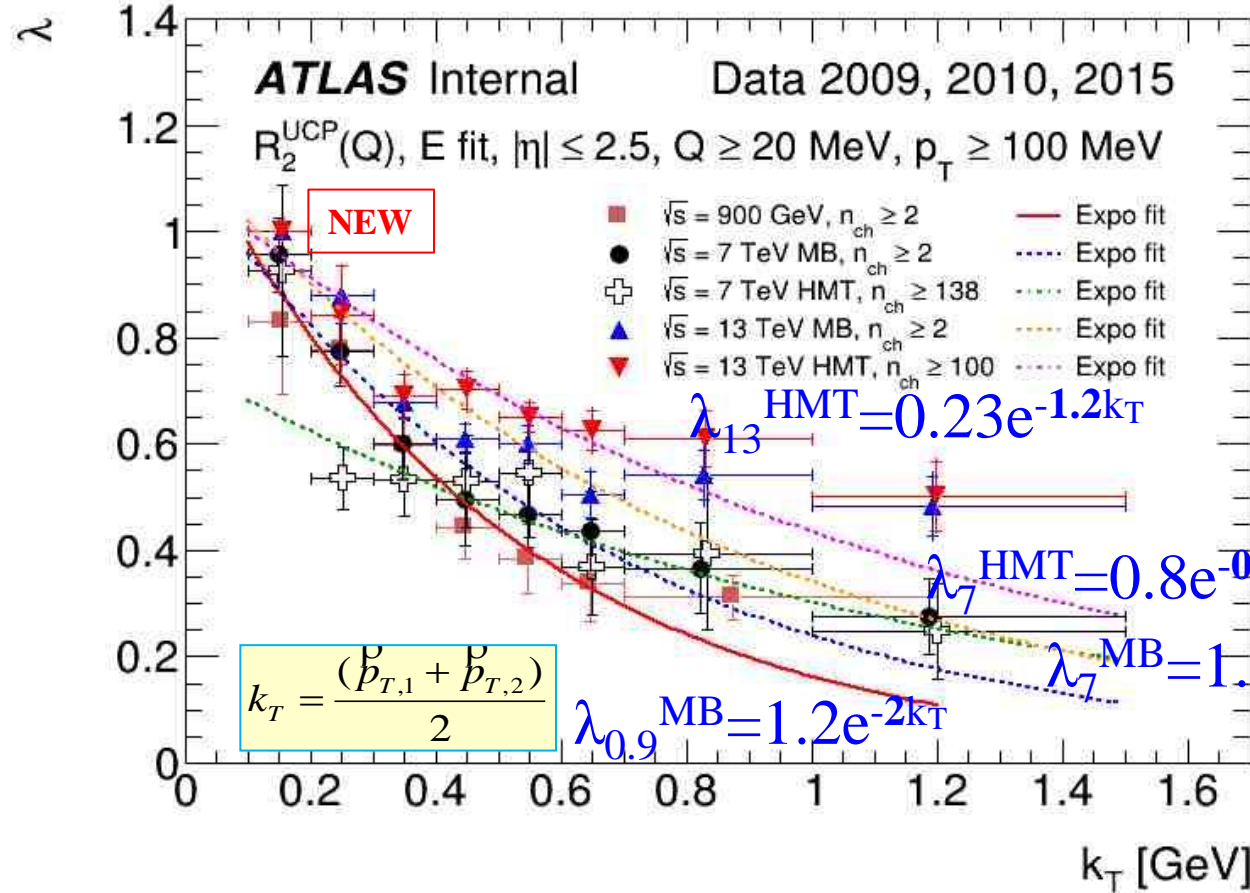


- The  $\lambda$  values are (trigger) multiplicity-independent within uncertainties at 13 TeV
- The  $R$  values increase with increasing (trigger) multiplicity region
- The  $\lambda$  values decrease exponentially with  $k_T$
- The  $R$  values decrease exponentially with  $k_T$  for HMT events



# $K_T$ DEPENDENCE OF BEC PARAMETERS AT 0.9 – 13 TEV

EPJC 75 (2015) 10, 466; ATL-COM-PHYS-2016-1621



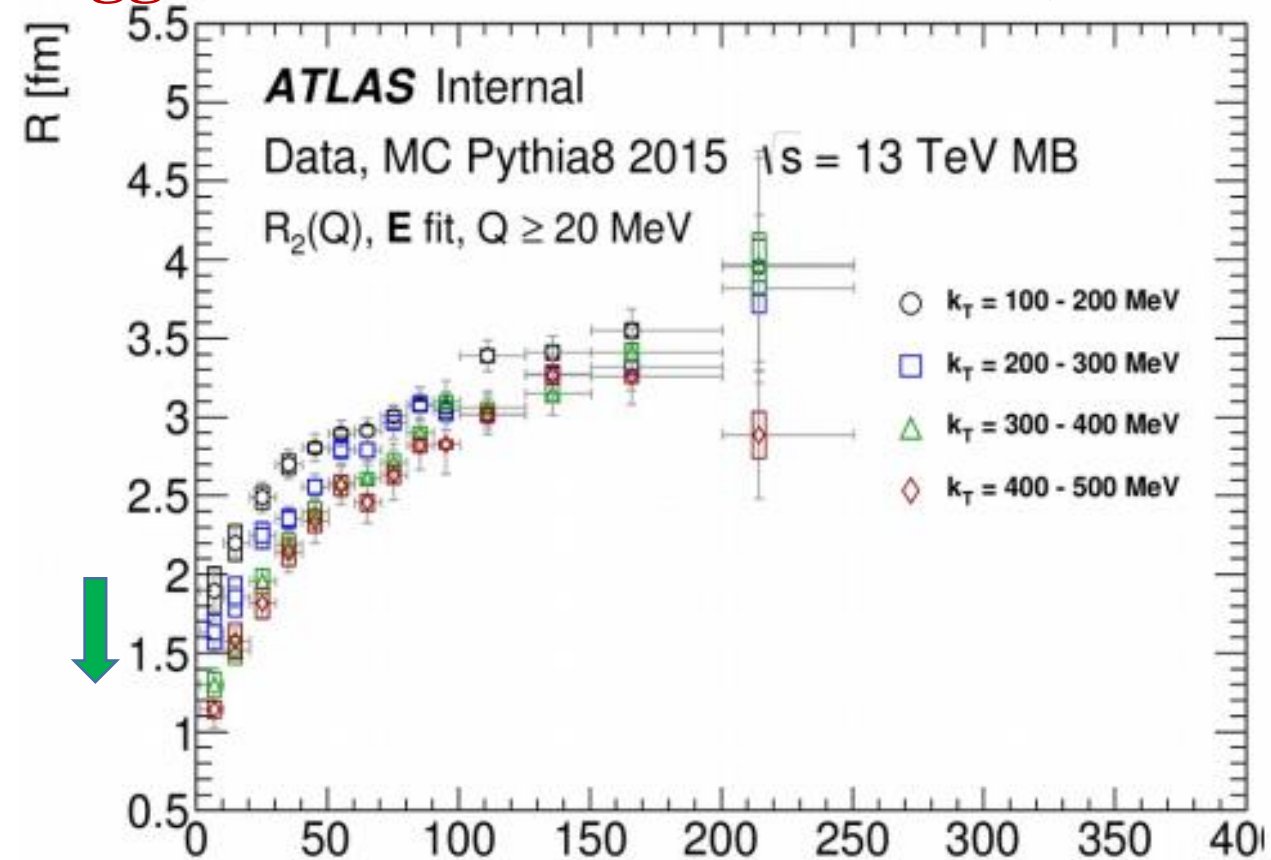
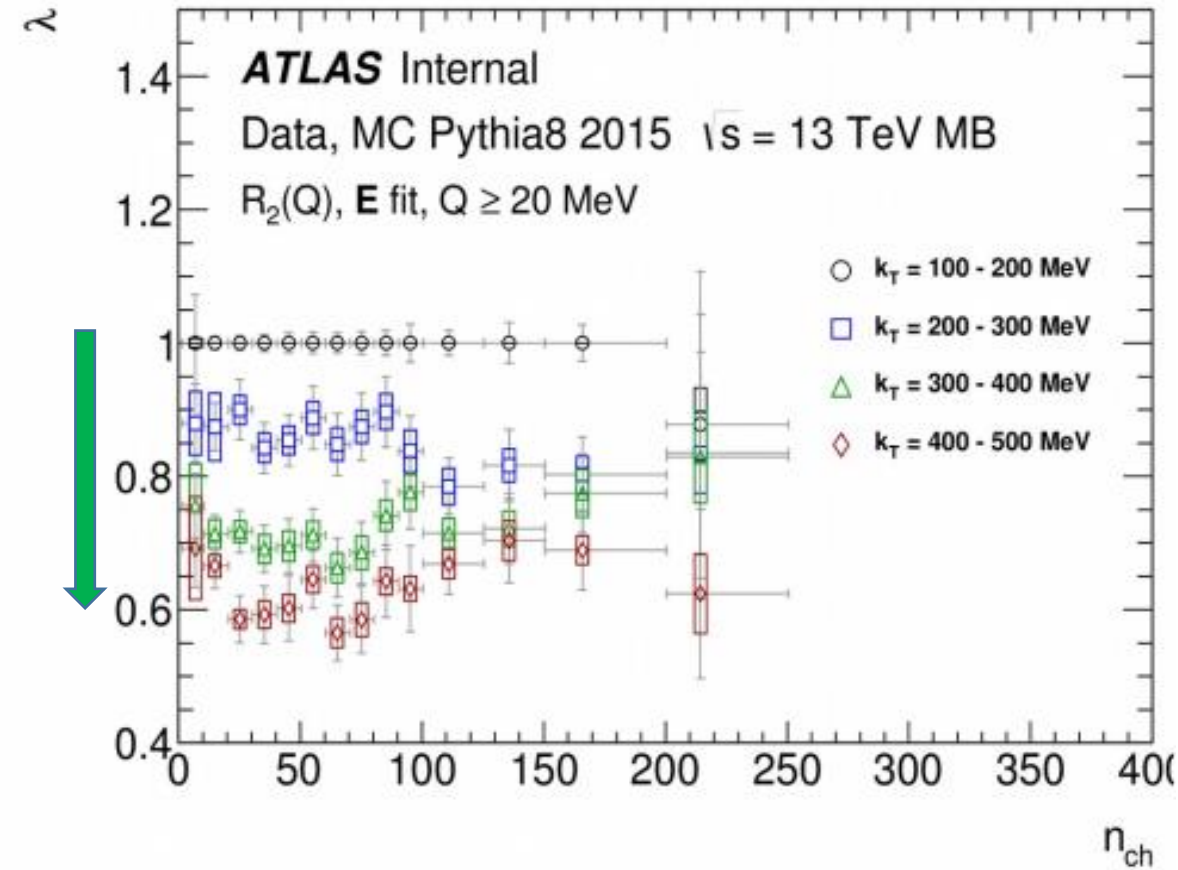
- The amplitude of fit of  $\lambda^{MB}$  vs  $k_T$  distributions is decrease from 1.2 to 0.23 with energy increasing.
- The slope of exponential fit of the  $R$  values vs  $k_T$  distributions decrease from 1.5 to 0.2 with increasing of energy.

# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS FOR $(N_{CH}, K_T)$ INTERVALS

## $\lambda$ distribution at 13 TeV

**MB trigger**

## R distribution at 13 TeV

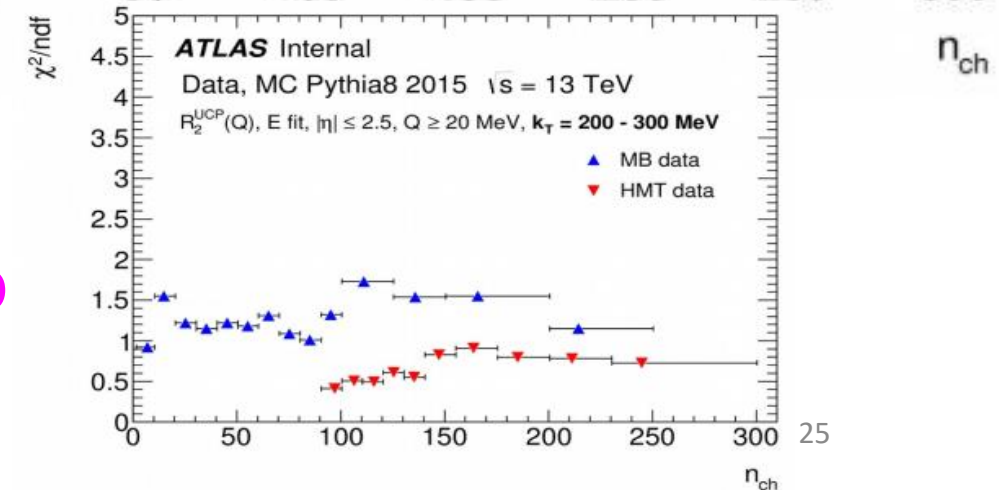
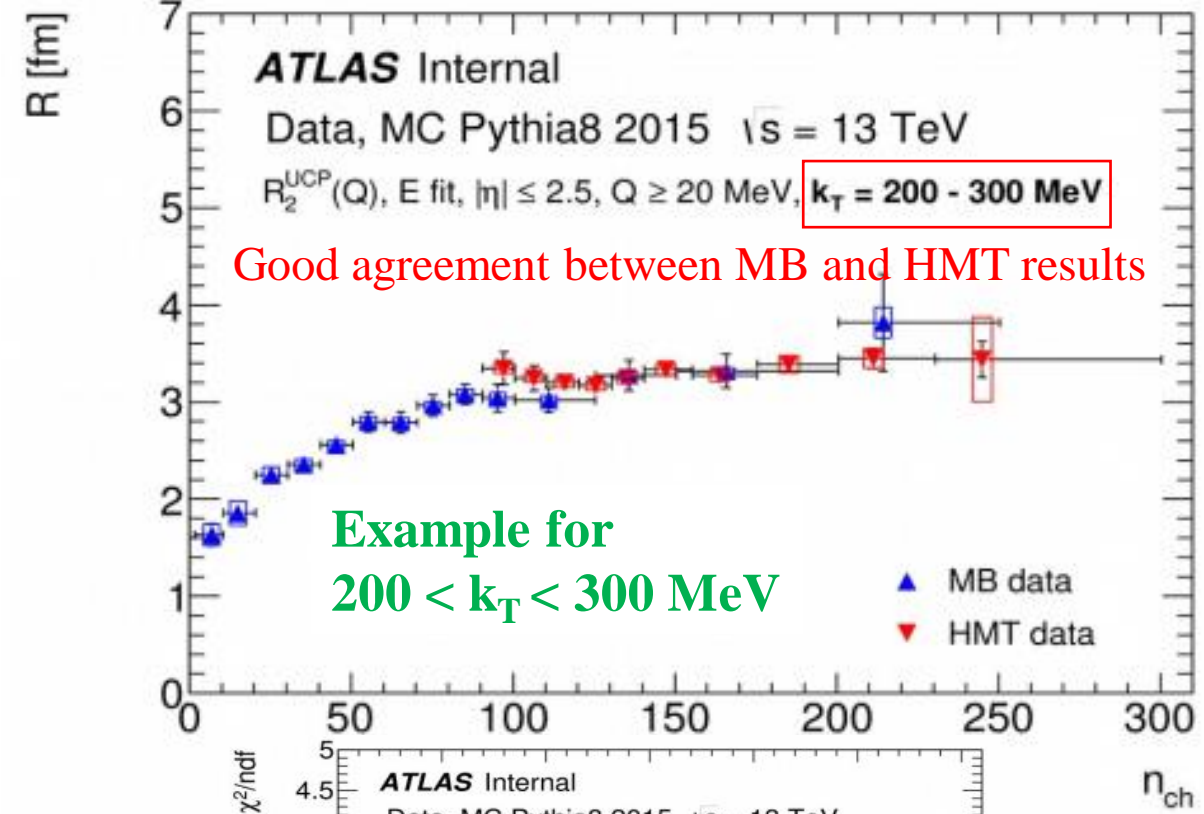
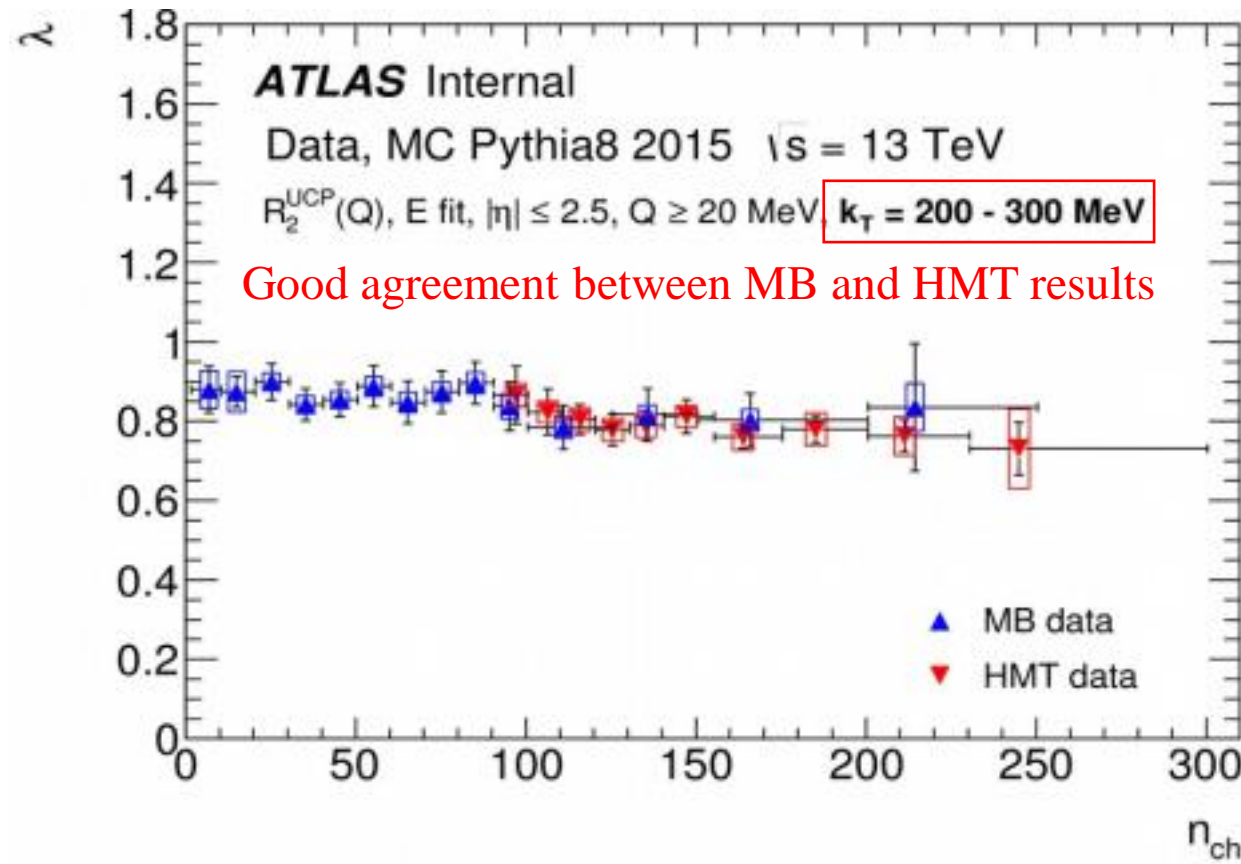


- The  $\lambda$  values are constant and decrease with  $k_T$  increasing
- The  $R$  values are decrease with  $k_T$  increasing
- The  $R$  values are increase as  $\sim \alpha \cdot n_{ch}^{1/3}$  with multiplicity



# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS

$\lambda$  distribution at 13 TeV MB + HMT triggers R distribution at 13 TeV



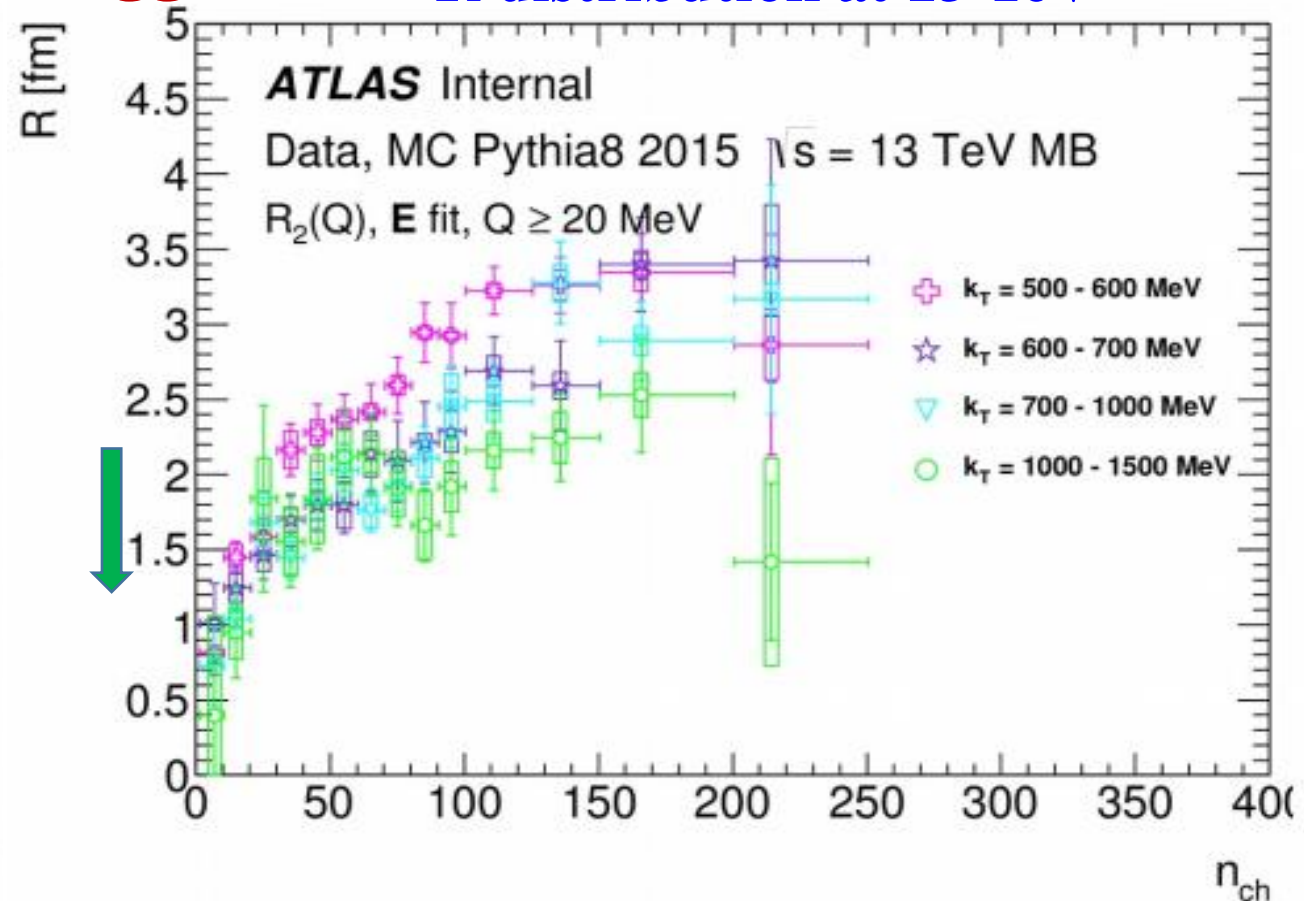
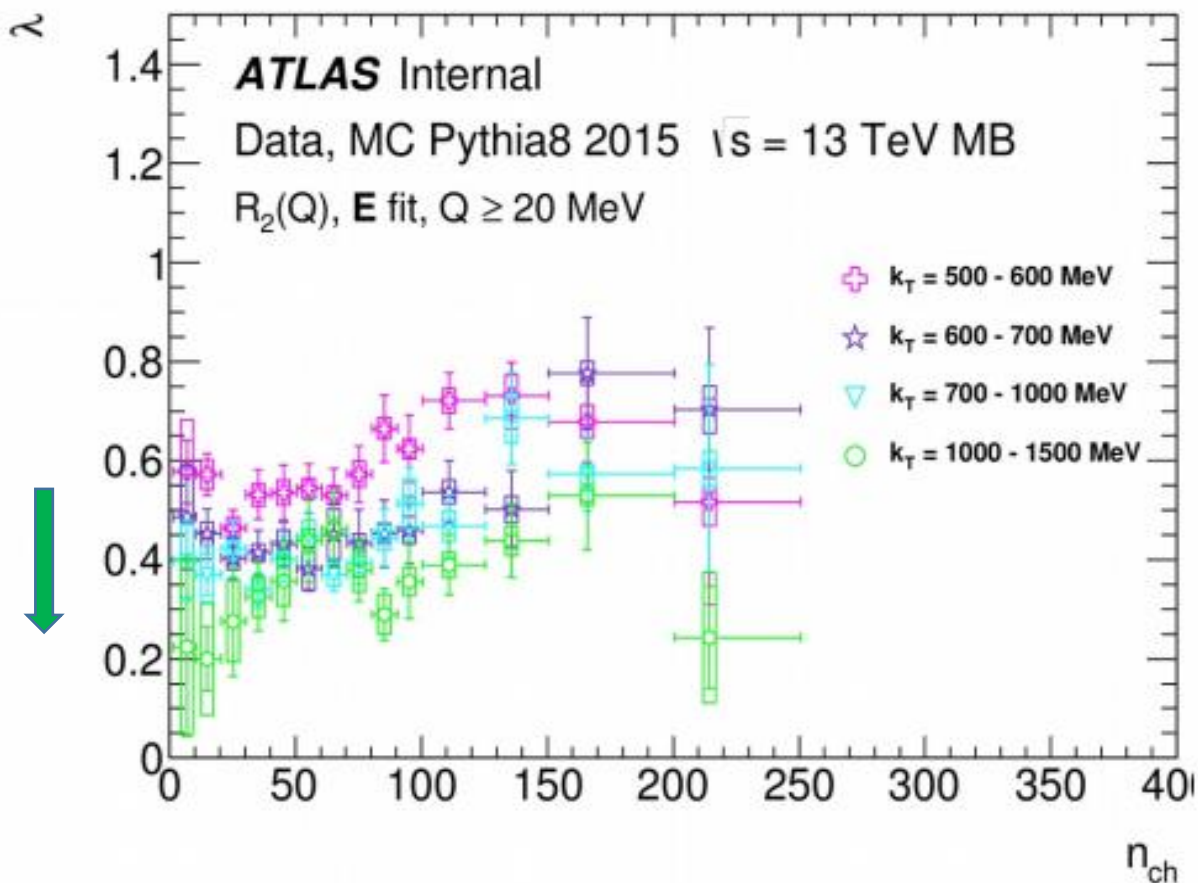
- Good connection between MB and HMT results
- The  $\lambda$  values slowly decreases in dependence of multiplicity
- The R values increases as  $\sim \alpha \cdot n_{ch}^{1/3}$  with multiplicity for  $n_{ch} < 100$
- The R are constant in dependence of multiplicity for  $n_{ch} > 100$

# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS FOR $(N_{ch}, K_T)$ INTERVALS - II

$\lambda$  distribution at 13 TeV

MB trigger

R distribution at 13 TeV



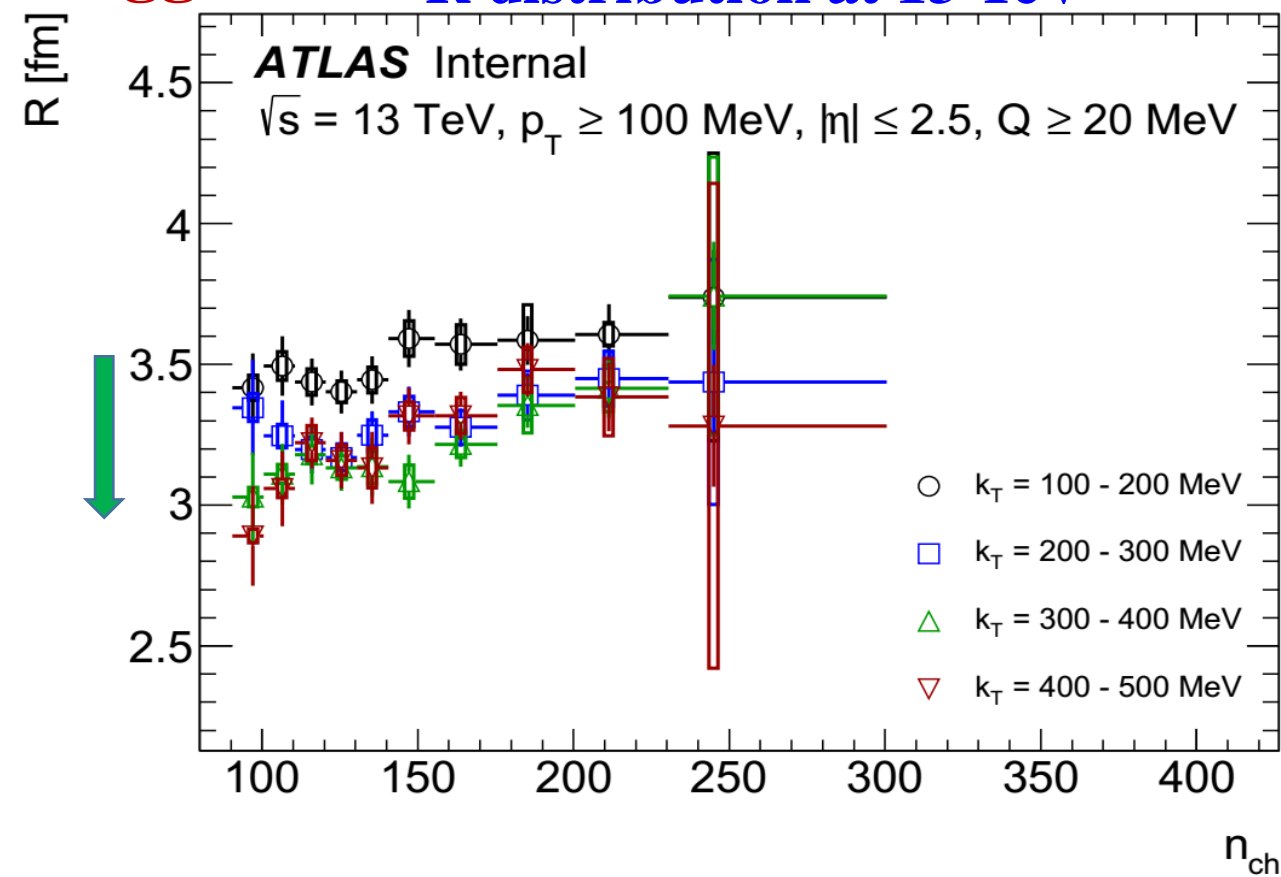
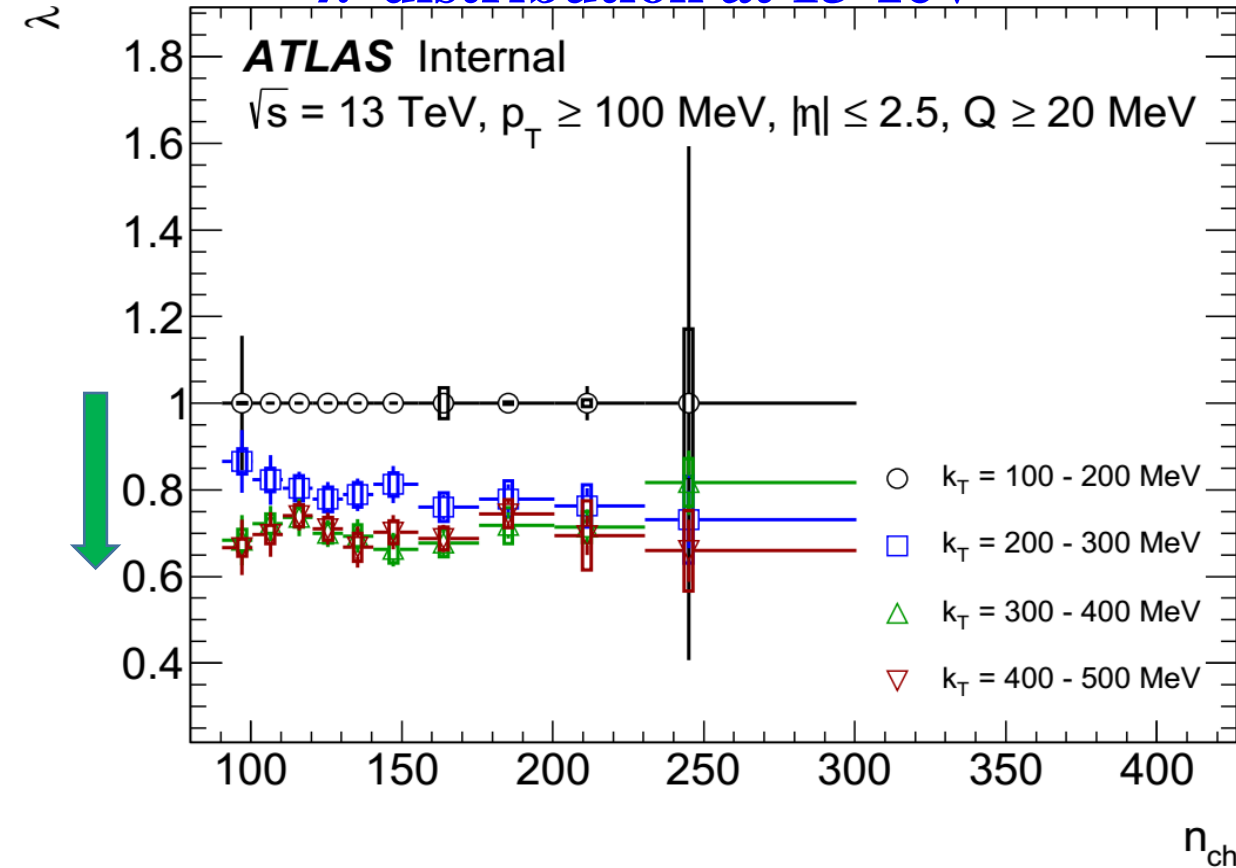
- The  $\lambda$  values are increase with multiplicity and decrease with  $k_T$  increasing
- The R values are decrease with  $k_T$  increasing
- The R values increases as  $\sim \alpha \cdot n_{ch}^{1/3}$  with multiplicity

# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS FOR $(N_{\text{CH}}, K_T)$ INTERVALS

## $\lambda$ distribution at 13 TeV

HMT trigger

## R distribution at 13 TeV



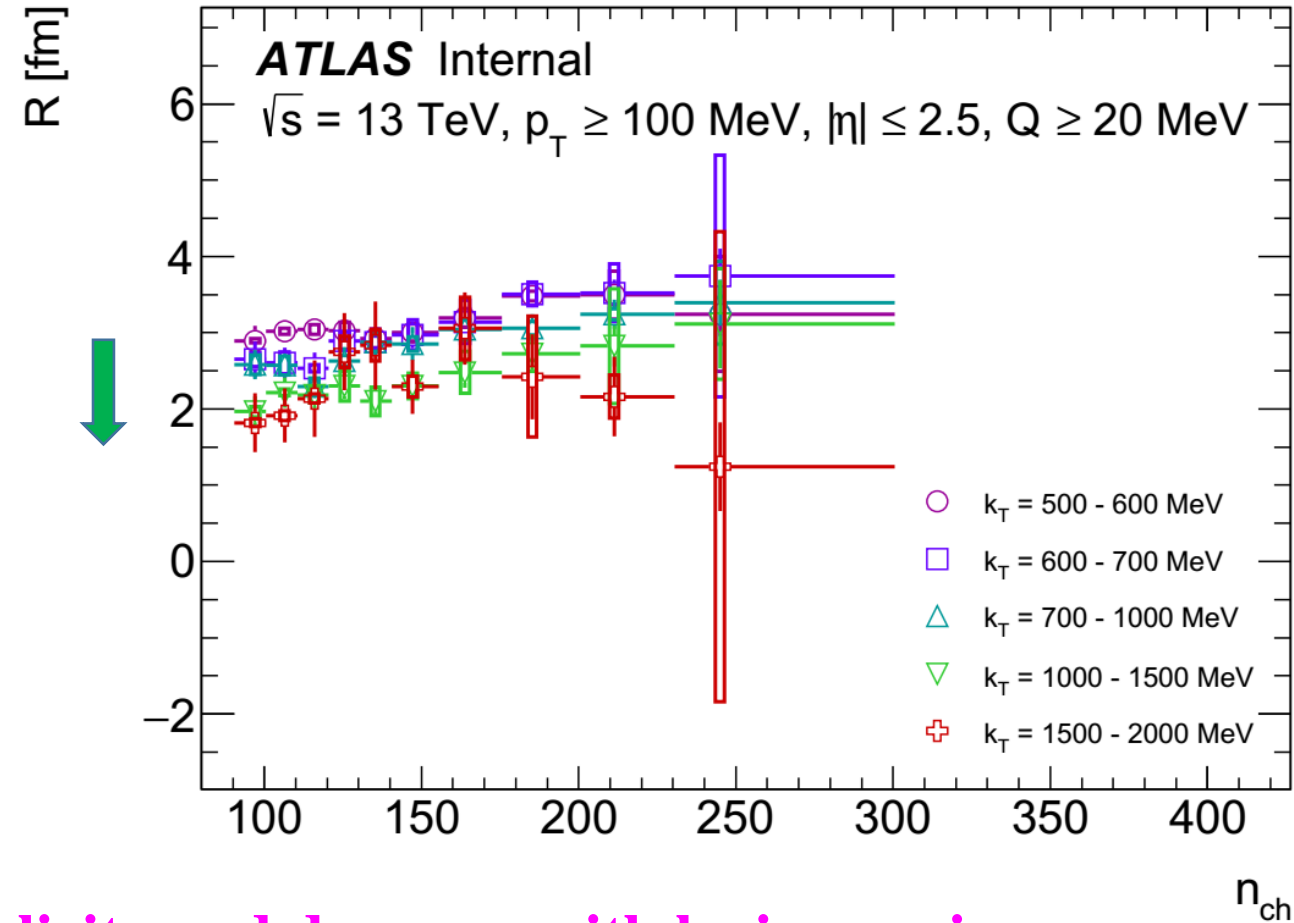
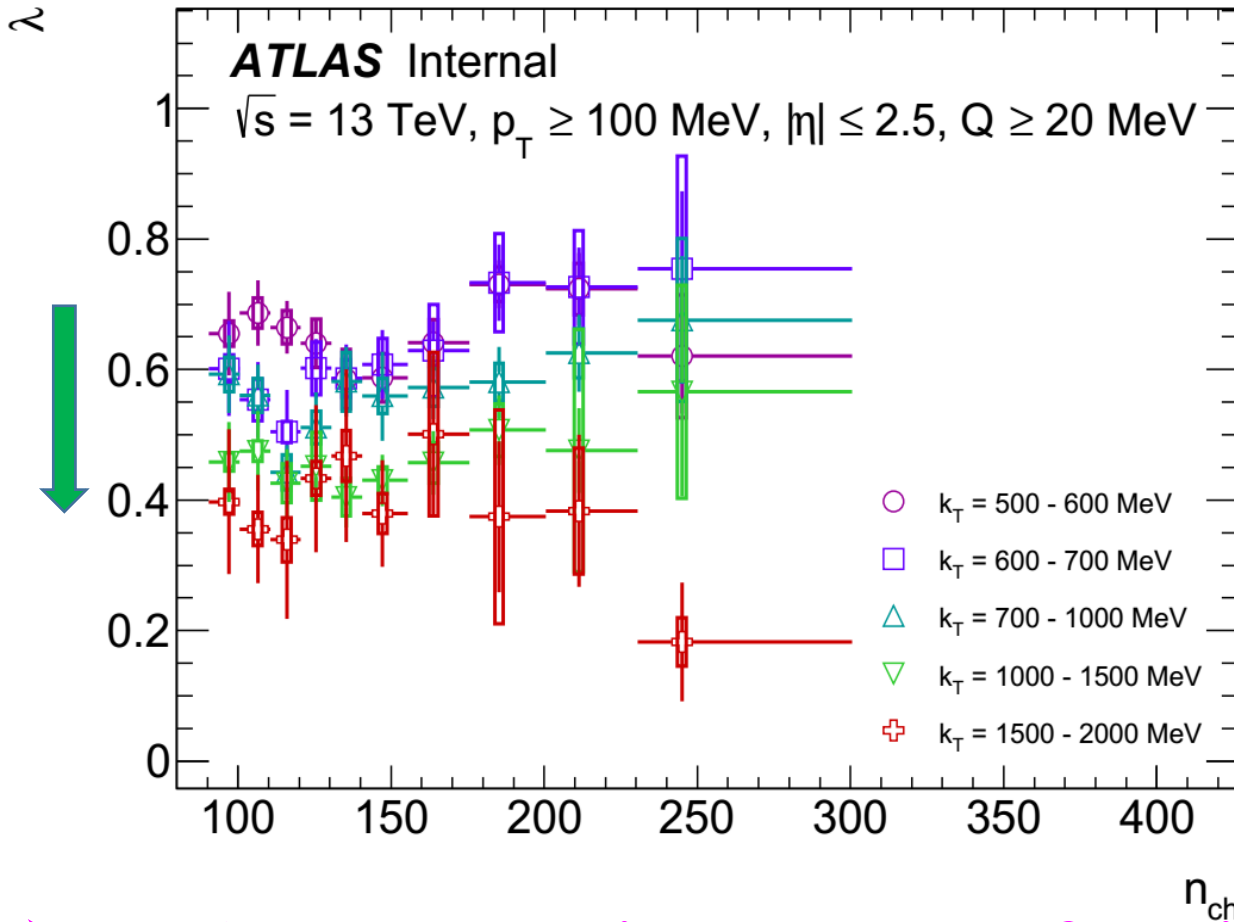
- The  $\lambda$  values are constant in dependence of multiplicity and decrease with  $k_T$  increasing
- The R values are constant in dependence of multiplicity and decrease with  $k_T$  increasing

# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS FOR $(N_{CH}, K_T)$ INTERVALS (CONT.)

$\lambda$  distribution at 13 TeV

HMT trigger

R distribution at 13 TeV



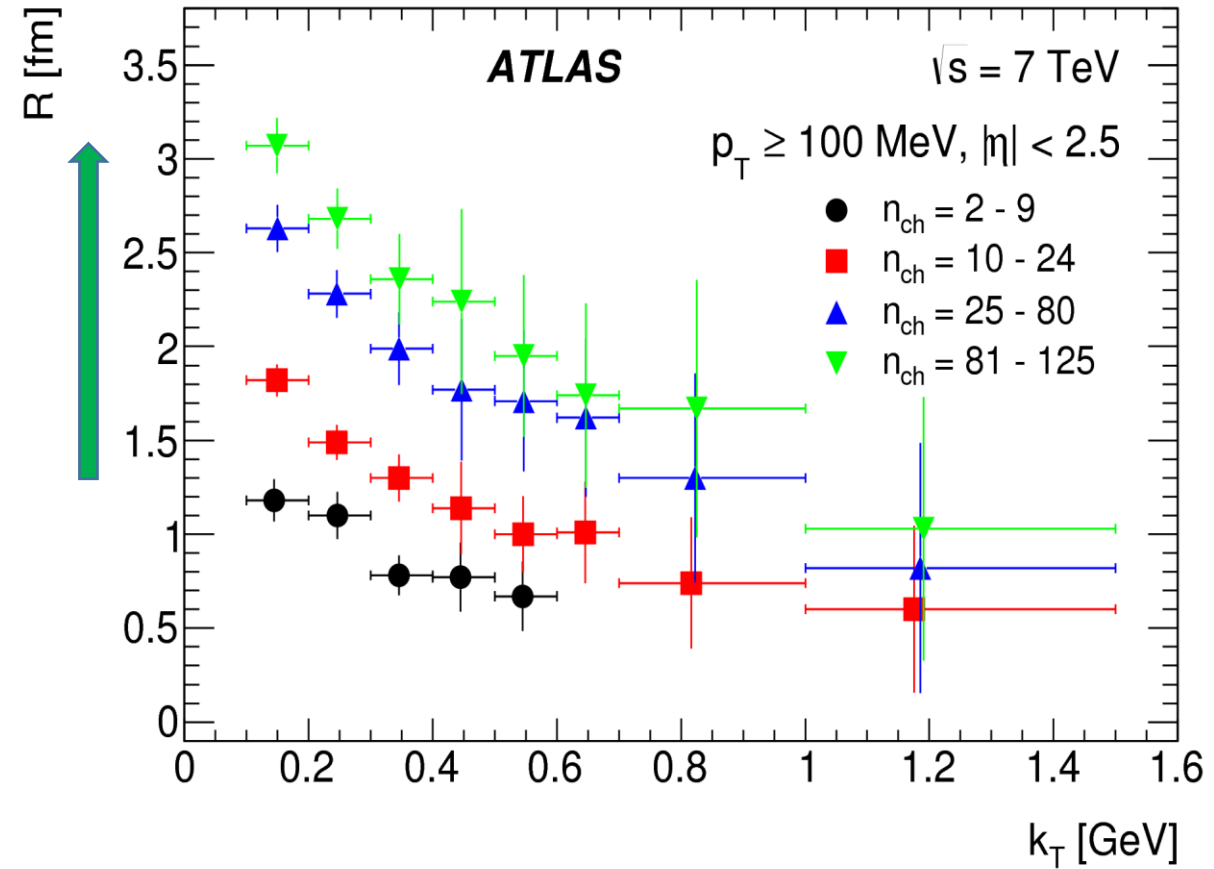
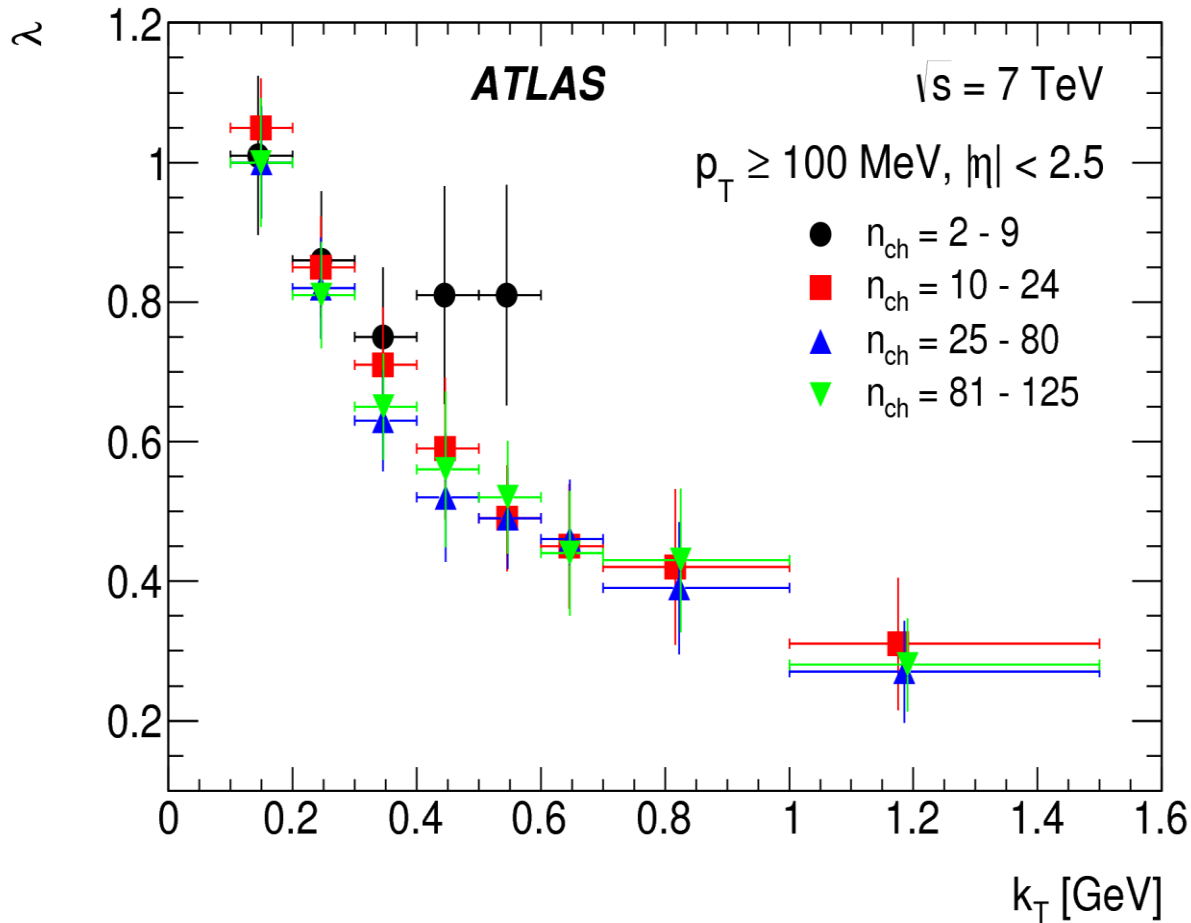
- The  $\lambda$  are constant in dependence of multiplicity and decrease with  $k_T$  increasing
- The R are constant in dependence of multiplicity and decrease with  $k_T$  increasing



$\lambda$  distribution at 7 TeV

MB trigger

R distribution at 7 TeV



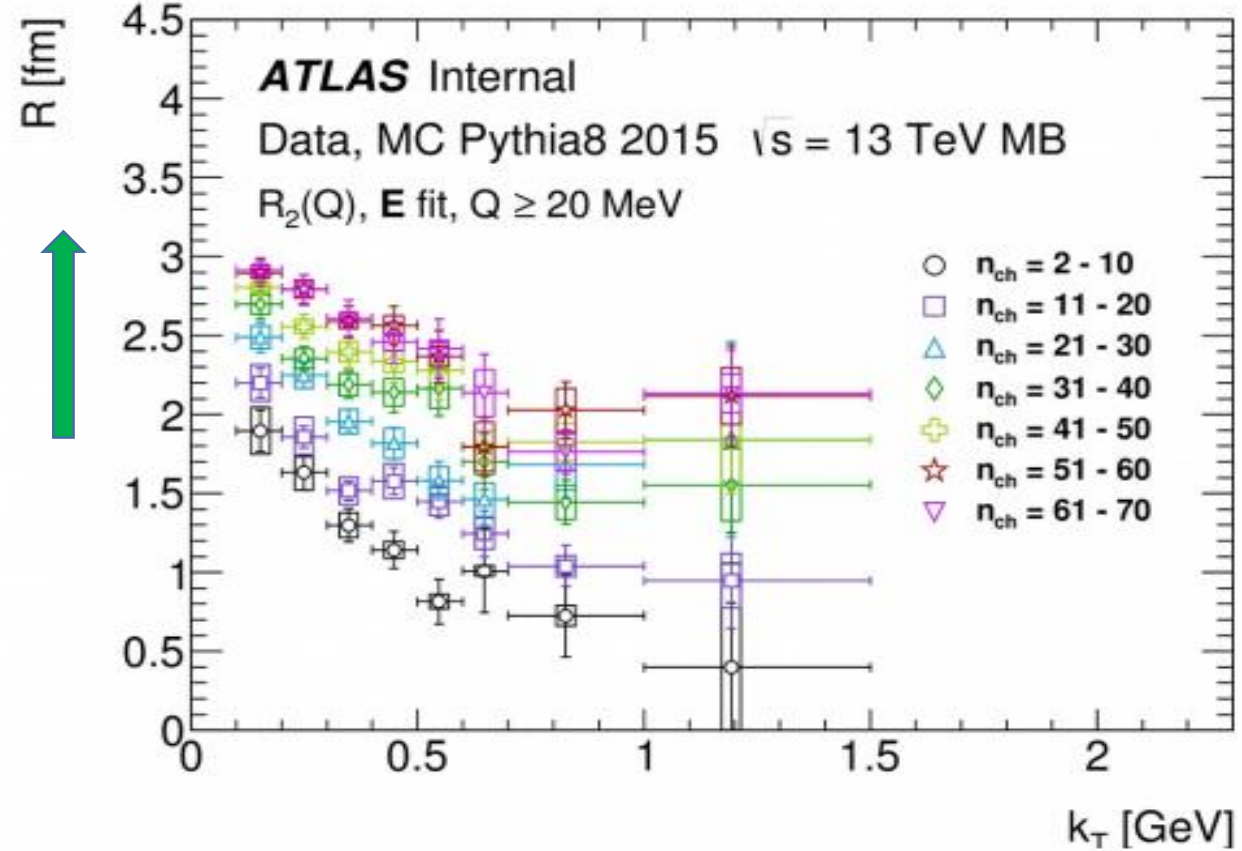
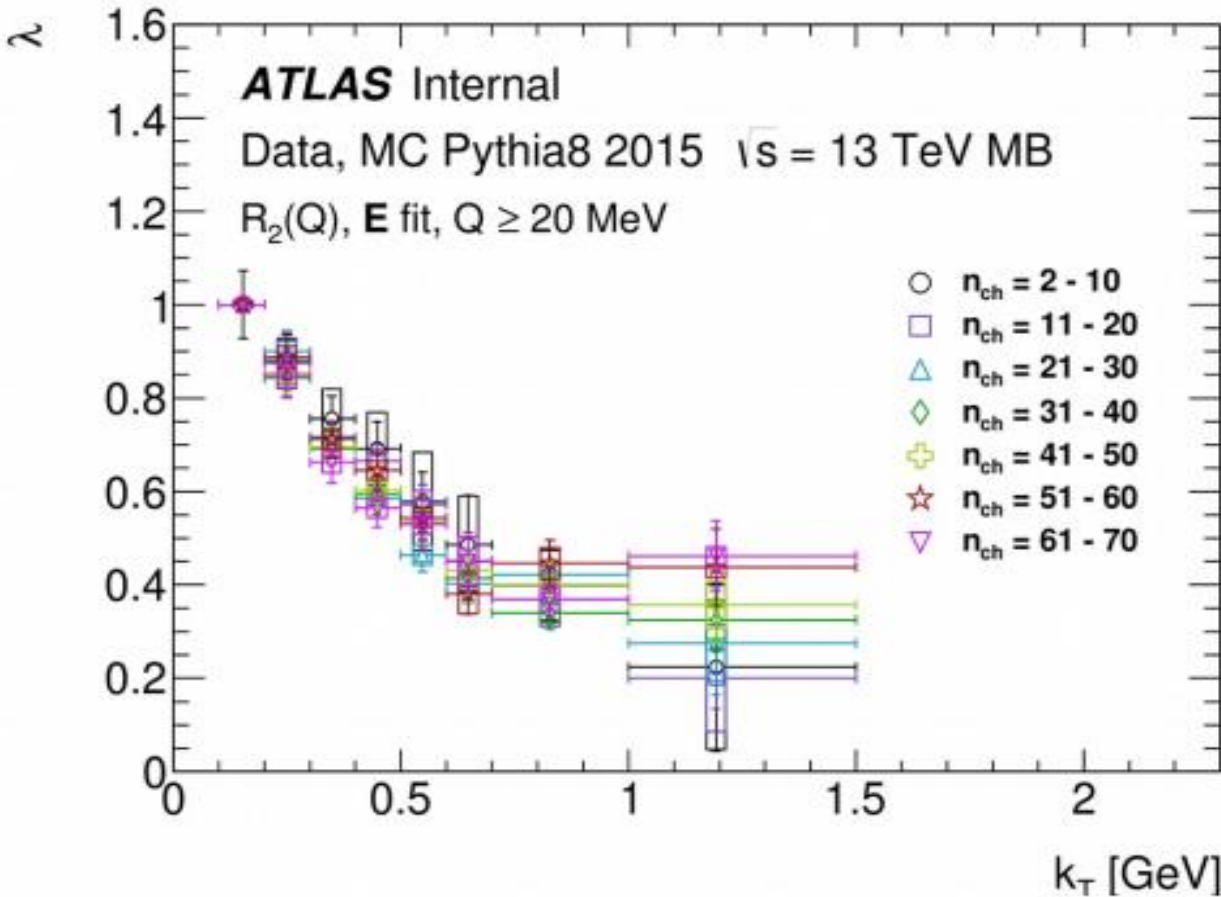
- The  $\lambda$  values are independent of multiplicity within uncertainties
- The R values increase with increasing multiplicity
- The  $\lambda$  and R values decrease exponentially with  $k_T$

# $K_T$ DEPENDENCE OF BEC PARAMETERS FOR $(N_{CH}, K_T)$ INTERVALS

$\lambda$  distribution at 13 TeV

MB trigger

R distribution at 13 TeV



- The  $\lambda$  values are independent of multiplicity within uncertainties
- The R values increase with increasing multiplicity
- The  $\lambda$  and R values decrease exponentially with  $k_T$

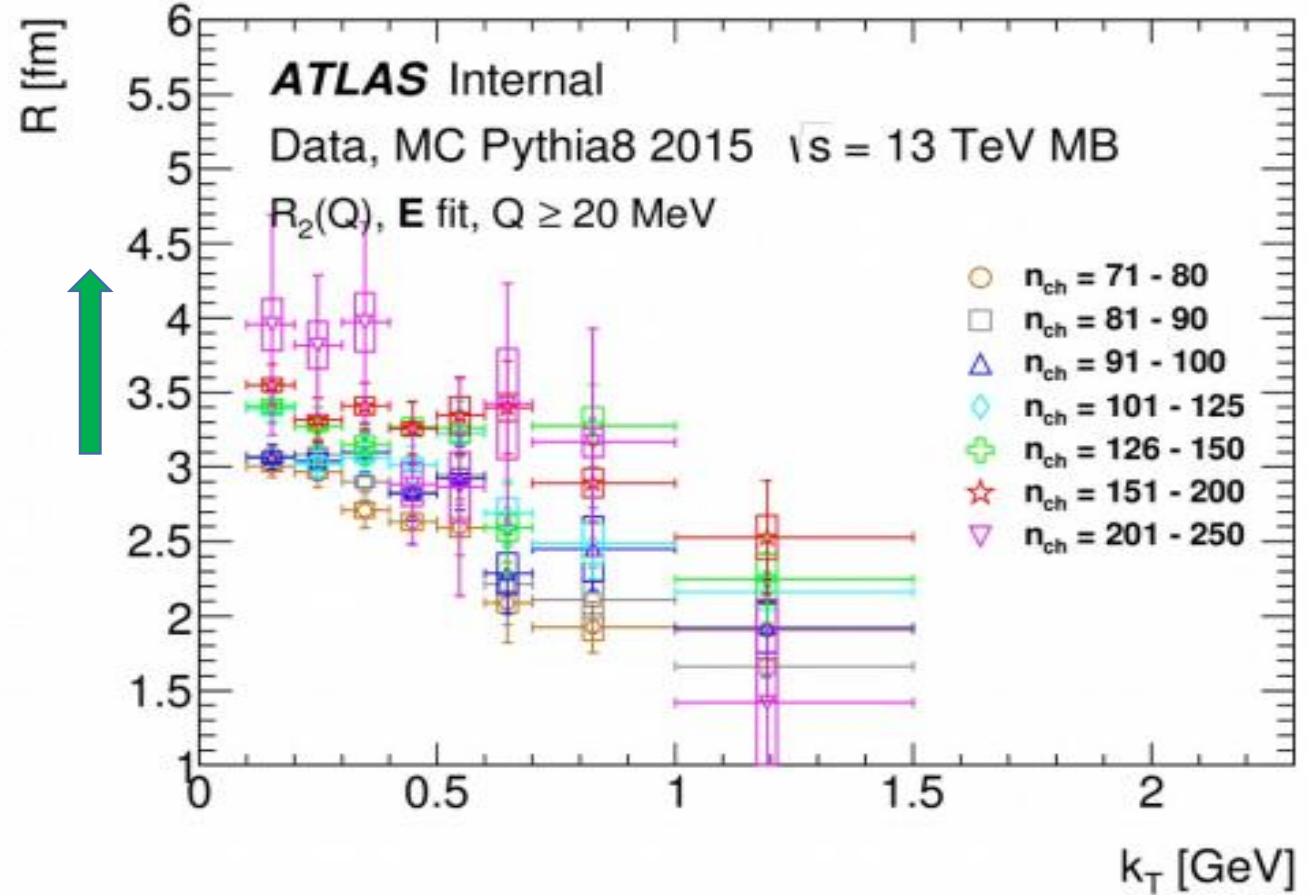
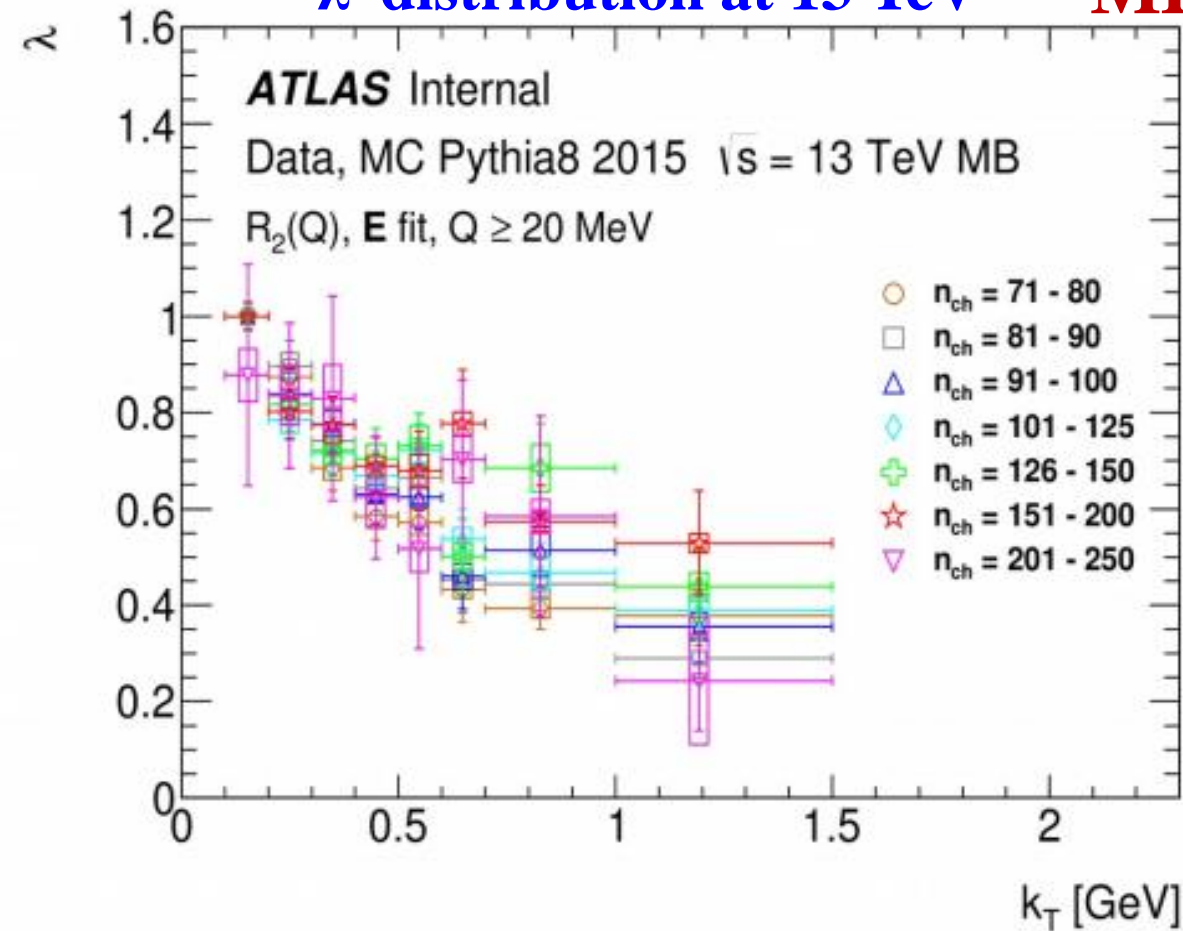
The effects similar to that observed for 7 TeV (see Backup slides)

# $K_T$ DEPENDENCE OF BEC PARAMETERS FOR $(N_{CH}, K_T)$ INTERVALS (CONT.)

$\lambda$  distribution at 13 TeV

MB trigger

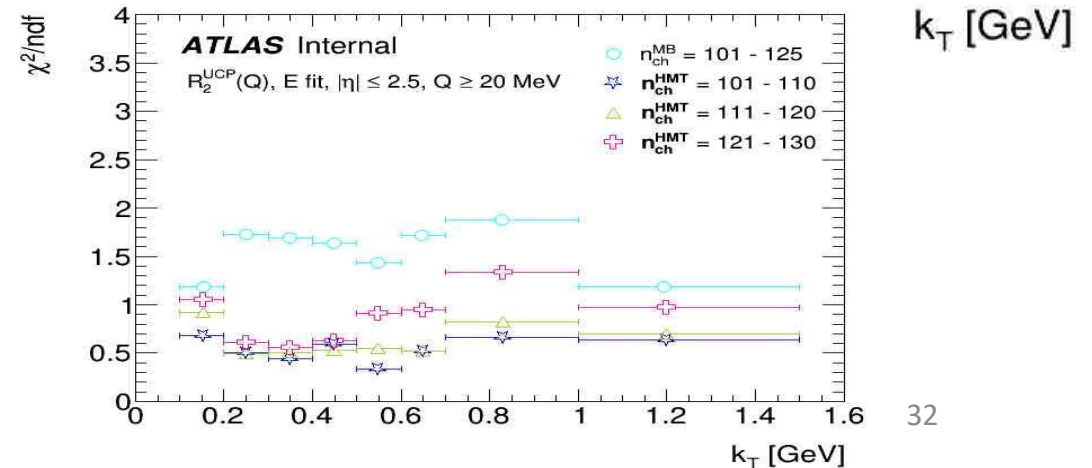
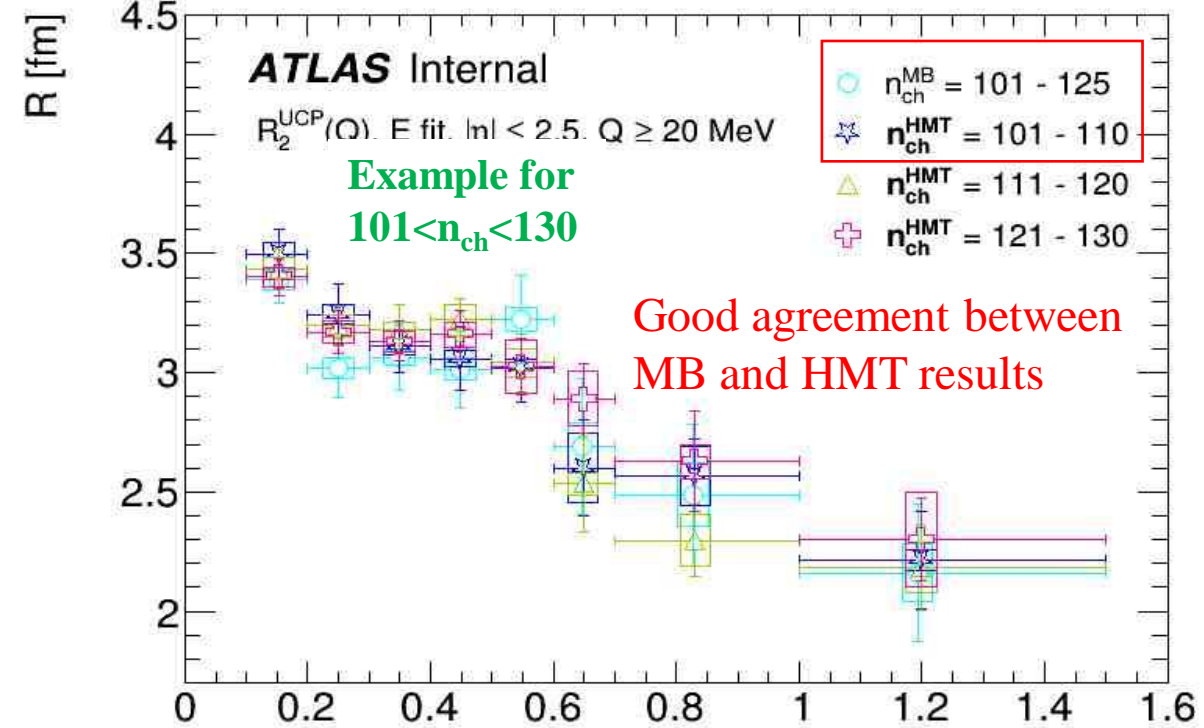
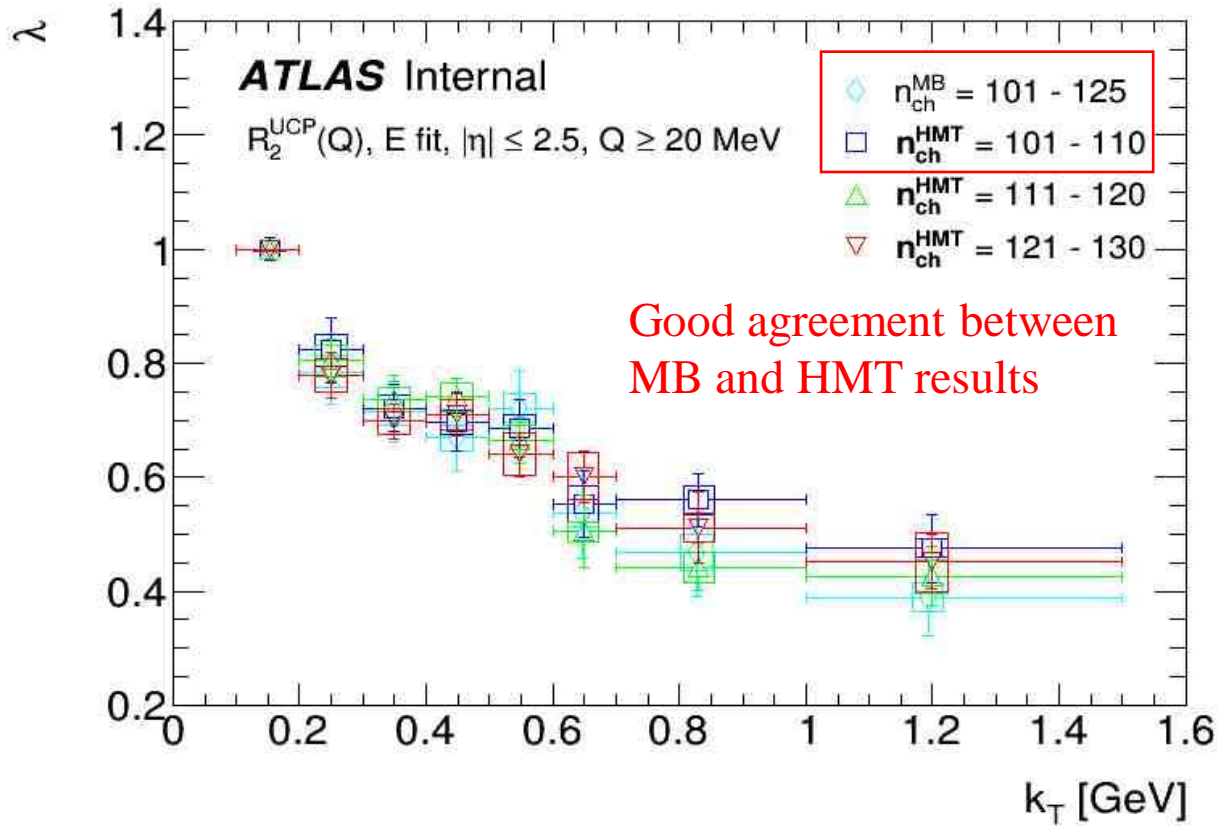
R distribution at 13 TeV



- The  $\lambda$  and R values are independent of multiplicity within uncertainties
- The  $\lambda$  and R values decrease exponentially with  $k_T$

# $K_T$ DEPENDENCE OF BEC PARAMETERS

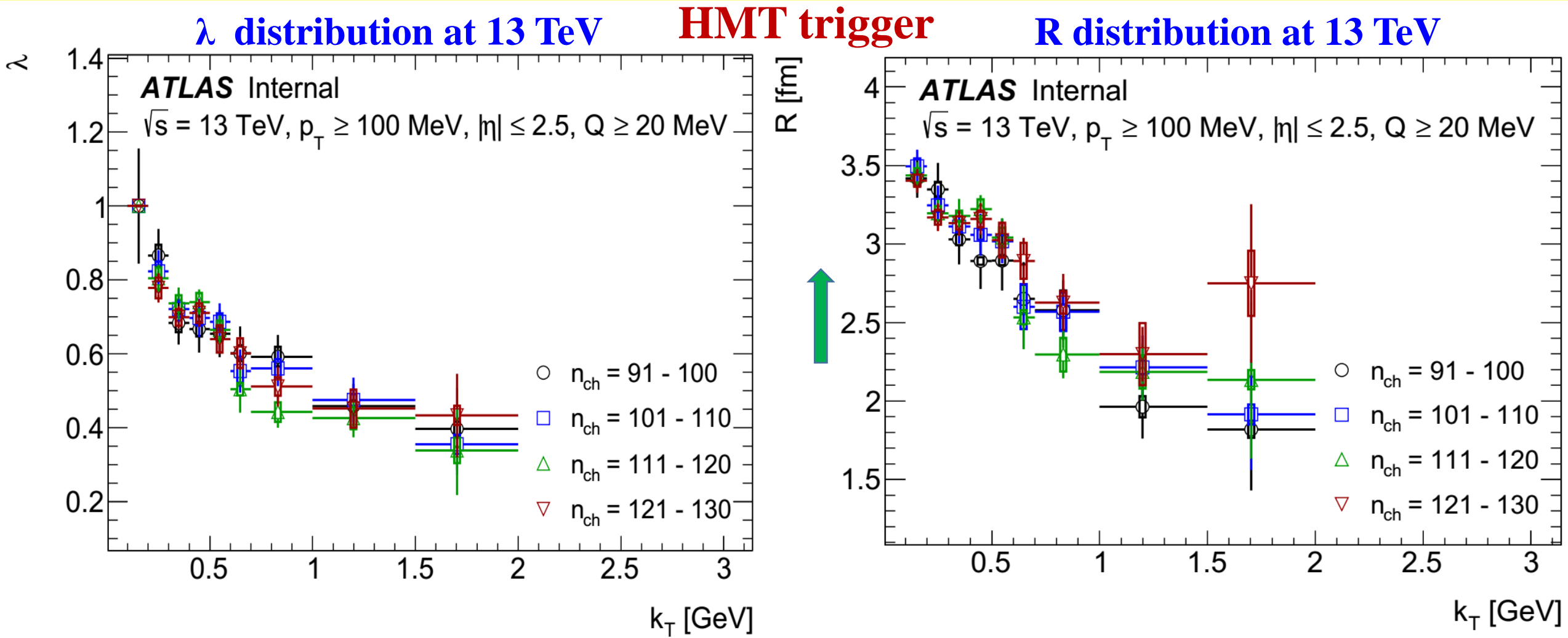
$\lambda$  distribution at 13 TeV MB + HMT triggers R distribution at 13 TeV



- Agreement between MB and HMT results
- The  $\lambda$  values decreases exponentially in dependence of  $k_T$
- The R values decreases exponentially in dependence of  $k_T$



# K<sub>T</sub> DEPENDENCE OF BEC PARAMETERS FOR (N<sub>CH</sub>, K<sub>T</sub>) INTERVALS



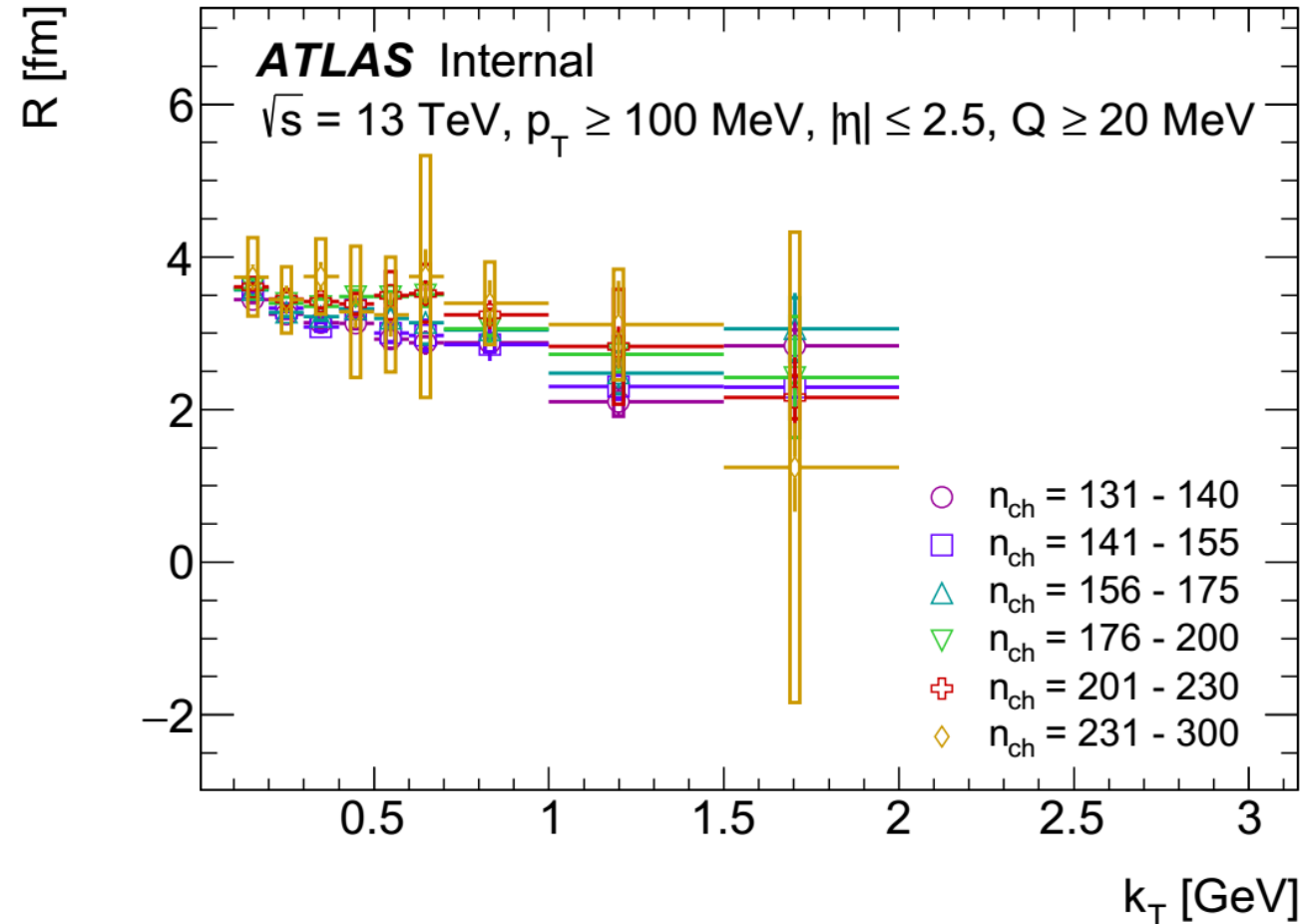
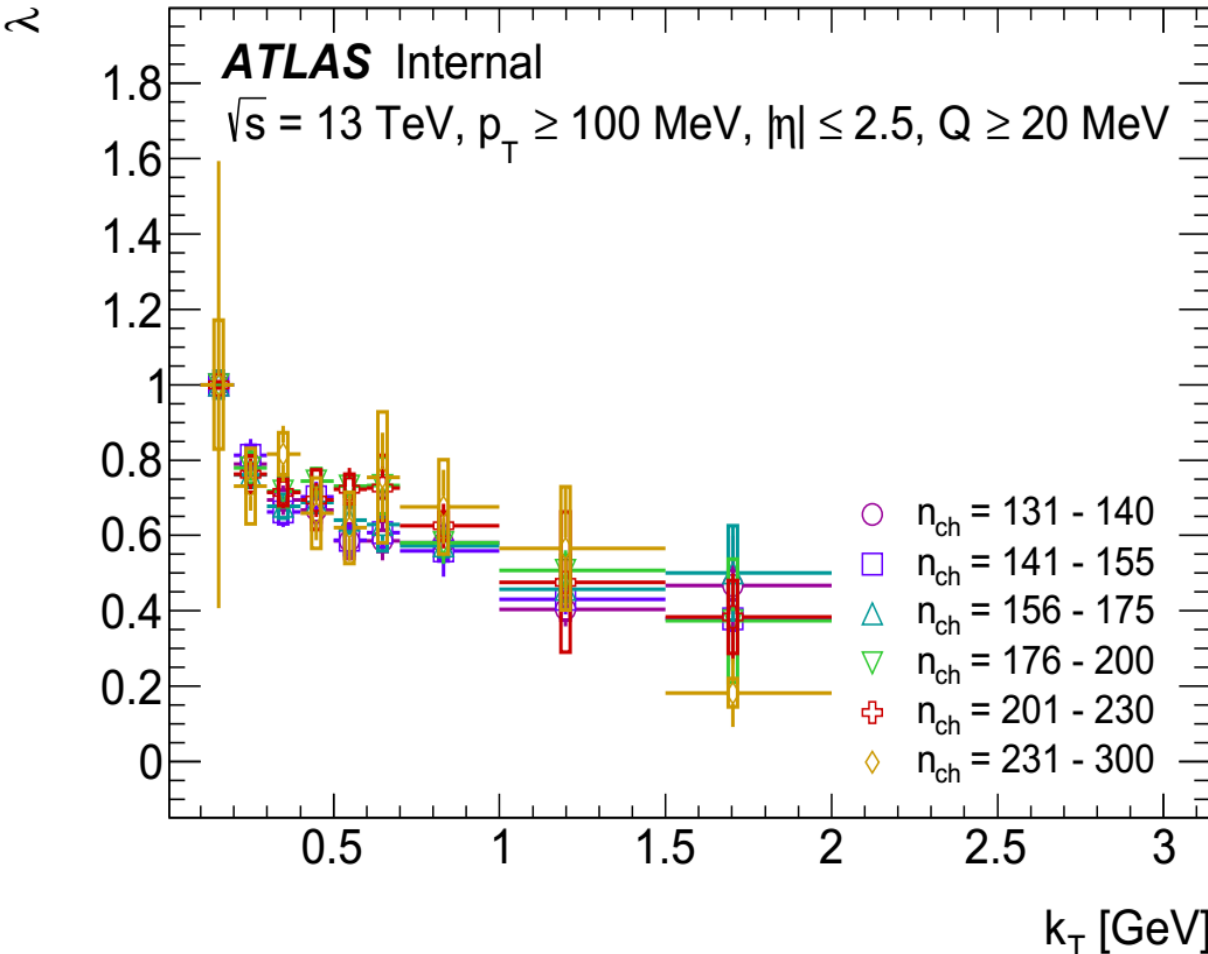
- The  $\lambda$  and  $R$  values are independent of multiplicity within uncertainties
- The  $\lambda$  and  $R$  values decrease exponentially with  $k_T$

# $K_T$ DEPENDENCE OF BEC PARAMETERS FOR $(N_{CH}, K_T)$ INTERVALS (CONT.)

$\lambda$  distribution at 13 TeV

HMT trigger

R distribution at 13 TeV



- The  $\lambda$  and R values are independent of multiplicity within uncertainties
- The  $\lambda$  and R values decrease exponentially with  $k_T$

# CONCLUSIONS

- ATLAS data taking in Run 2 is going very well. Good data taking efficiency and Data Quality
- A rich set of new physics results:
  - ❖ Evidence for Higgs bosons decaying to a pair of b-quarks
  - ❖ Evidence for Zt production, rate consistent with the SM
  - ❖ A first study of Bose-Einstein correlations in pp collisions at 13 TeV
    - ✓ Confirmation of the BEC radius saturation effect for high multiplicity region. The feature for the first time observed by ATLAS at 7 TeV
    - ✓ Energy dependence of  $\lambda$

**THANK YOU VERY MUCH FOR ATTENTION!**



# BACKUP SLIDES

# EXOTIC SEARCHES

## ATLAS Exotics Searches\* - 95% CL Upper Exclusion Limits

Status: July 2017

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 37.0) \text{ fb}^{-1}$$

$$\sqrt{s} = 8, 13 \text{ TeV}$$

	Model	$\ell, \gamma$	Jets <sup>†</sup>	$E_{\text{T}}^{\text{miss}}$	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit		Reference	
Extra dimensions	ADD $G_{KK} + g/q$	$0 e, \mu$	1-4 j	Yes	36.1	$M_D$	7.75 TeV	$n = 2$	ATLAS-CONF-2017-060
	ADD non-resonant $\gamma\gamma$	$2 \gamma$	-	-	36.7	$M_S$	8.6 TeV	$n = 3$ HLZ NLO	CERN-EP-2017-132
	ADD QBH	-	2 j	-	37.0	$M_{\text{th}}$	8.9 TeV	$n = 6$	1703.09217
	ADD BH high $\sum p_T$	$\geq 1 e, \mu$	$\geq 2 j$	-	3.2	$M_{\text{th}}$	8.2 TeV	$n = 6, M_D = 3 \text{ TeV}$ , rot BH	1606.02265
	ADD BH multijet	-	$\geq 3 j$	-	3.6	$M_{\text{th}}$	9.55 TeV	$n = 6, M_D = 3 \text{ TeV}$ , rot BH	1512.02586
	RS1 $G_{KK} \rightarrow \gamma\gamma$	$2 \gamma$	-	-	36.7	$G_{KK}$ mass	4.1 TeV	$k/\bar{M}_{\text{Pl}} = 0.1$	CERN-EP-2017-132
	Bulk RS $G_{KK} \rightarrow WW \rightarrow qq\ell\nu$	$1 e, \mu$	1 J	Yes	36.1	$G_{KK}$ mass	1.75 TeV	$k/\bar{M}_{\text{Pl}} = 1.0$	ATLAS-CONF-2017-051
	2UED / RPP	$1 e, \mu$	$\geq 2 b, \geq 3 j$	Yes	13.2	KK mass	1.6 TeV	Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow tt) = 1$	ATLAS-CONF-2016-104
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2 e, \mu$	-	-	36.1	$Z'$ mass	4.5 TeV		ATLAS-CONF-2017-027
	SSM $Z' \rightarrow \tau\tau$	$2 \tau$	-	-	36.1	$Z'$ mass	2.4 TeV		ATLAS-CONF-2017-050
	Leptophobic $Z' \rightarrow bb$	-	2 b	-	3.2	$Z'$ mass	1.5 TeV		1603.08791
	Leptophobic $Z' \rightarrow tt$	$1 e, \mu$	$\geq 1 b, \geq 1J/2j$	Yes	3.2	$Z'$ mass	2.0 TeV	$\Gamma/m = 3\%$	ATLAS-CONF-2016-014
	SSM $W' \rightarrow \ell\nu$	$1 e, \mu$	-	Yes	36.1	$W'$ mass	5.1 TeV		1706.04786
	HVT $V' \rightarrow WV \rightarrow qq\bar{q}q$ model B	$0 e, \mu$	2 J	-	36.7	$V'$ mass	3.5 TeV	$g_V = 3$	CERN-EP-2017-147
CI	HVT $V' \rightarrow WH/ZH$ model B	multi-channel	-	-	36.1	$V'$ mass	2.93 TeV	$g_V = 3$	ATLAS-CONF-2017-055
	LRSM $W'_R \rightarrow tb$	$1 e, \mu$	2 b, 0-1 j	Yes	20.3	$W'_R$ mass	1.92 TeV		1410.4103
	LRSM $W'_R \rightarrow tb$	$0 e, \mu$	$\geq 1 b, 1 J$	-	20.3	$W'_R$ mass	1.76 TeV		1408.0886
	CI $qq\bar{q}q$	-	2 j	-	37.0	$\Lambda$	21.8 TeV	$\eta_{LL}^-$	1703.09217
DM	CI $\ell\ell q\bar{q}$	$2 e, \mu$	-	-	36.1	$\Lambda$	40.1 TeV	$\eta_{LL}^-$	ATLAS-CONF-2017-027
	CI $u\bar{u}t\bar{t}$	$2(SS)/\geq 3 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	20.3	$\Lambda$	4.9 TeV	$ C_{RR}  = 1$	1504.04605
	Axial-vector mediator (Dirac DM)	$0 e, \mu$	1-4 j	Yes	36.1	$m_{\text{med}}$	1.5 TeV	$g_\gamma = 0.25, g_t = 1.0, m(\chi) < 400 \text{ GeV}$	ATLAS-CONF-2017-060
LQ	Vector mediator (Dirac DM)	$0 e, \mu, 1 \gamma$	$\leq 1 j$	Yes	36.1	$m_{\text{med}}$	1.2 TeV	$g_\gamma = 0.25, g_t = 1.0, m(\chi) < 480 \text{ GeV}$	1704.03848
	$VV_{\chi\chi}$ EFT (Dirac DM)	$0 e, \mu$	1 J, $\leq 1 j$	Yes	3.2	$M_\chi$	700 GeV	$m(\chi) < 150 \text{ GeV}$	1608.02372
	Scalar LQ 1 <sup>st</sup> gen	$2 e$	$\geq 2 j$	-	3.2	LQ mass	1.1 TeV	$\beta = 1$	1605.06035
Heavy quarks	Scalar LQ 2 <sup>nd</sup> gen	$2 \mu$	$\geq 2 j$	-	3.2	LQ mass	1.05 TeV	$\beta = 1$	1605.06035
	Scalar LQ 3 <sup>rd</sup> gen	$1 e, \mu$	$\geq 1 b, \geq 3 j$	Yes	20.3	LQ mass	640 GeV	$\beta = 0$	1508.04735
	VLQ $TT \rightarrow Ht + X$	$0$ or $1 e, \mu$	$\geq 2 b, \geq 3 j$	Yes	13.2	T mass	1.2 TeV	$\mathcal{B}(T \rightarrow Ht) = 1$	ATLAS-CONF-2016-104
Excited fermions	VLQ $TT \rightarrow Zt + X$	$1 e, \mu$	$\geq 1 b, \geq 3 j$	Yes	36.1	T mass	1.16 TeV	$\mathcal{B}(T \rightarrow Zt) = 1$	1705.10751
	VLQ $TT \rightarrow Wb + X$	$1 e, \mu$	$\geq 1 b, \geq 1J/2j$	Yes	36.1	T mass	1.35 TeV	$\mathcal{B}(T \rightarrow Wb) = 1$	CERN-EP-2017-094
	VLQ $BB \rightarrow Hb + X$	$1 e, \mu$	$\geq 2 b, \geq 3 j$	Yes	20.3	B mass	700 GeV	$\mathcal{B}(B \rightarrow Hb) = 1$	1505.04306
	VLQ $BB \rightarrow Zb + X$	$2/\geq 3 e, \mu$	$\geq 2/\geq 1 b$	-	20.3	B mass	790 GeV	$\mathcal{B}(B \rightarrow Zb) = 1$	1409.5500
	VLQ $BB \rightarrow Wt + X$	$1 e, \mu$	$\geq 1 b, \geq 1J/2j$	Yes	36.1	B mass	1.25 TeV	$\mathcal{B}(B \rightarrow Wt) = 1$	CERN-EP-2017-094
	VLQ $QQ \rightarrow WqWq$	$1 e, \mu$	$\geq 4 j$	Yes	20.3	Q mass	690 GeV		1509.04261
Other	Excited quark $q^* \rightarrow qg$	-	2 j	-	37.0	$q^*$ mass	6.0 TeV	only $u^*$ and $d^*$ , $\Lambda = m(q^*)$	1703.09127
	Excited quark $q^* \rightarrow q\gamma$	$1 \gamma$	1 j	-	36.7	$q^*$ mass	5.3 TeV	only $u^*$ and $d^*$ , $\Lambda = m(q^*)$	CERN-EP-2017-148
	Excited quark $b^* \rightarrow bg$	-	1 b, 1 j	-	13.3	$b^*$ mass	2.3 TeV		ATLAS-CONF-2016-060
	Excited quark $b^* \rightarrow Wt$	$1$ or $2 e, \mu$	1 b, 2-0 j	Yes	20.3	$b^*$ mass	1.5 TeV	$f_E = f_L = f_R = 1$	1510.02664
	Excited lepton $\ell^*$	$3 e, \mu$	-	-	20.3	$\ell^*$ mass	3.0 TeV	$\Lambda = 3.0 \text{ TeV}$	1411.2921
	Excited lepton $\nu^*$	$3 e, \mu, \tau$	-	-	20.3	$\nu^*$ mass	1.6 TeV	$\Lambda = 1.6 \text{ TeV}$	1411.2921
Other	LRSM Majorana $\nu$	$2 e, \mu$	2 j	-	20.3	$N^0$ mass	2.0 TeV	$m(W_R) = 2.4 \text{ TeV}$ , no mixing	1506.06020
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$	$2, 3, 4 e, \mu$ (SS)	-	-	36.1	$H^{\pm\pm}$ mass	870 GeV	DY production	ATLAS-CONF-2017-053
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$	$3 e, \mu, \tau$	-	-	20.3	$H^{\pm\pm}$ mass	400 GeV	DY production, $\mathcal{B}(H_L^{\pm\pm} \rightarrow \ell\tau) = 1$	1411.2921
	Monotop (non-res prod)	$1 e, \mu$	1 b	Yes	20.3	spin-1 invisible particle mass	657 GeV	$a_{\text{non-res}} = 0.2$	1410.5404
	Multi-charged particles	-	-	-	20.3	multi-charged particle mass	785 GeV	DY production, $ q  = 5e$	1504.04188
	Magnetic monopoles	-	-	-	7.0	monopole mass	1.34 TeV	DY production, $ g  = 1g_0$ , spin 1/2	1509.08059

$\sqrt{s} = 8 \text{ TeV}$        $\sqrt{s} = 13 \text{ TeV}$

Mass scale [TeV]

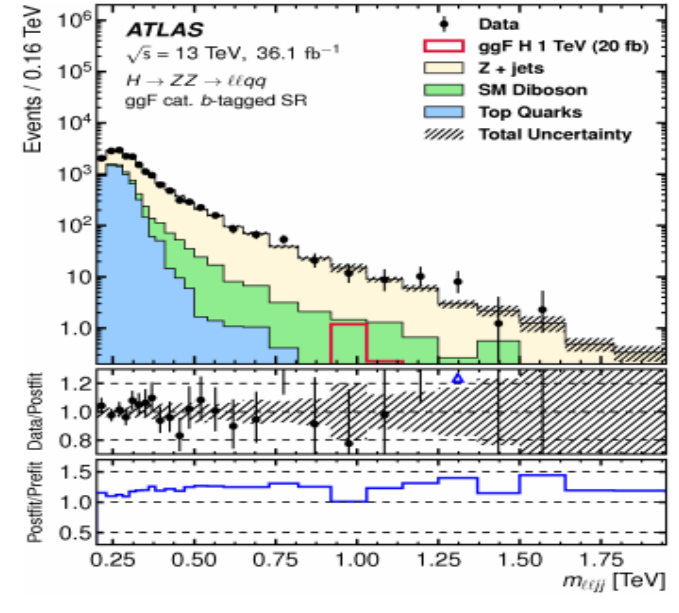
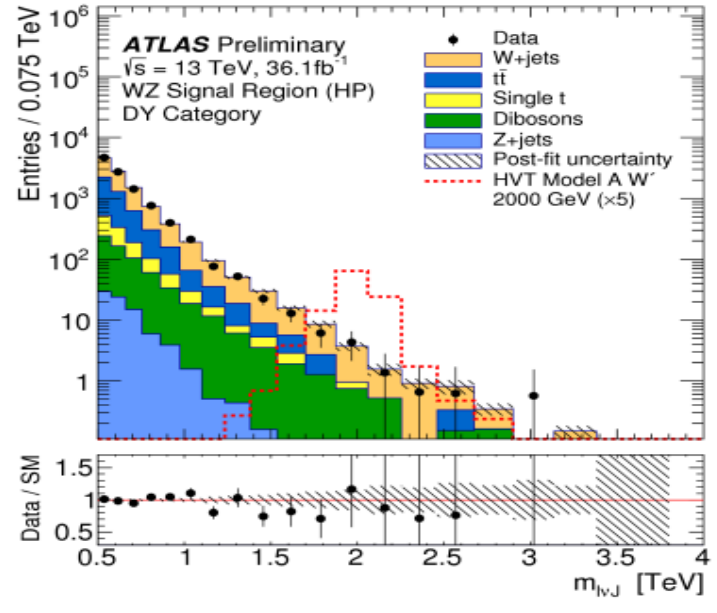
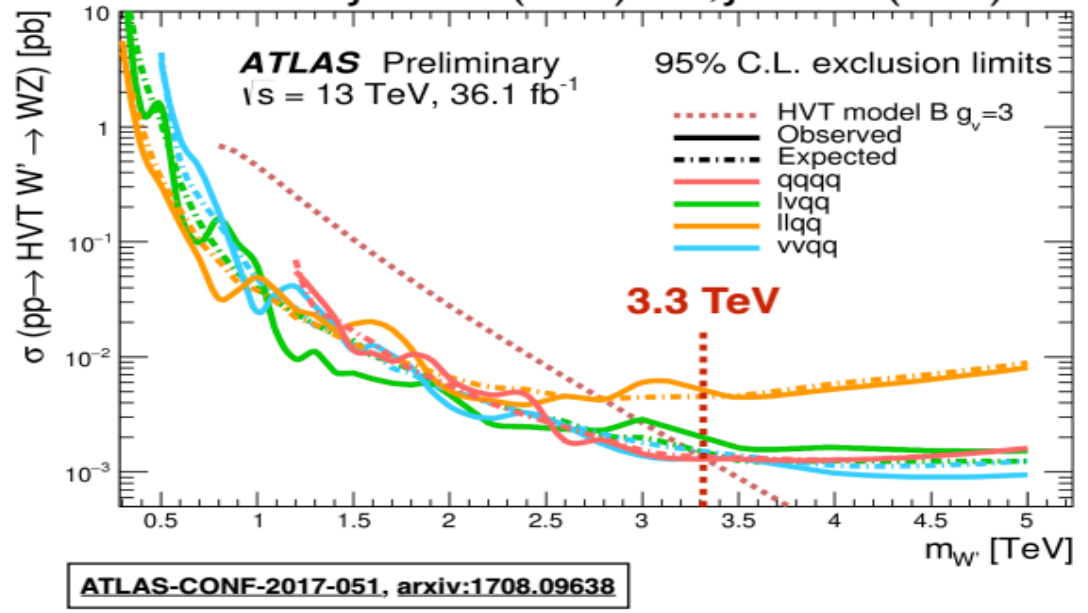
\*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).

# RESONANCE SEARCHES (DIBOSON)

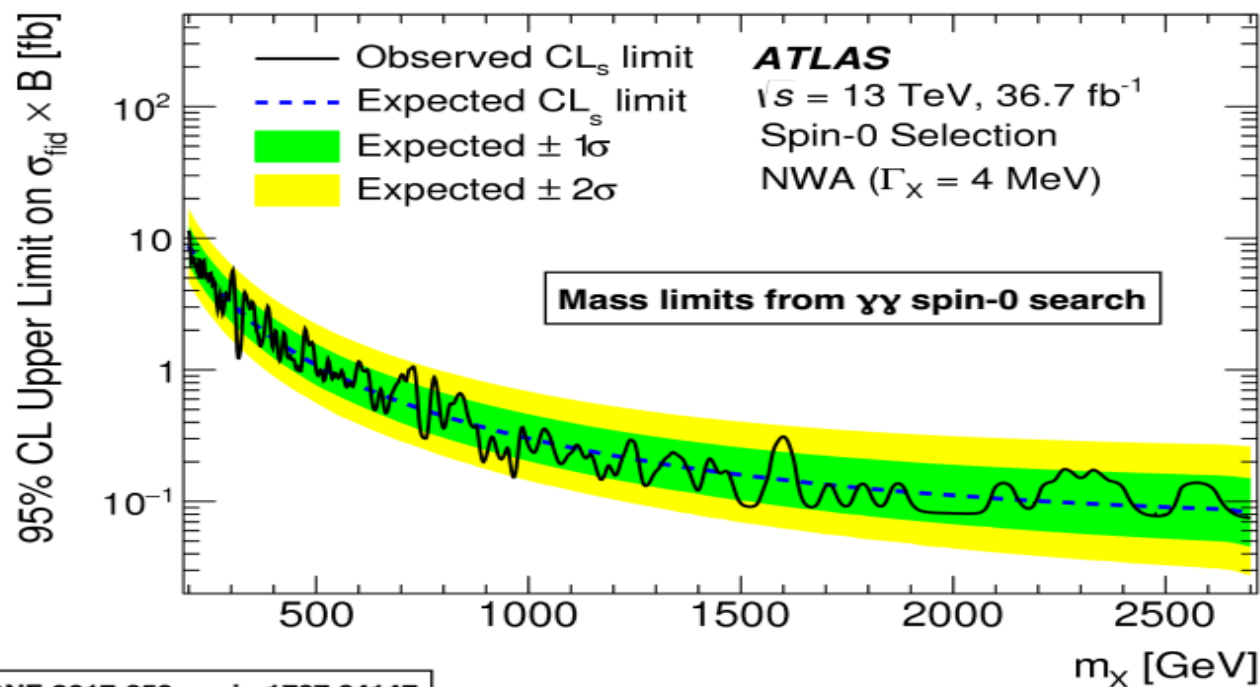
- Searches for new resonances decaying to a pair of W, Z or H bosons. At the time of the **LHCC in May**, only the  $VH \rightarrow qqbb$  result was ready. Now:
  - $VV \rightarrow qqqq/qq\ell\nu/qq\ell\ell/qq\nu\nu, VH \rightarrow qqbb/\ell\nu bb/\ell\ell bb$ .
- Reconstruct merged jets at high  $p_T$  for resonances at the highest mass, using substructure techniques ('boson-tagging').
- No significant excess observed over the SM expectations.
- Set limits in framework of Heavy Vector Triplet (HVT) model\*.

\* JHEP 09 (2014) 060, JHEP 01 (2013) 166

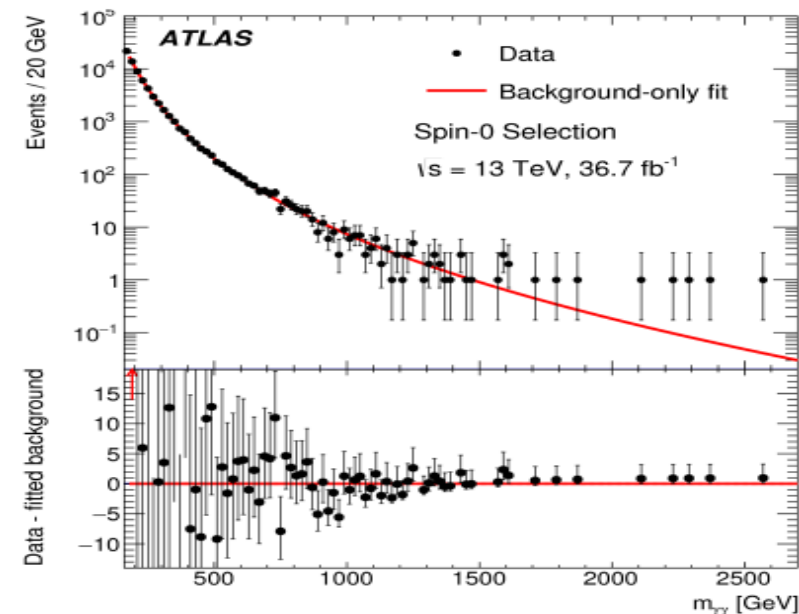
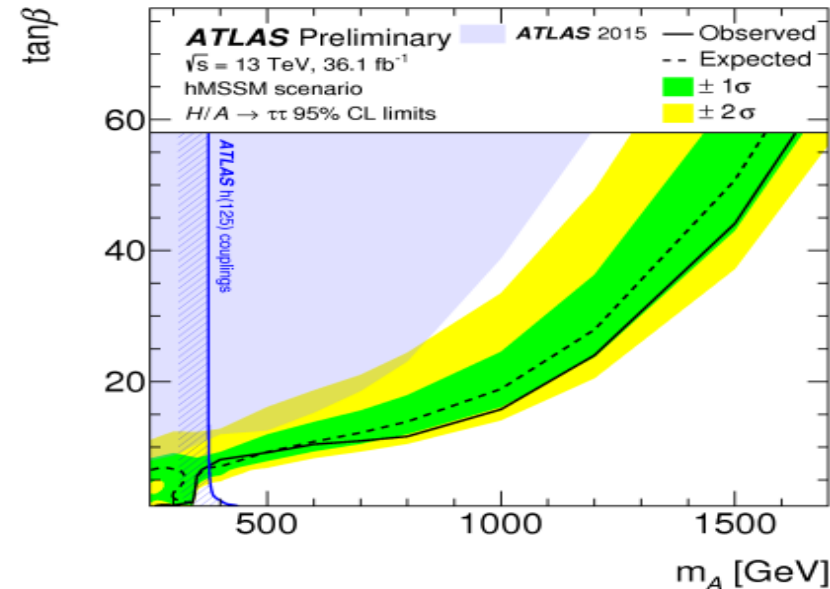


# RESONANCE SEARCHES ( $\gamma\gamma$ , $\tau\tau$ )

- Diphoton and ditau searches sensitive to new heavy scalars, e.g. Higgs bosons.
  - $\gamma\gamma$  search also targets spin-2 (graviton) production with a dedicated selection.
  - $\tau\tau$  searches sensitive to SUSY Higgs (H/A).
- No significant excess over the SM expectation.



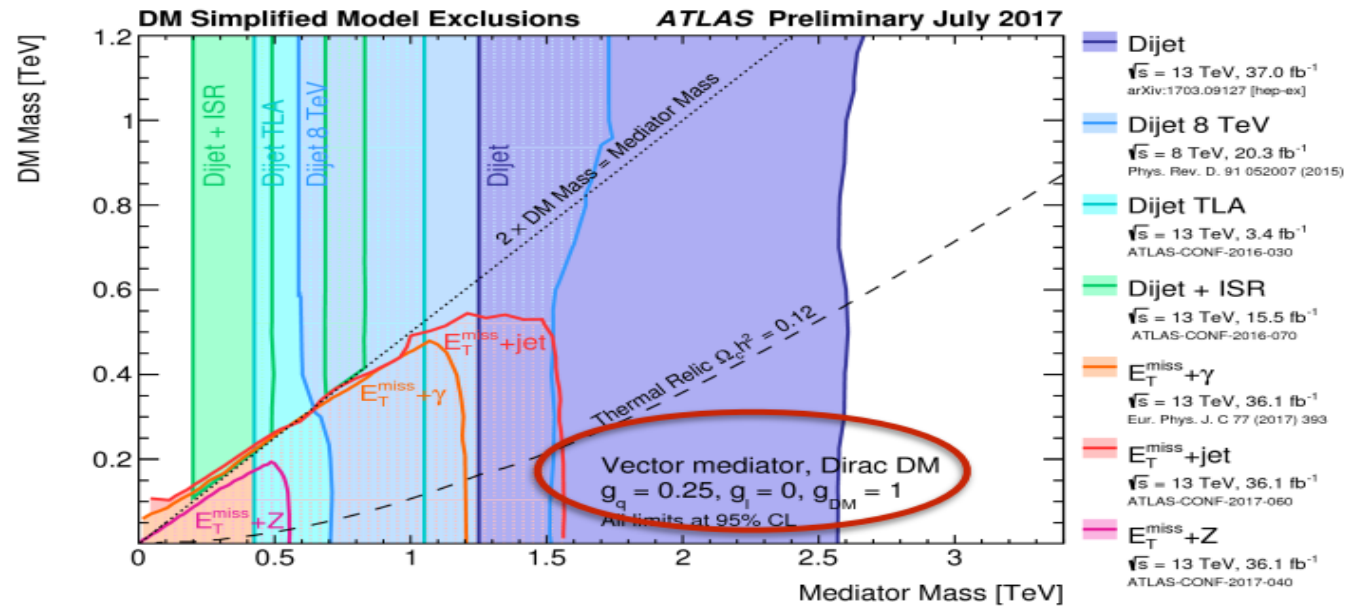
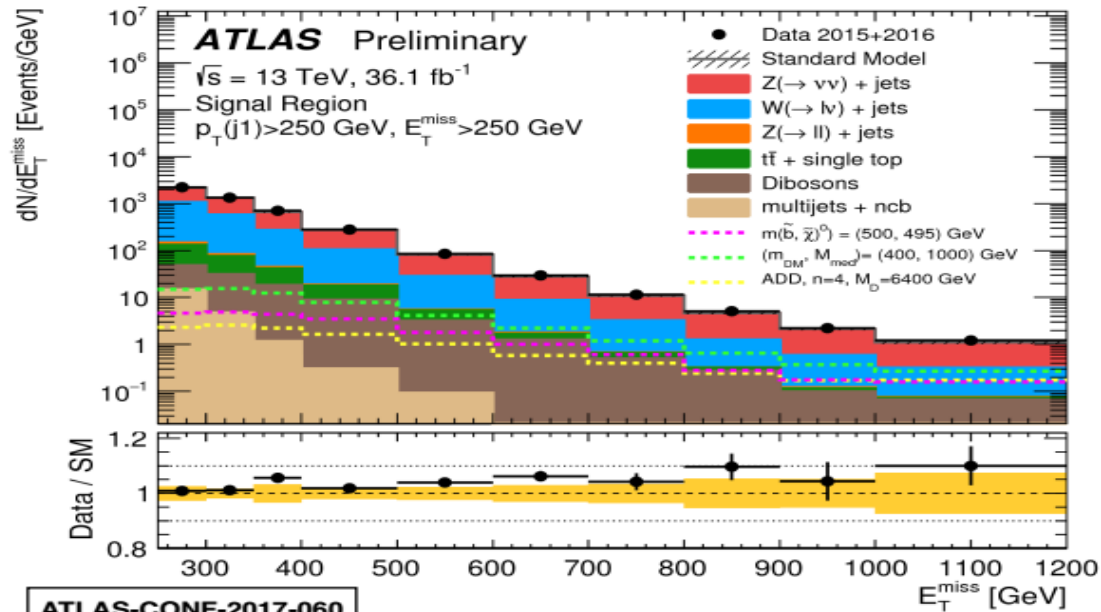
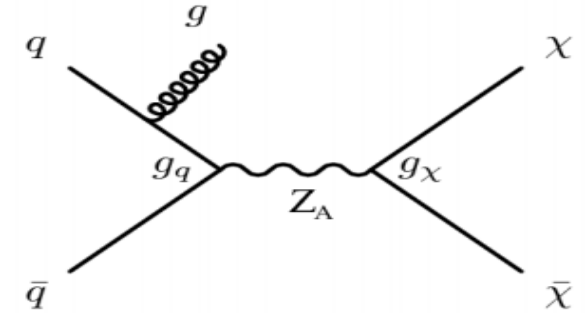
ATLAS-CONF-2017-050, arxiv:1707.04147





# SEARCHES FOR DARK MATTER

- Generic dark matter models tested with searches for mono-jet/ $\gamma$ / $Z$ / $H(\rightarrow\gamma\gamma/bb)+E_T^{\text{miss}}$ , with recoil against invisible dark matter particle(s).
  - Complementary to direct dark matter searches, and direct searches for the mediator decaying to e.g. a pair of jets.
- Phenomenology depends on mass of DM, mass of heavy mediator and value and type of couplings.
- No significant excesses over the SM predictions.



# SEARCHES FOR SUPERSYMMETRY

## ATLAS SUSY Searches\* - 95% CL Lower Limits

May 2017

ATLAS Preliminary

$\sqrt{s} = 7, 8, 13$  TeV

	Model	$e, \mu, \tau, \gamma$	Jets	$E_T^{\text{miss}}$	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference	
							$\sqrt{s} = 7, 8$ TeV	$\sqrt{s} = 13$ TeV
Inclusive Searches	MSUGRA/CMSSM	0-3 $e, \mu$ /1-2 $\tau$	2-10 jets/3 $b$	Yes	20.3	$\tilde{q}, \tilde{g}$	1.85 TeV	$m(\tilde{g})=m(\tilde{g})$
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	36.1	$\tilde{q}$	1.57 TeV	$m(\tilde{\chi}_1^0) < 200$ GeV, $m(1^{\text{st}} \text{ gen. } \tilde{q})=m(2^{\text{nd}} \text{ gen. } \tilde{q})$
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed)	mono-jet	1-3 jets	Yes	3.2	$\tilde{q}$	608 GeV	$m(\tilde{g})-m(\tilde{\chi}_1^0) < 5$ GeV
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	36.1	$\tilde{g}$	2.02 TeV	$m(\tilde{\chi}_1^0) < 200$ GeV
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0 \rightarrow qqW^\pm\tilde{\chi}_1^0$	0	2-6 jets	Yes	36.1	$\tilde{g}$	2.01 TeV	$m(\tilde{\chi}_1^0) < 200$ GeV, $m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\nu\nu)\tilde{\chi}_1^0$	3 $e, \mu$	4 jets	-	36.1	$\tilde{g}$	1.825 TeV	$m(\tilde{\chi}_1^0) < 400$ GeV
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	0	7-11 jets	Yes	36.1	$\tilde{g}$	1.8 TeV	$m(\tilde{\chi}_1^0) < 400$ GeV
	GMSB ( $\tilde{\ell}$ NLSP)	1-2 $\tau$ + 0-1 $\ell$	0-2 jets	Yes	3.2	$\tilde{g}$	2.0 TeV	$c\tau(\text{NLSP}) < 0.1$ mm
	GGM (bino NLSP)	2 $\gamma$	-	Yes	3.2	$\tilde{g}$	1.65 TeV	$m(\tilde{\chi}_1^0) < 950$ GeV, $c\tau(\text{NLSP}) < 0.1$ mm, $\mu < 0$
	GGM (higgsino-bino NLSP)	$\gamma$	1 $b$	Yes	20.3	$\tilde{g}$	1.37 TeV	$m(\tilde{\chi}_1^0) > 680$ GeV, $c\tau(\text{NLSP}) < 0.1$ mm, $\mu > 0$
	GGM (higgsino-bino NLSP)	$\gamma$	2 jets	Yes	13.3	$\tilde{g}$	1.8 TeV	$m(\text{NLSP}) > 430$ GeV
	GGM (higgsino NLSP)	2 $e, \mu$ (Z)	2 jets	Yes	20.3	$\tilde{g}$	900 GeV	1503.03290
Gravitino LSP	0	mono-jet	Yes	20.3	$F^{1/2}$ scale	865 GeV	$m(\tilde{G}) > 1.8 \times 10^{-4}$ eV, $m(\tilde{g})=m(\tilde{g})=1.5$ TeV	
3 <sup>rd</sup> gen. $\tilde{g}$ med.	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 $b$	Yes	36.1	$\tilde{g}$	1.92 TeV	$m(\tilde{\chi}_1^0) < 600$ GeV
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0-1 $e, \mu$	3 $b$	Yes	36.1	$\tilde{g}$	1.97 TeV	$m(\tilde{\chi}_1^0) < 200$ GeV
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow b\tilde{t}\tilde{\chi}_1^0$	0-1 $e, \mu$	3 $b$	Yes	20.1	$\tilde{g}$	1.37 TeV	$m(\tilde{\chi}_1^0) < 300$ GeV
3 <sup>rd</sup> gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 $b$	Yes	36.1	$\tilde{b}_1$	950 GeV	$m(\tilde{\chi}_1^0) < 420$ GeV
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^0$	2 $e, \mu$ (SS)	1 $b$	Yes	36.1	$\tilde{b}_1$	275-700 GeV	$m(\tilde{\chi}_1^0) < 200$ GeV, $m(\tilde{\chi}_1^\pm) = m(\tilde{\chi}_1^0) + 100$ GeV
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^0$	0-2 $e, \mu$	1-2 $b$	Yes	4.7/13.3	$\tilde{t}_1$	117-170 GeV	$m(\tilde{\chi}_1^0) = 2m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^\pm) = 55$ GeV
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$	0-2 $e, \mu$	0-2 jets/1-2 $b$	Yes	20.3/36.1	$\tilde{t}_1$	90-198 GeV	$m(\tilde{\chi}_1^0) = 1$ GeV
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet	Yes	3.2	$\tilde{t}_1$	90-323 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0) = 5$ GeV
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 $e, \mu$ (Z)	1 $b$	Yes	20.3	$\tilde{t}_1$	150-600 GeV	$m(\tilde{\chi}_1^0) > 150$ GeV
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 $e, \mu$ (Z)	1 $b$	Yes	36.1	$\tilde{t}_2$	290-790 GeV	$m(\tilde{\chi}_1^0) = 0$ GeV
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$	1-2 $e, \mu$	4 $b$	Yes	36.1	$\tilde{t}_2$	320-880 GeV	$m(\tilde{\chi}_1^0) = 0$ GeV
EW direct	$\tilde{\ell}_{L,R}, \tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 $e, \mu$	0	Yes	36.1	$\tilde{\ell}$	90-440 GeV	$m(\tilde{\chi}_1^0) = 0$
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow \tilde{\nu}(\ell\tilde{\nu})$	2 $e, \mu$	0	Yes	36.1	$\tilde{\chi}_1^\pm$	710 GeV	$m(\tilde{\chi}_1^0) = 0, m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^+\tilde{\chi}_1^+/\tilde{\chi}_2^+, \tilde{\chi}_1^+ \rightarrow \tilde{\tau}\nu(\tau\tilde{\nu}), \tilde{\chi}_2^+ \rightarrow \tilde{\tau}\tau(\nu\tilde{\nu})$	2 $\tau$	-	Yes	36.1	$\tilde{\chi}_1^\pm$	760 GeV	$m(\tilde{\chi}_1^0) = 0, m(\tilde{\tau}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_L\nu\tilde{\ell}_L(\ell\tilde{\nu}), \tilde{\ell}\tilde{\nu}\tilde{\ell}_L(\ell\tilde{\nu})$	3 $e, \mu$	0	Yes	36.1	$\tilde{\chi}_1^+, \tilde{\chi}_2^0$	1.16 TeV	$m(\tilde{\chi}_1^0) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^\pm) = 0, m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0Z\tilde{\chi}_1^0$	2-3 $e, \mu$	0-2 jets	Yes	36.1	$\tilde{\chi}_1^+, \tilde{\chi}_2^0$	580 GeV	$m(\tilde{\chi}_1^0) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^\pm) = 0, \tilde{\ell}$ decoupled
	$\tilde{\chi}_1^+\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0h\tilde{\chi}_1^0, h \rightarrow b\tilde{b}/WW/\tau\tau/\gamma\gamma$	$e, \mu, \gamma$	0-2 $b$	Yes	20.3	$\tilde{\chi}_1^+, \tilde{\chi}_2^0$	270 GeV	$m(\tilde{\chi}_1^0) = m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^\pm) = 0, \tilde{\ell}$ decoupled
	$\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_2^0 \rightarrow \tilde{\ell}_R\ell$	4 $e, \mu$	0	Yes	20.3	$\tilde{\chi}_2^0, \tilde{\chi}_3^0$	635 GeV	$m(\tilde{\chi}_2^0) = m(\tilde{\chi}_3^0), m(\tilde{\chi}_1^\pm) = 0, m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_2^0) + m(\tilde{\chi}_1^0))$
	GGM (wino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$	1 $e, \mu + \gamma$	-	Yes	20.3	$\tilde{W}$	115-370 GeV	$c\tau < 1$ mm
	GGM (bino NLSP) weak prod., $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$	2 $\gamma$	-	Yes	20.3	$\tilde{W}$	590 GeV	$c\tau < 1$ mm
	Long-lived particles	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	36.1	$\tilde{\chi}_1^\pm$	430 GeV
Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$		dE/dx trk	-	Yes	18.4	$\tilde{\chi}_1^\pm$	495 GeV	$m(\tilde{\chi}_1^0) - m(\tilde{\chi}_1^\pm) \sim 160$ MeV, $\tau(\tilde{\chi}_1^\pm) < 15$ ns
Stable, stopped $\tilde{g}$ R-hadron		0	1-5 jets	Yes	27.9	$\tilde{g}$	850 GeV	$m(\tilde{\chi}_1^0) = 100$ GeV, $10 \mu\text{s} < \tau(\tilde{g}) < 1000$ s
Stable $\tilde{g}$ R-hadron		trk	-	-	3.2	$\tilde{g}$	1.58 TeV	1606.05129
Metastable $\tilde{g}$ R-hadron		dE/dx trk	-	-	3.2	$\tilde{g}$	1.57 TeV	1604.04520
GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{\ell}, \tilde{\mu}) + \tau(e, \mu)$		1-2 $\mu$	-	-	19.1	$\tilde{\chi}_1^0$	537 GeV	$10 < \tan\beta < 50$
GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$ , long-lived $\tilde{\chi}_1^0$		2 $\gamma$	-	Yes	20.3	$\tilde{\chi}_1^0$	440 GeV	$1 < \tau(\tilde{\chi}_1^0) < 3$ ns, SPS8 model
$\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow e\tilde{e}\nu/\mu\tilde{\mu}\nu$		displ. $e\tilde{e}/\mu\tilde{\mu}\nu$	-	-	20.3	$\tilde{\chi}_1^0$	1.0 TeV	$7 < c\tau(\tilde{\chi}_1^0) < 740$ mm, $m(\tilde{g}) = 1.3$ TeV
GGM $\tilde{g}\tilde{g}, \tilde{\chi}_1^0 \rightarrow Z\tilde{G}$	displ. vtx + jets	-	-	20.3	$\tilde{\chi}_1^0$	1.0 TeV	$6 < c\tau(\tilde{\chi}_1^0) < 480$ mm, $m(\tilde{g}) = 1.1$ TeV	
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu/\mu\tau$	$e\mu, e\tau, \mu\tau$	-	-	3.2	$\tilde{\nu}_\tau$	1.9 TeV	$\lambda_{111}^2 = 0.11, \lambda_{132/133/233} = 0.07$
	Bilinear RPV CMSSM	2 $e, \mu$ (SS)	0-3 $b$	Yes	20.3	$\tilde{q}, \tilde{g}$	1.45 TeV	$m(\tilde{g}) = m(\tilde{g}), c\tau_{LSP} < 1$ mm
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow e\tilde{e}\nu, \mu\tilde{\mu}\nu$	4 $e, \mu$	-	Yes	13.3	$\tilde{\chi}_1^\pm$	1.14 TeV	$m(\tilde{\chi}_1^0) > 400$ GeV, $\lambda_{12k} \neq 0$ ( $k = 1, 2$ )
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\nu\tau, e\nu\tau$	3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{\chi}_1^\pm$	450 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{133} \neq 0$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{q}$	0	4-5 large-R jets	-	14.8	$\tilde{g}$	1.08 TeV	$BR(\tilde{t}) = BR(\tilde{b}) = BR(\tilde{c}) = 0\%$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$	0	4-5 large-R jets	-	14.8	$\tilde{g}$	1.55 TeV	$m(\tilde{\chi}_1^0) = 800$ GeV
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$	1 $e, \mu$	8-10 jets/0-4 $b$	-	36.1	$\tilde{g}$	2.1 TeV	$m(\tilde{\chi}_1^0) = 1$ TeV, $\lambda_{112} \neq 0$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$	1 $e, \mu$	8-10 jets/0-4 $b$	-	36.1	$\tilde{g}$	1.65 TeV	$m(\tilde{t}_1) = 1$ TeV, $\lambda_{323} \neq 0$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$	0	2 jets + 2 $b$	-	15.4	$\tilde{t}_1$	410 GeV	ATLAS-CONF-2016-022, ATLAS-CONF-2016-084
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\ell}$	2 $e, \mu$	2 $b$	-	36.1	$\tilde{t}_1$	0.4-1.45 TeV	$BR(\tilde{t}_1 \rightarrow b\tilde{\ell}/\mu) > 20\%$
Other	Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$	0	2 $c$	Yes	20.3	$\tilde{c}$	510 GeV	$m(\tilde{\chi}_1^0) < 200$ GeV

\*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10<sup>-1</sup>

1

Mass scale [TeV]

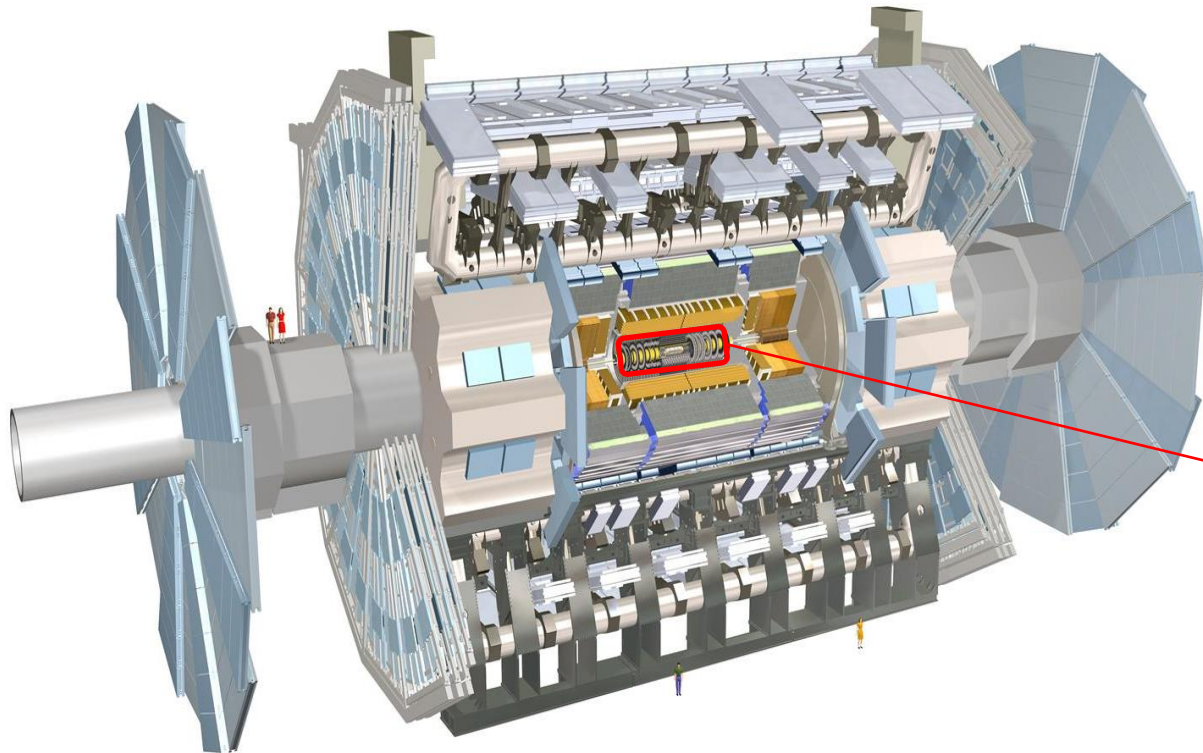
Mass reach searches for Supersymmetry. Only a representative selection of the available results is shown.



# OVERVIEW

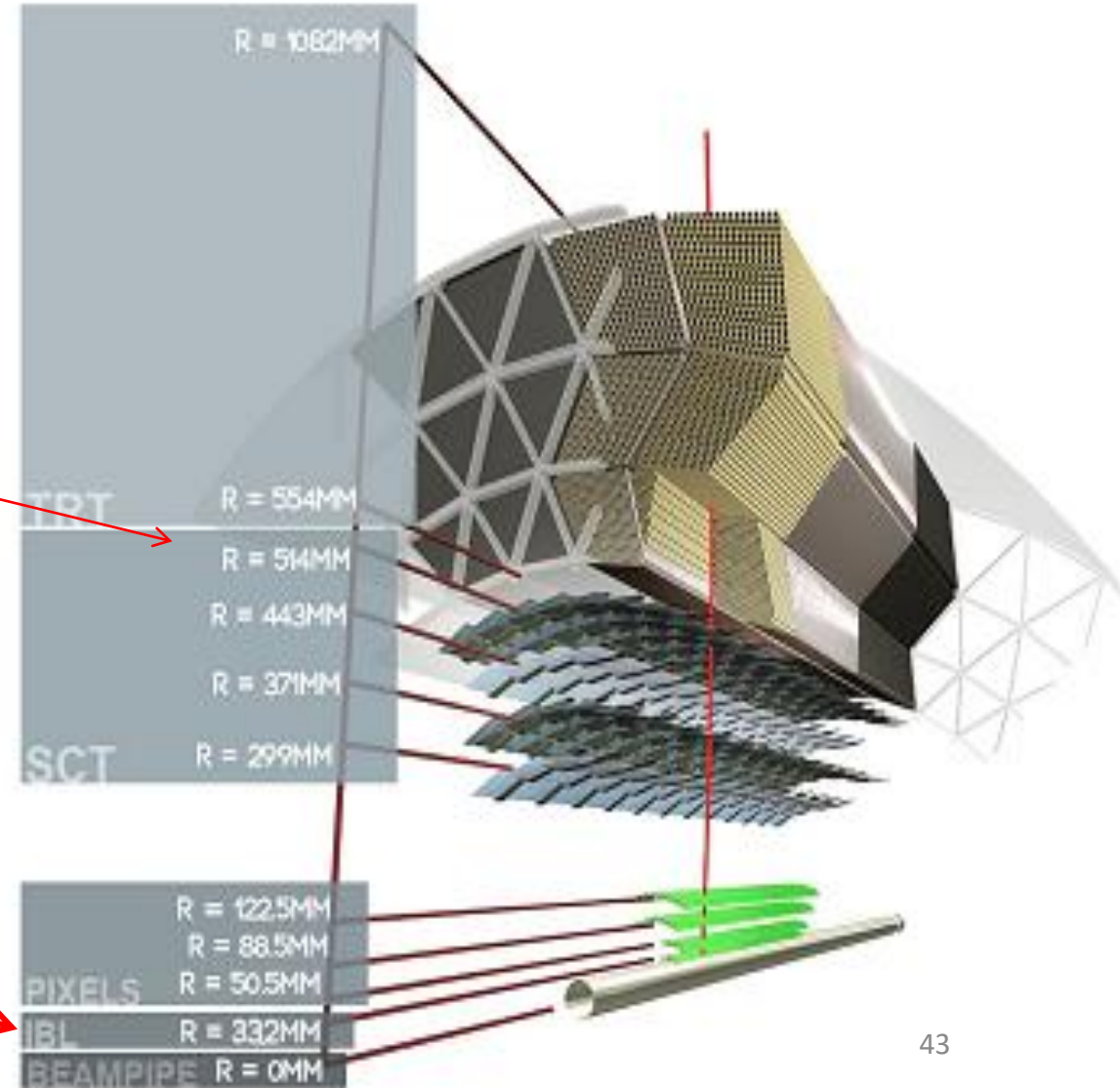
The focus of ATLAS is high- $p_T$  physics, and also provides a window onto important soft QCD processes. These have intrinsic interest but also the understanding of underpins searches for new physics.

## Task ► Bose Einstein Correlations



**ATLAS Inner Detector (ID) main tracking device:**

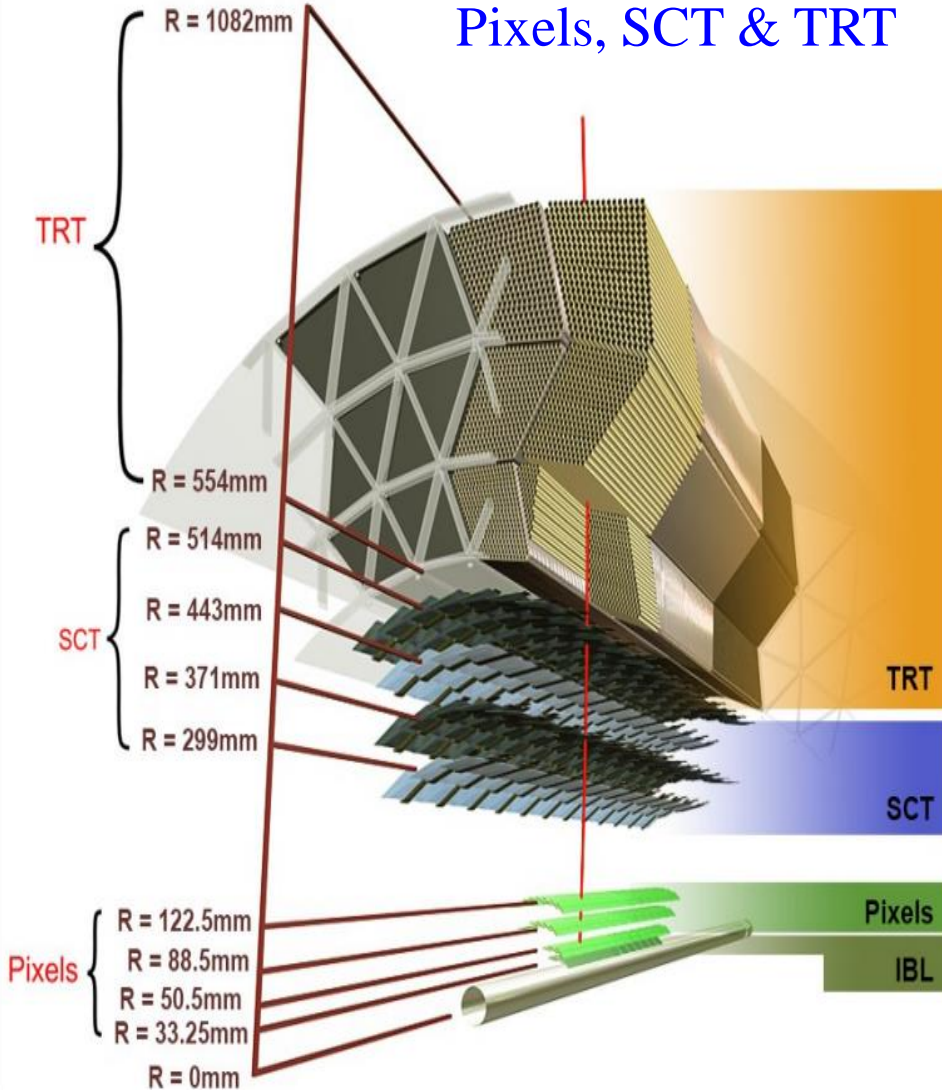
Consists of **Pixel**, **Silicon strip (SCT)** and **drift tube (TRT)** detectors. Single hit resolution between  $10\ \mu\text{m}$  (Pixel) and  $1\ \mu\text{m}$  (TRT). **New: Insertable B-Layer (IBL)** in the Pixel



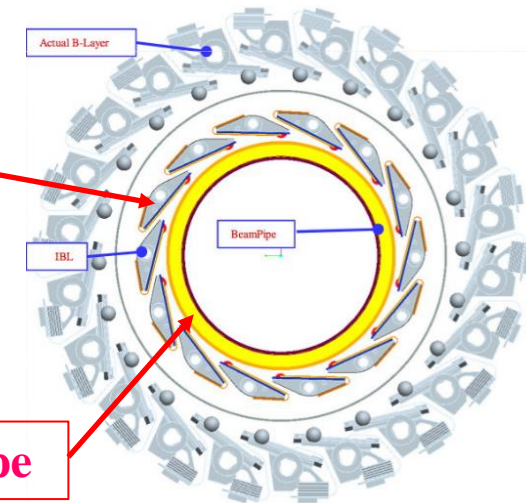


# INNER DETECTORS (ID)

## ATLAS tracking detectors: Pixels, SCT & TRT



- ❑ **New innermost 4-th layer** for the Pixel detector [IBL = Insertable B-Layer]
- ❑ Required complete removal of the ATLAS Pixel volume
- ❑ IBL fully operational



Two times better tracks impact parameters resolution at 13 TeV!



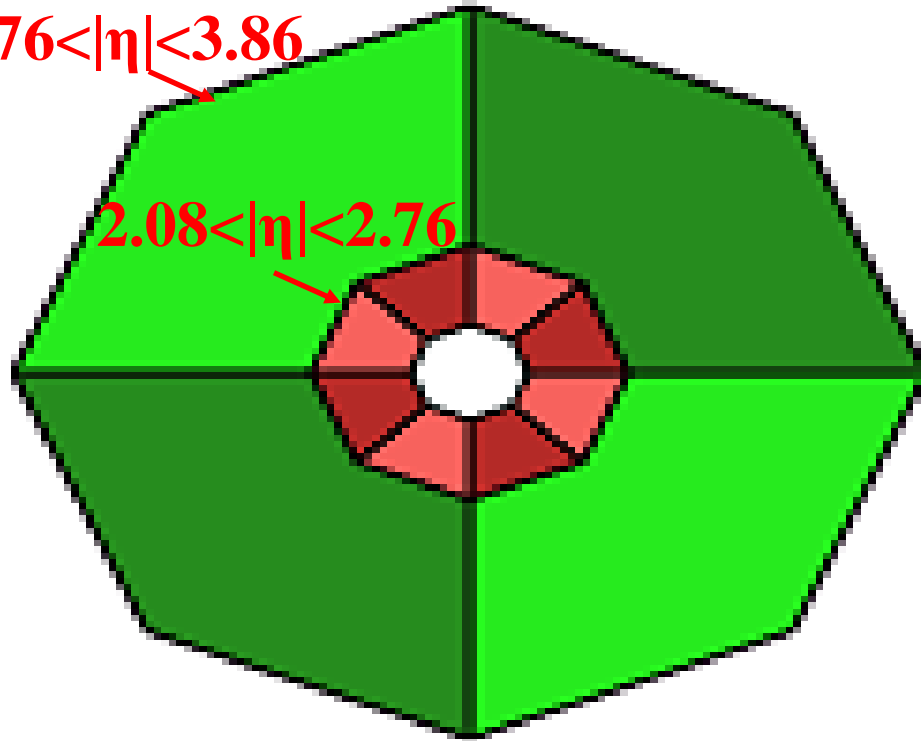
# MINIMUM BIAS TRIGGER SCINTILLATOR

24 independent wedge-shaped plastic scintillators (12 per side) read out by PMTs,

$$2.08 < |\eta| < 3.86^*$$

$$2.76 < |\eta| < 3.86$$

$$2.08 < |\eta| < 2.76$$



\* Pseudorapidity is defined as  $\eta = -\frac{1}{2} \ln(\tan(\theta/2))$ ,  $\theta$  is the polar angle with respect to the beam.

- Designed for triggering on min bias events, >99% efficiency
- **MBTS** timing used to veto halo and beam gas events
- Also being used as gap trigger for various diffractive subjects

# MINIMUM-BIAS AND HIGH MULTIPLICITY TRACK TRIGGERS

➤ For these analysis the events collected with **Minimum-bias** (MB) trigger named as *HLT\_noalg\_mb\_L1MBTS\_1* were used.

□ This trigger required at least one hit in one of the 12+12 sectors (A and C sides) of the MBTS detector.

✓ Integral Luminosity  $\sim 151 \mu\text{b}^{-1}$ ; Statistic:  $9.6 \times 10^6$  events with  $2.8 \times 10^8$  tracks

➤ For these analysis the events collected with **High multiplicity track** (HMT) trigger named as *HLT\_mb\_sp900\_trk60\_hmt\_L1MBTS\_1\_1* were used.

□ High-multiplicity track (HMT) events were collected at 13 TeV using a dedicated high-multiplicity track trigger:

❖ requires more than 900 SCT space-points,

❖ more than 60 reconstructed good quality charged tracks with  $p_T > 0.4$  GeV associated with the primary vertex.

✓ Integral Luminosity  $\sim 8.4 \text{ nb}^{-1}$ ; Statistic:  $9.1 \times 10^6$  events with  $9.8 \times 10^8$  tracks

# MOTIVATION FOR BOSE-EINSTEIN CORRELATIONS

- Bose-Einstein correlations (BEC) represent a unique probe of the *space-time geometry* of the *hadronization region* and allow the *determination the size and shape of the source* from which particles are emitted.
- Studies of the dependence of BEC on *particle multiplicity* and *transverse momentum* are of special interest. They help in the understanding of multiparticle production mechanisms.
- High-multiplicity data in proton interactions can serve as a reference for studies in nucleus-nucleus collisions. The effect is reproduced in hydrodynamical and Pomeron-based approaches for hadronic interactions where high multiplicities play a crucial role.

# BOSE-EINSTEIN CORRELATIONS AND HANBURY BROWN – TWISS INTERFEROMETRY

Bose-Einstein correlations (BEC) are often considered to be the analogue of the **Hanbury Brown and Twiss effect** in astronomy, describing the interference of **incoherently-emitted identical bosons**.

Intensity interferometry of photons in radio-astronomy: measures angular diameter of two stars, so the physical size of the source

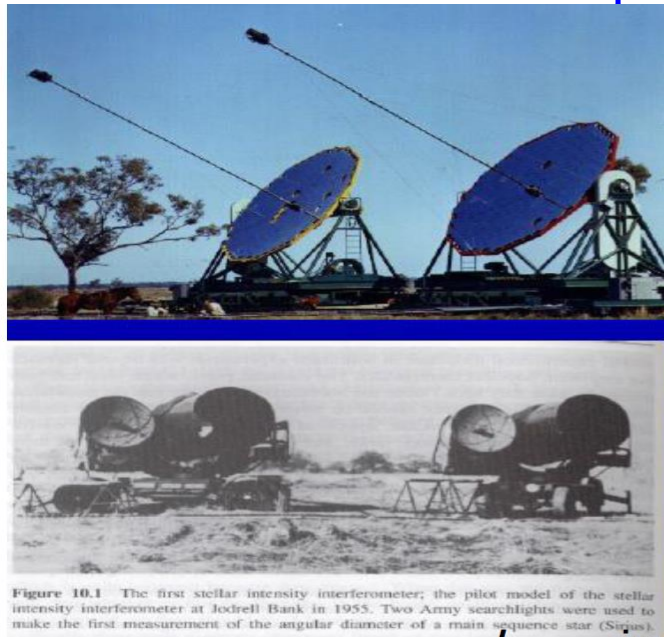
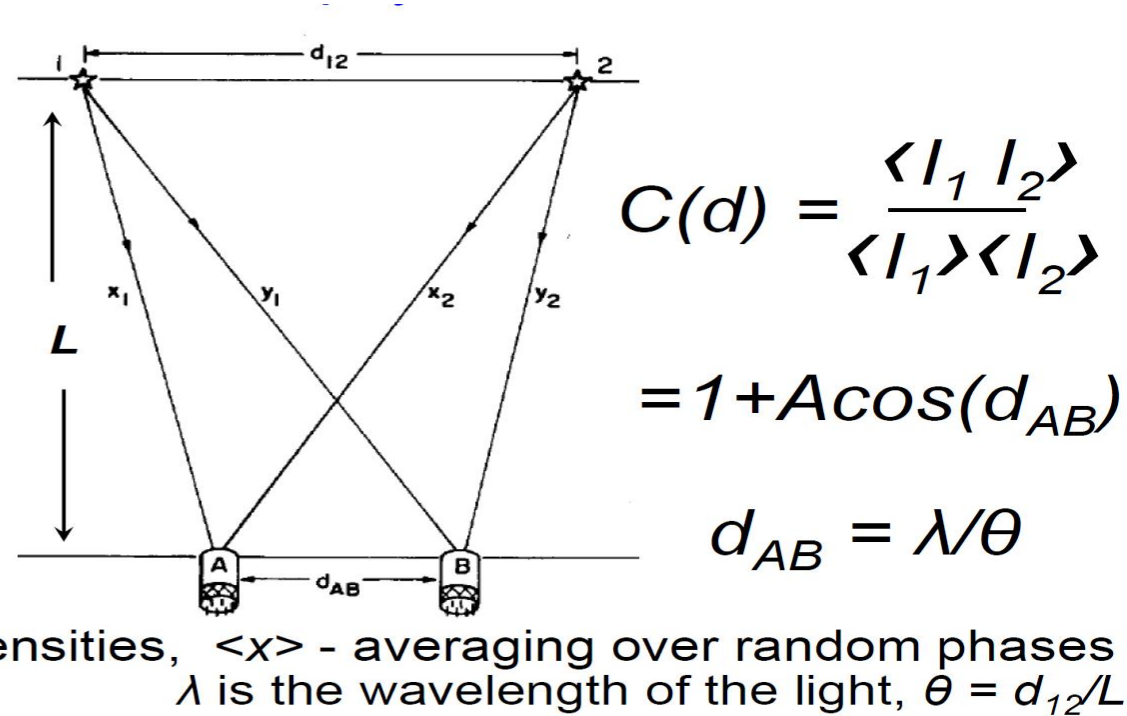


Figure 10.1 The first stellar intensity interferometer; the pilot model of the stellar intensity interferometer at Jodrell Bank in 1955. Two Army searchlights were used to make the first measurement of the angular diameter of a main sequence star (Sargas).



Varying  $d_{AB}$  one learns the angle, and using the individual wave vectors, the physical size of the source



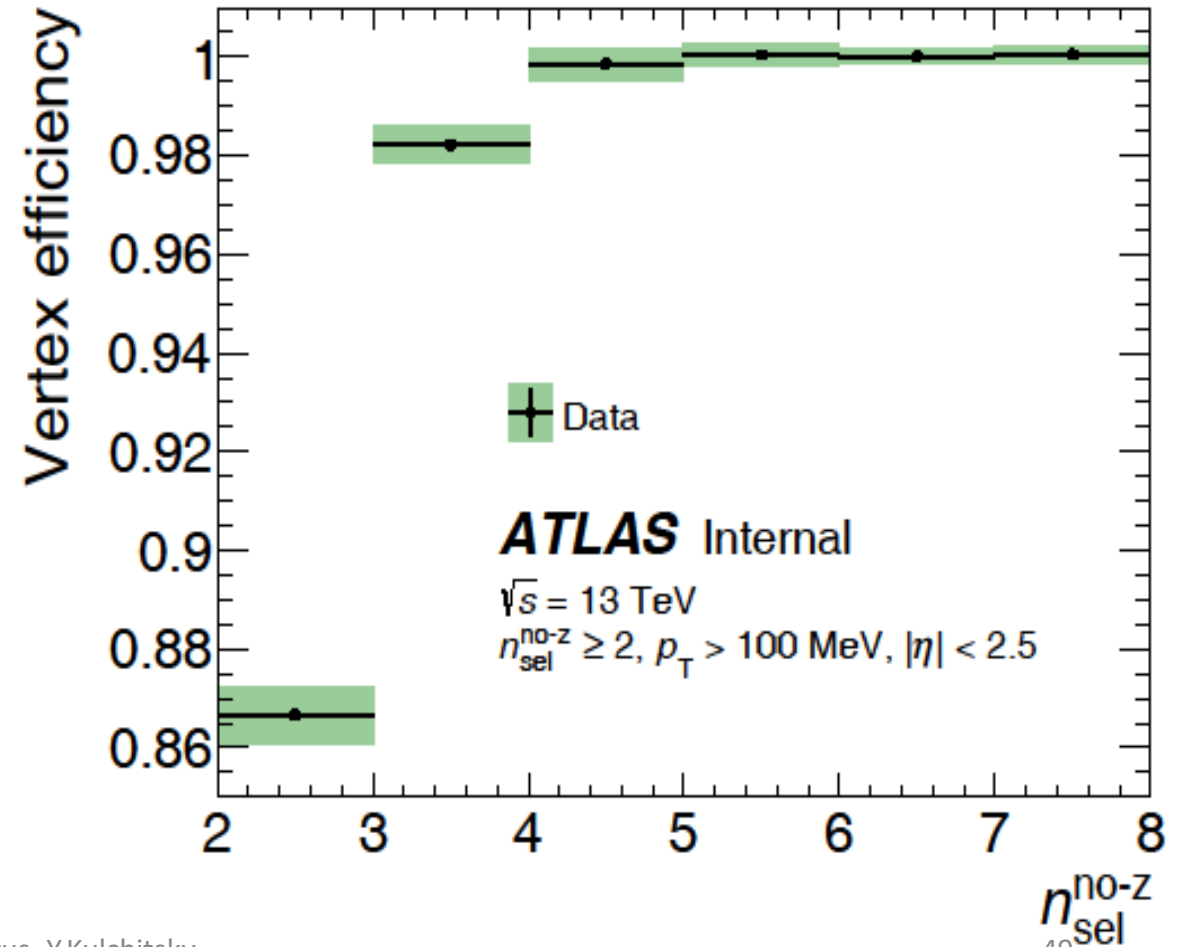
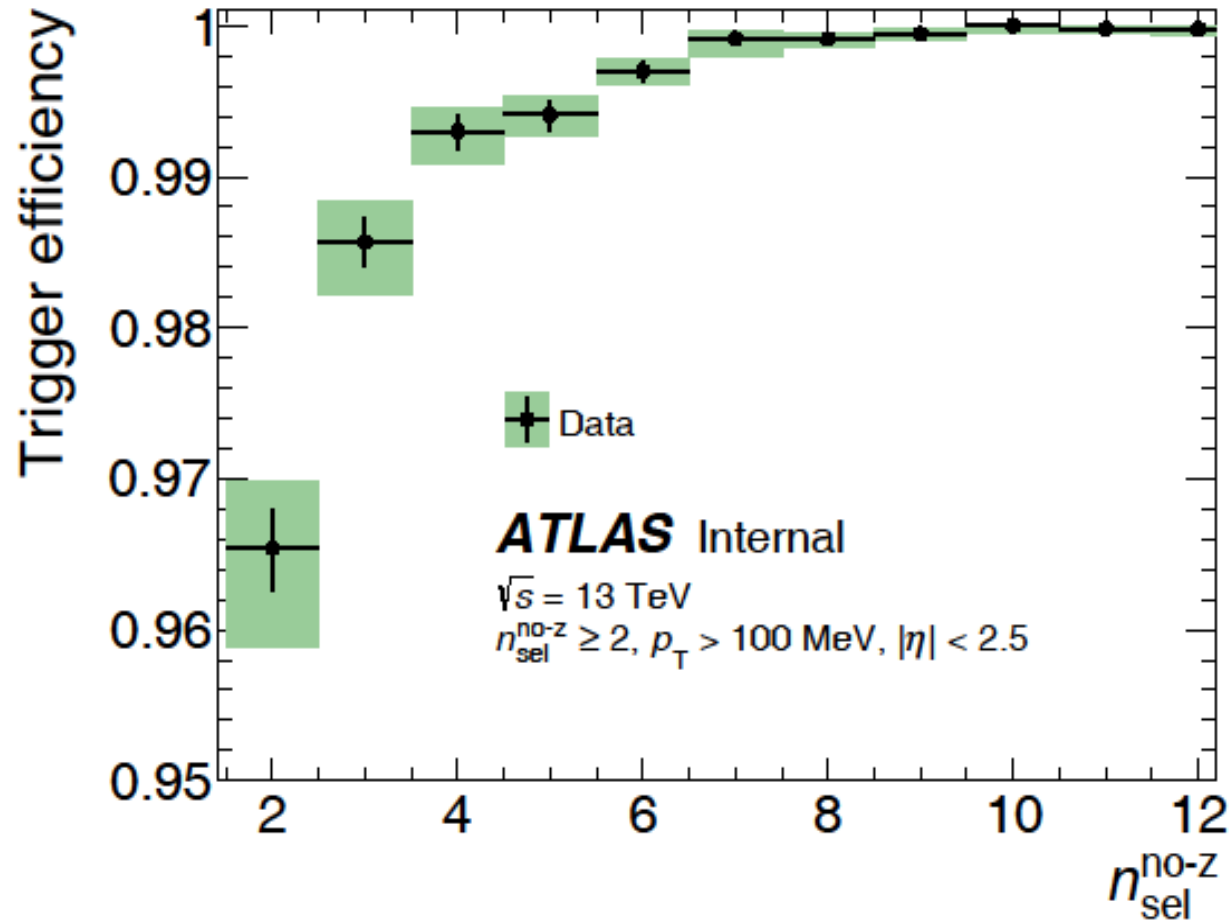
# EVENT CORRECTIONS

We correct the events on:

$$w(n) = \frac{1}{\epsilon_{trig}(n) \cdot \epsilon_{vert}(n)}$$

The trigger efficiency –  $\epsilon_{trig}(n)$ ,

The vertex reconstruction efficiency –  $\epsilon_{vert}(n)$



# DATA SAMPLES

Information from  
ATL-COM-PHYS-2015-1379 p.51

Run	$\mu$	Events in MinBias Stream	L1MBTS_1_1 Prescales	sp900_trk60 Prescales
267358	0.003	<b>Low <math>\mu</math></b> 8.6 20.9 millions of triggers	2	<b>Minimum-Bias trigger</b> <b>Events: <math>9.6 \times 10^6</math>;</b> <b>Tracks: <math>2.8 \times 10^8</math></b>
267359	0.007		2	
267360	0.03-0.04	12.5	LB:294-296, 100 LB:297-444, 1	1
267367	0.03	<b>High <math>\mu</math></b> 17.1 206.4 millions of triggers	LB:328-329, 100 LB:330-393, 1 LB:393-396, 200 LB:397-399, 2 LB:400-405, 10 LB:406-420, 4 LB:465-500, 3 LB:502-505, 3	1 1 1
267385	0.04		71.2	LB:241-584, 3 LB:870-908, 1 LB:910, 1
267599	0.01 0.03 0.03 0.03	105.6	LB:289-482, 1 LB:483-495, 1.9 LB:496-500, 1.4 LB:501-1130, 1	1 1 1 1

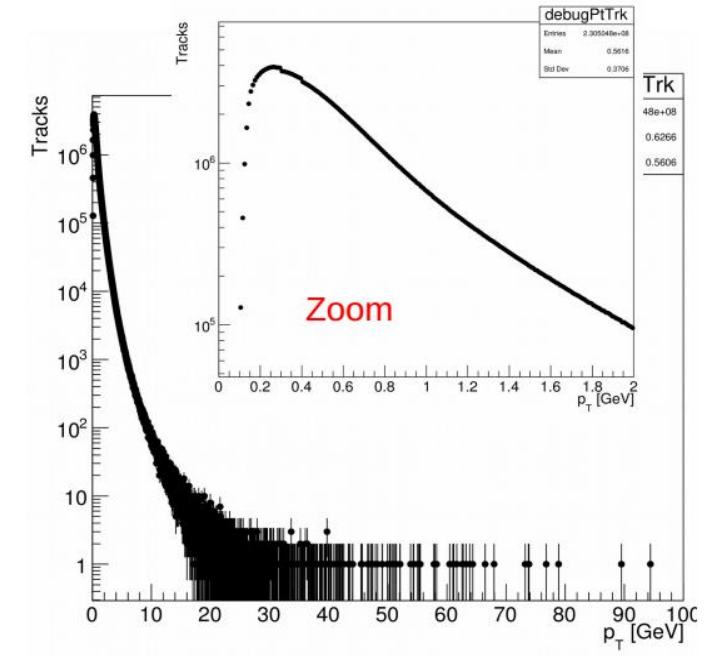
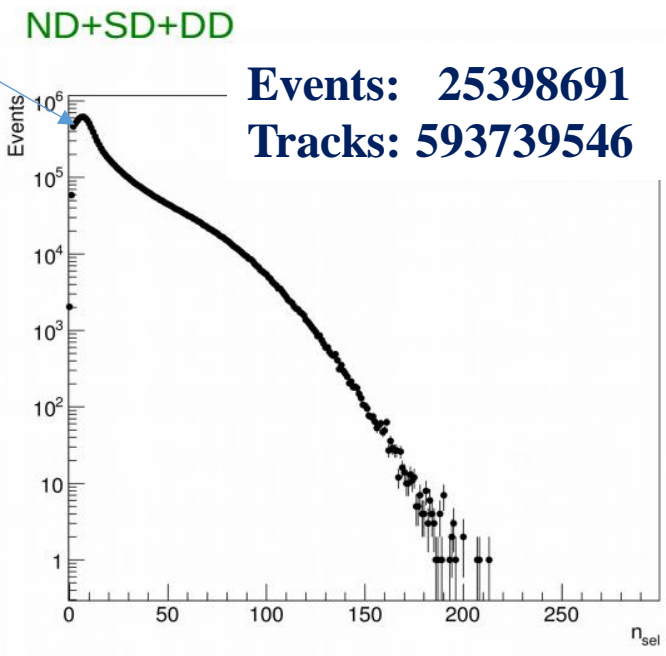
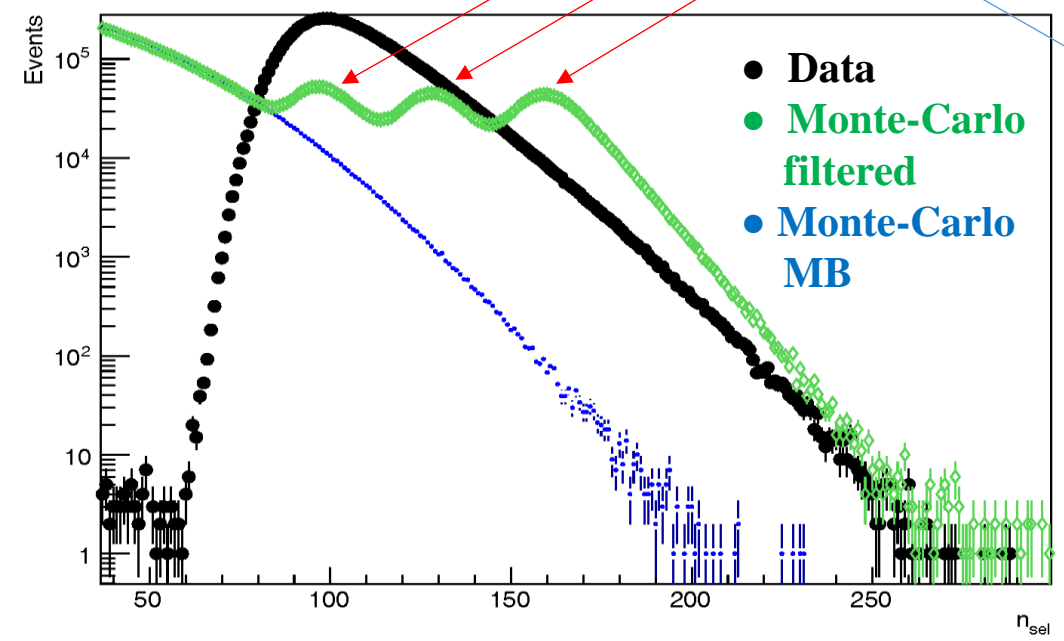
This data can be used for pile-up study for runs with higher  $\mu$

# MONTE-CARLO SAMPLES

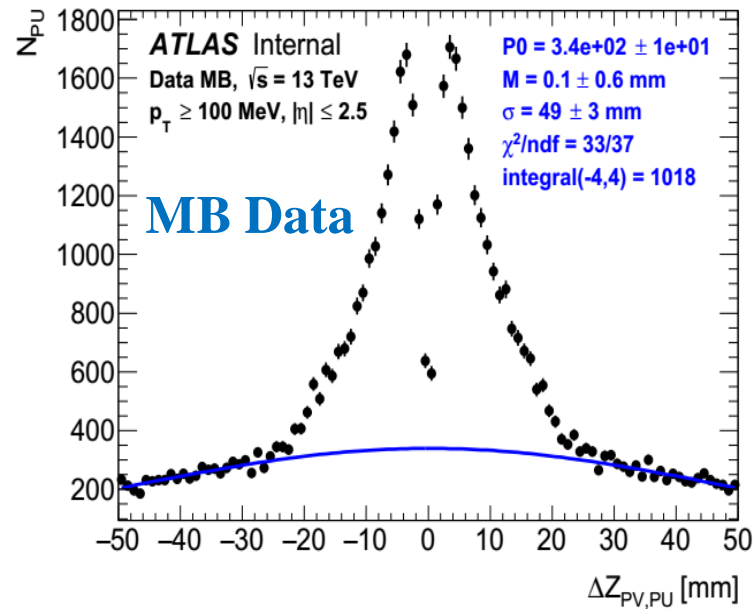
Generator	Version	Tune	PDF
Pythia8	8.185	A2	MSTW2008LO
Pythia8	8.186	Monash	NNPDF2.3LO
EPOS	LHCv3400	LHC	—

Stat.  
[mil]

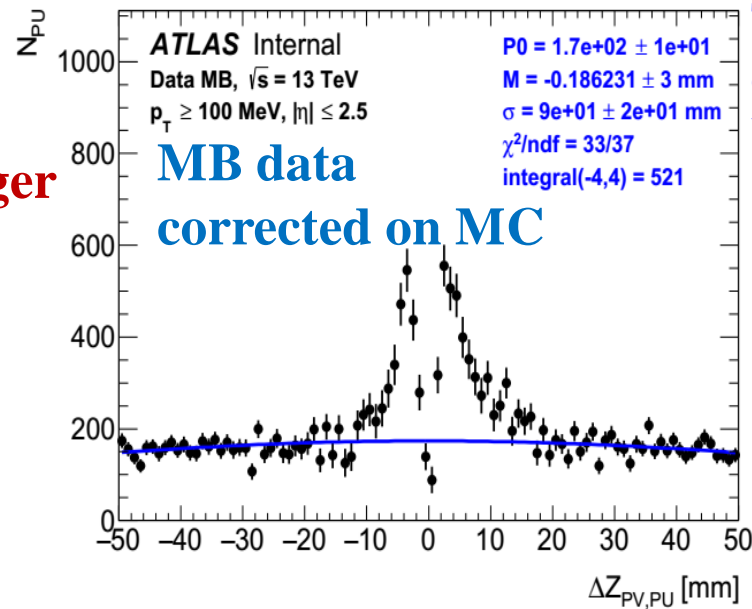
Monte-Carlo	Name of dataset:	Stat. [mil]
Pythia8 A2	361203.Pythia8_A2_MSTW2008LO_ND_minbias.merge.AOD.e3639_a782_a787_r6264	30
	361204.Pythia8_A2_MSTW2008LO_SD_minbias.merge.AOD.e3639_a782_a787_r6264	
	361205.Pythia8_A2_MSTW2008LO_DD_minbias.merge.AOD.e3639_a782_a787_r6264	
Pythia8 A2	361214.Pythia8_A2MSTW2008LO_minbias_NDnch120.merge.AOD.e3908_a782_a787_r6264	1
	361215.Pythia8_A2MSTW2008LO_minbias_NDnch160.merge.AOD.e3908_a782_a787_r6264	1
Filtered	361216.Pythia8_A2MSTW2008LO_minbias_NDnch200.merge.AOD.e3908_a782_a787_r6264	1
Pythia8 Monash	361200.Pythia8_Monash_NNPDF23LO_ND_minbias.merge.AOD.e3639_s2601_s2132_r6616_r6270	16
	361201.Pythia8_Monash_NNPDF23LO_SD_minbias.merge.AOD.e3639_s2601_s2132_r6616_r6264	
	361202.Pythia8_Monash_NNPDF23LO_DD_minbias.merge.AOD.e3639_s2601_s2132_r6616_r6264	
EPOS	361224.Epos_minbias_inelastic.merge.AOD.e3908_a782_s2183_a787_r6264mc15	20



# PILE-UP FOR HMT AND MB EVENTS

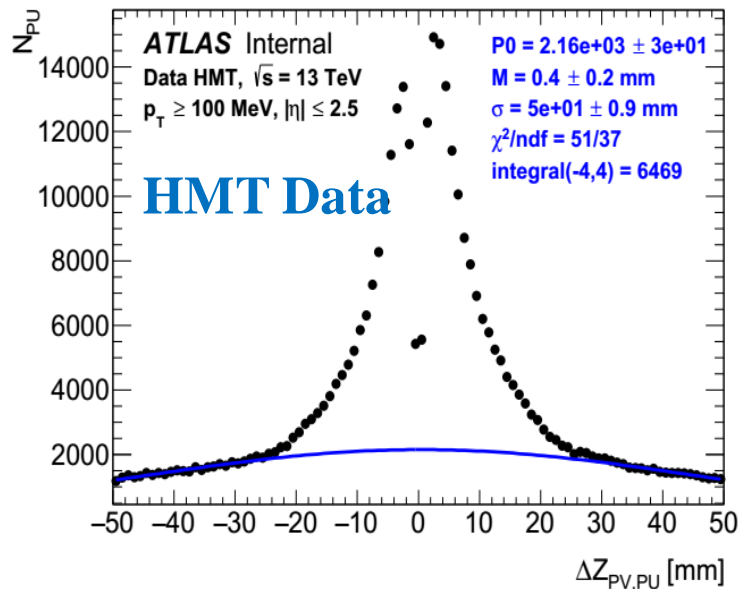


**MB trigger**

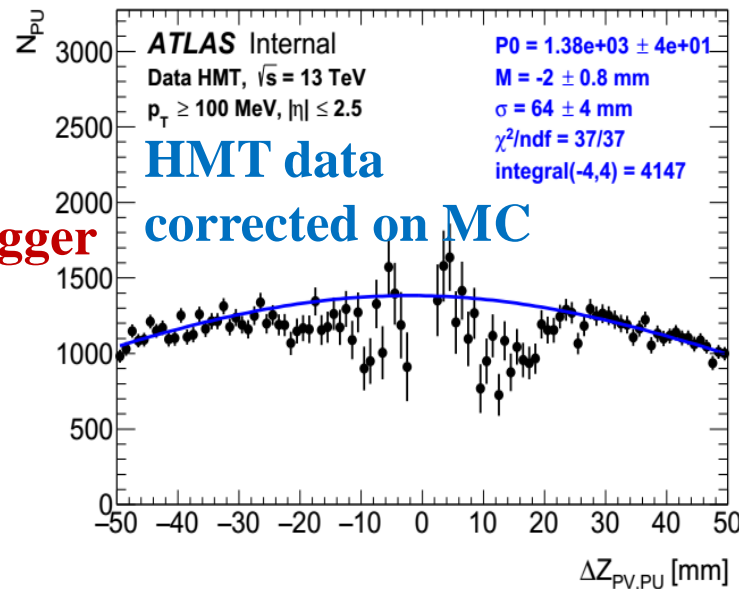


The distribution of the distance between Z coordinates of Primary Vertex and Pile-Up Vertexes for MB and HMT events for Data (left) and Data corrected on MC (right)

For MB events the number of pile-up vertexes in the Primary Vertex (PV) region  $\pm 4$  mm is  $\sim 520$  after correction on MC, and the number of tracks in Pile-up vertex is 9.4. Therefore the fraction of pile-up tracks in MB events is **0.002%**



**HMT trigger**



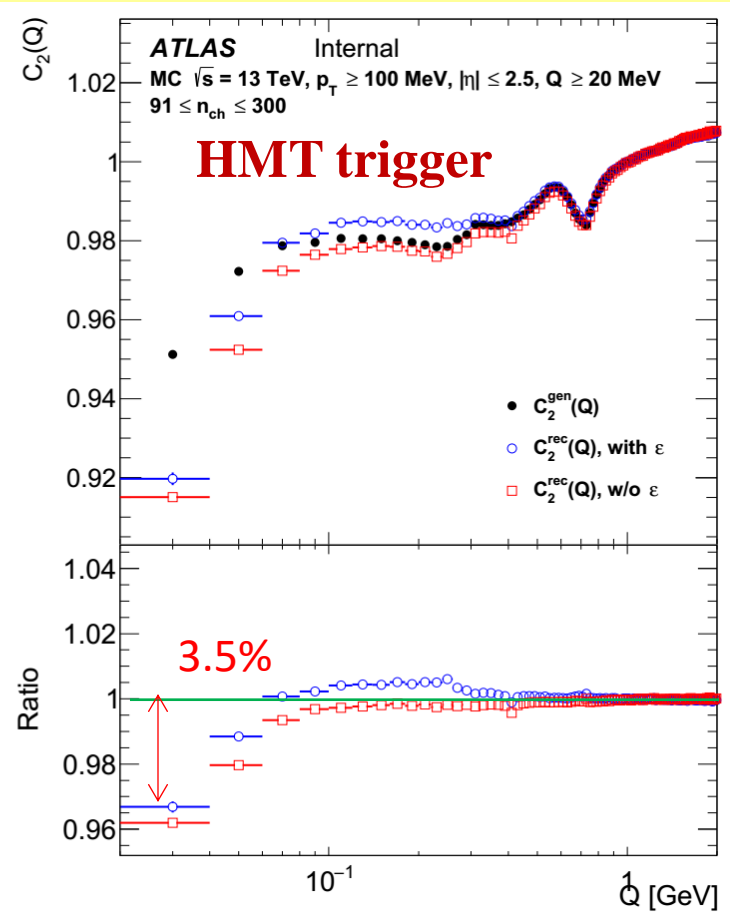
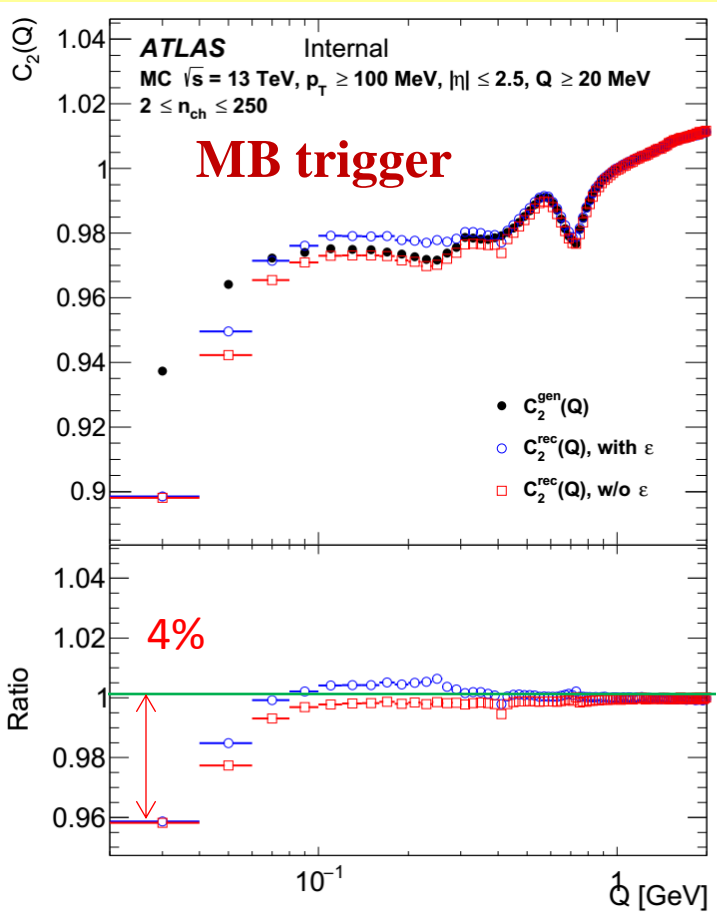
For HMT events the number of pile-up vertexes in the Primary Vertex (PV) region  $\pm 4$  mm is  $\sim 4150$ , after correction on MC, and the number of tracks in Pile-up vertex is 23. Therefore the fraction of pile-up tracks in MB events is **0.01%**

**We can conclude that mean number of pile-up tracks per MB or HMT event is negligible**

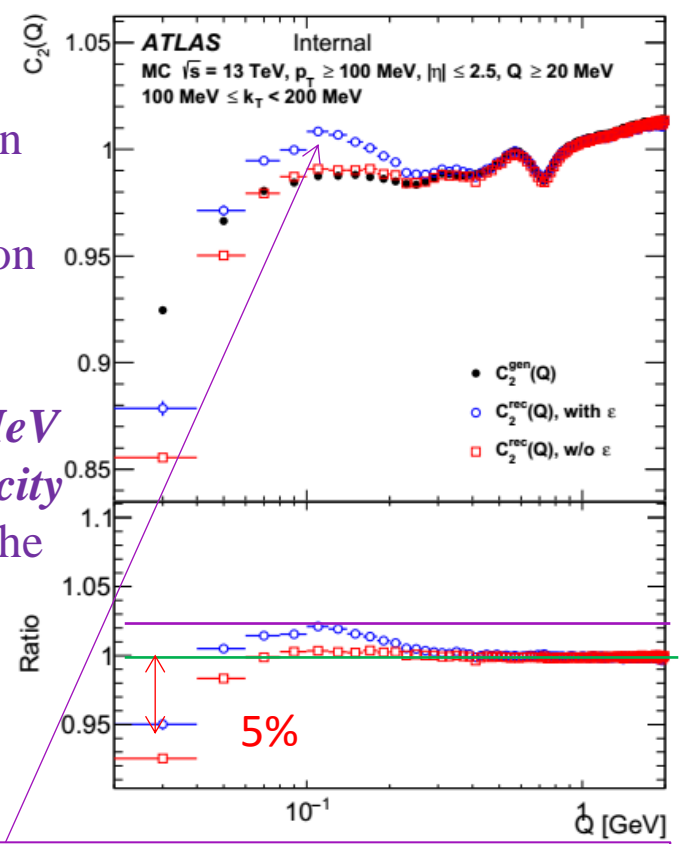
Mean number of tracks (pile-up tracks) per event: MB – 26 (0.0005) tracks/event; HMT – 108 (0.01) tracks/event



# CLOSURE TEST FOR TWO-PARTICLES $C_2$ CORRELATION FUNCTION



Un-closure between for reconstruction efficiency correction important only for smallest  $k_T$  region from 100 to 200 MeV and high multiplicity region  $n_{ch} > 60$  on the level up to 2%.



**Only for  $100 < k_T < 200$  MeV and  $n_{ch} > 60$**

The closure tests for  $C_2(Q)$  correlation function of Pythia8 A2 datasets at 13 TeV MB events for multiplicity region  $2 \leq n_{ch} \leq 250$  and HMT events for multiplicity region  $91 \leq n_{ch} \leq 300$  built using the unlike-sign pair reference sample.

The differences between the particle level distributions and the reconstructed distributions after unfolding are assigned to the full error for bins of  $R_2(Q)$  correlation functions and included in the final fitting error of the BEC parameters

# MINIMUM-BIAS EVENT SELECTION CRITERIA

**Events pass the data quality criteria (“good events”): all ID sub-systems nominal cond., stable beam, defined beam spot)**

- Accept on signal-arm Minimum Bias Trigger scintillator,
- Primary vertex (2 tracks with  $p_T > 100$  MeV),
- Veto to any additional vertices with  $\geq 4$  tracks,
- At least 2 tracks with  $p_T > 100$  MeV,  $|\eta| < 2.5$ ,
- At least 1 first Pixel layer hit & 2, 4, or 6 SCT hits for  $p_T > 100, 300, 400$  MeV respectively,
- IBL hit required if expected (if not expected, next to innermost hit required if expected),
- Cuts on the transverse impact parameter:  $|d_0^{BL}| < 1.5$  mm (w.r.t beam line),
- Cuts on the longitudinal impact parameter:  $|\Delta z_0 \sin\Theta| < 1.5$  mm ( $\Delta z_0$  is difference between tracks  $z_0$  and vertex  $z$  position),
- Track fit  $\chi^2$  probability  $> 0.01$  for tracks with  $p_T > 10$  GeV.

## **Correct distributions for detector effects:**

- where possible the data used to reduce the MC dependencies
- Monte Carlo derived corrections for tracking

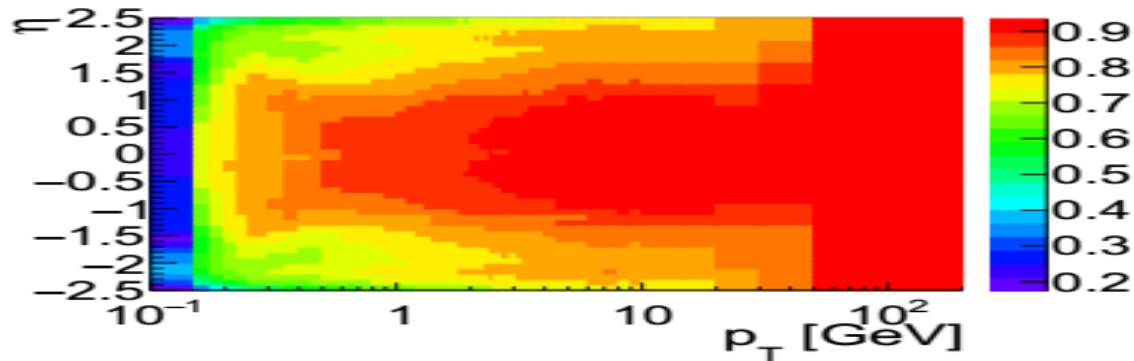
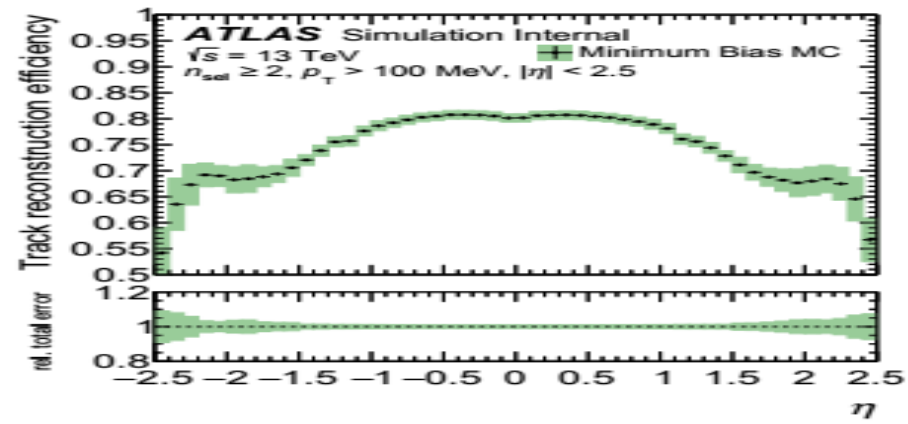
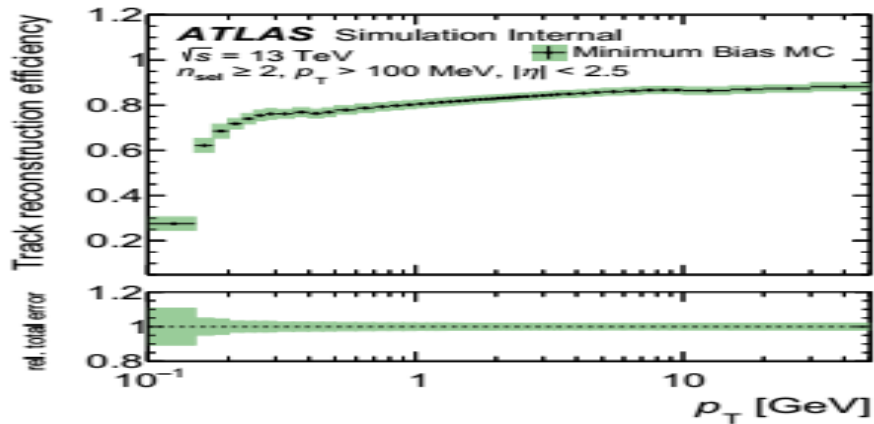
# TRACK RECONSTRUCTION CORRECTIONS

Performed corrections on:

1. The reconstruction track efficiency –  $\varepsilon(p_t, \eta)$ ,
2. The fraction of non-primary (secondaries and fake) tracks –  $f_{\text{nonp}}(p_t, \eta)$ ,
3. The fraction of tracks for which the corresponding primary particles are outside the kinematic range –  $f_{\text{okr}}(p_t, \eta)$ ,
4. The strange baryon tracks –  $f_{\text{sb}}(p_t, \eta)$ ,

We use the formula, as earlier and as in MB studies:

$$w_i(p_t, \eta) = \frac{(1 - f_{\text{nonp}}(p_t, \eta) - f_{\text{okr}}(p_t, \eta) - f_{\text{sb}}(p_t, \eta))}{\varepsilon(p_t, \eta)}$$



# THE PHASE SPACE CORRECTION

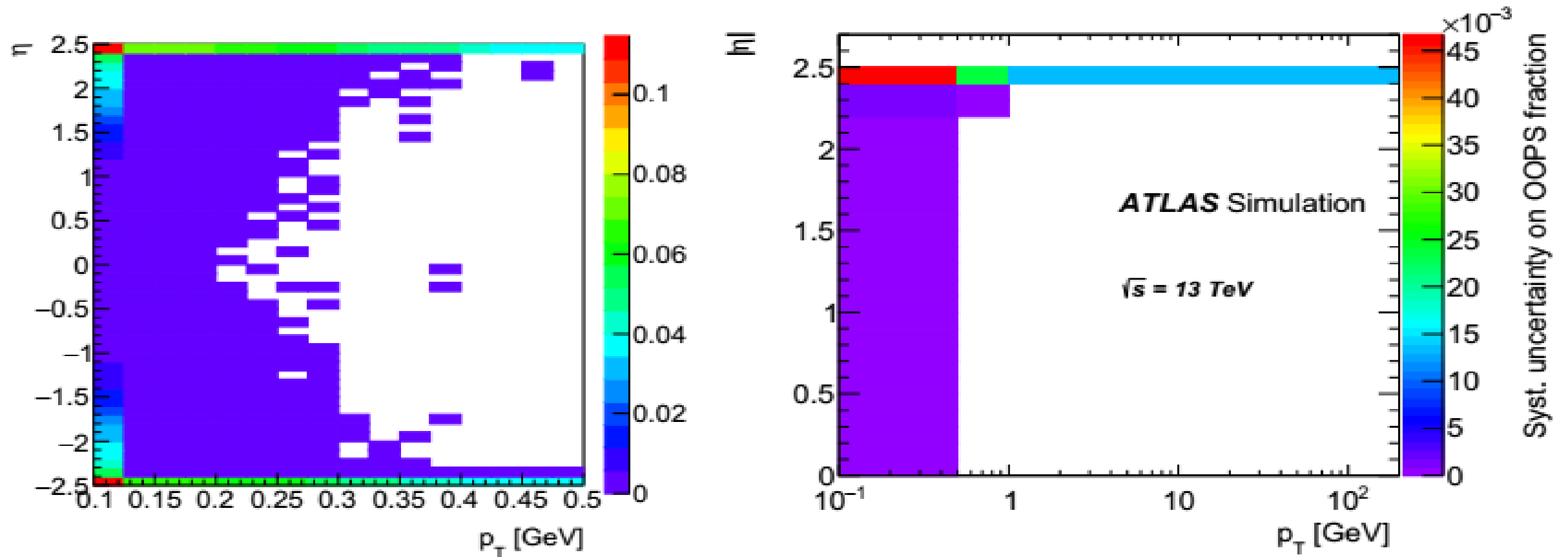


Figure 16: The out of phase space correction (OOPS) in  $p_T$  and  $\eta$  bins (left) and the systematic uncertainty on the out of phase space correction fractions (b). The systematic is made up of several contributions added up in quadrature, where each contribution is calculated as the difference in migration fractions between samples (see body text for further explanation).



# FAKE TRACK CORRECTION

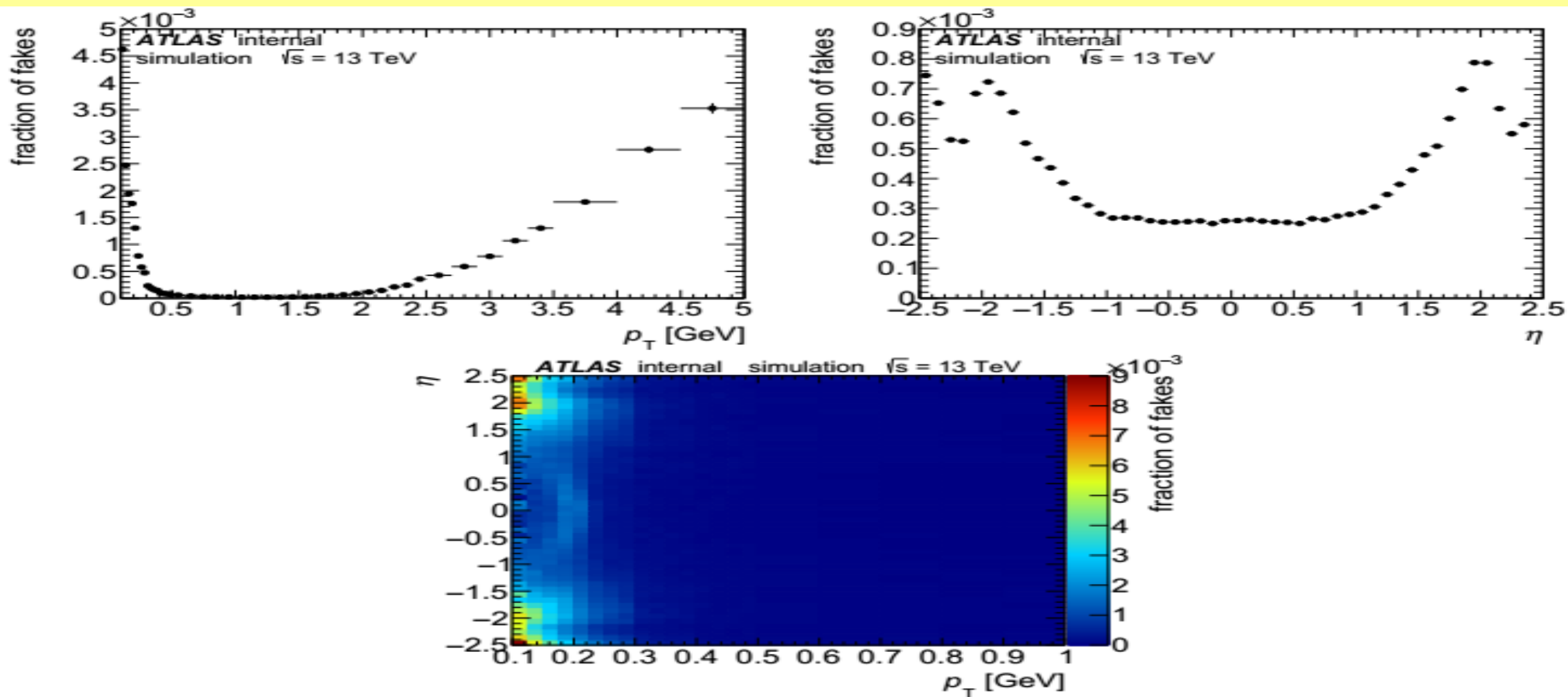
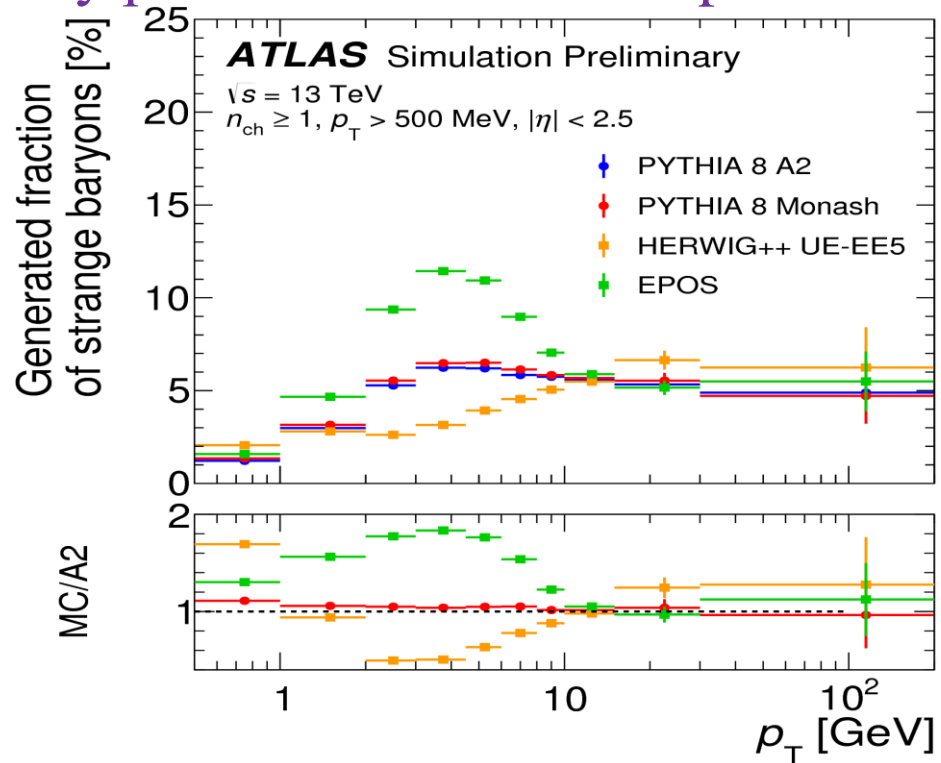


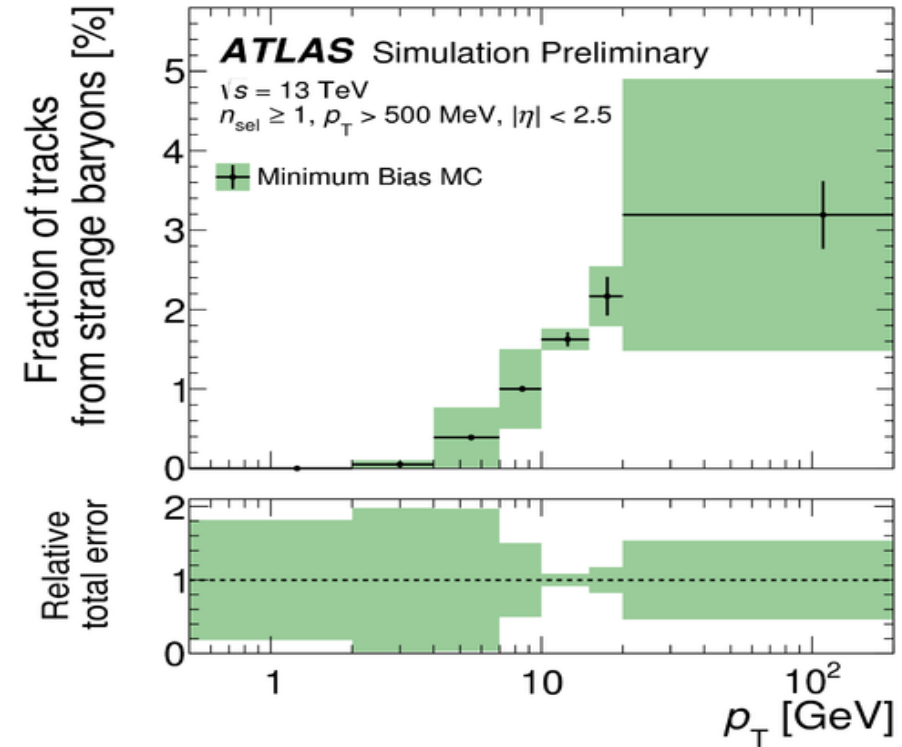
Figure 9: The fraction of fakes after applying the full event selection (see Section 3) as a function of  $p_T$  (left top) or  $\eta$  (right top) and the two-dimensional dependency of  $p_T$  and  $\eta$  (bottom). The fraction is below 1% and therefore negligible for the analysis.

# STRANGE BARYONS

- Particles with lifetime  $30 \text{ ps} < \tau < 300 \text{ ps}$  are no longer considered primary particles in the analysis, decay products are treated like secondary particles.
- All of these particles were **strange baryons**: with low reconstruction efficiency ( $< 0.1\%$ ) and large variations in predicted rates lead to a model dependence
- Primary particles have  $\tau > 300 \text{ ps}$



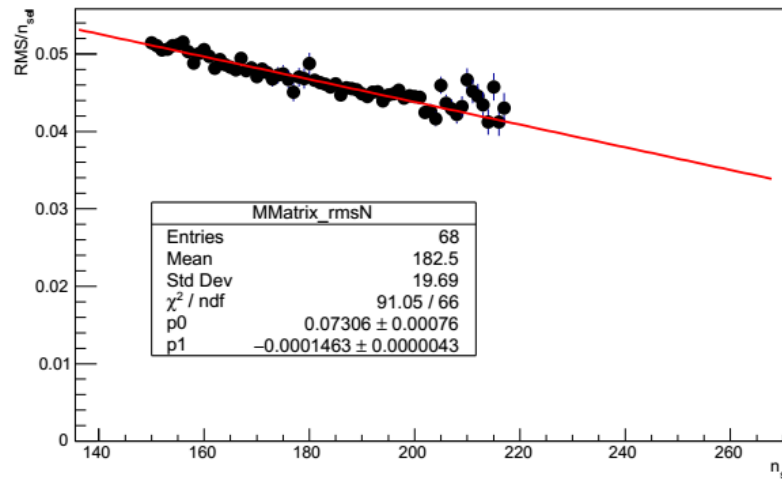
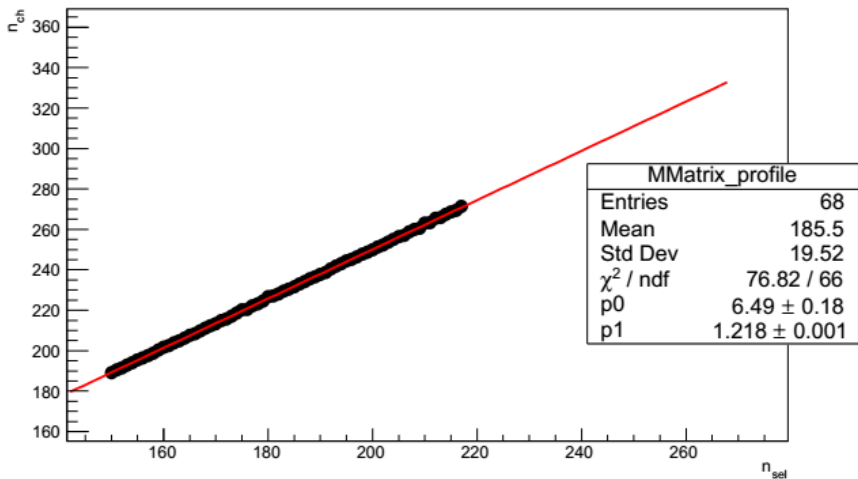
The fraction of strange baryons in generated particles as a function of particle  $p_{\text{T}}$  as predicted by various generators.



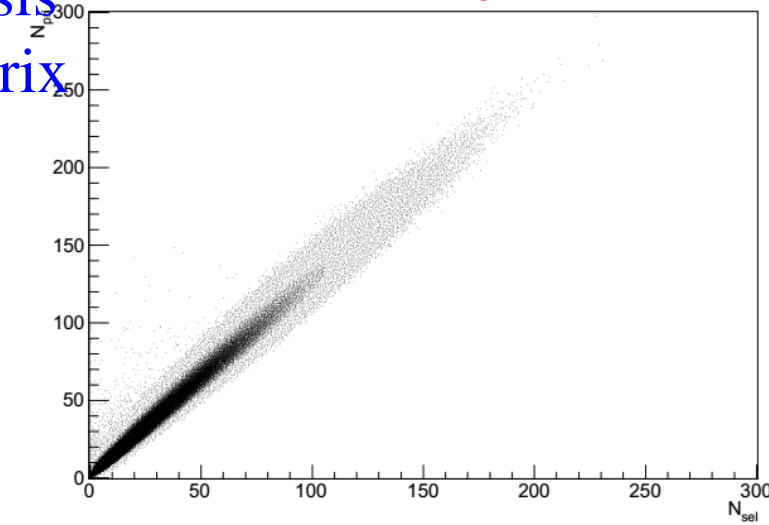
The fraction of reconstructed tracks coming from strange baryons as a function of track  $p_{\text{T}}$  as predicted by PYTHIA8 A2.

# MULTIPLICITY UNFOLDING

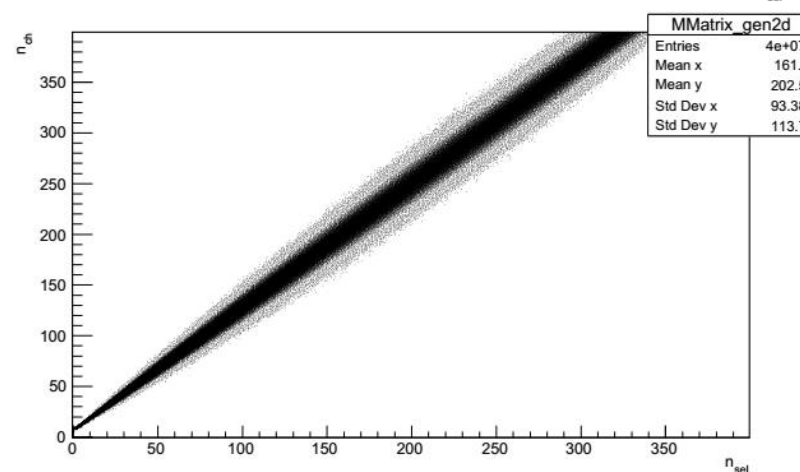
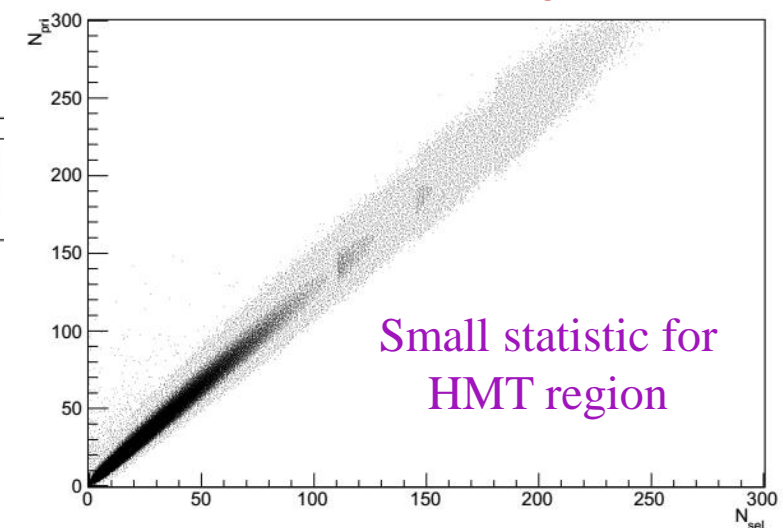
- **Multiplicity unfolding** is the same like for Minimum-bias analysis
- For High multiplicity track events the Multiplicity unfolding matrix additionally included the MC with filters: more 120, 160, 200



MB unfolding matrix



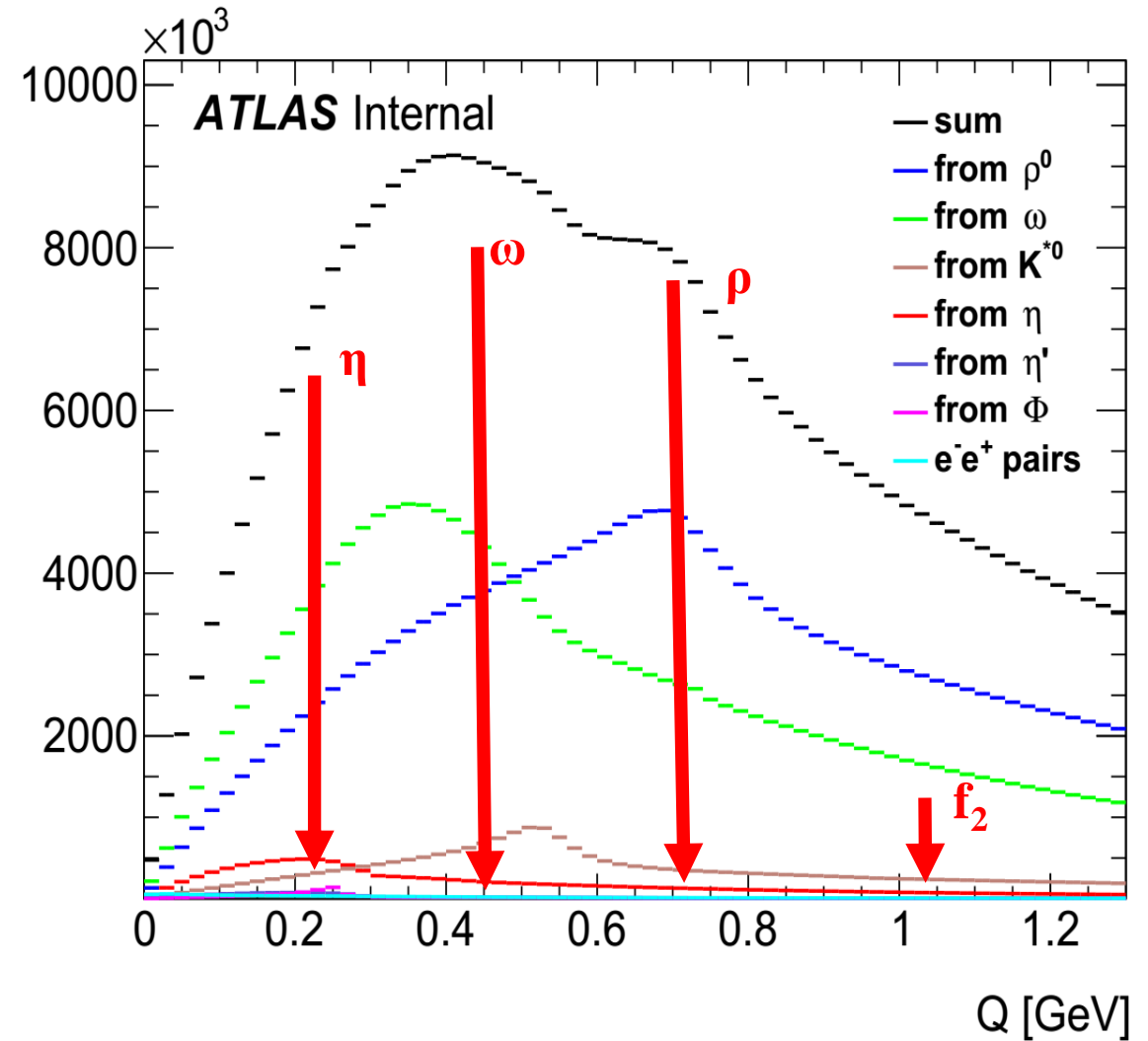
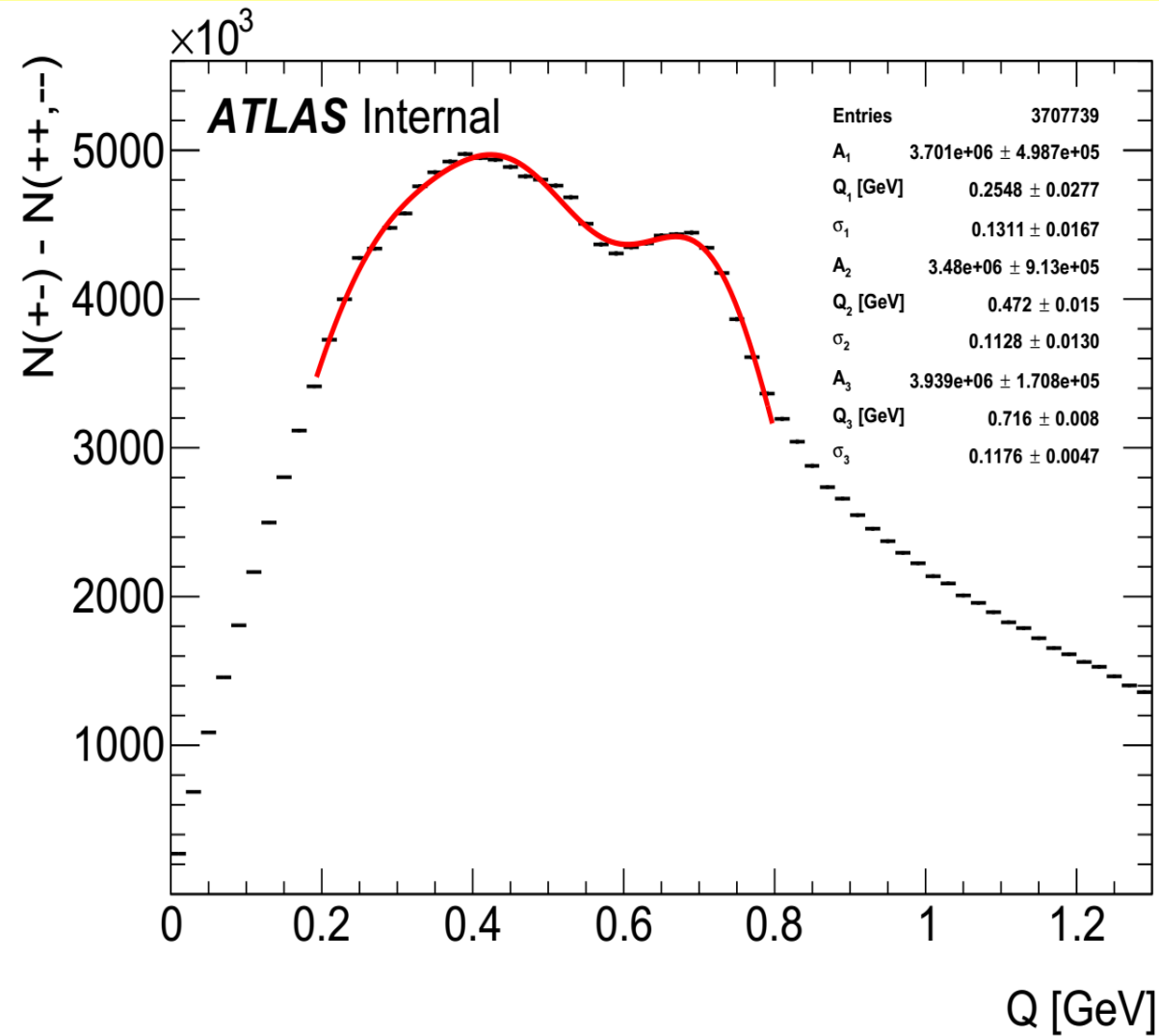
MB + HMT unfolding matrix



LHC days in Belarus, 11.04.2015

For high multiplicity region  $n_{\text{sel}} > 140$  with small statistic the calculation of mean and  $\text{RMS}/n_{\text{sel}}$  of unfolding matrix were done and used for calculation of unfolding matrix for high multiplicity region

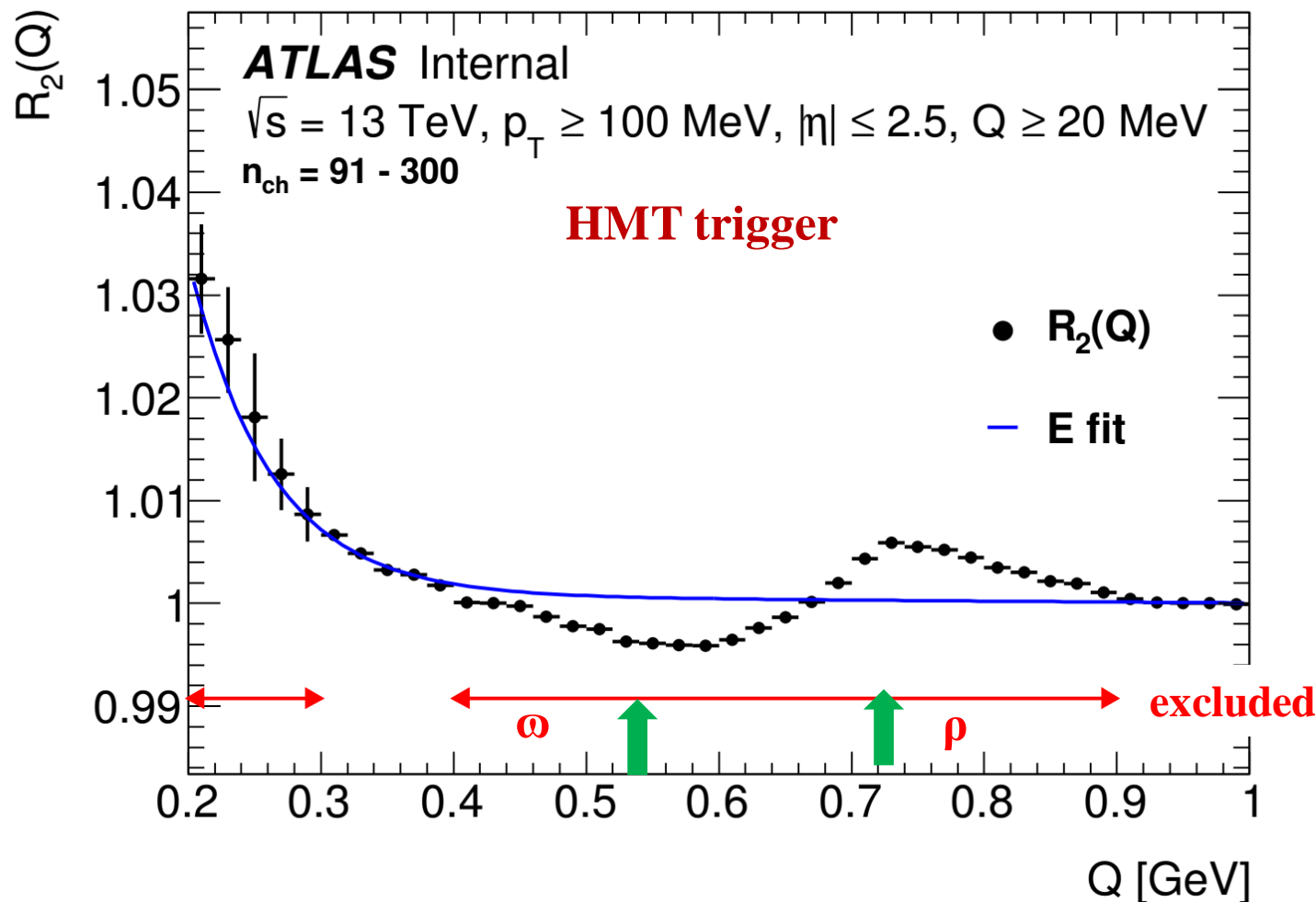
# RESONANCES STUDY



The Q spectrum generated by Pythia8 and the decomposition of its resonant part into leading contributions.



# ZOOM OF INCLUSIVE $R_2$ DISTRIBUTION



Inclusive HMT two particles  $R_2$  correlation function

$$R_2 = C_2^{\text{data}} / C_2^{\text{MC}} = N_{\text{data}}^{\pm\pm} / N_{\text{MC}}^{\pm\pm} \cdot N_{\text{MC}}^{+-} / N_{\text{data}}^{+-}$$

Three bump regions because MC underestimated or overestimates:

- 1)  $\eta \rightarrow \pi^+ \pi^- \pi^0$  and  $\eta' \rightarrow \pi^+ \pi^- \gamma$ ;
- 2)  $\omega \rightarrow \pi^+ \pi^- \pi^0$  and  $\rho \rightarrow \pi^+ \pi^-$ ,
- 3)  $f_2 \rightarrow \pi^+ \pi^-$ ;

The excluded regions at 13 TeV

- 1) 0.2–0.3 GeV – not important after non-closure correction;
- 2) 0.4–0.9 GeV – important;
- 3) 1.0–1.16 GeV (only for  $2 \leq n_{ch} \leq 40$  and  $100 \leq k_T \leq 200 \text{ MeV}$ ) – not important for BEC.

The excluded region at 7 TeV was 0.5–0.9 GeV.

## HOW COMES THAT THE $\rho$ - AND $\omega$ -MESONS ARE SO BADLY MODELED?

1. There are not MC, which is in agreement with all charged-particle distributions at 0.9 – 13 TeV.
2. The MC description of resonances are not so badly in scale of MC description of charged-particle distributions.
3. In  $R_2$  correlation functions the ratio of experimental  $N_{\text{data}}^{+-}$  to MC  $N_{\text{MC}}^{+-}$  unlike sign particles distribution are used:

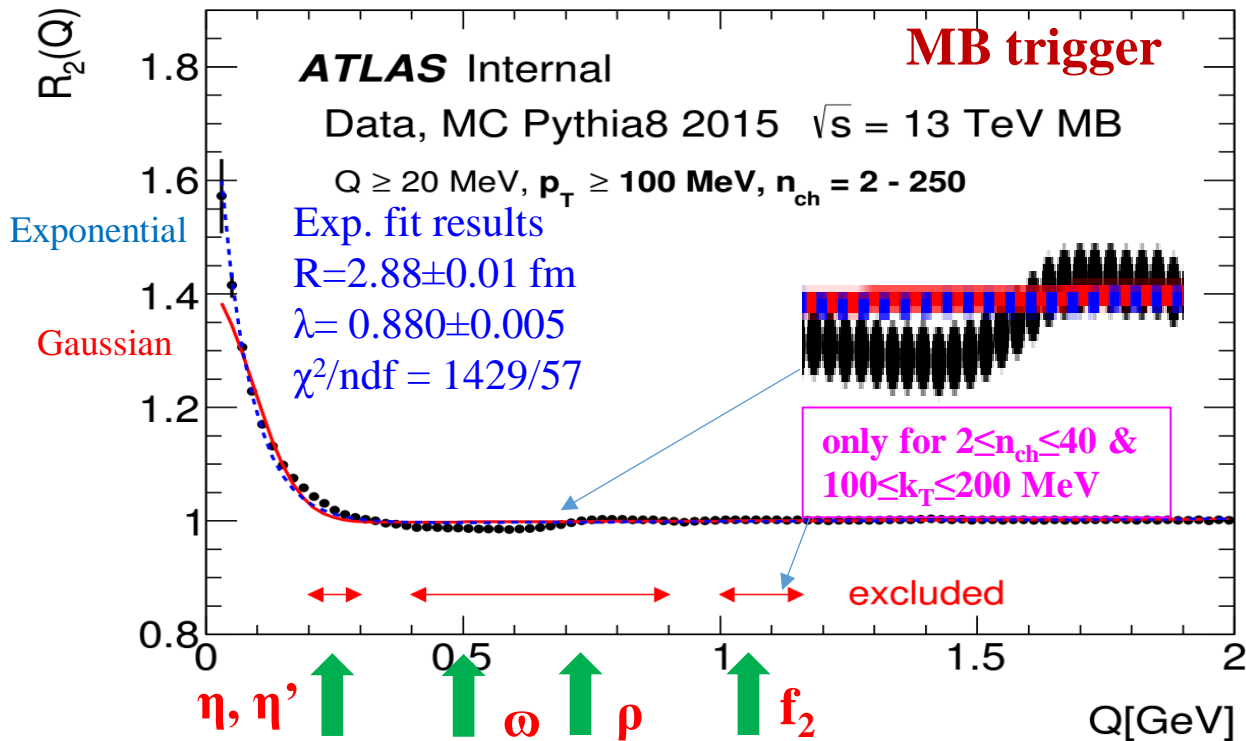
$$R_2 = C_2^{\text{data}} / C_2^{\text{MC}} = N_{\text{data}}^{++} / N_{\text{MC}}^{++} \cdot N_{\text{MC}}^{+-} / N_{\text{data}}^{+-}$$

4. One can see in Figure on previous page that MC little bit underestimate  $\omega$ -meson and little bit overestimate  $\rho$ -meson. We excluded region of these mesons because of small statistical errors and therefore large  $\chi^2/\text{ndf}$  for this region.
5. The situation as we have in our analysis is standard for similar analysis at other energies.

# INCLUSIVE $R_2$ DISTRIBUTIONS

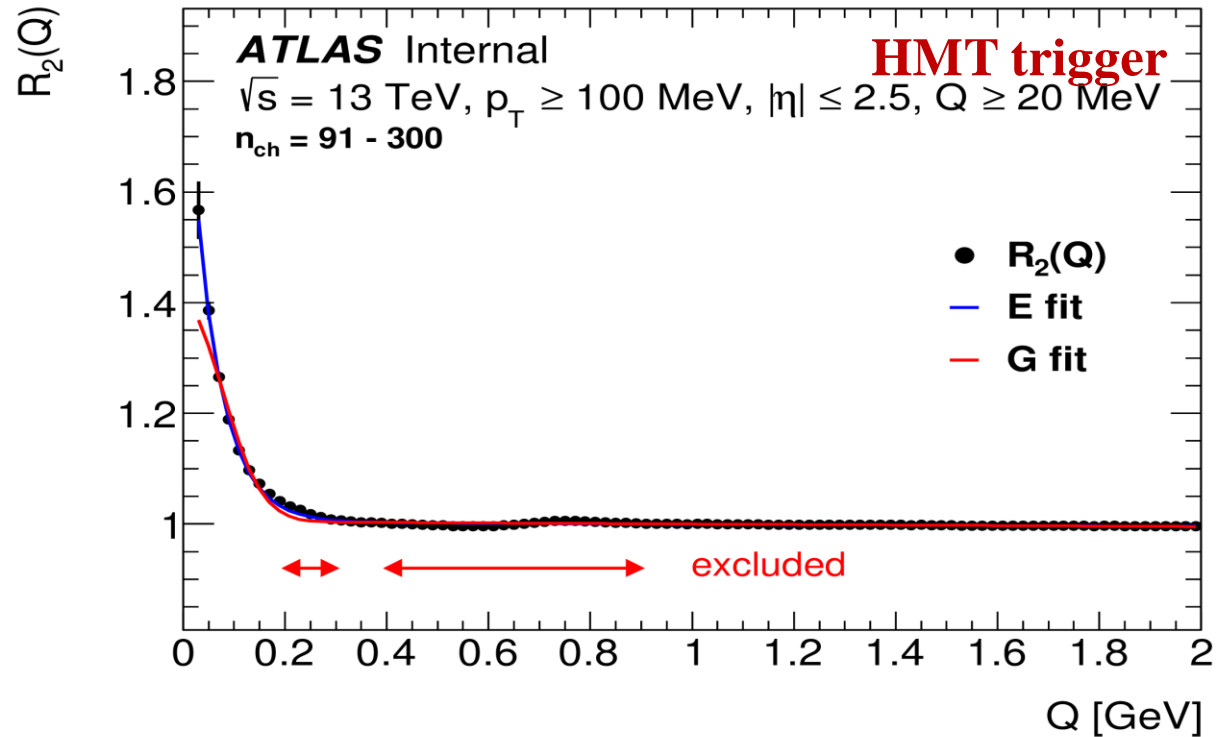
Inclusive MB two particles  $R_2$  correlation function

Lum.  $\sim 151 \mu\text{b}^{-1}$ ; Statistic:  $9.6 \times 10^6$  events with  $2.8 \times 10^8$  tracks



Inclusive HMT two particles  $R_2$  correlation function

Lum.  $\sim 8.4 \text{nb}^{-1}$ ; Statistic:  $9.1 \times 10^6$  events with  $9.8 \times 10^8$  tracks



Fit to extract strength and source size. Goldhaber spherical shape with a Gaussian distribution of the source. Exponential, radial Lorentzian distribution of the source -> much better at low Q. Three bumps regions because MC overestimates: 1)  $\eta \rightarrow \pi^+ \pi^- \pi^0$  or  $\eta \rightarrow \pi^+ \pi^- \gamma$ ; 2)  $\omega \rightarrow \pi^+ \pi^- \pi^0$  and  $\rho \rightarrow \pi^+ \pi^-$ , 3)  $f_2 \rightarrow \pi^+ \pi^-$ ; Therefore regions 0.2–0.3 GeV; 0.4–0.9 GeV and 1.0–1.16 GeV (only for  $2 \leq n_{\text{ch}} \leq 40$  and  $100 \leq k_T \leq 200$  MeV) excluded from the fit. Q region is from 0.02 to 2 GeV.

**Fit function:**  $R_2(Q) = C_0 [1 + \lambda \Omega(QR)] (1 + \epsilon Q)$ ,  $\epsilon$ -term counts for the long-range correlations

Studies of one-dimensional BEC effects in pp collisions for  $p_T > 100$  MeV and  $|\eta| < 2.5$  at 13 TeV

# COULOMB CORRECTION

The measured  $N(Q)$  distribution for like or unlike signed particle (track) pairs in presence of the Coulomb interaction is given by:

where  $N_{meas}(Q)$  is the measured distribution,  $N(Q)$  is the distribution free of Coulomb correlations.

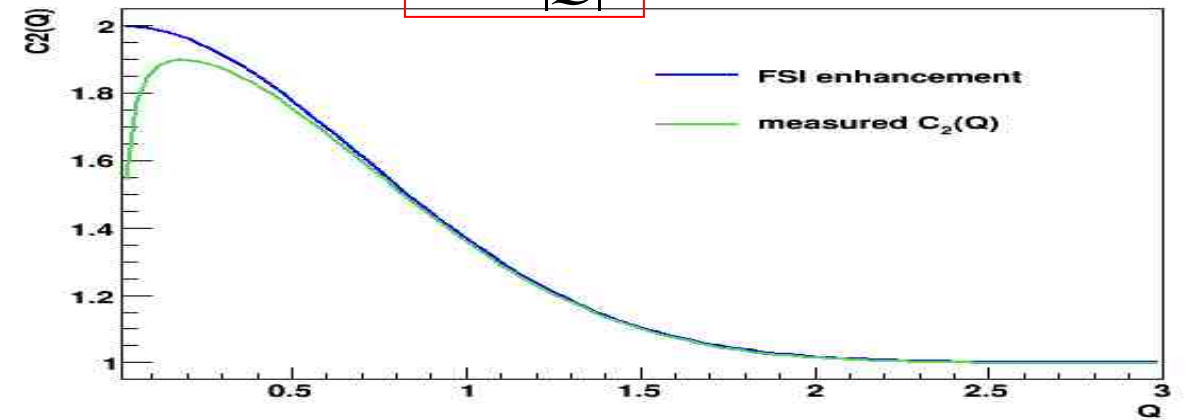
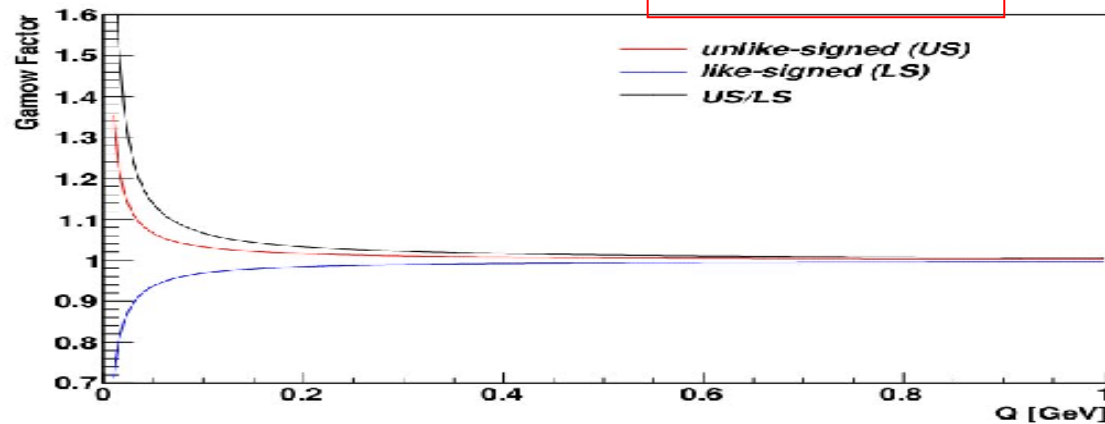
$$N_{meas}(Q) = G(Q)N(Q)$$

Gamow penetration  $G(Q)$  factor

$$G(Q) = \frac{2\pi\eta}{e^{2\pi\eta} - 1}$$

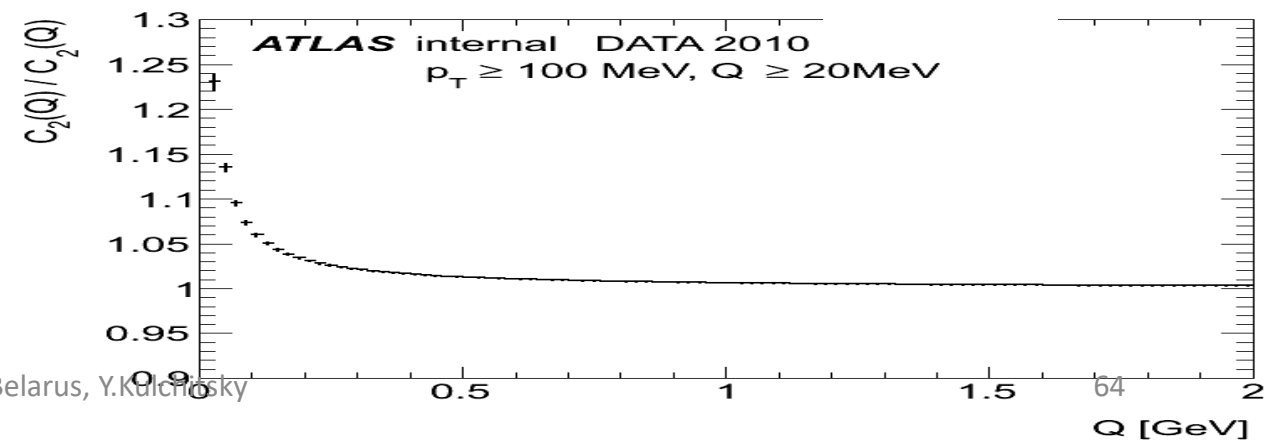
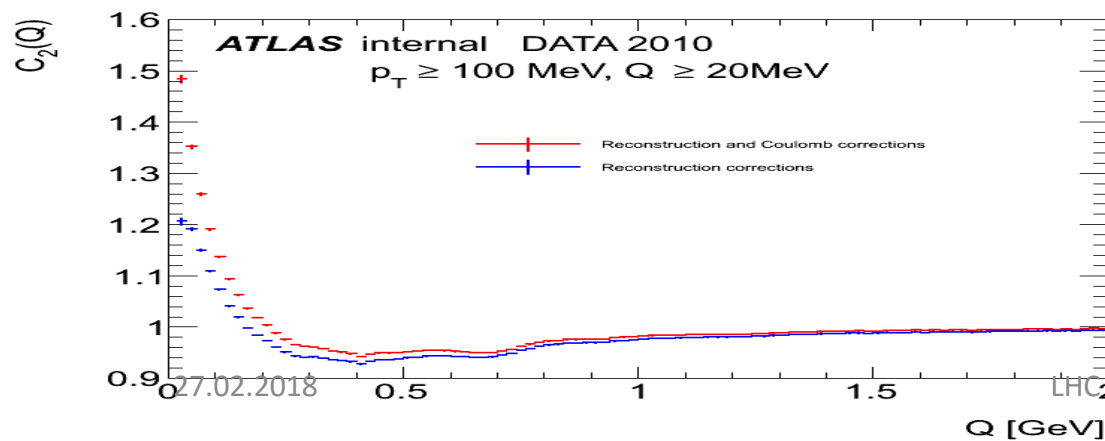
Sommerfeld parameter  $\eta$

$$\eta = \frac{\pm \alpha_m}{|Q|}$$



Comparison of the same  $C_2(Q)$  function: **blue** – primary data  $C_2(Q)$ , **red** –  $C_2(Q)$  corrected to Coulomb correction

**Ratio** of the same  $C_2(Q)$  function – with and without correction on Coulomb correction

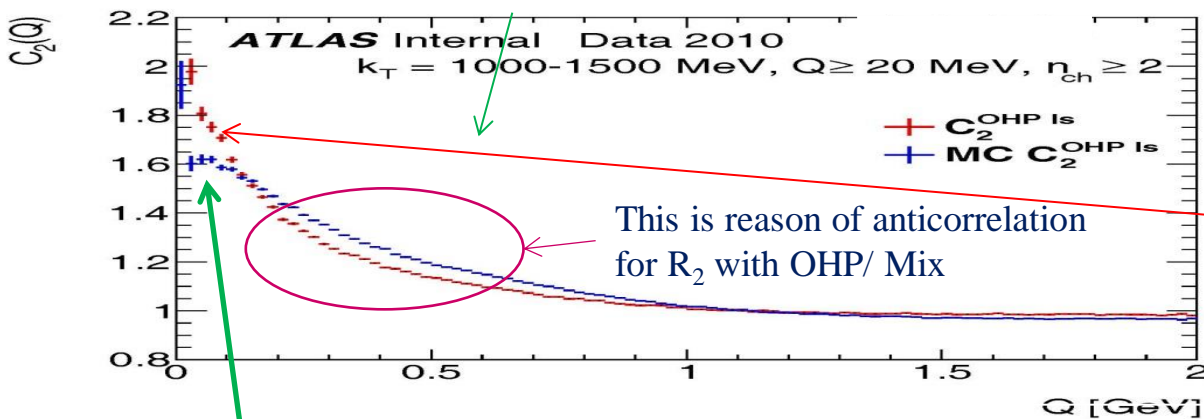




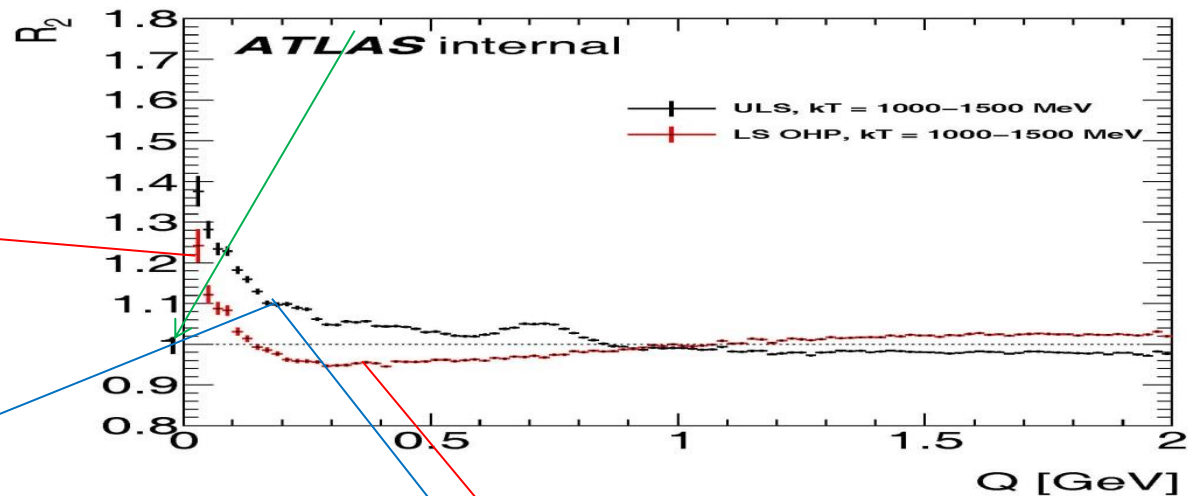
# OHP (MIX) AND UCP REFERENCE SAMPLES: $K_T$

The two-particle correlation function  $C_2(Q)$  for different  $k_T$  intervals using the opposite-hemisphere reference sample for data (red) and MC (blue)

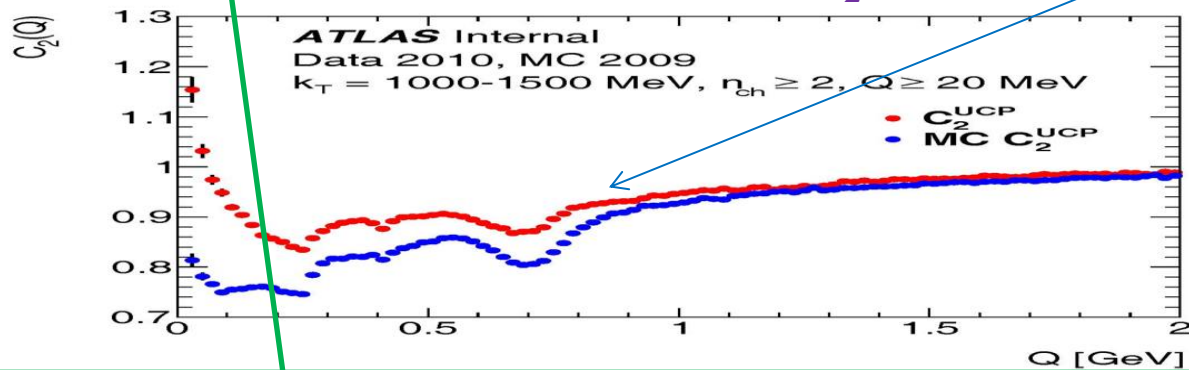
Artificial peak in  $C_2$  in BEC region



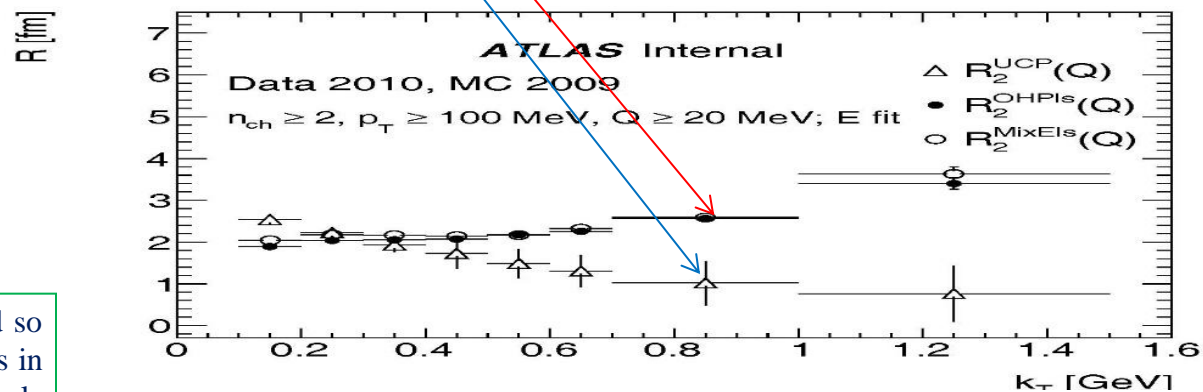
Small BEC peak for  $R_2$ , OHP



Reflection of resonances in  $C_2$  UCP



$k_T$  dependences for  $R$  for  $R_2$  with different reference

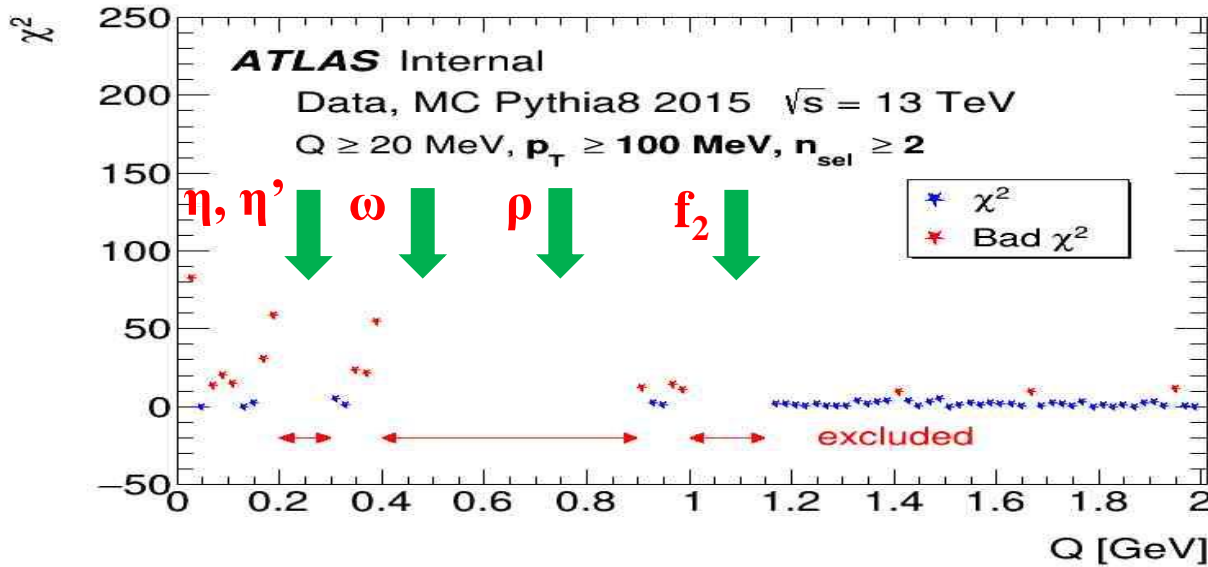


**Disadvantage for OHP/MIX:** violation of energy-momentum constraint, event topology, destroying other features such as non-BEC etc.

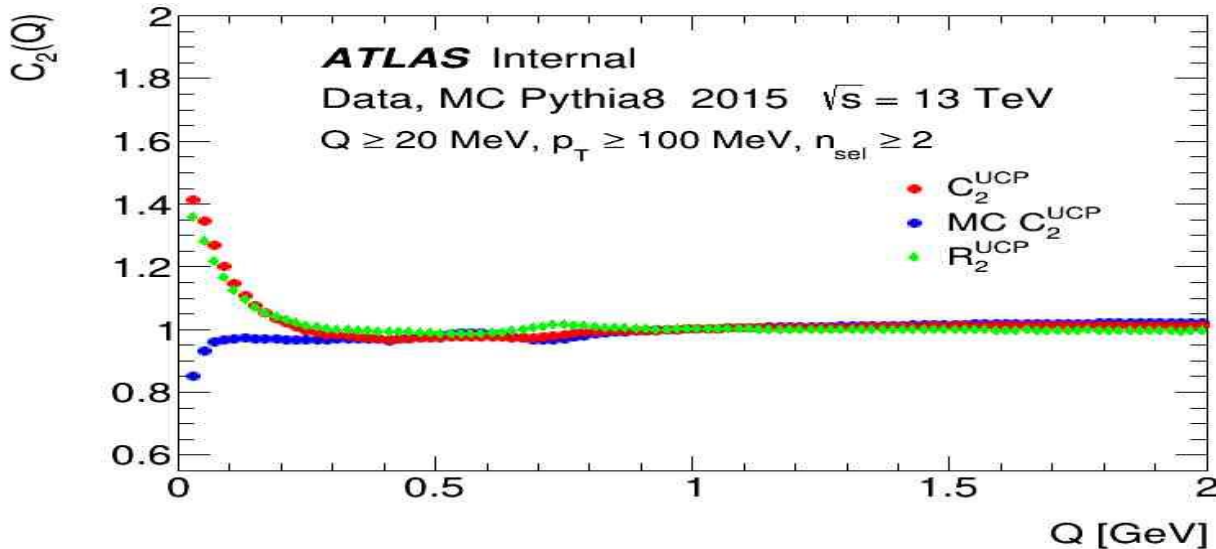
To note is that the slope in the MC can be explained by the fact that MC is tuned to the data and so reflects different dynamical constraints, but MC has no possibility to reproduce a peak at small  $Q$  as in the data but shows a broad enhancement.. The additional correlations in large multiplicity intervals seem to be due to multi-jet events in MC where the correlations between particles within the same jet can contribute to the region of low- $Q$ s. In this case, the single-ratio  $C_2(Q)$  correlation function numerator contains contributions from multi-jets, while the denominator does not have this effect as no correlations are expected in randomly paired particles.

# INCLUSIVE $R_2$ DISTRIBUTION FOR $P_T > 100$ MEV: $2 \leq N_{CH}$

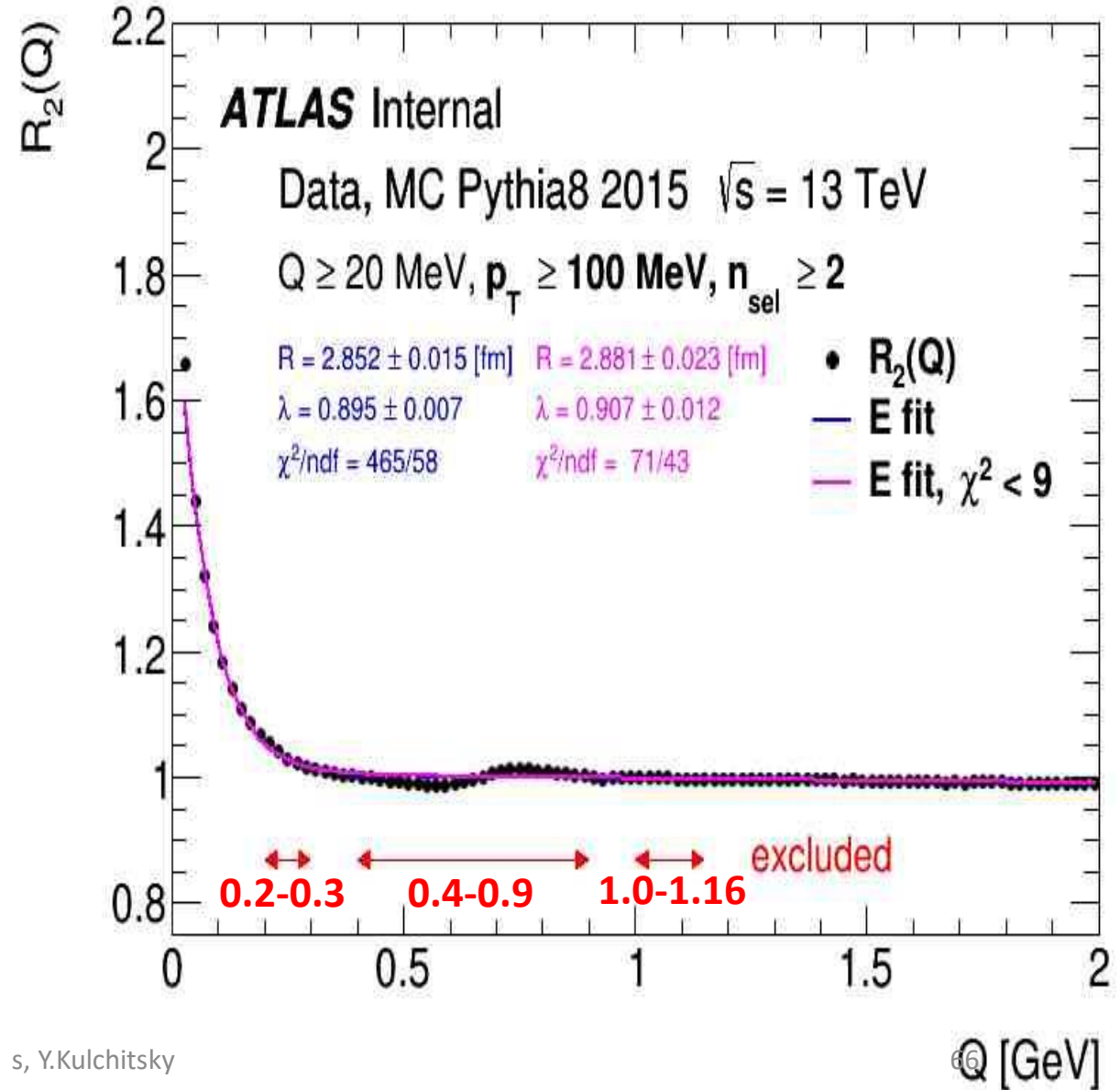
## $\chi^2$ vs Q distribution



## Comparison of $C_2$ and $R_2$ distributions

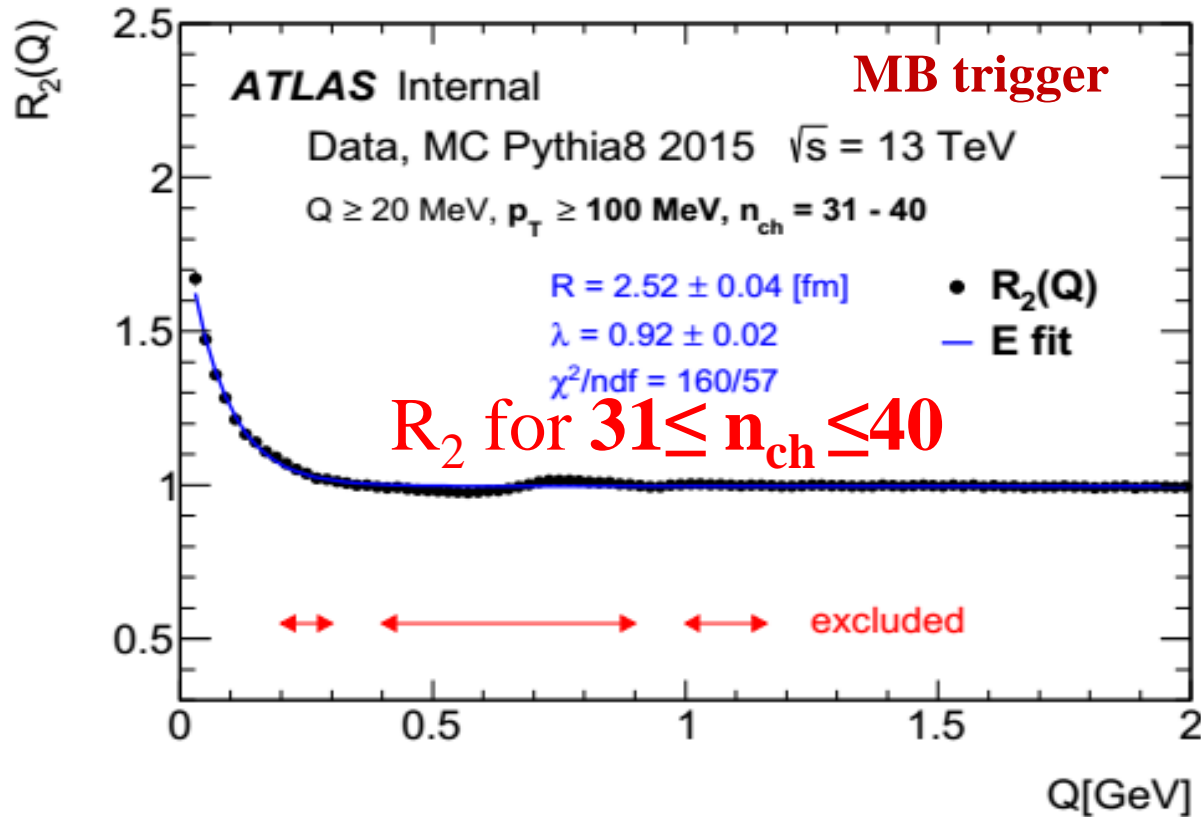


## Exponential fit results comparison for $R_2$

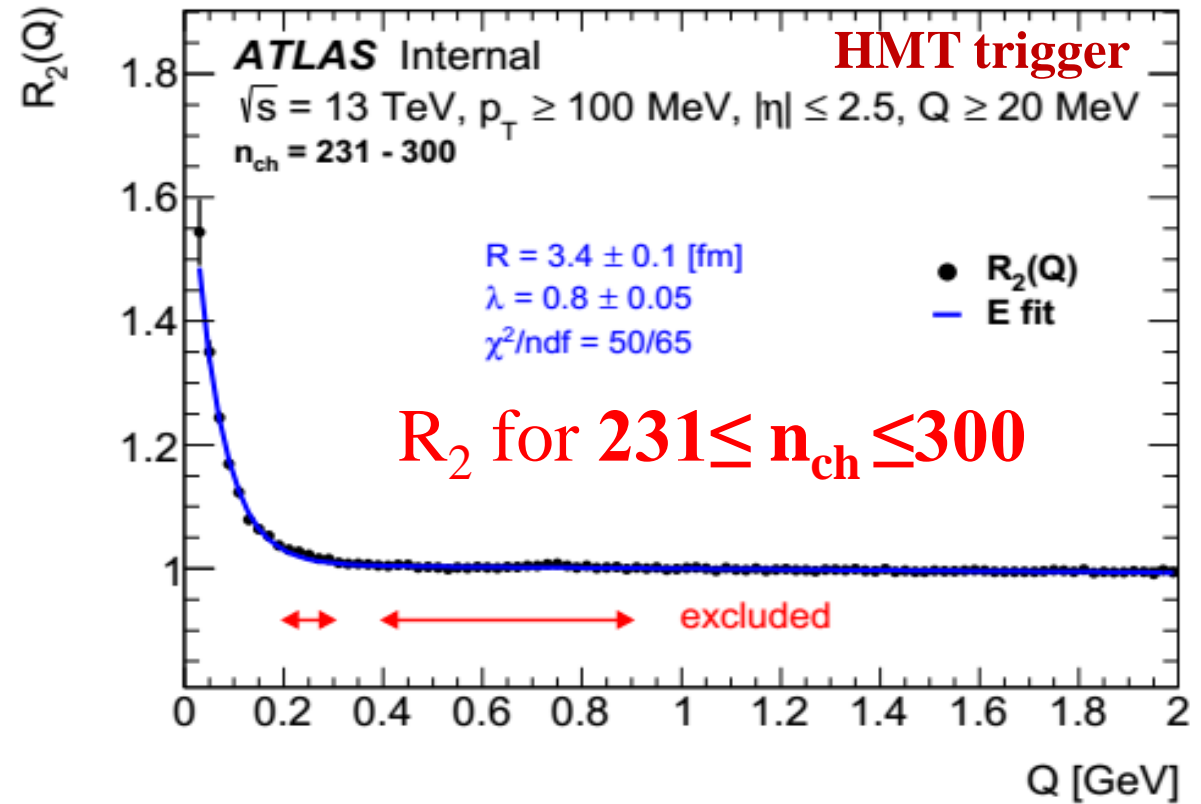


# TYPICAL TWO-PARTICLE $R_2$ CORRELATION FUNCTIONS

MB two particles  $R_2$  correlation function



HMT two particles  $R_2$  correlation function

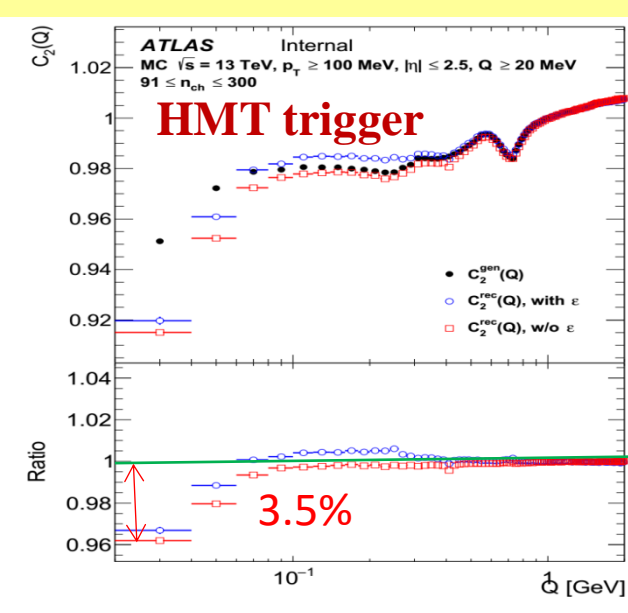
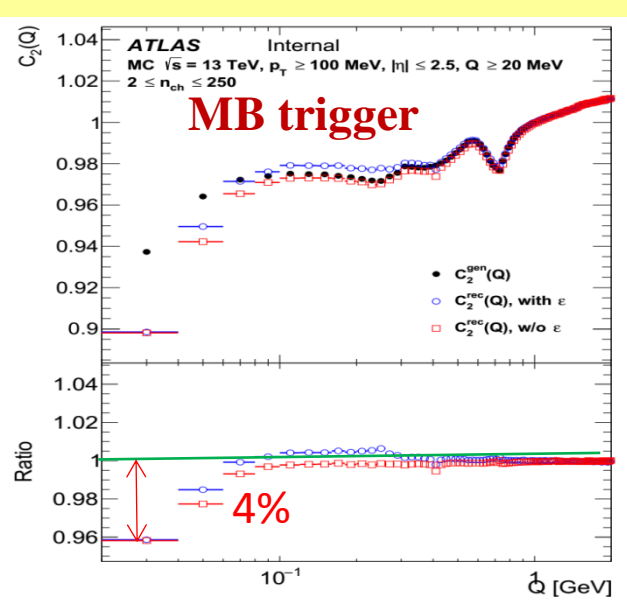
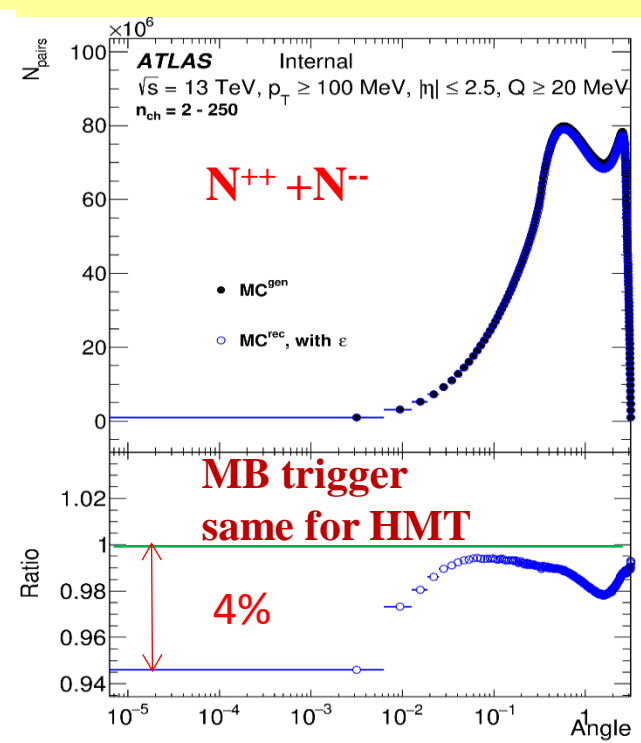
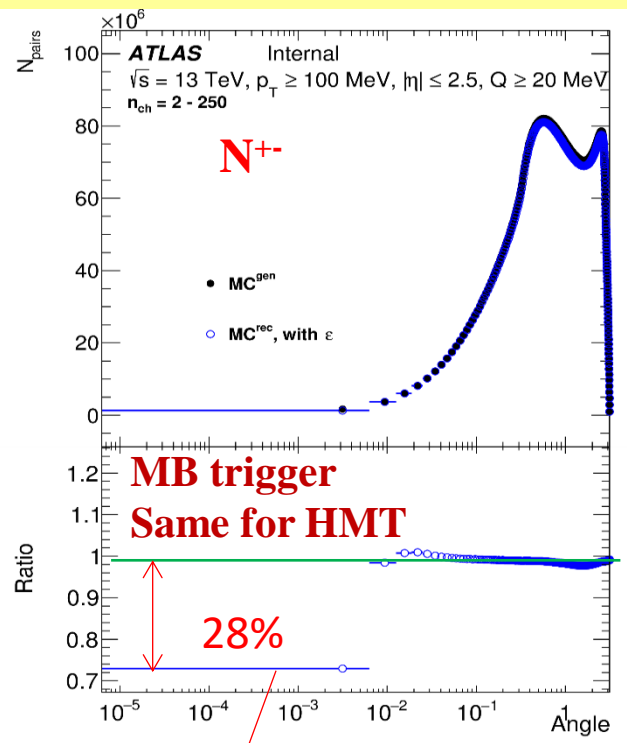


Fit to extract strength and source size. Goldhaber spherical shape with a Gaussian distribution of the source. Exponential, radial Lorentzian distribution of the source -> much better at low Q. Three bumps regions because MC overestimates: 1)  $\eta \rightarrow \pi^+ \pi^- \pi^0$  or  $\eta \rightarrow \pi^+ \pi^- \gamma$ ; 2)  $\omega \rightarrow \pi^+ \pi^- \pi^0$  and  $\rho \rightarrow \pi^+ \pi^-$ , 3)  $f_2 \rightarrow \pi^+ \pi^-$ ; Therefore regions 0.2–0.3 GeV; 0.4–0.9 GeV and 1.0–1.16 GeV (only for  $2 \leq n_{ch} \leq 40$  and  $100 \leq k_T \leq 200$  MeV) excluded from the fit. Q region is from 0.02 to 2 GeV.

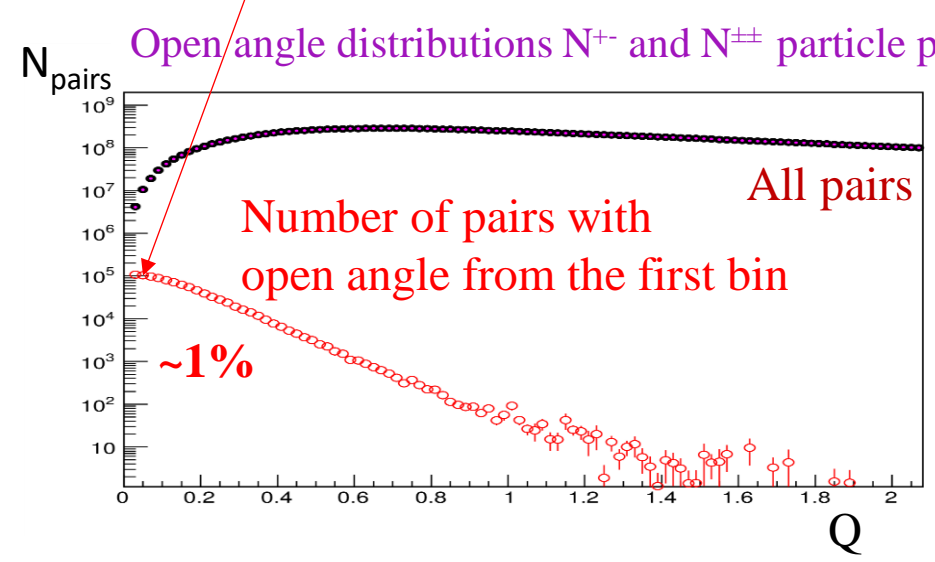
**Fit function:**  $R_2(Q) = C_0 [1 + \lambda \Omega(QR)] (1 + \epsilon Q)$ ,  $\epsilon$ -term counts for the long-range correlations

Studies of one-dimensional BEC effects in pp collisions for  $p_T > 100$  MeV and  $|\eta| < 2.5$  at 13 TeV

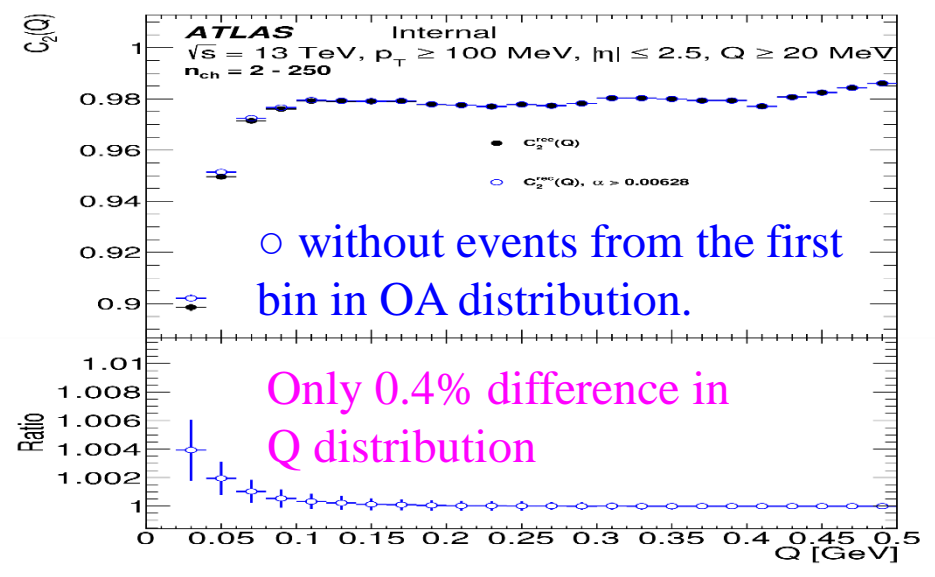
# OPEN ANGLE AND NON-CLOSURE FOR TWO PARTICLES



Closure test for  $C_2$  correlation functions for MB/HMT events

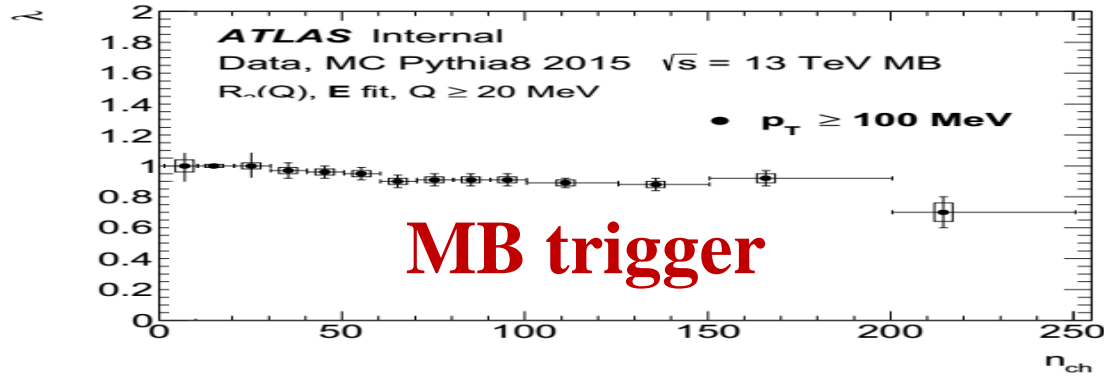


Q-distributions of  $N^{++}$  and  $N^{+-}$  particle pairs with all open angles (black circles), OA, and with open angles from the first bin (red circles). The fraction of the last is on the level of  $\sim 1\%$  for the  $0.02 < Q < 0.06 \text{ GeV}$  region

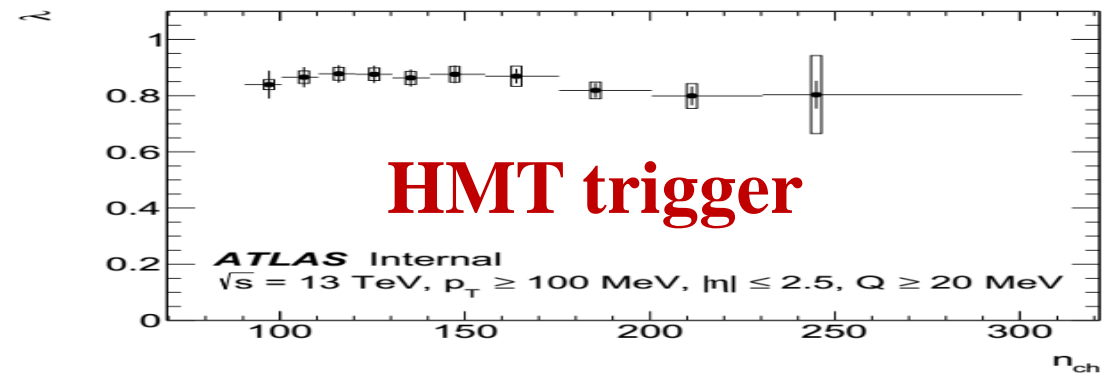




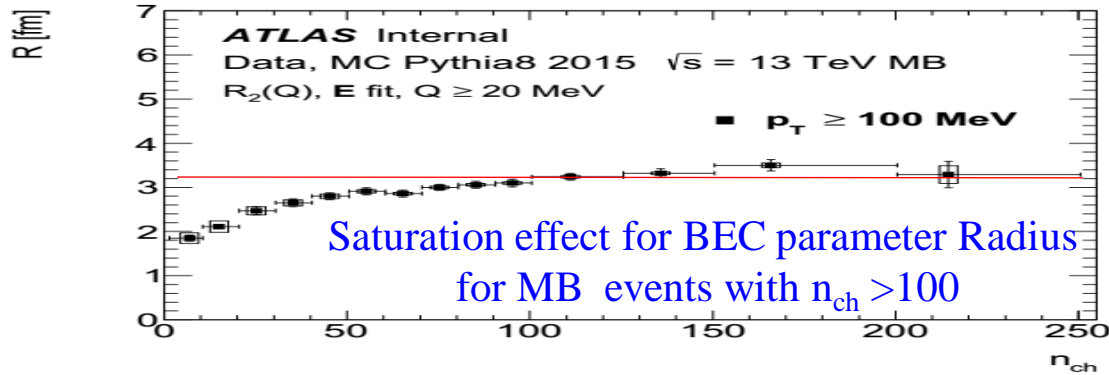
# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS



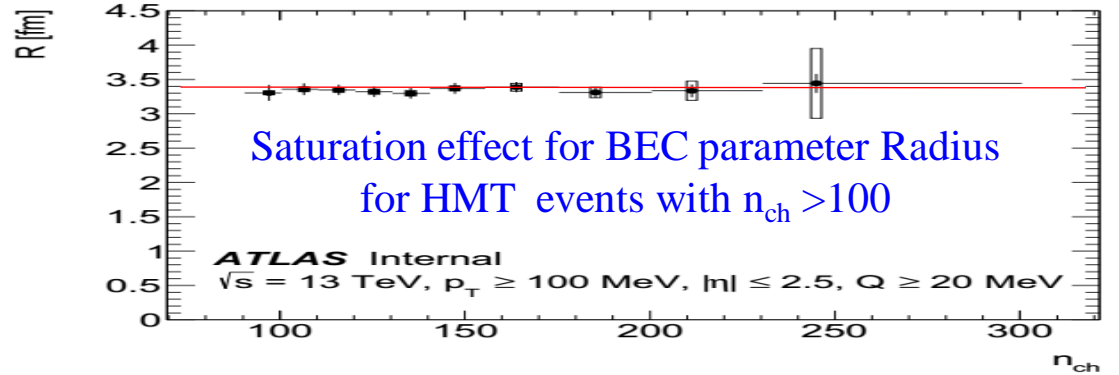
(a)



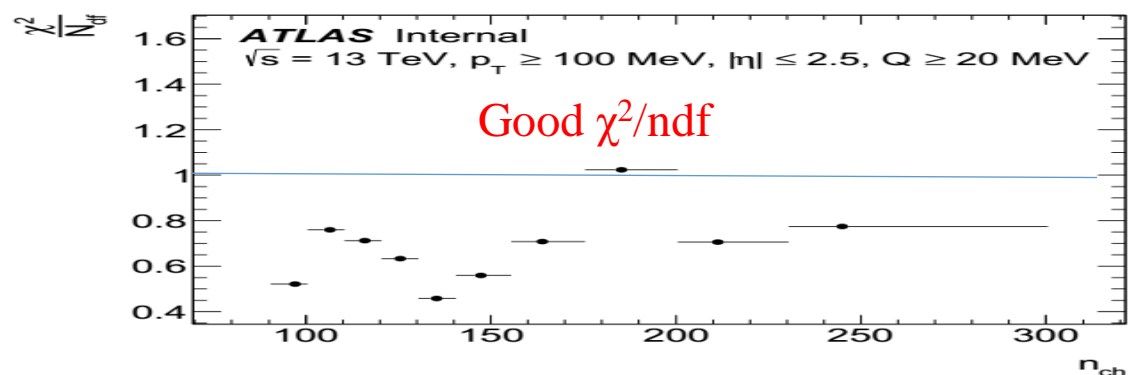
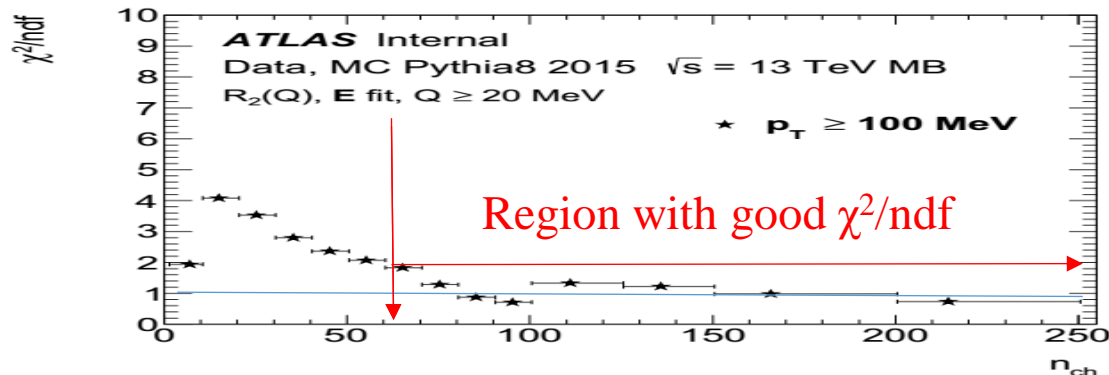
(b)



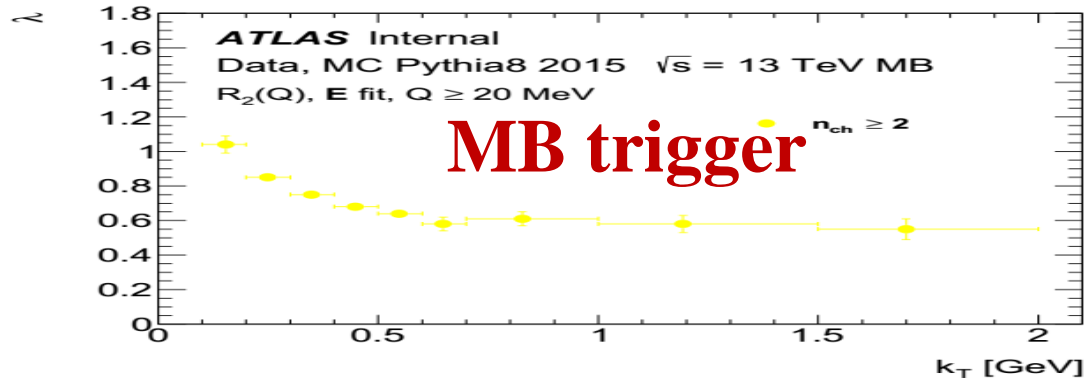
(c)



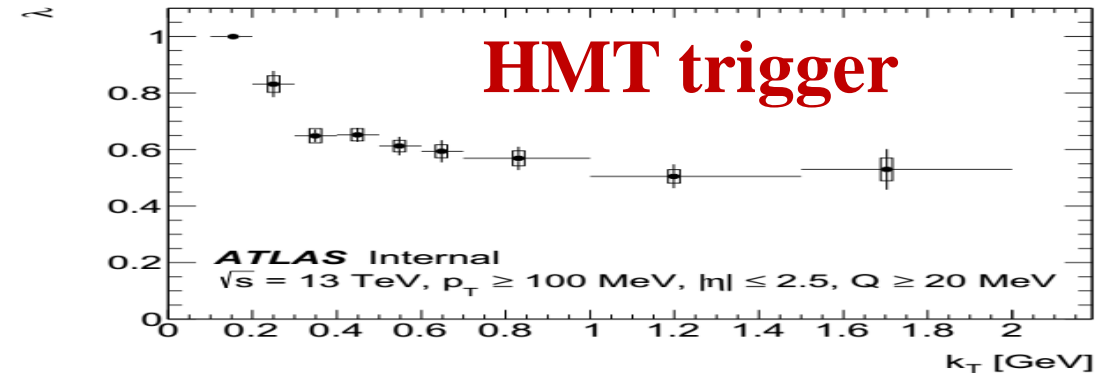
(d)



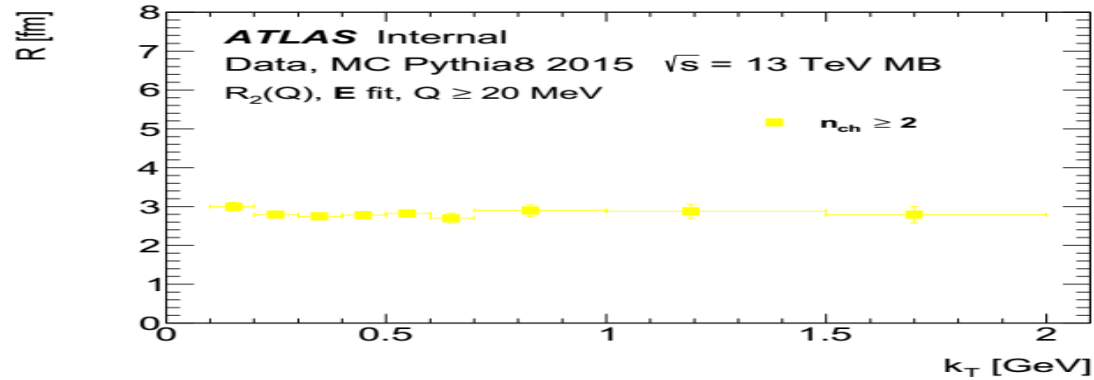
# $K_T$ DEPENDENCE OF BEC PARAMETERS



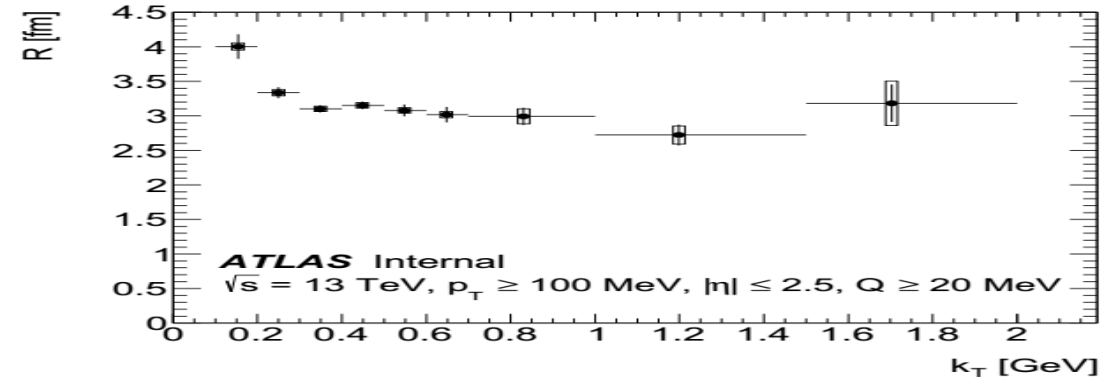
(a)



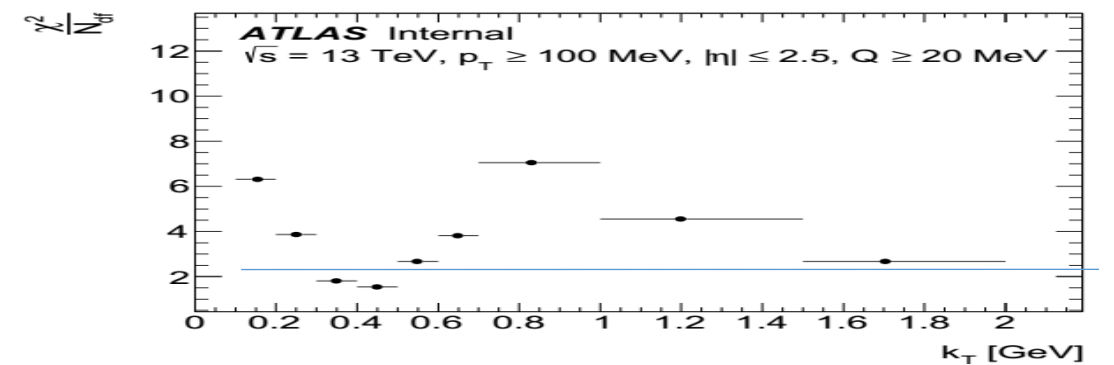
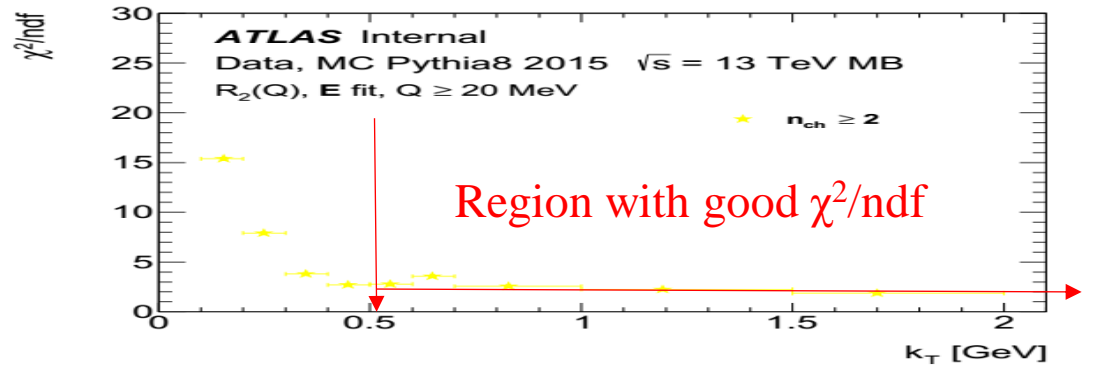
(b)



(c)



(d)



# FIT RESULTS FOR BEC PARAMETERS AT 13 TEV

The fit results of the BEC parameters  $R$  and  $\lambda$  dependence on the multiplicity,  $n_{\text{ch}}$ , and the transverse momentum of the pair,  $k_{\text{T}}$ , for different functional forms and for different data samples at 13 TeV. The errors represent the statistical and systematical uncertainties.

BEC param.	Fit function	13 TeV	
		Minimum-bias events	High-multiplicity track events
$R(n_{\text{ch}})$	$\alpha \sqrt[3]{n_{\text{ch}}}$ $\beta$	$\alpha = 0.77 \pm 0.02 \pm 0.02 \text{ fm } (n_{\text{ch}} \leq 70)$	— $\beta = 3.35 \pm 0.03 \pm 0.07 \text{ fm } (n_{\text{ch}} \geq 100)$
$\lambda(n_{\text{ch}})$	$\gamma e^{-\delta n_{\text{ch}}}$	$\gamma = 1.01 \pm 0.01 \pm 0.02$ $\delta = 0.0011 \pm 0.0001 \pm 0.0002$	
$R(k_{\text{T}})$	$\xi e^{-\kappa k_{\text{T}}}$	$\xi = 2.95 \pm 0.31 \pm 0.10 \text{ fm}$ $\kappa = 0.21 \pm 0.24 \pm 0.13 \text{ GeV}^{-1}$	$\xi = 3.42 \pm 0.19 \pm 0.01 \text{ fm}$ $\kappa = 0.20 \pm 0.12 \pm 0.03 \text{ GeV}^{-1}$
$\lambda(k_{\text{T}})$	$\mu e^{-\nu k_{\text{T}}}$	$\mu = 1.15 \pm 0.09 \pm 0.01$ $\nu = 1.22 \pm 0.19 \pm 0.07 \text{ GeV}^{-1}$	$\mu = 1.23 \pm 0.10 \pm 0.01$ $\nu = 1.15 \pm 0.21 \pm 0.09 \text{ GeV}^{-1}$

# FIT RESULTS FOR BEC PARAMETERS AT 0.9 – 7 TEV

The fit results of the BEC parameters  $R$  and  $\lambda$  dependence on the multiplicity,  $n_{\text{ch}}$ , and the transverse momentum of the pair,  $k_{\text{T}}$ , for different functional forms and for different data samples at 0.9 – 7 TeV. The errors represent the quadratic sum of the statistical and systematical uncertainties.

BEC param.	Fit function	0.9 TeV		7 TeV	
				Minimum-bias events	High-multiplicity events
$R(n_{\text{ch}})$	$\alpha \sqrt[3]{n_{\text{ch}}}$ $\beta$	$\alpha = 0.64 \pm 0.07 \text{ fm } (n_{\text{ch}} \leq 82)$		$\alpha = 0.63 \pm 0.05 \text{ fm } (n_{\text{ch}} < 55)$	—
		—		$\beta = 2.28 \pm 0.32 \text{ fm } (n_{\text{ch}} \geq 55)$	
$\lambda(n_{\text{ch}})$	$\gamma e^{-\delta n_{\text{ch}}}$	$\gamma = 1.06 \pm 0.10$		$\gamma = 0.96 \pm 0.07$	
		$\delta = 0.011 \pm 0.004$		$\delta = 0.0038 \pm 0.0008$	
$R(k_{\text{T}})$	$\xi e^{-\kappa k_{\text{T}}}$	$\xi = 2.64 \pm 0.33 \text{ fm}$		$\xi = 2.88 \pm 0.27 \text{ fm}$	$\xi = 3.39 \pm 0.54 \text{ fm}$
		$\kappa = 1.48 \pm 0.67 \text{ GeV}^{-1}$		$\kappa = 1.05 \pm 0.58 \text{ GeV}^{-1}$	$\kappa = 0.92 \pm 0.73 \text{ GeV}^{-1}$
$\lambda(k_{\text{T}})$	$\mu e^{-\nu k_{\text{T}}}$	$\mu = 1.20 \pm 0.18$		$\mu = 1.12 \pm 0.10$	$\mu = 0.75 \pm 0.10$
		$\nu = 2.00 \pm 0.35 \text{ GeV}^{-1}$		$\nu = 1.54 \pm 0.26 \text{ GeV}^{-1}$	$\nu = 0.91 \pm 0.45 \text{ GeV}^{-1}$



- The systematic uncertainties of the inclusive Bose-Einstein correlation parameters,  $\mathbf{R}$  (the effective radius parameter of the source) and  $\lambda$  (the strength of the effect parameter), of the *fit of  $R_2(Q)$  correlation functions with exponential model* are summarized in the Table.
- The systematic uncertainties are combined by adding them in quadrature and the resulting values are given in the bottom row.
- The same sources of uncertainty are considered for the differential measurements in  $n_{ch}$  and the average transverse momentum  $k_T$  of a pair, and their impact on the fit parameters is found to be similar in size.

Source	0.9 TeV		7 TeV		7 TeV (HM)	
	$\lambda$	$R$	$\lambda$	$R$	$\lambda$	$R$
Track reconstruction efficiency	0.6%	0.7%	0.3%	0.2%	1.3%	0.3%
Track splitting and merging	negligible		negligible		negligible	
Monte Carlo samples	14.5%	12.9%	7.6%	10.4%	5.1%	8.4%
Coulomb correction	2.6%	0.1%	5.5%	0.1%	3.7%	0.5%
Fitted range of $Q$	1.0%	1.6%	1.6%	2.2%	5.5%	6.0%
Starting value of $Q$	0.4%	0.3%	0.9%	0.6%	0.5%	0.3%
Bin size	0.2%	0.2%	0.9%	0.5%	4.1%	3.4%
Exclusion interval	0.2%	0.2%	1%	0.6%	0.7%	1.1%
Total	14.8%	13.0%	9.6%	10.7%	9.4%	10.9%

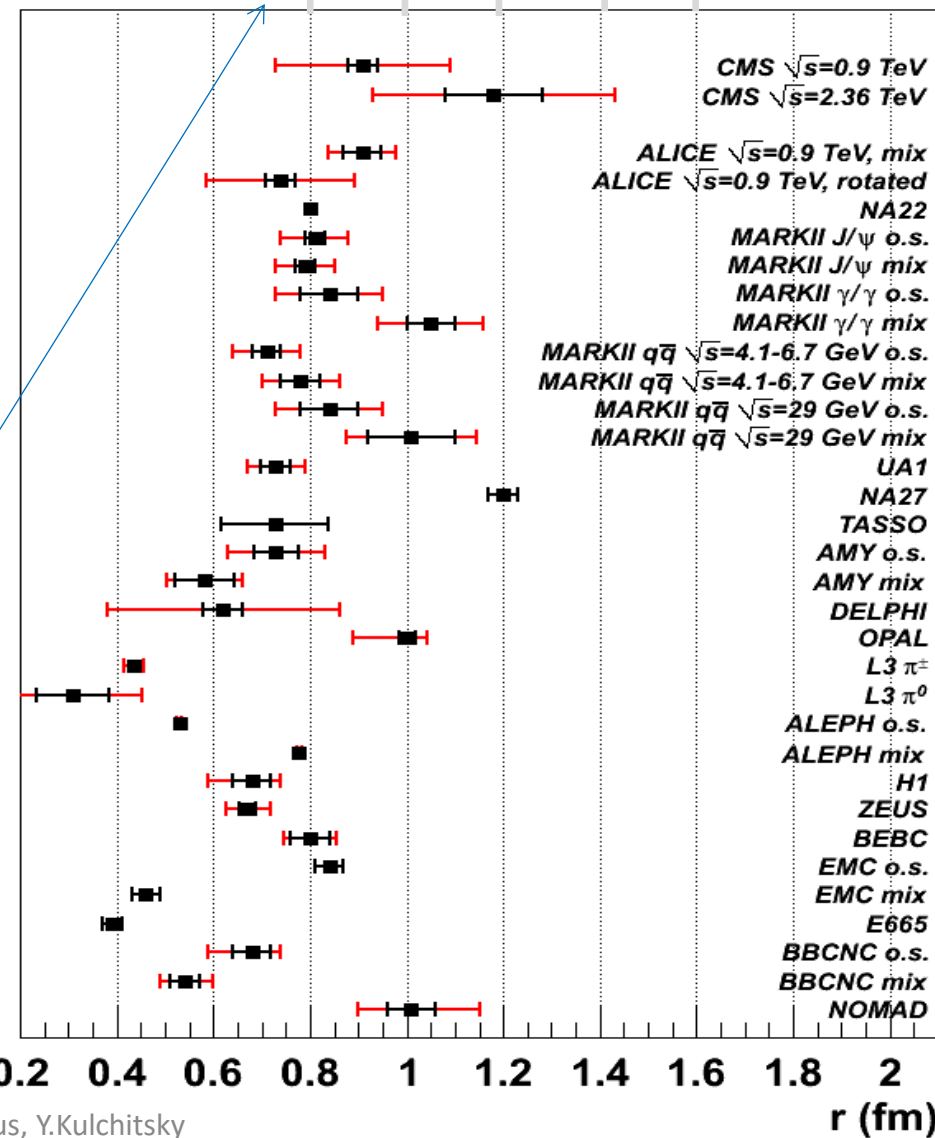
# COMPARISON WITH OTHER EXPERIMENTS

The results of BEC parameters for  
**Exponential fits of  $R_2$  used total uncertainties**  
 Statistical uncertainties are below 2-4 %

Energy [GeV]	$n_{ch}$	$\lambda$	$R^{(E)}$ [fm]
0.9	$\geq 2$	$0.74 \pm 0.10$	$1.83 \pm 0.25$
7	$\geq 2$	$0.71 \pm 0.07$	$2.06 \pm 0.22$
7 (HMT)	$\geq 150$	$0.74 \pm 0.06$	$2.36 \pm 0.30$
13	$\geq 2$	$0.88 \pm 0.05$	$2.88 \pm 0.05$
13 (HMT)	$\geq 93$	$0.86 \pm 0.02$	$3.31 \pm 0.04$

ATLAS  $\sqrt{s} = 13$  TeV  
 ATLAS  $\sqrt{s} = 7$  TeV HMT  
 ATLAS  $\sqrt{s} = 7$  TeV  
 ATLAS  $\sqrt{s} = 0.9$  TeV

Larger kinematic region

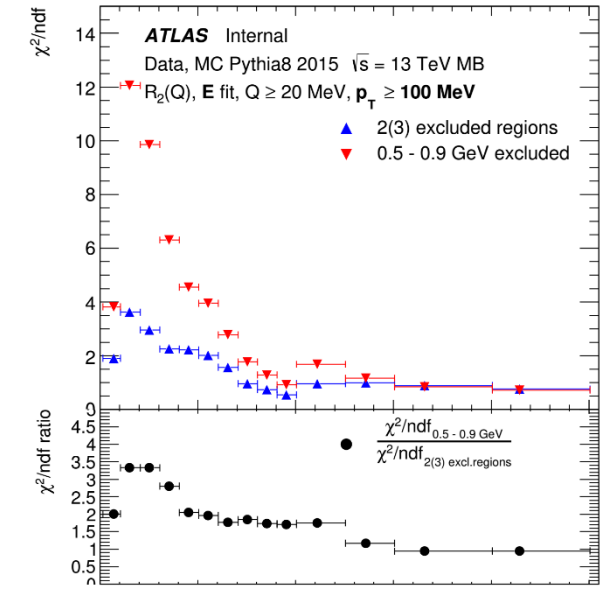
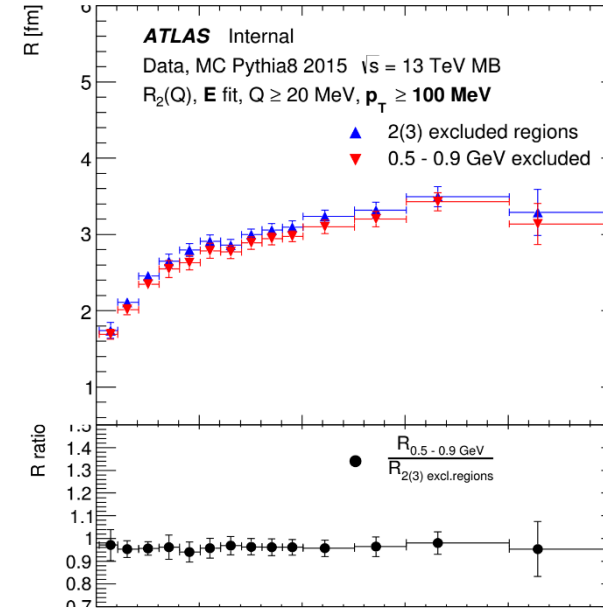
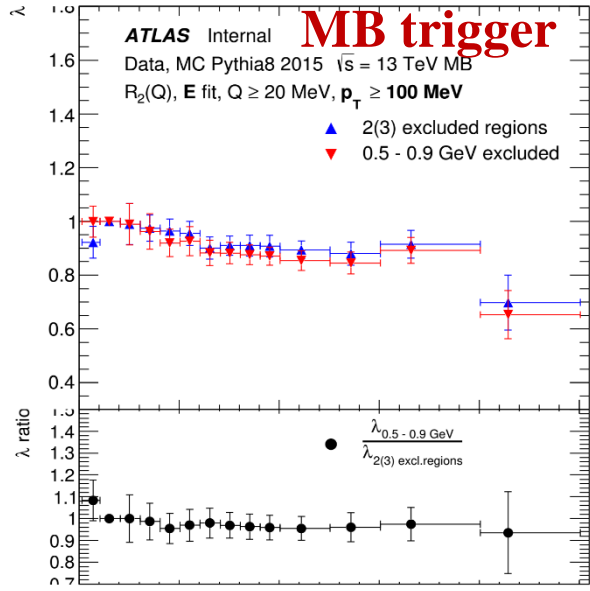


Comparison with results of previous experiments

$$R^{(G)} = R^{(E)} / \sqrt{\pi}$$

Energy [GeV]	$n_{ch}$	$R^{(G)}$ [fm]
0.9	$\geq 2$	$1.03 \pm 0.14$
7	$\geq 2$	$1.16 \pm 0.12$
7 (HMT)	$\geq 150$	$1.33 \pm 0.17$
13	$\geq 2$	$1.58 \pm 0.03$
13 (HMT)	$\geq 93$	$1.87 \pm 0.02$

# COMPARISON OF MULTIPLICITY RESULTS WITH ONE EXCLUDED REGION 0.5 – 0.9 GeV



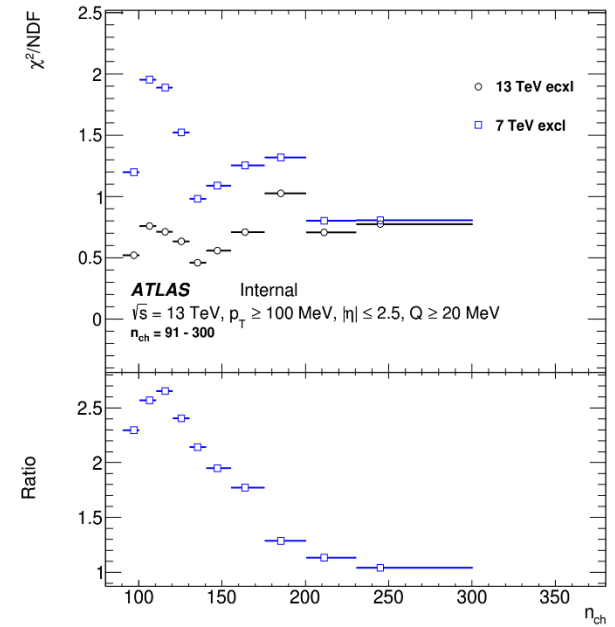
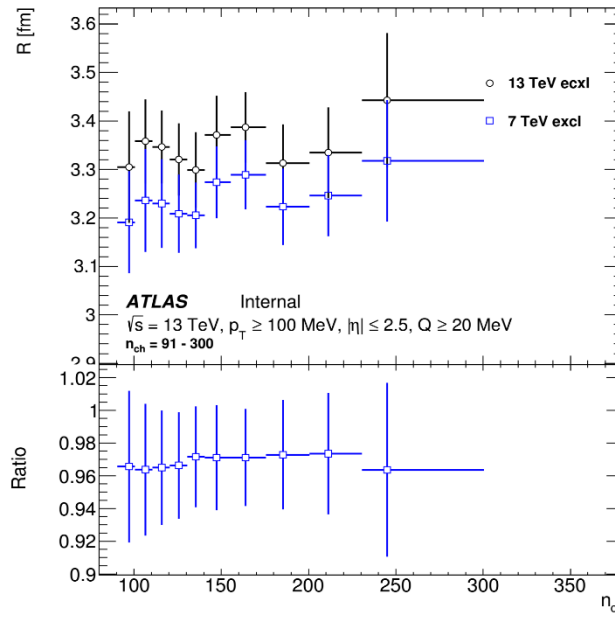
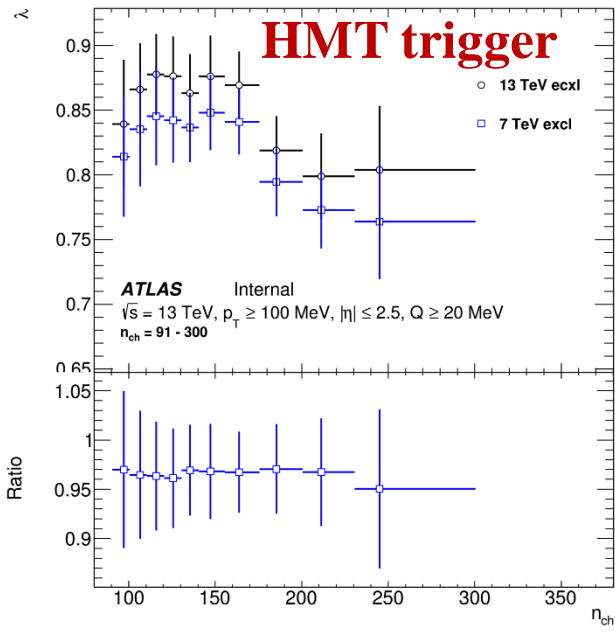
Three bump regions because MC underestimated or overestimates:

1.  $\eta \rightarrow \pi^+ \pi^- \pi^0$  and  $\eta' \rightarrow \pi^+ \pi^- \gamma$ ;
2.  $\omega \rightarrow \pi^+ \pi^- \pi^0$  and  $\rho \rightarrow \pi^+ \pi^-$ ;
3.  $f_2 \rightarrow \pi^+ \pi^-$ ;

The excluded regions at 13 TeV are

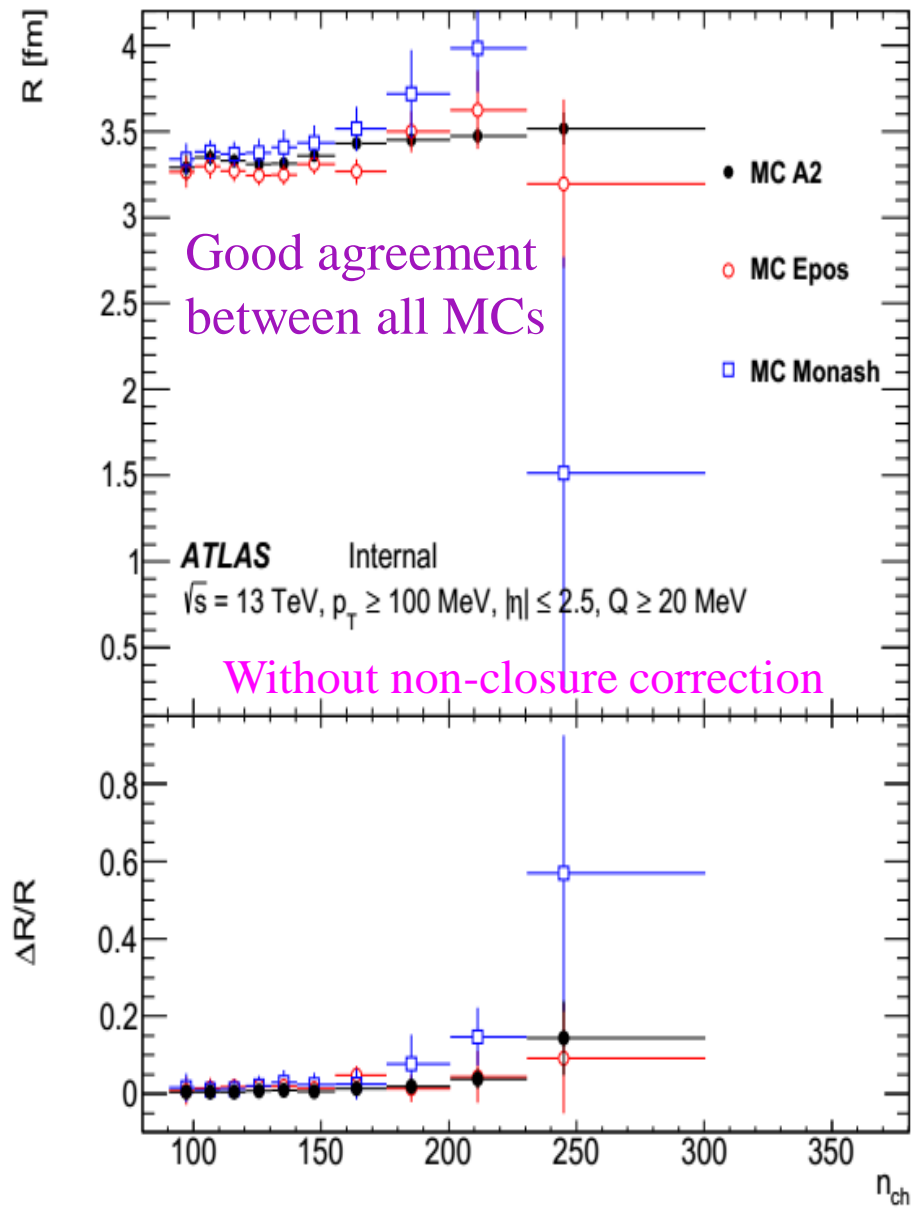
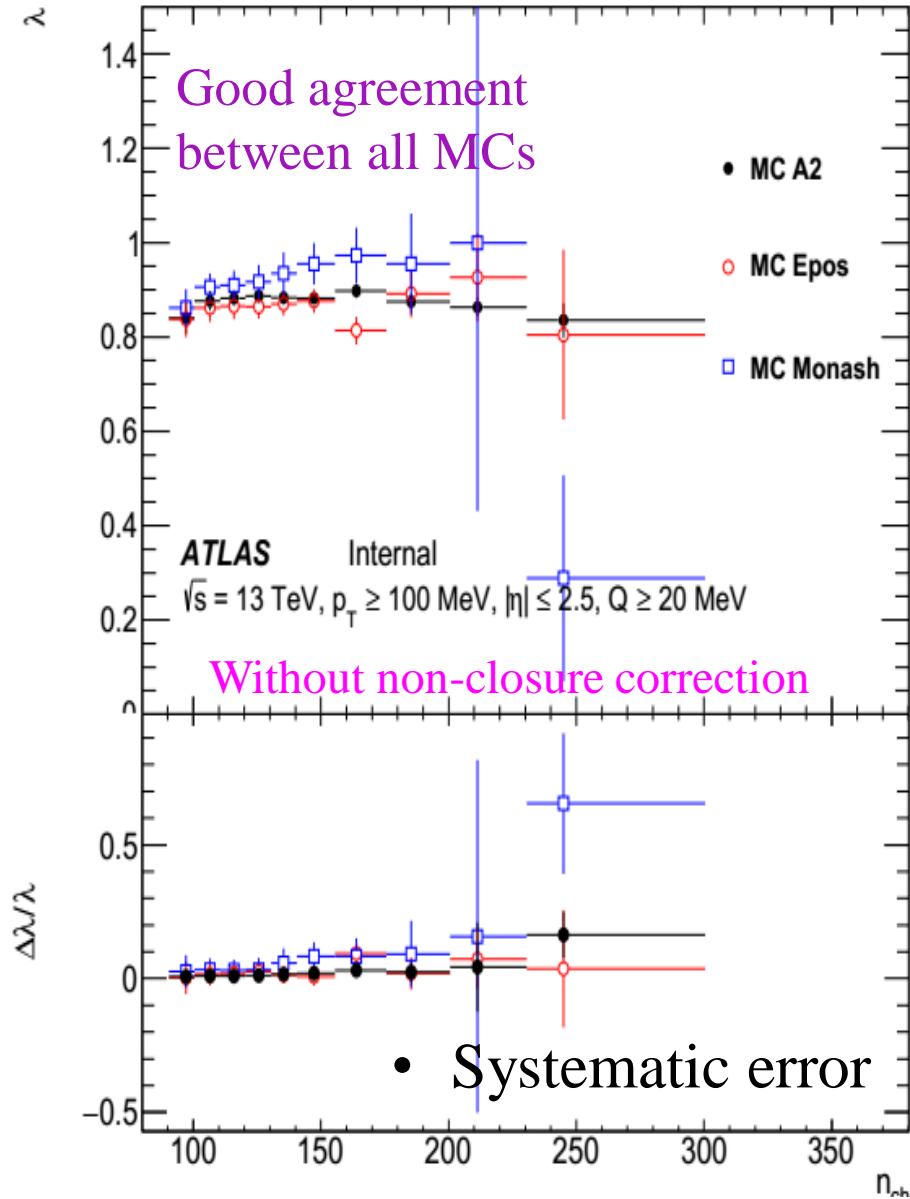
1. 0.2–0.3 GeV – not important after non-closure correction;
2. 0.4–0.9 GeV – important;
3. 1.0–1.16 GeV (only for  $2 \leq n_{ch} \leq 40$  and  $100 \leq k_T \leq 200$  MeV) – not important for BEC.

The excluded region at 7 TeV is 0.5–0.9 GeV.

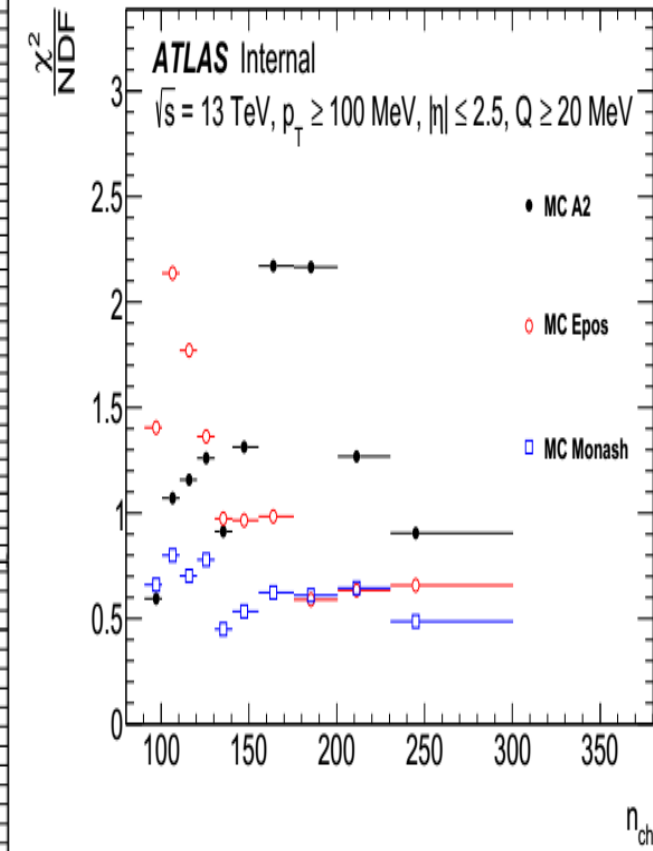


BEC parameters are in a good agreement in error bars but  $\chi^2/ndf$  more higher

# MULTIPLICITY DEPENDENCE OF BEC PARAMETERS FOR DIFFERENT MC



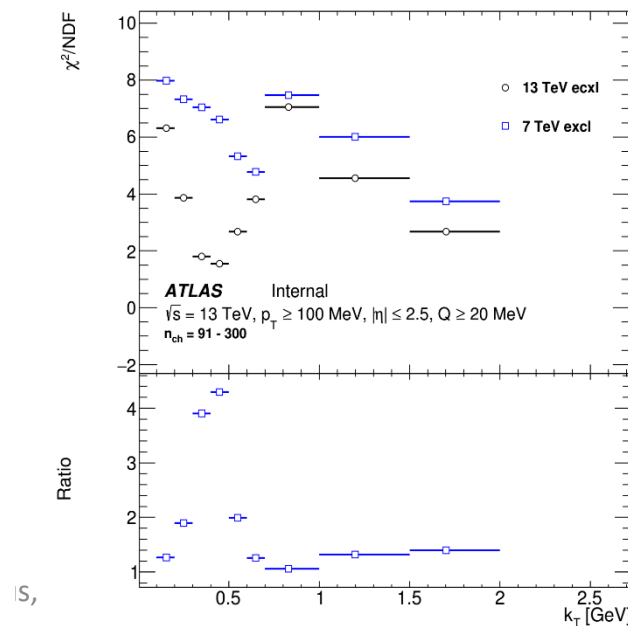
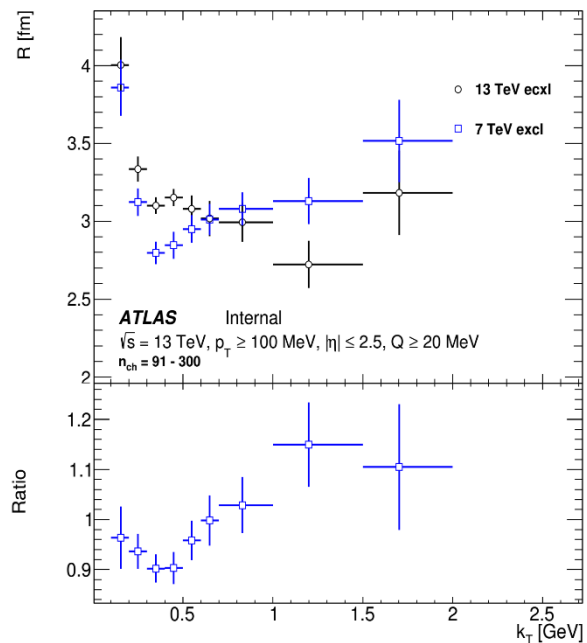
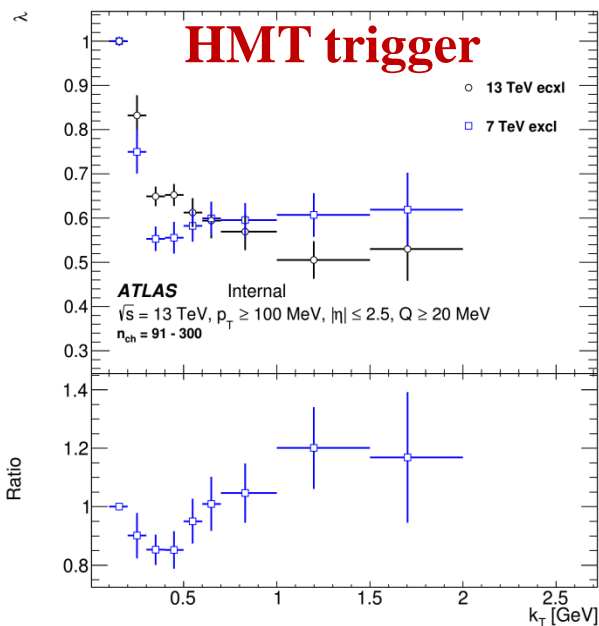
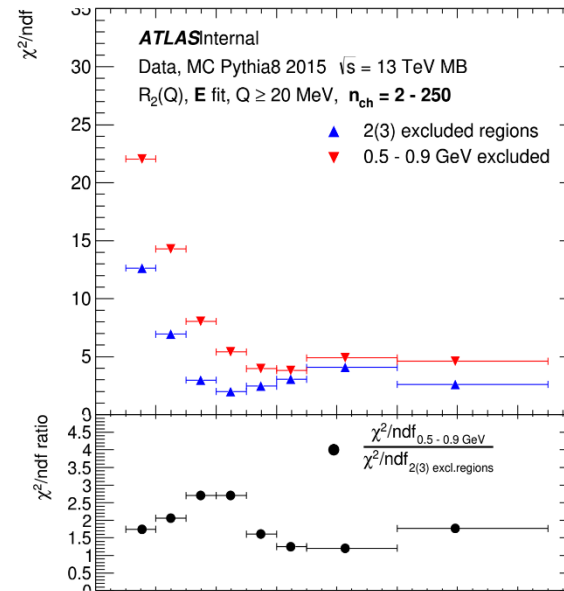
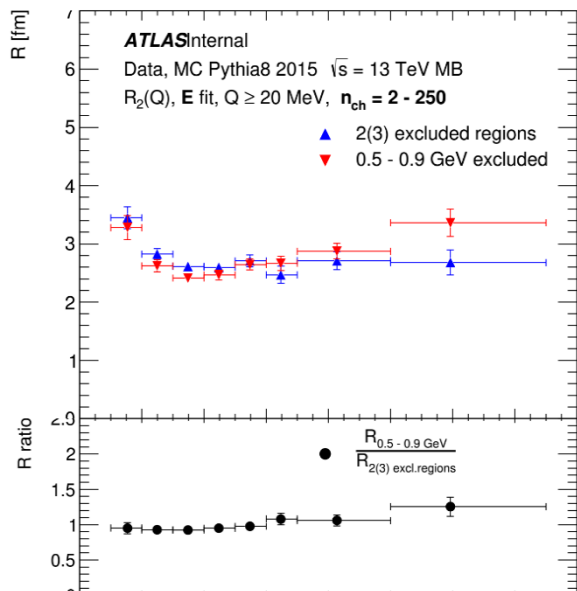
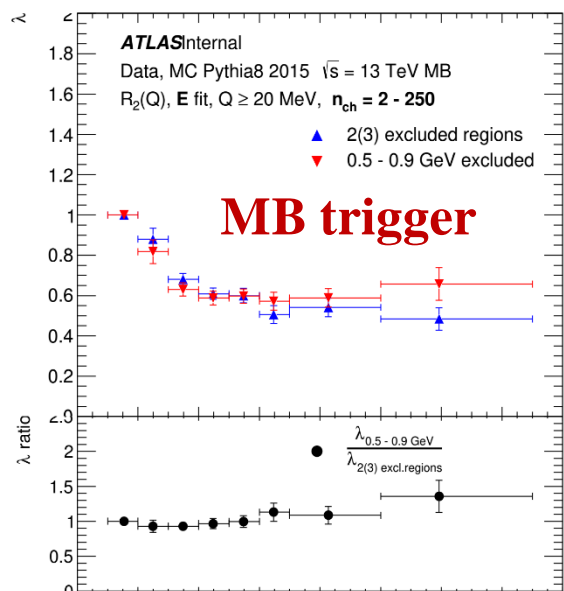
## HMT trigger



Radius saturation effect for all MCs for high multiplicity region



# COMPARISON OF $k_T$ RESULTS WITH ONE EXCLUDED REGION 0.5 – 0.9 GeV



Three bump regions because MC underestimated or overestimates:

1.  $\eta \rightarrow \pi^+ \pi^- \pi^0$  and  $\eta' \rightarrow \pi^+ \pi^- \gamma$ ;
2.  $\omega \rightarrow \pi^+ \pi^- \pi^0$  and  $\rho \rightarrow \pi^+ \pi^-$ ;
3.  $f_2 \rightarrow \pi^+ \pi^-$ ;

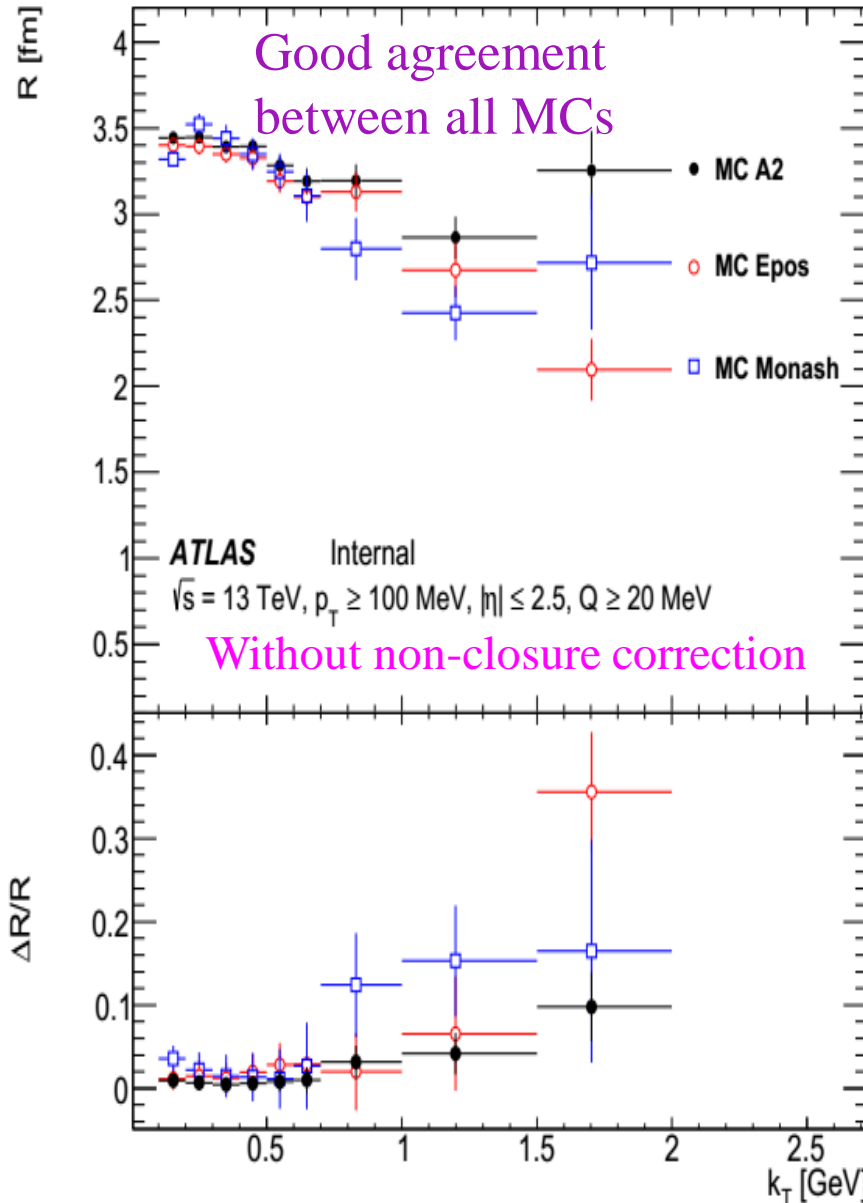
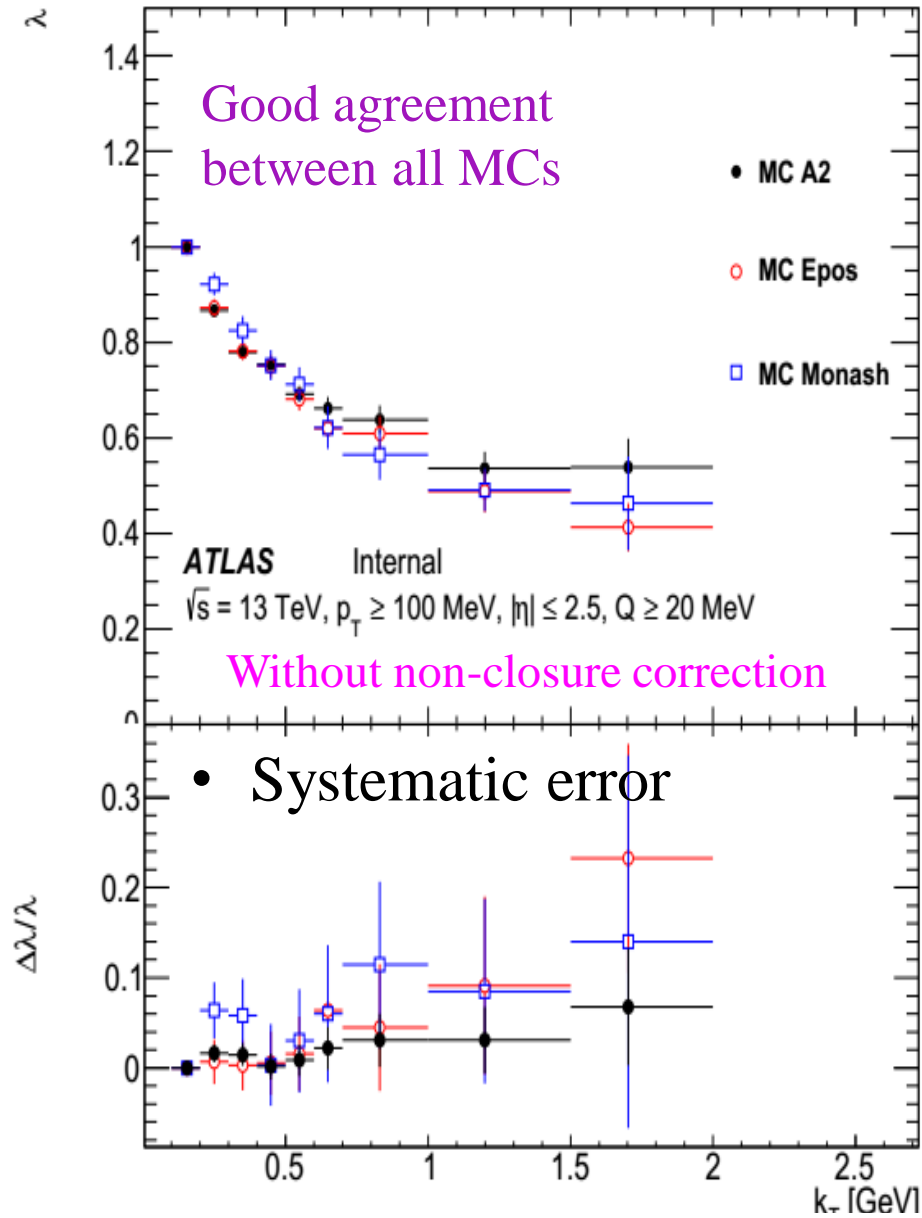
The excluded regions at 13 TeV are

1. 0.2–0.3 GeV – not important after non-closure correction;
2. 0.4–0.9 GeV – important;
3. 1.0–1.16 GeV (only for  $2 \leq n_{ch} \leq 40$  and  $100 \leq k_T \leq 200$  MeV) – not important for BEC.

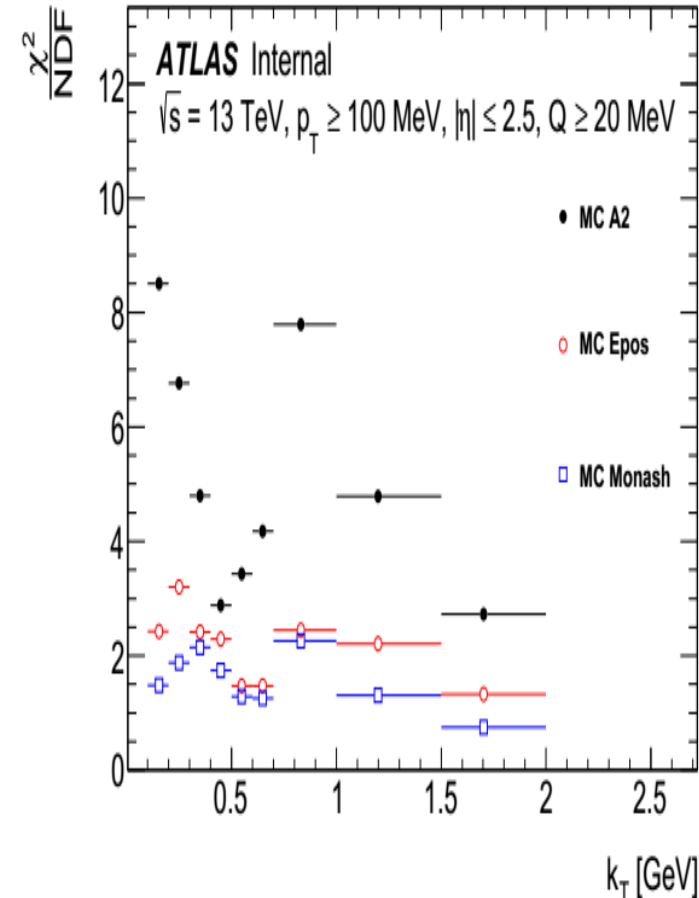
The excluded region at 7 TeV is 0.5–0.9 GeV.

BEC parameters are in a good agreement in error bars but  $\chi^2/\text{ndf}$  more higher

# K<sub>T</sub> DEPENDENCE OF BEC PARAMETERS FOR DIFFERENT MC



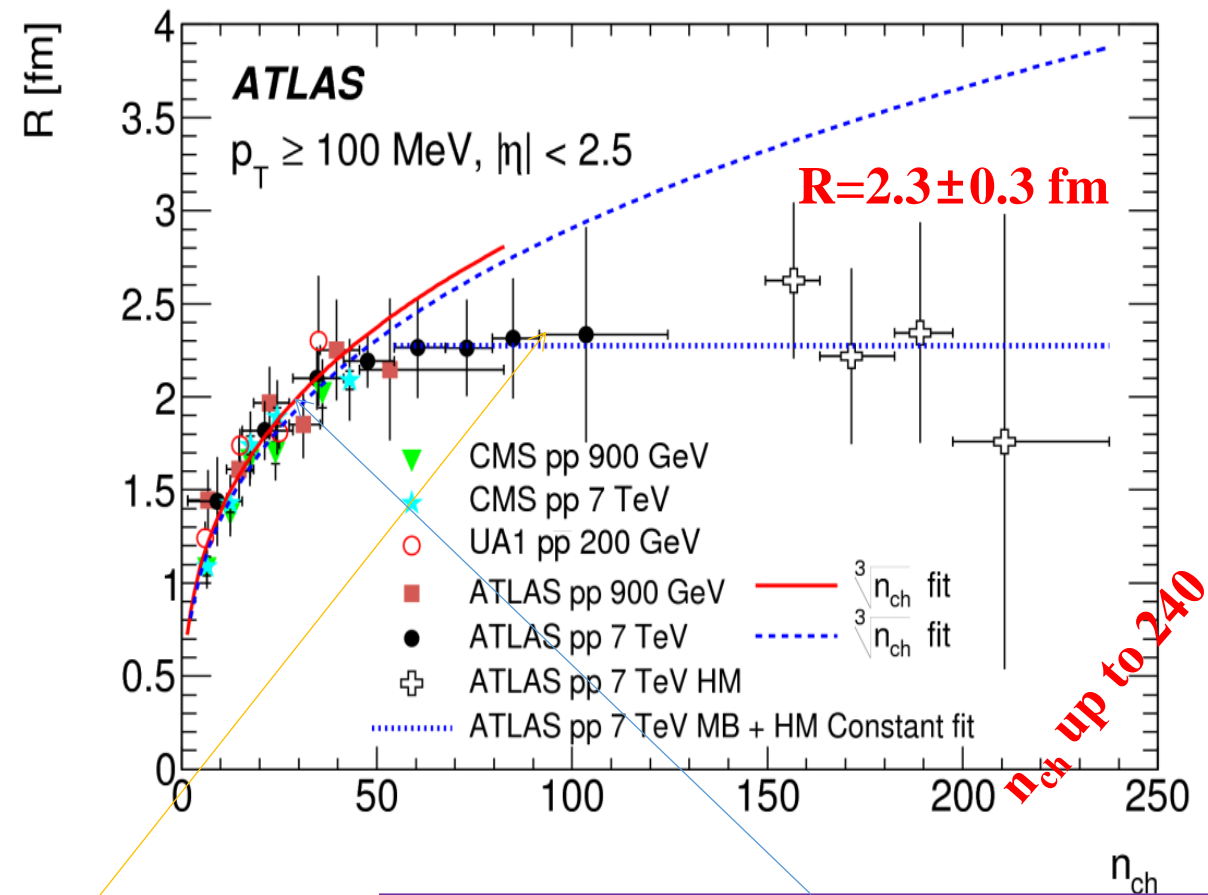
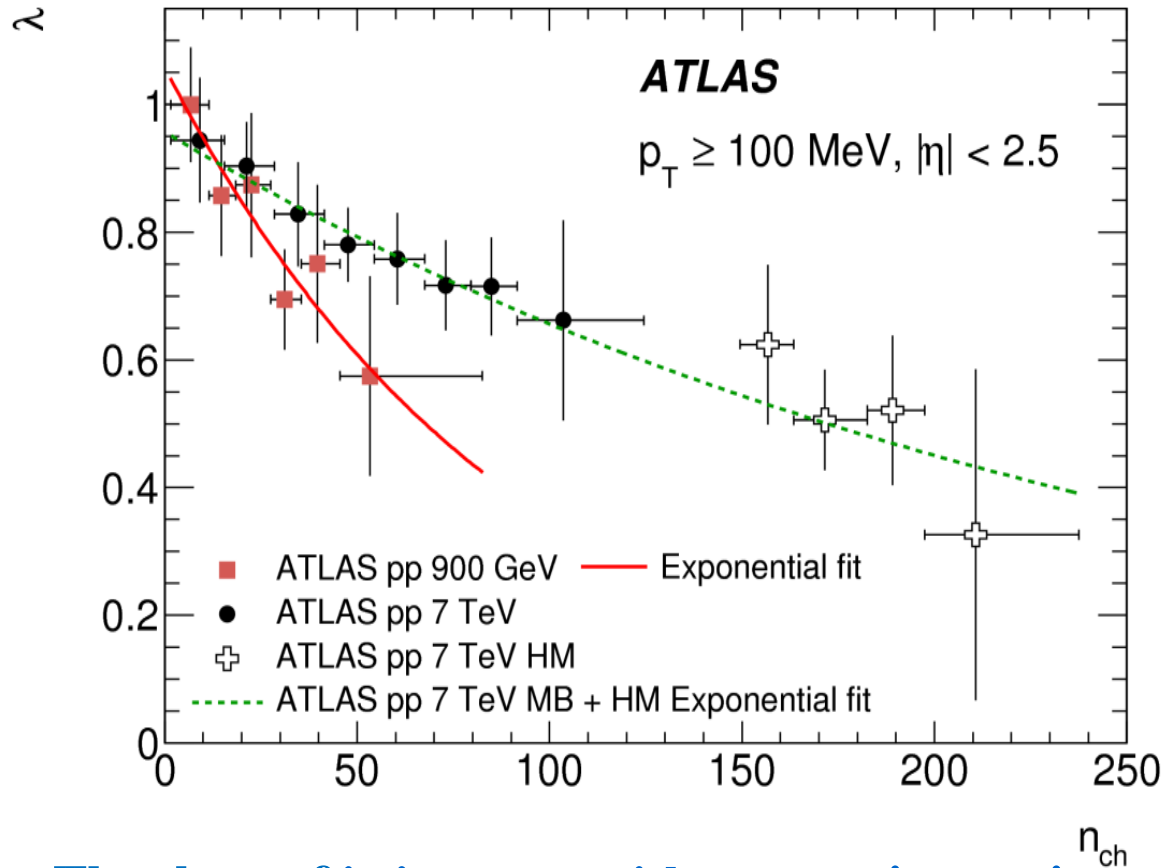
## HMT trigger



EPOS and Pythia8 Monash – more exponential distribution for BEC radius for high multiplicity region

# MULTIPLICITY DEPENDENCE OF $\lambda$ AND R BEC PARAMETERS AT 0.9, 7 TEV

EPJC 75 (2015) 466



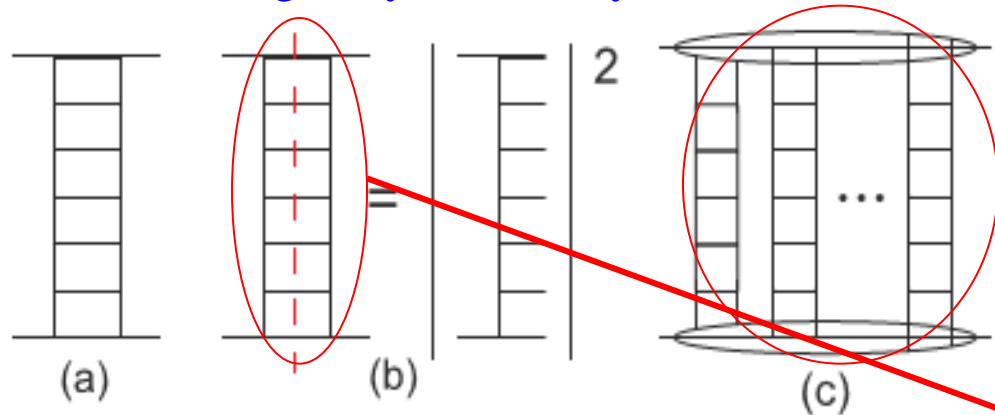
- ▶ The slope of  $\lambda$  increase with energy increasing
- ▶ The R are energy independent within the uncertainties
- ▶ The  $\lambda$  exponentially decrease with multiplicity

Good Agreement with  
CMS & UA1

- ▶ R of the  $\alpha \cdot n_{ch}^{1/3}$  fit for  $n_{ch} \leq 55$ : 0.9 TeV is  $\alpha = 0.64 \pm 0.07$  fm, 7 TeV is  $\alpha = 0.63 \pm 0.05$  fm
- ▶ R is a **Constant** for  $n_{ch} > 55$  at 7 TeV  $R = 2.28 \pm 0.32$  fm observed for the first time

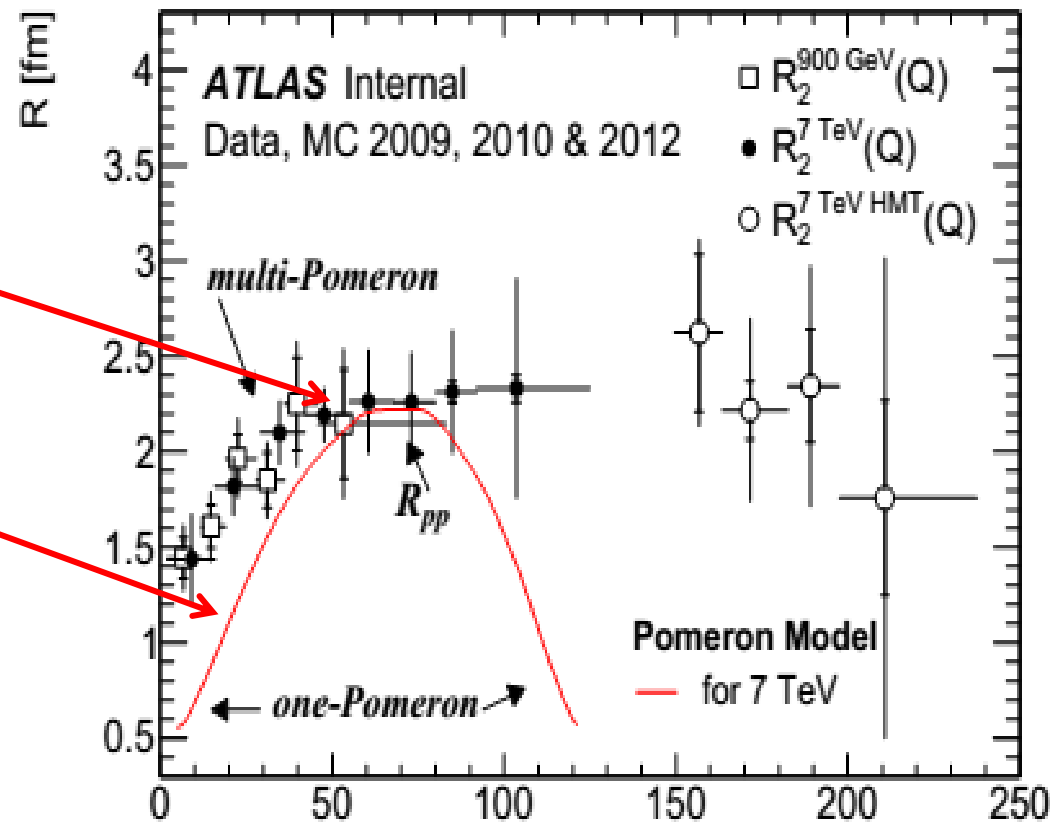
# THEORY PREDICTION FOR R PARAMETER OF BEC

V.A.Shegelsky, et al, Pomeron universality from identical pion correlations at the LHC, Phys.Letter B703 (2011) 288.  
 M.G.Ryskin, V.A.Shegelsky, Nucl.Phys B219 (2011) 10.



The ladder diagram for one-Pomeron exchange; (b) cutting one-Pomeron exchange leads to the multiperipheral chain of final state particles; (c) a multi-Pomeron exchange diagram.

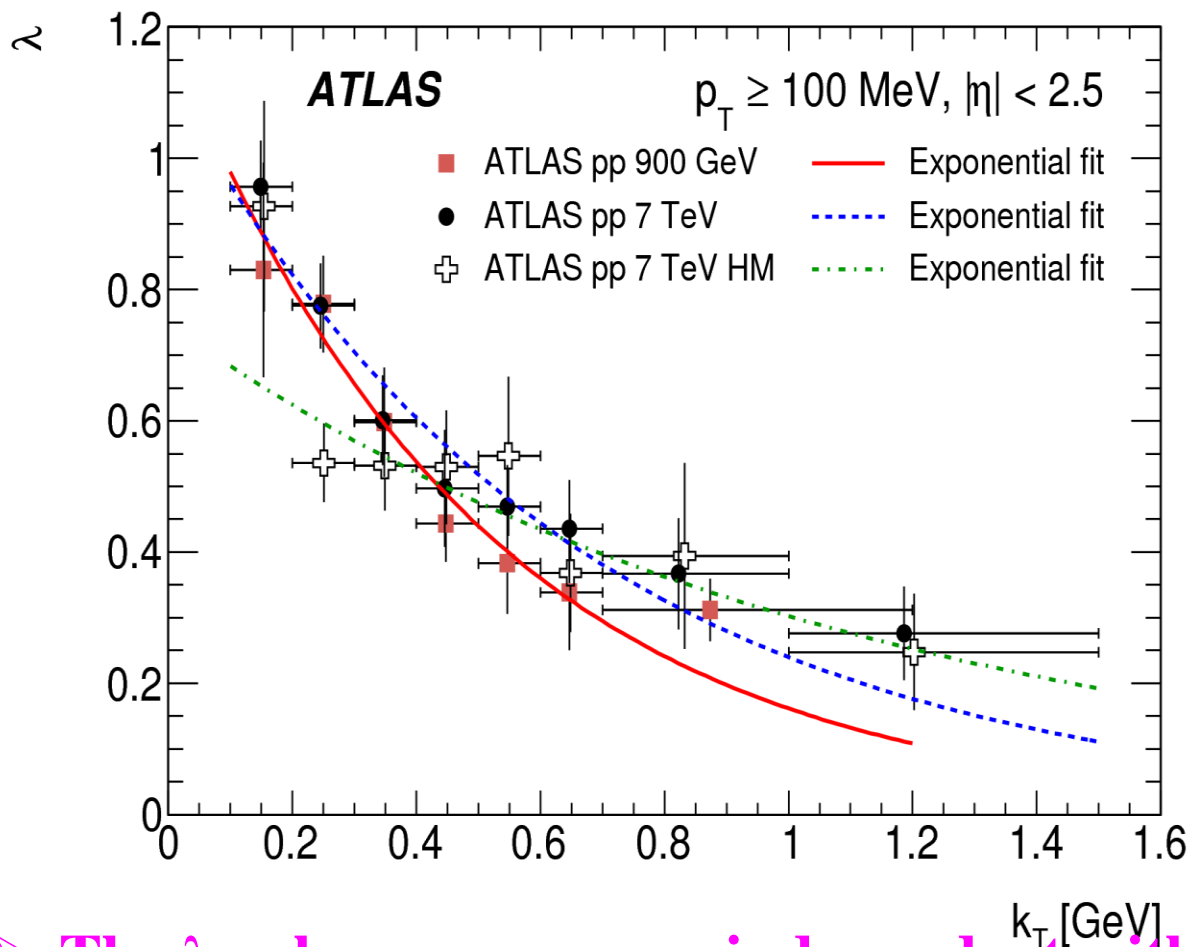
**Interpretation:** ► The BEC radius for one parton-parton interaction (underline events, cut Pomeron) is  $\sim 1$  fm, like for smallest multiplicity. ► For high multiplicity events we see BEC signal from some parton-parton interactions. ► The radius for high multiplicity can be interpret as an average distance between separate parton-parton interactions is  $\sim 2$  fm.



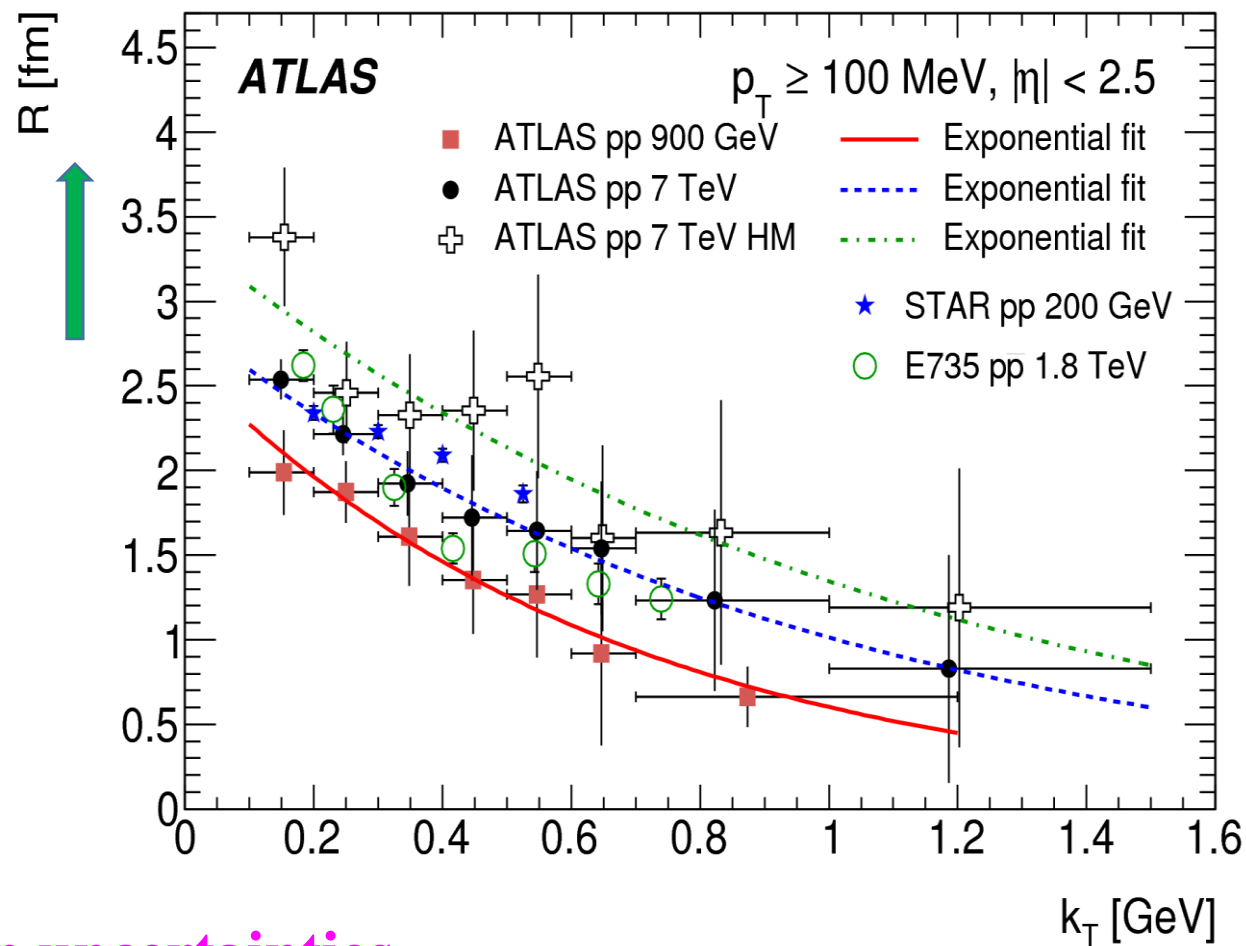
The prediction of Pomeron model  $R=2.2$  fm is in agreement with saturated radius  $R=2.3$  fm at 7 TeV for middle multiplicity region. There is not agreement with data for  $n_{ch} > 80$ .



$\lambda$  distribution at 7 TeV



$R$  distribution at 7 TeV



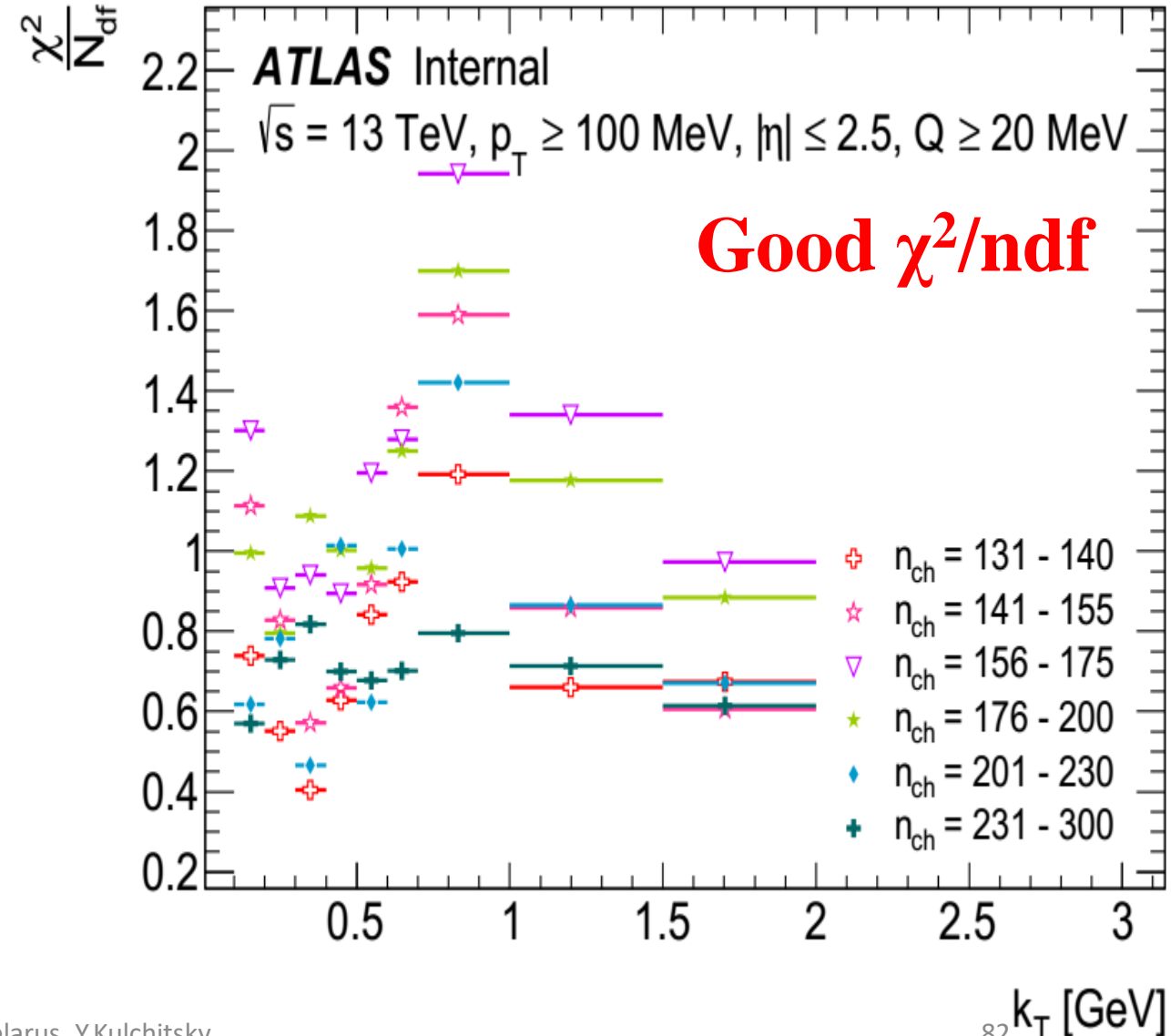
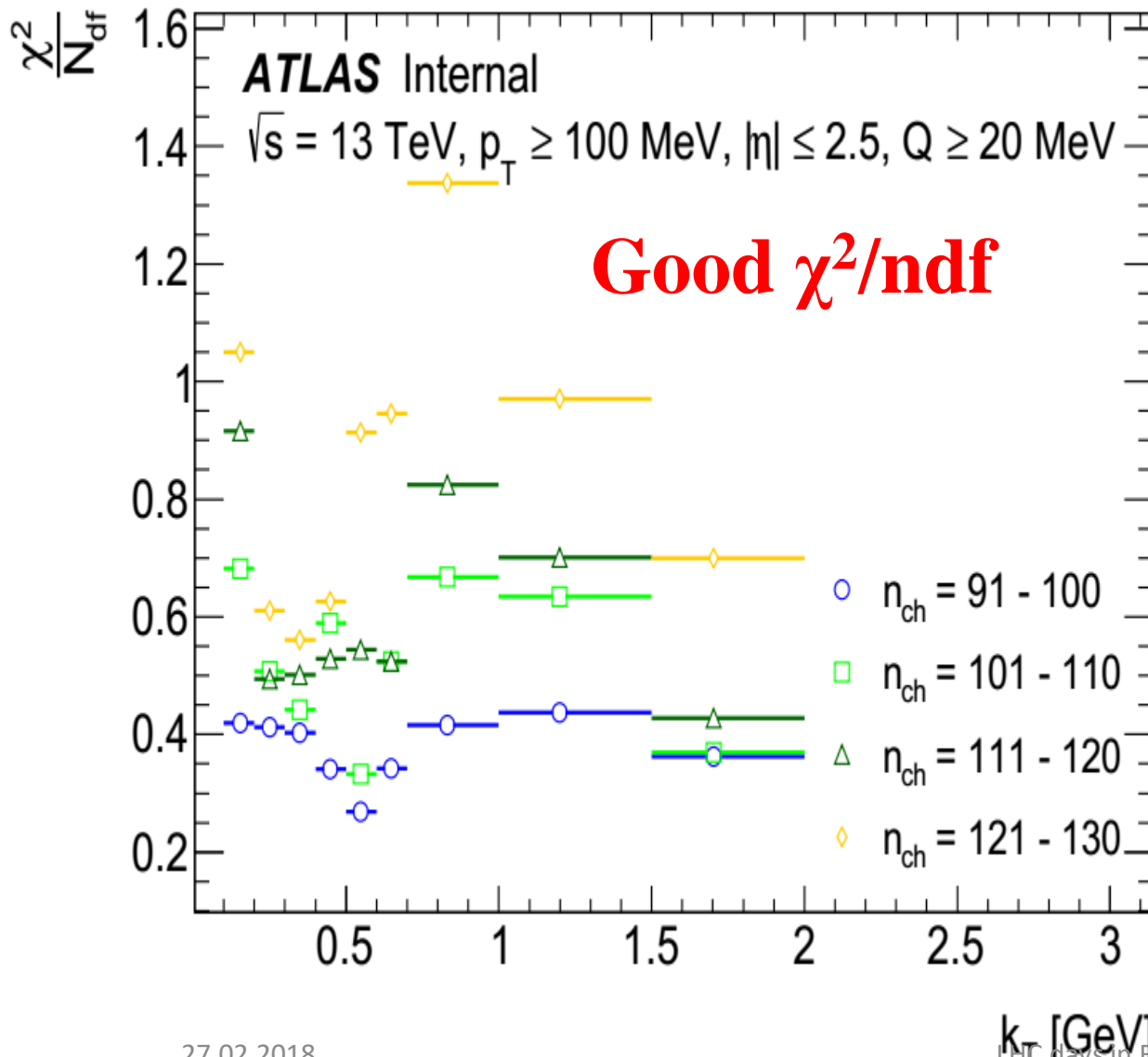
- The  $\lambda$  values are energy-independent within uncertainties
- The  $R$  values increase with increasing multiplicity
- The  $\lambda$  and  $R$  values decrease exponentially with  $k_T$

# K<sub>T</sub> DEPENDENCE OF $\chi^2/\text{NDF}$ FOR (N<sub>CH</sub>, K<sub>T</sub>) INTERVALS

$\chi^2/\text{ndf}$  distribution at 13 TeV

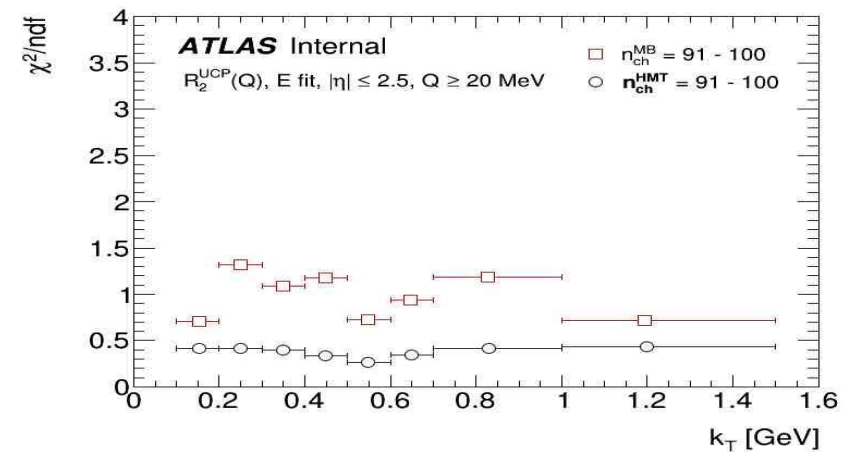
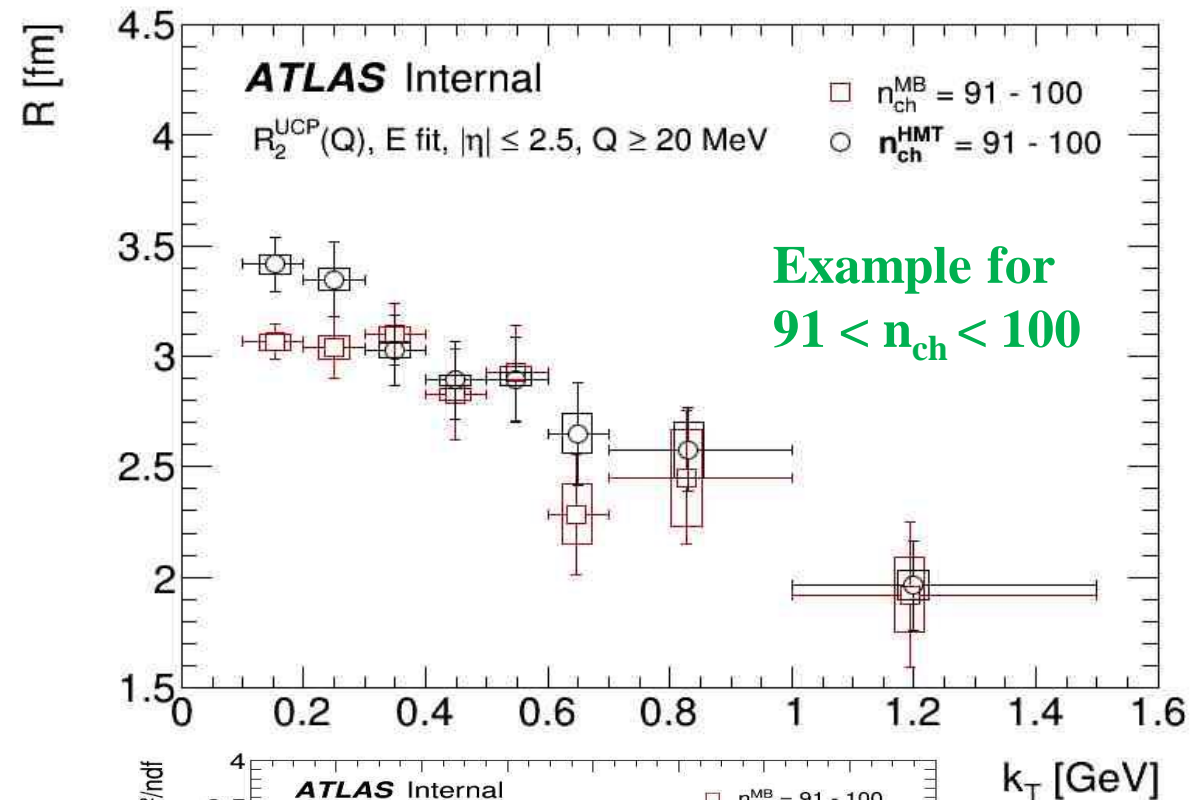
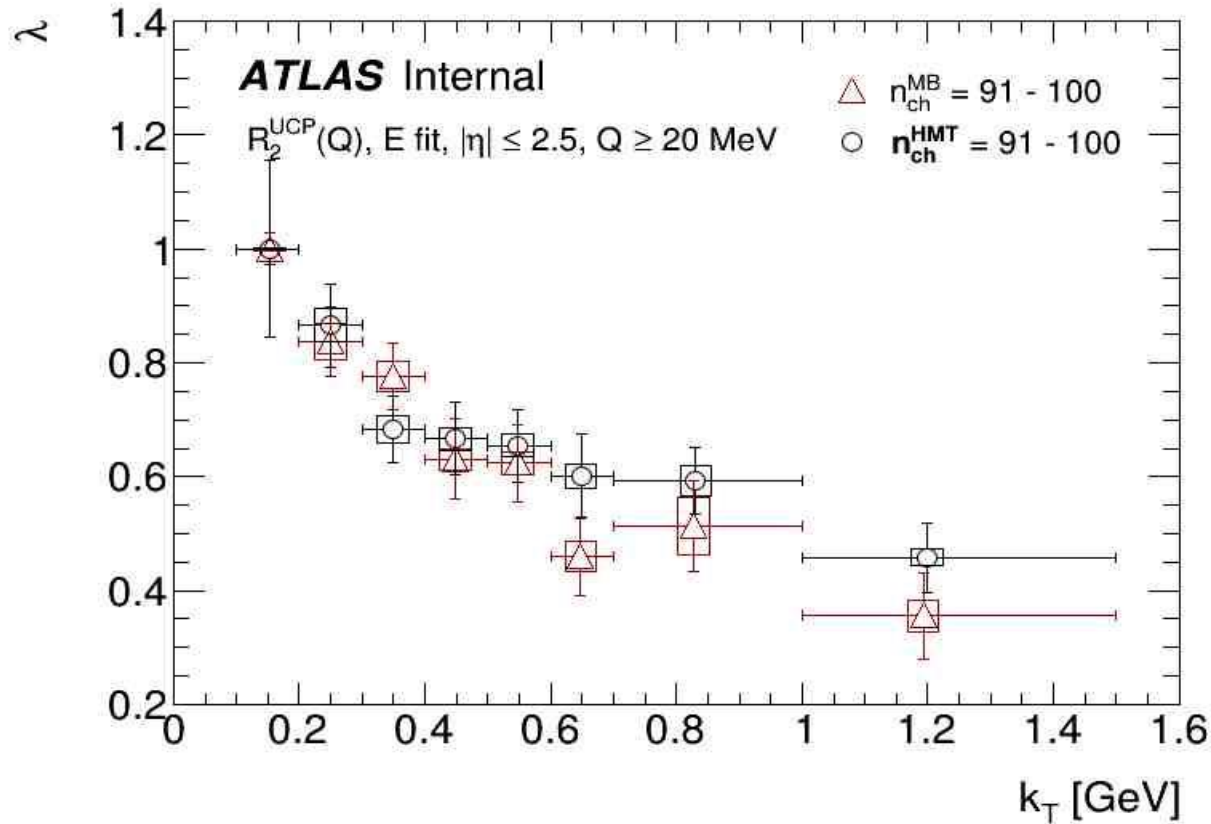
HMT trigger

$\chi^2/\text{ndf}$  distribution at 13 TeV



# $K_T$ DEPENDENCE OF BEC PARAMETERS

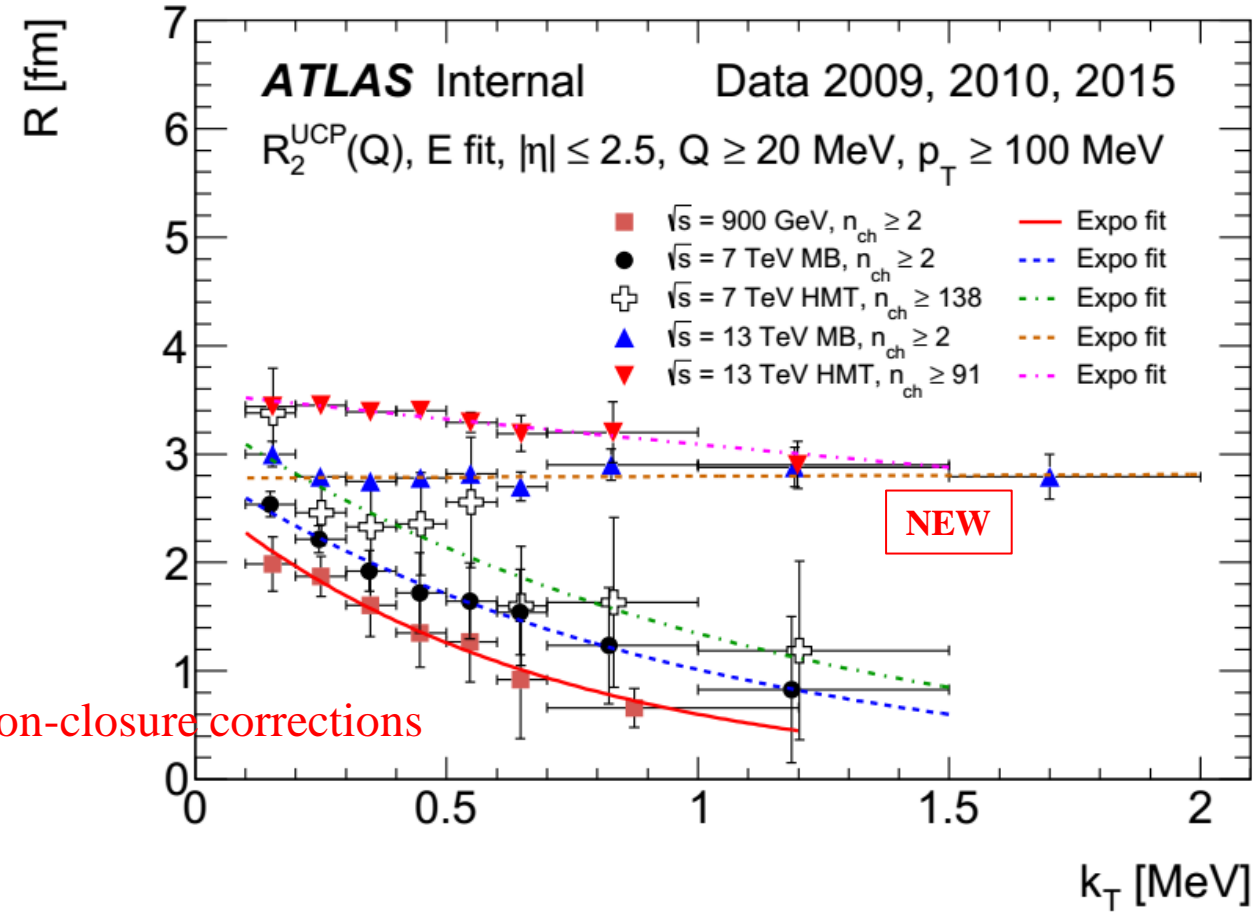
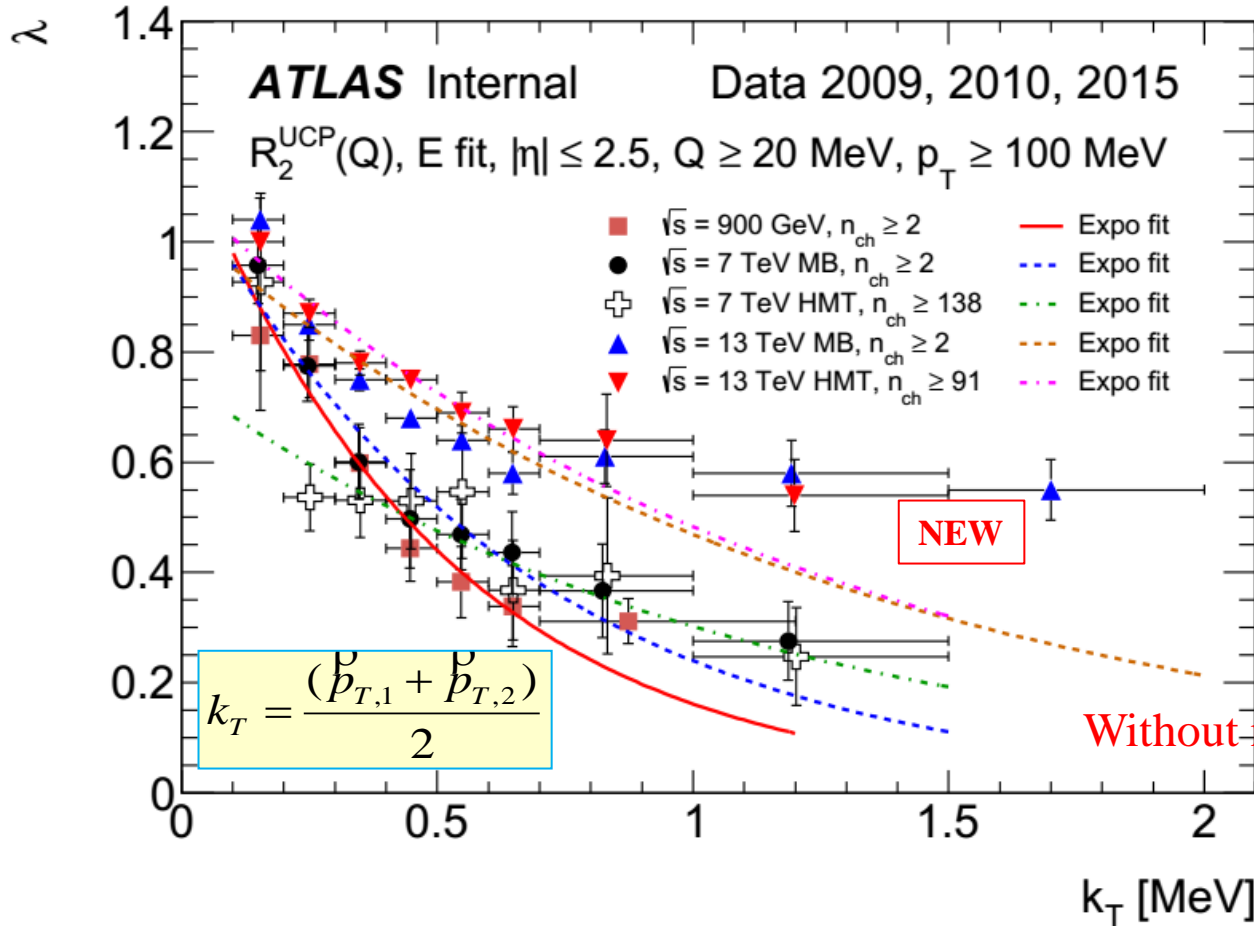
$\lambda$  distribution at 13 TeV MB + HMT triggers R distribution at 13 TeV



- Agreement between MB and HMT results
- The  $\lambda$  values decreases exponentially in dependence of  $k_T$
- The  $R$  values decreases exponentially in dependence of  $k_T$

# $K_T$ DEPENDENCE OF BEC PARAMETERS AT 0.9 – 13 TEV

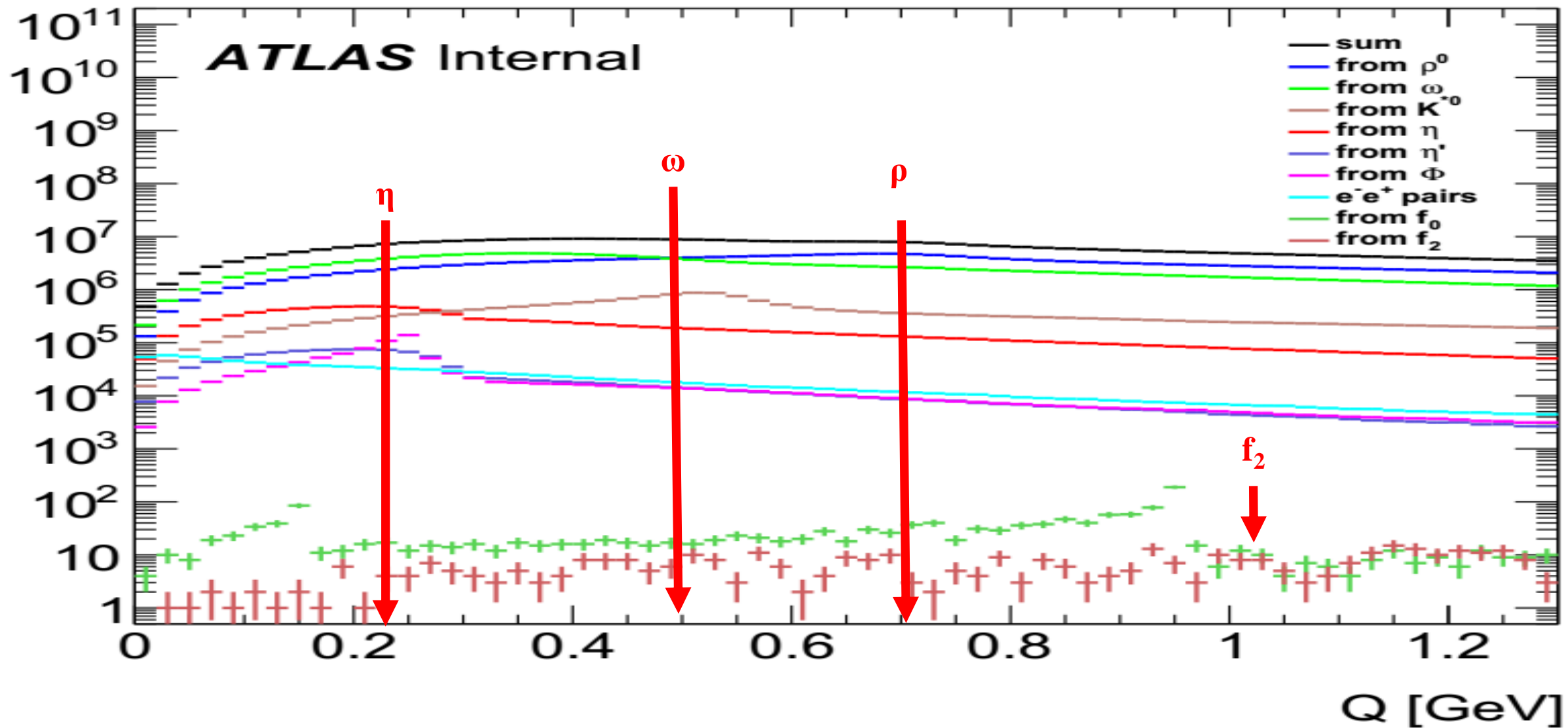
EPJC 75 (2015) 10, 466; ATL-COM-PHYS-2016-1621



- The  $\lambda$  values are (trigger) multiplicity-independent within uncertainties at 13 TeV
- The R values increase with increasing multiplicity region (trigger)
- The  $\lambda$  values decrease exponentially with  $k_T$
- The R values decrease exponentially with  $k_T$  for HMT events



# RESONANCES STUDY AT 13 TEV



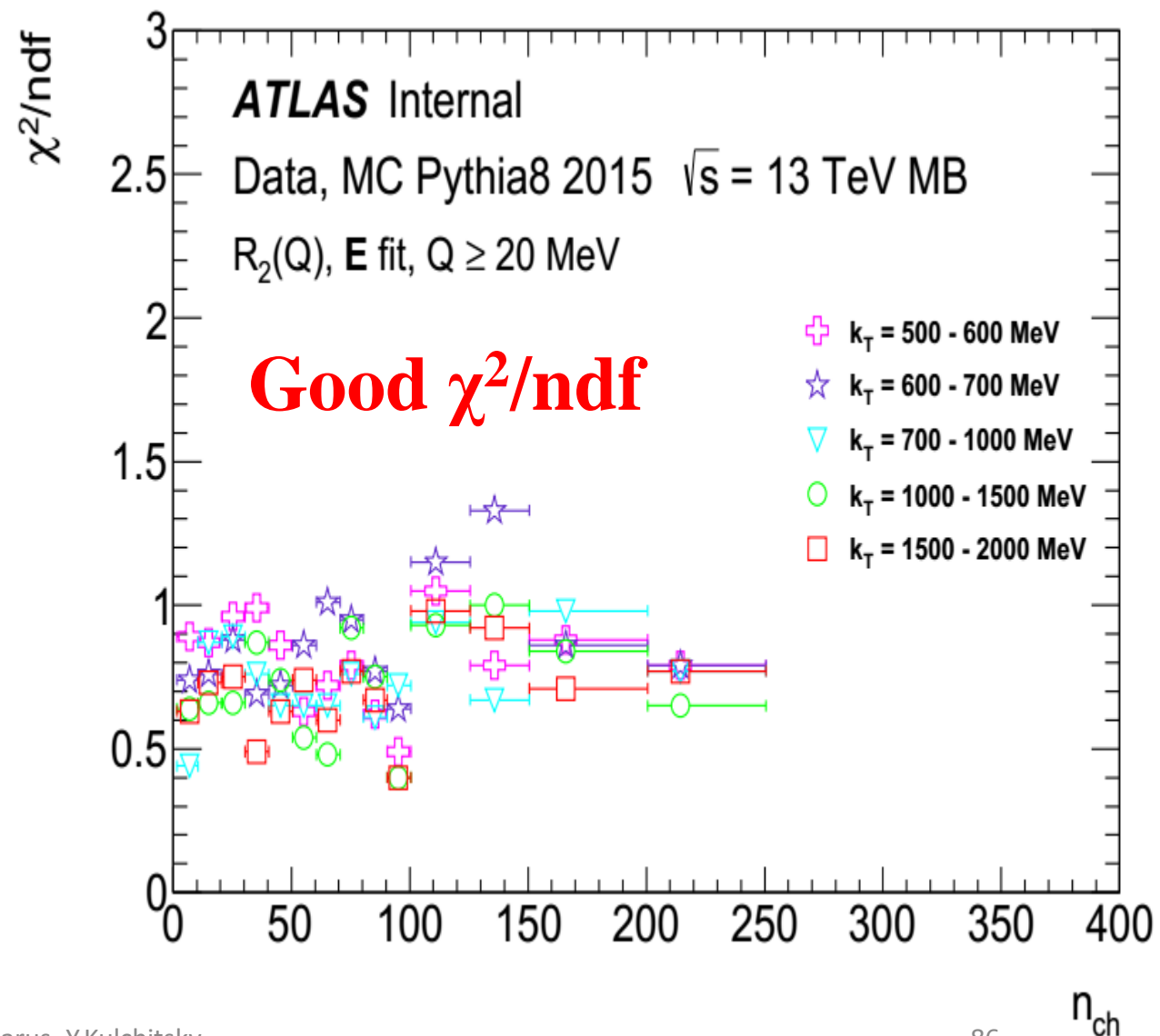
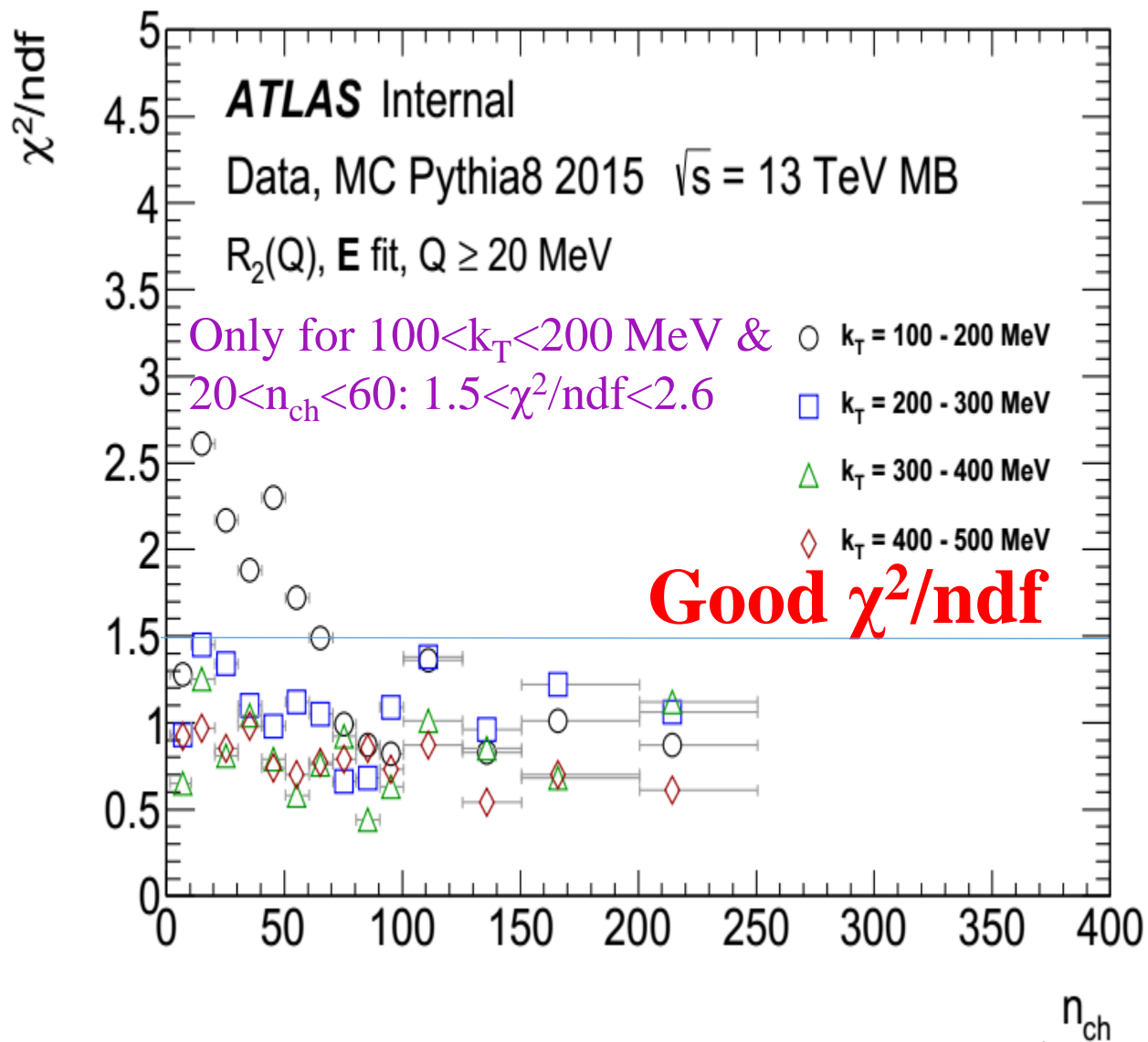
The Q spectrum generated by Pythia8 and the decomposition of its resonant part into leading contributions.

# MULTIPLICITY DEPENDENCE OF $\chi^2/\text{NDF}$ FOR $(N_{\text{CH}}, K_T)$ INTERVALS

$\chi^2/\text{ndf}$  distribution at 13 TeV

**MB trigger**

$\chi^2/\text{ndf}$  distribution at 13 TeV

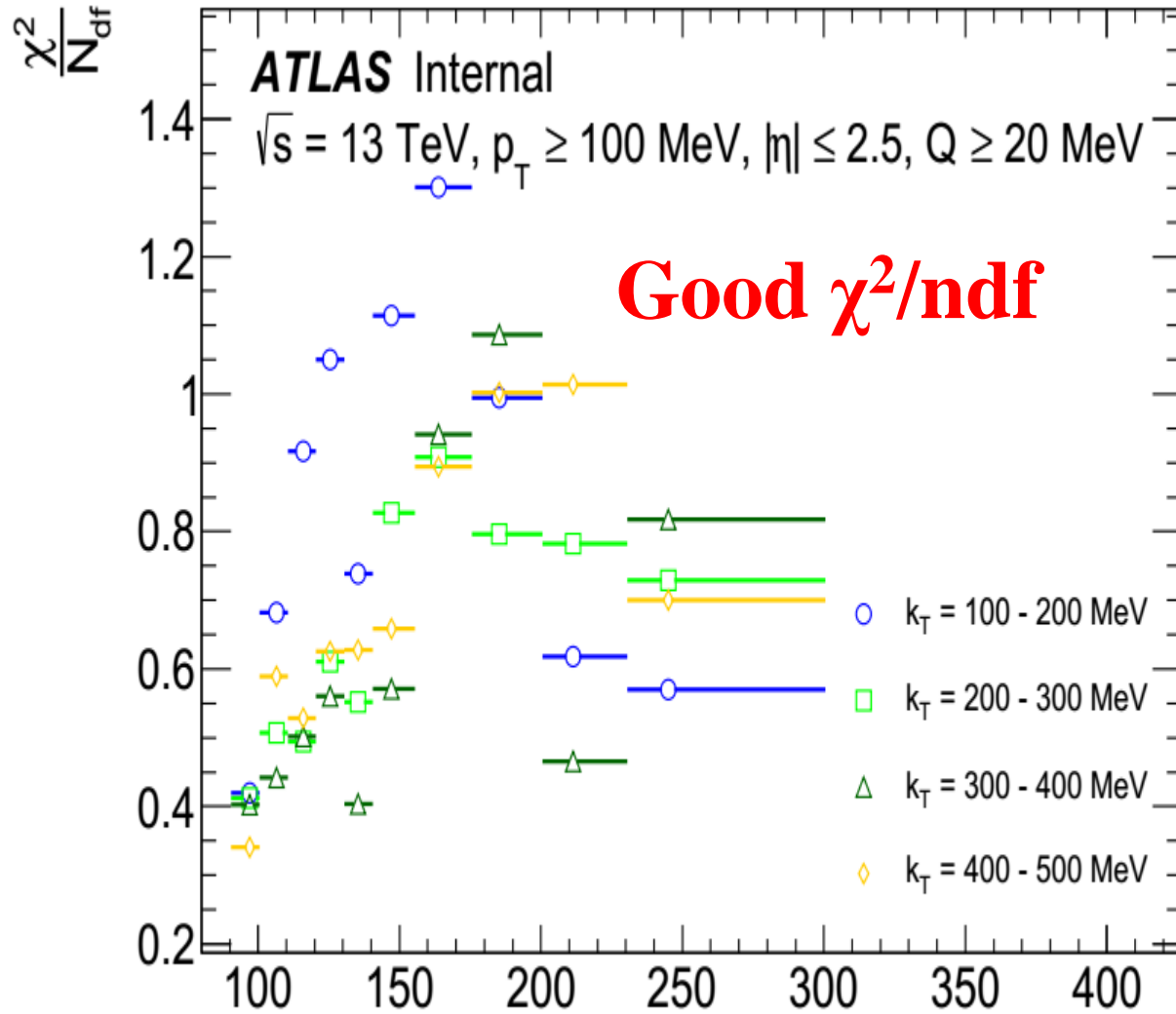


# MULTIPLICITY DEPENDENCE OF $\chi^2/\text{NDF}$ FOR $(N_{\text{CH}}, K_T)$ INTERVALS

$\chi^2/\text{ndf}$  distribution at 13 TeV

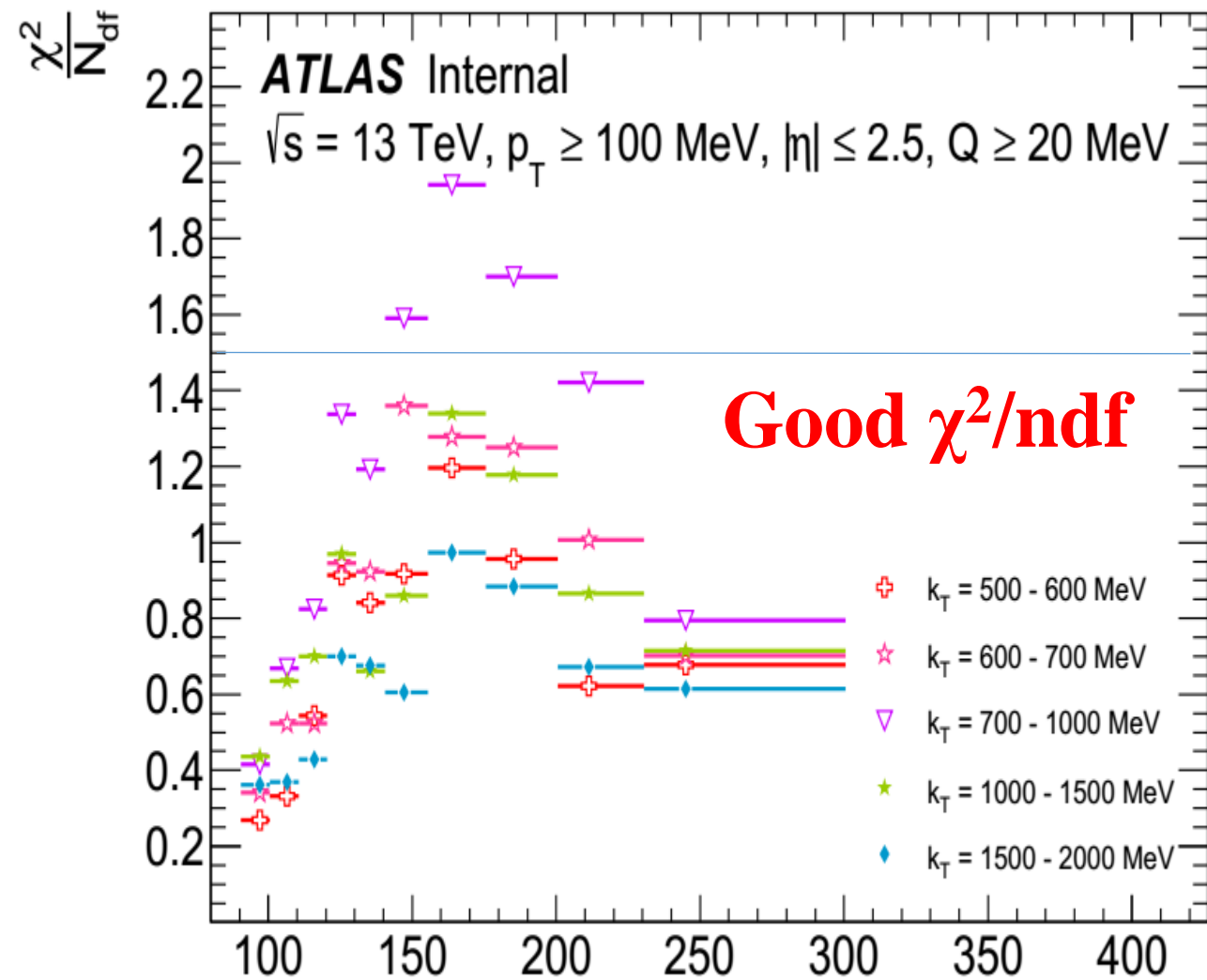
**HMT trigger**

$\chi^2/\text{ndf}$  distribution at 13 TeV



$n_{\text{ch}}$

Iarus, Y.Kulchitsky



$n_{\text{ch}}$