

Extracting (un)polarized cross sections for pion photoproduction off the neutron from deuteron data

arXiv:1804.04757

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Introduction

Meson productions induced by photon and electron beams

Interests

- Information on N^* properties (mass, width, missing resonances, etc.)
complementary to πN scattering data ← “complete” measurements
- Electromagnetic $N \rightarrow N^*$ transition form factors (Q^2 –dependence)
→ important information for hadron structures

Measurements of neutron target data (incl. polarization obs.) are recently active

$\gamma n \rightarrow \pi^0 n$ $d\sigma/d\Omega$ A2@MAMI PRL 112, 142001 (2014)

E A2@MAMI PLB 770, 523 (2017)

$\gamma n \rightarrow \pi^- p$ $d\sigma/d\Omega$ CLAS@JLab PRC 96, 035204 (2017)

E CLAS@JLab PRL 118, 242002 (2017)

$\gamma n \rightarrow \eta n$ $d\sigma/d\Omega$ A2@MAMI PRC 90, 015205 (2014)

E A2@MAMI PRC 95, 055201 (2017)

More data are available, and more are coming soon

Why neutron target data ?

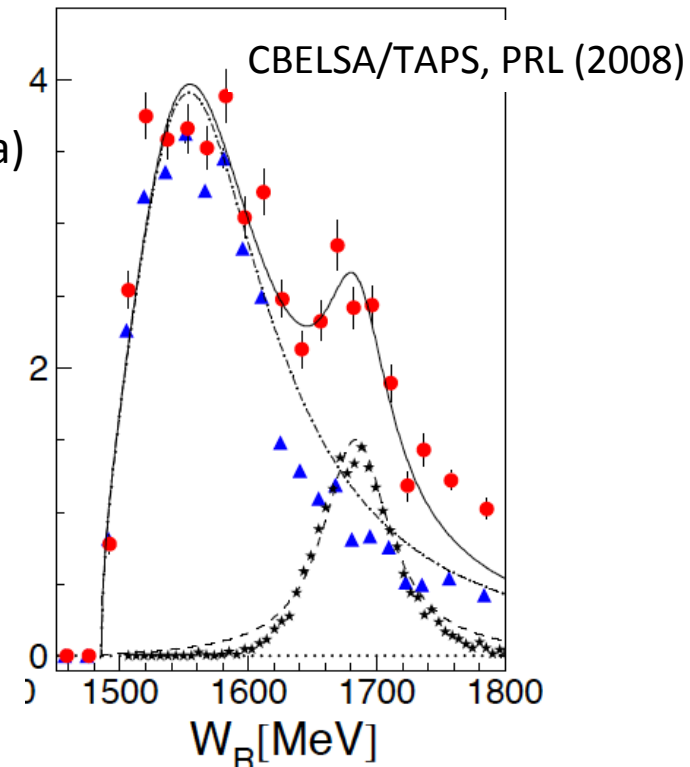
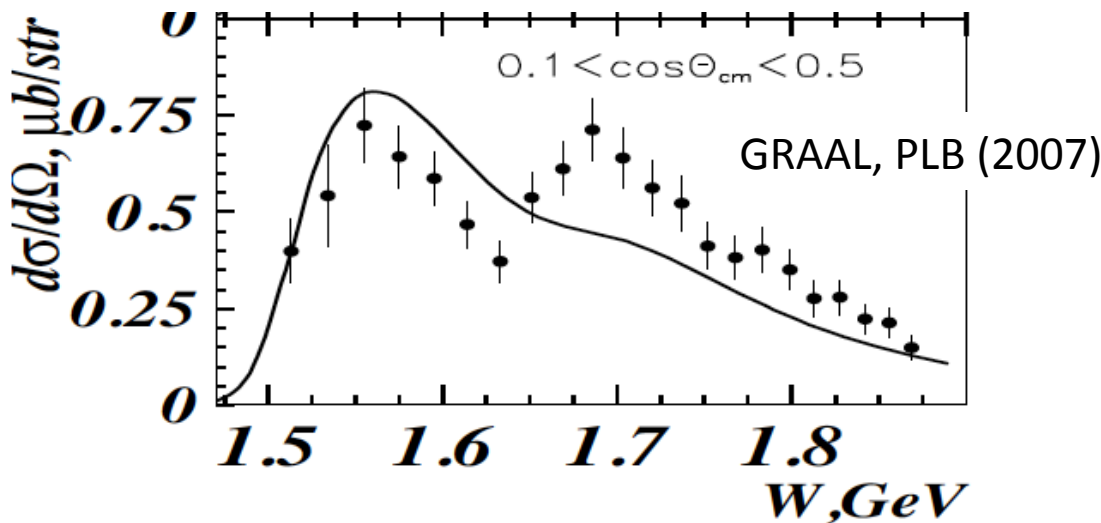
First motivation

$\gamma^{(*)} p \rightarrow N^*$ AND $\gamma^{(*)} n \rightarrow N^*$ form factors \rightarrow Isospin structure of $\gamma^{(*)} N \rightarrow N^*$ form factors

- Interesting quantities for understanding hadron structures
- Necessary for application to neutrino-induced reactions
 \rightarrow needed for analyzing data from neutrino-oscillation experiments SXN et al. PRD (2015)

Unexpected surprises

Narrow bump at $W \sim 1.68$ GeV in $\gamma n \rightarrow \eta n$ (not in πN data)

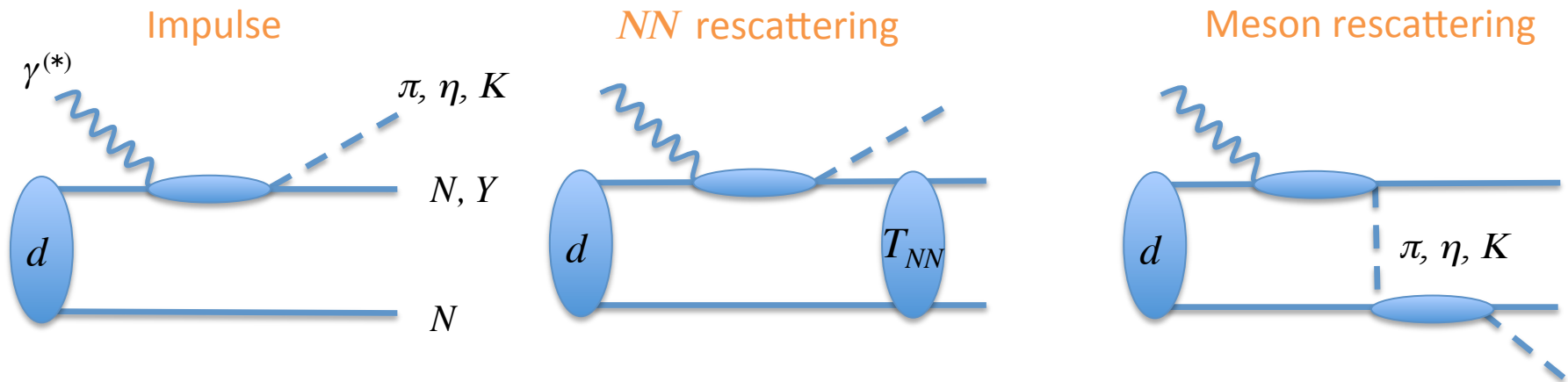


Photon and electron-induced meson productions off deuteron

To obtain $\gamma^{(*)} 'n' \rightarrow \pi N, \eta N, K Y$ data, **Deuteron** is the primary target; we need to understand:

How to extract $\gamma^{(*)} 'n' \rightarrow MB$ cross sections from $\gamma^{(*)} d \rightarrow MBN$ data

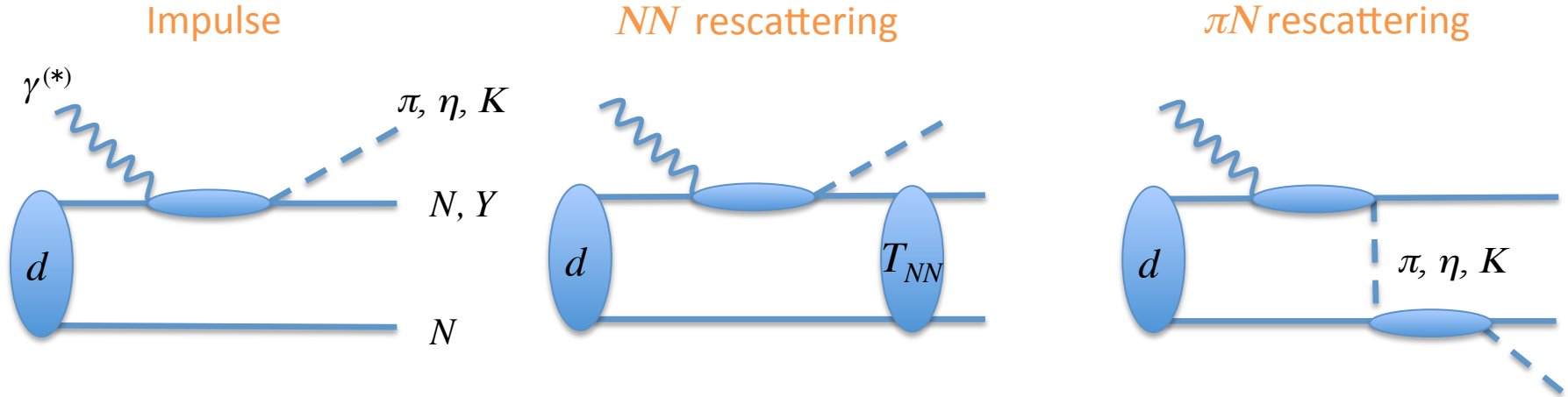
To address this question theoretically, we first need understand meson productions on deuteron



Features

- Photo- and electro-excitations of bound nucleons
- initial nucleons are in Fermi motion
- Final state interactions (**FSI**)

Photon and electron-induced meson productions off deuteron



Q: How to extract $\sigma(\gamma^{(*)} n' \rightarrow MB)$ from $\sigma(\gamma^{(*)} d \rightarrow MBN)$?

- Common (and practical) procedure

Kinematical cuts \rightarrow quasi-free $\gamma^{(*)} n' \rightarrow MB$ events selected

Concerns: FSI and/or cuts could distort $\sigma(\gamma^{(*)} n' \rightarrow MB)$ from true one ?

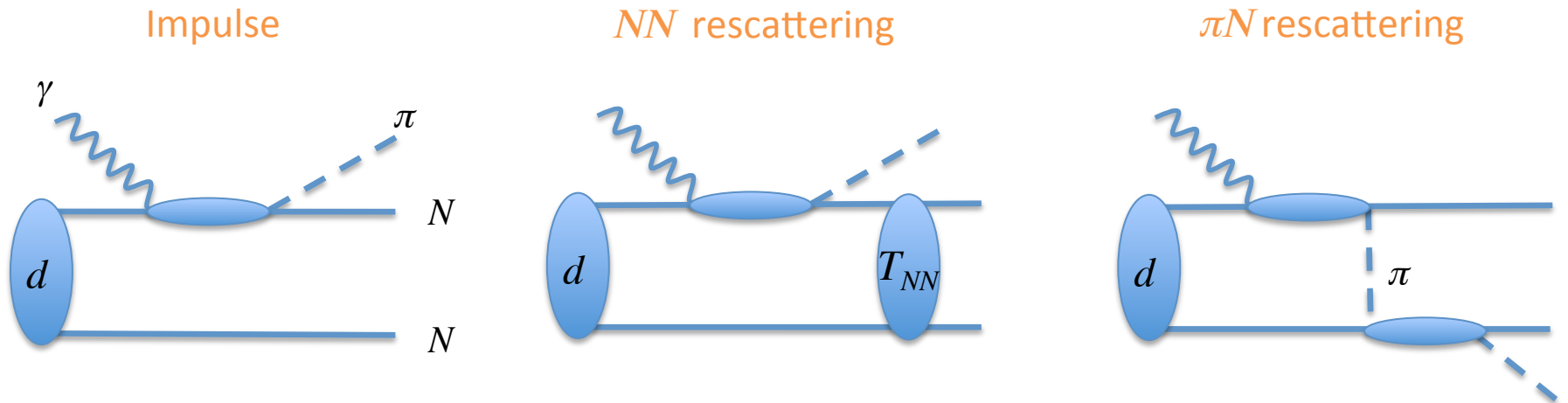
This work

- A dynamical model for $\gamma d \rightarrow \pi NN$ with FSI is developed
- $d\sigma/d\Omega_\pi$ and Σ, E, G for $\gamma d \rightarrow \pi NN$ are calculated; FSI effects examined
- $d\sigma/d\Omega_\pi$ and Σ, E, G for $\gamma 'n' \rightarrow \pi^0 n, \pi^- p$ are extracted from $\gamma d \rightarrow \pi NN$
- We address how FSI and cuts could distort $\gamma 'n'$ observables

$\gamma d \rightarrow \pi NN$ reaction model based on
dynamical coupled-channels model

Model for $\gamma d \rightarrow \pi N N$

Multiple scattering theory truncated at the first-order rescattering

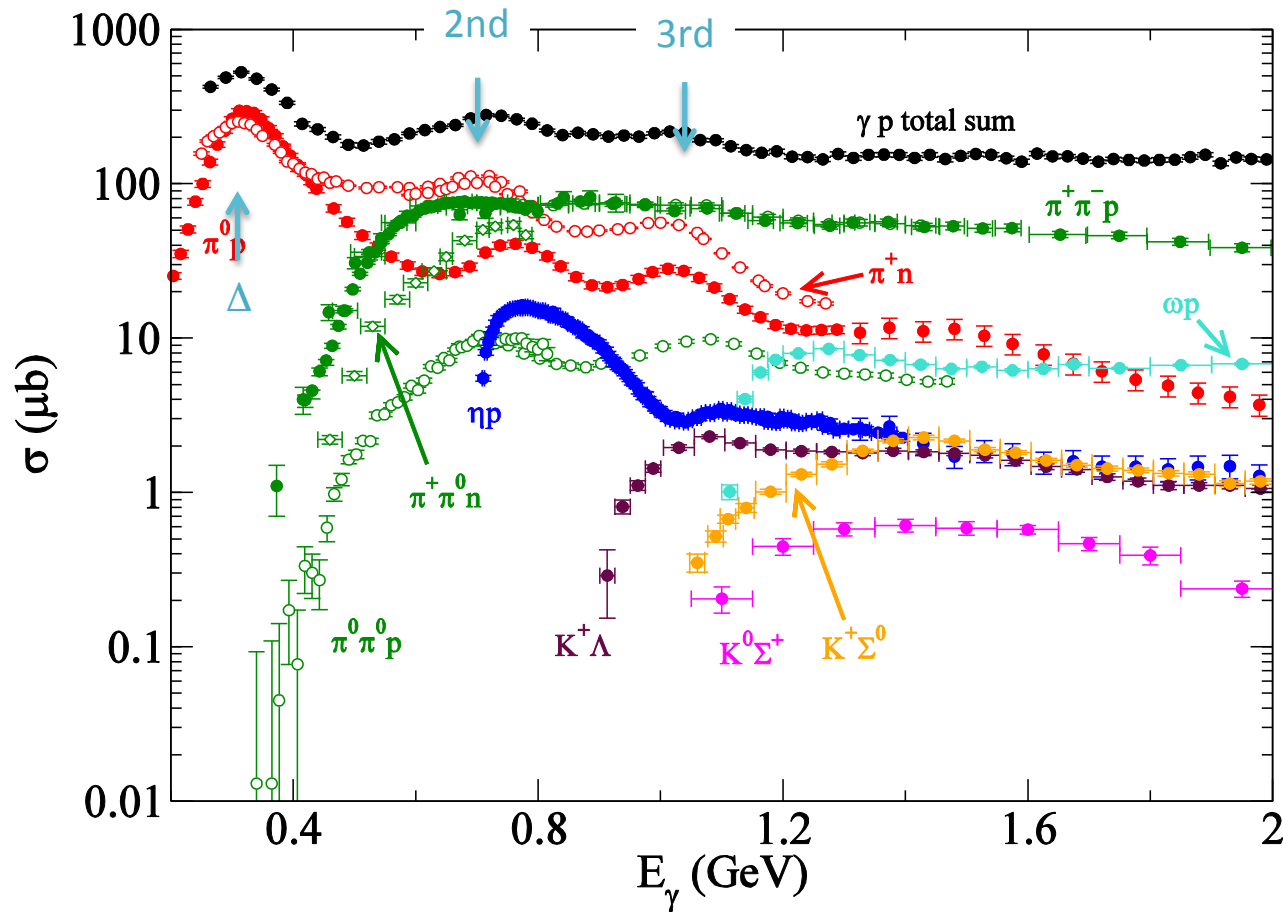


Elementary amplitudes

- $\gamma N \rightarrow \pi N, \pi N \rightarrow \pi N$ amplitude \leftarrow DCC model (Kamano et al., PRC94 (2016))
- T_{NN} , deuteron w.f. \leftarrow CD-Bonn potential (Machleidt et al., PRC 63 (2001))

3-dim. loop integral with off-shell amplitudes are numerically evaluated

Dynamical coupled-channels model for meson productions in resonance region



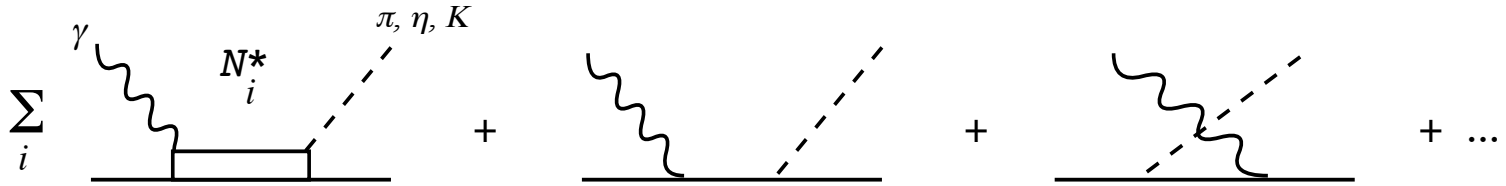
Data for $\gamma p \rightarrow X$

Need develop a model to describe these reactions

- Several **nucleon resonances** form characteristic peaks
- 2π production is comparable to 1π
- η, K productions (**multi-channel couplings are important physics**)

Dynamical coupled-channels model for resonance region

Resonance excitation + non-resonant meson-exchange mechanisms



Theoretically sound model should also account for:

- **Channel-couplings** required by unitarity ($\pi N, \eta N, K\Lambda, K\Sigma$ stable channels)
- **2 π production** mechanisms ($\rho N, \sigma N, \pi\Delta \leftrightarrow \pi\pi N$ channels)

Dynamical Coupled-Channels (DCC) model accounts for these features

developed through analyzing data for $\gamma N, \pi N \rightarrow \pi N, \eta N, K\Lambda, K\Sigma$

~ 26,000 data points

DCC (Dynamical Coupled-Channel) model

Matsuyama et al., Phys. Rep. **439**, 193 (2007)

Kamano et al., PRC 88, 035209 (2013)

Coupled-channel Lippmann-Schwinger equation for meson-baryon scattering

$$T_{ab} = V_{ab} + \sum_c V_{ac} G_c T_{cb}$$

$$\{a, b, c\} = \pi N, \eta N, \pi\pi N, \pi\Delta, \sigma N, \rho N, K\Lambda, K\Sigma$$

Coupled-channels and hadron rescattering required by unitarity
is fully taken into account

In addition, γN channel is included perturbatively

DCC (Dynamical Coupled-Channel) model

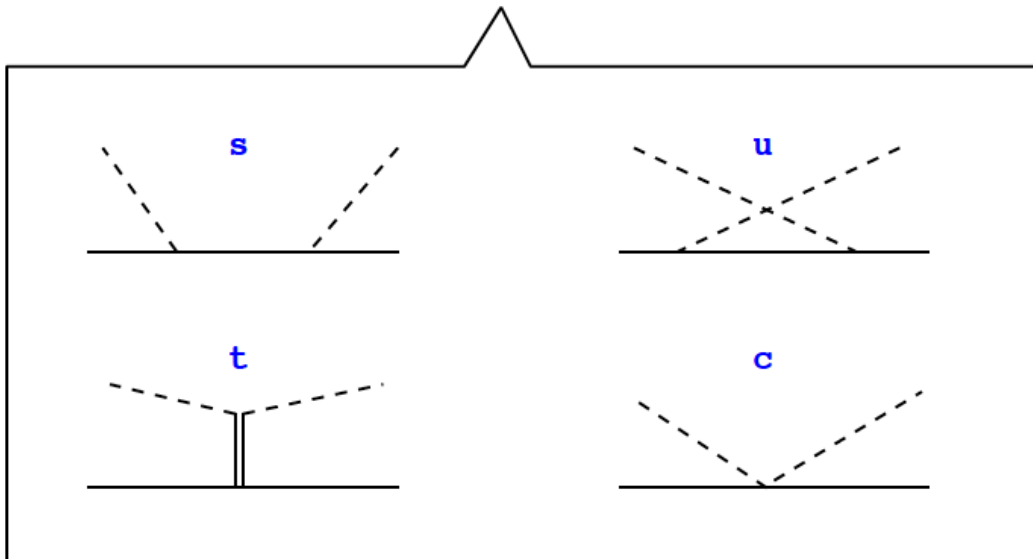
Matsuyama et al., Phys. Rep. **439**, 193 (2007)

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Coupled-channel Lippmann-Schwinger equation for meson-baryon scattering

$$T_{ab} = V_{ab} + \sum_c V_{ac} G_c T_{cb}$$

$$\mathbf{V}_{ab} = \text{[diagram 1]} + \text{[diagram 2]} + \mathbf{Z}$$



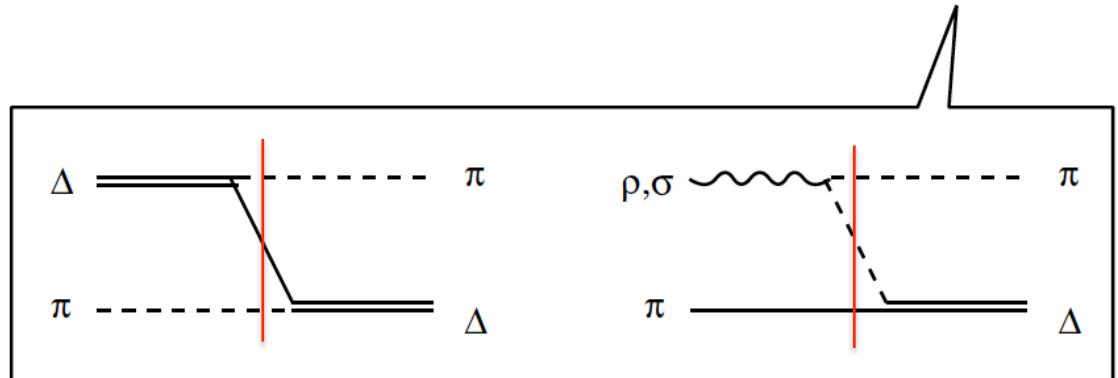
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$$T_{ab} = V_{ab} + \sum_c V_{ac} G_c T_{cb}$$



essential for three-body unitarity

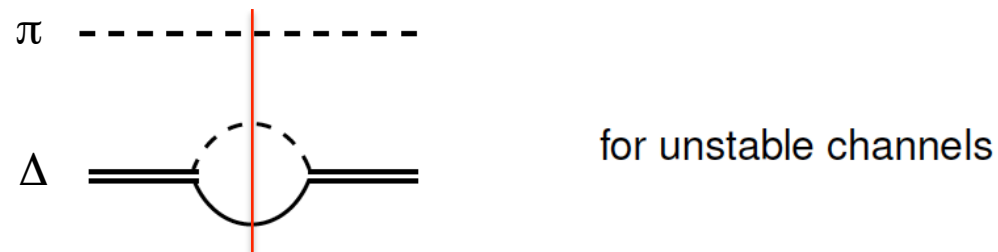
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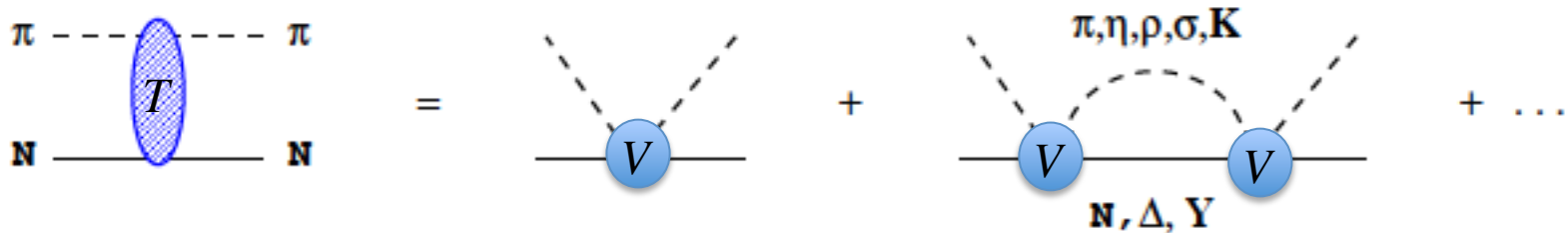
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Kamano et al., PRC **88**, 035209 (2013)

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In addition, γN channel is included perturbatively

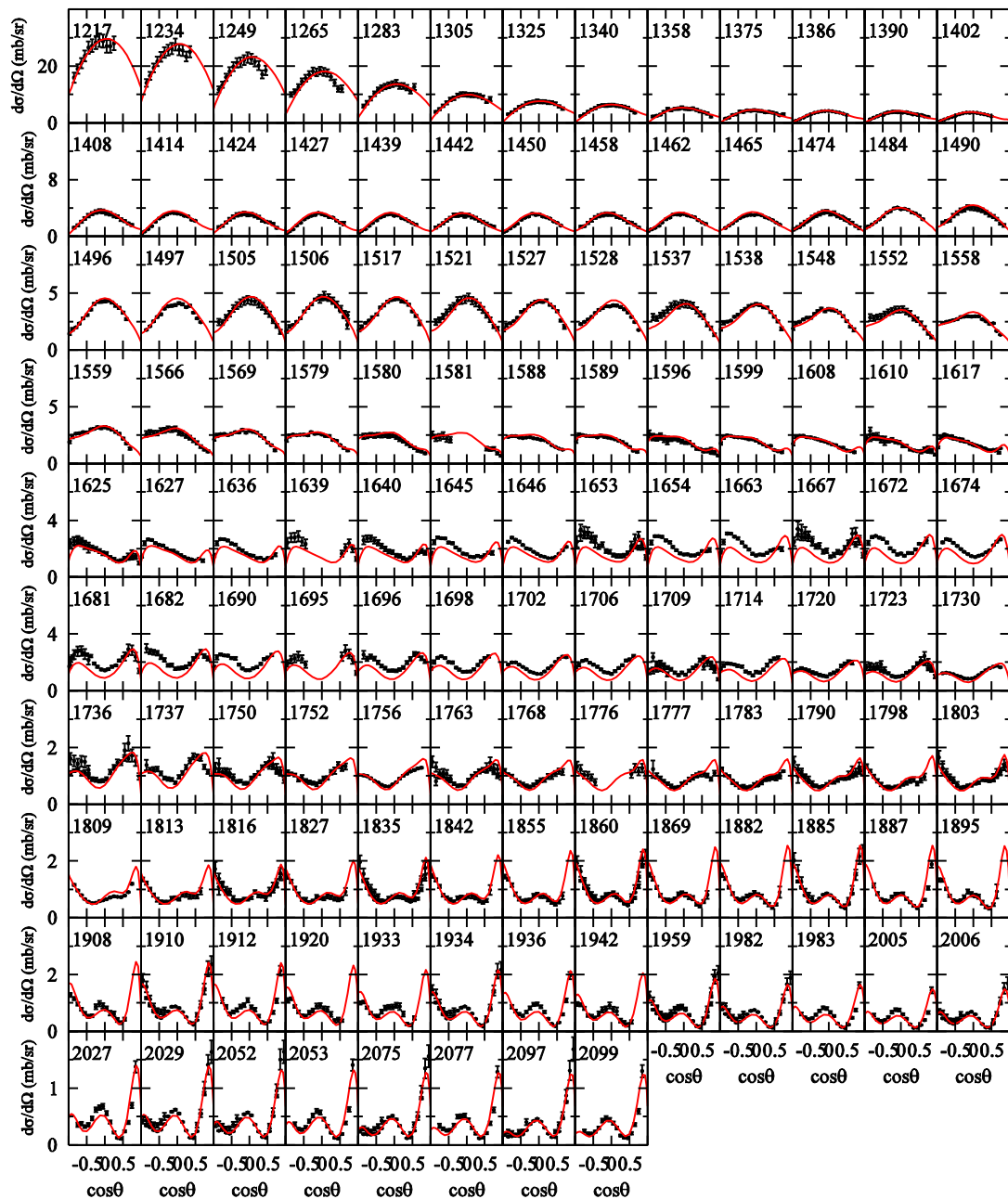


$\gamma p \rightarrow \pi^0 p$ $d\sigma/d\Omega$ for $W < 2.1$ GeV

Comparison of DCC model with data

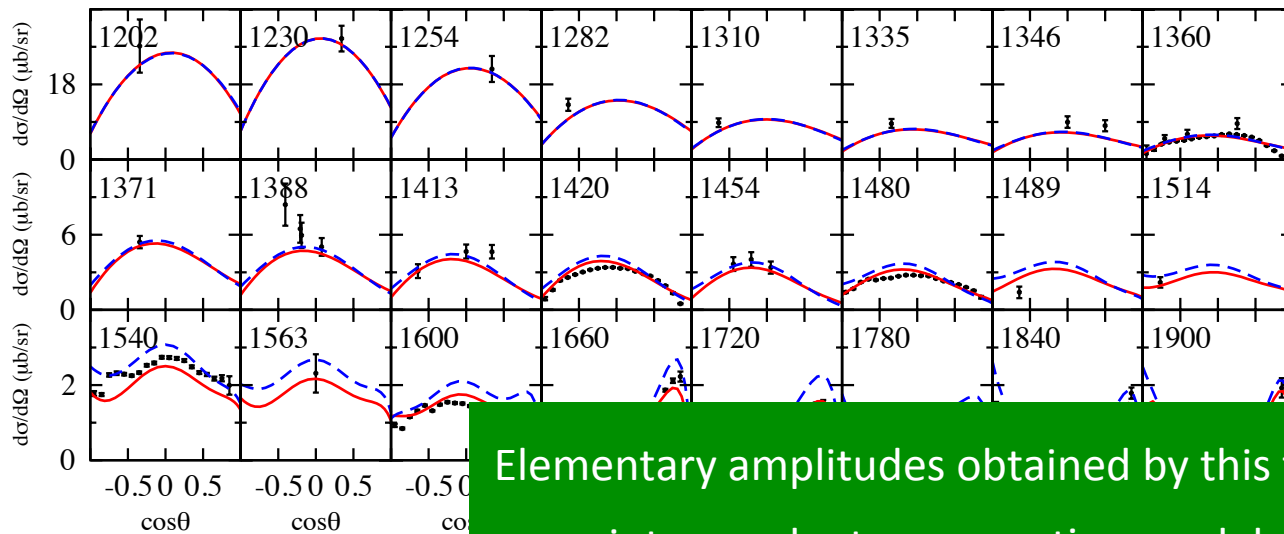
Kamano, Nakamura, Lee, Sato, PRC 88 (2013)

Reasonable fit to data
in the whole resonance region



$\gamma n \rightarrow \pi^0 n$ $d\sigma/d\Omega$ for $W < 2$ GeV

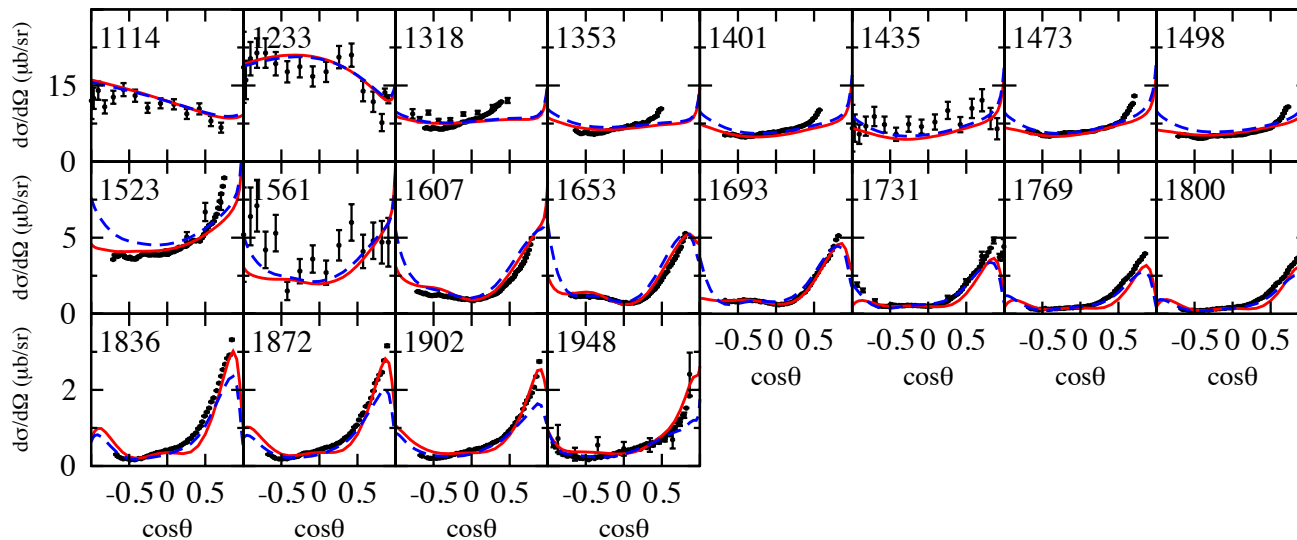
SXN, Kamano, Lee, Sato, arXiv:1804.04757



— : new fit
- - - : previous fit
PRC 88 (2013)

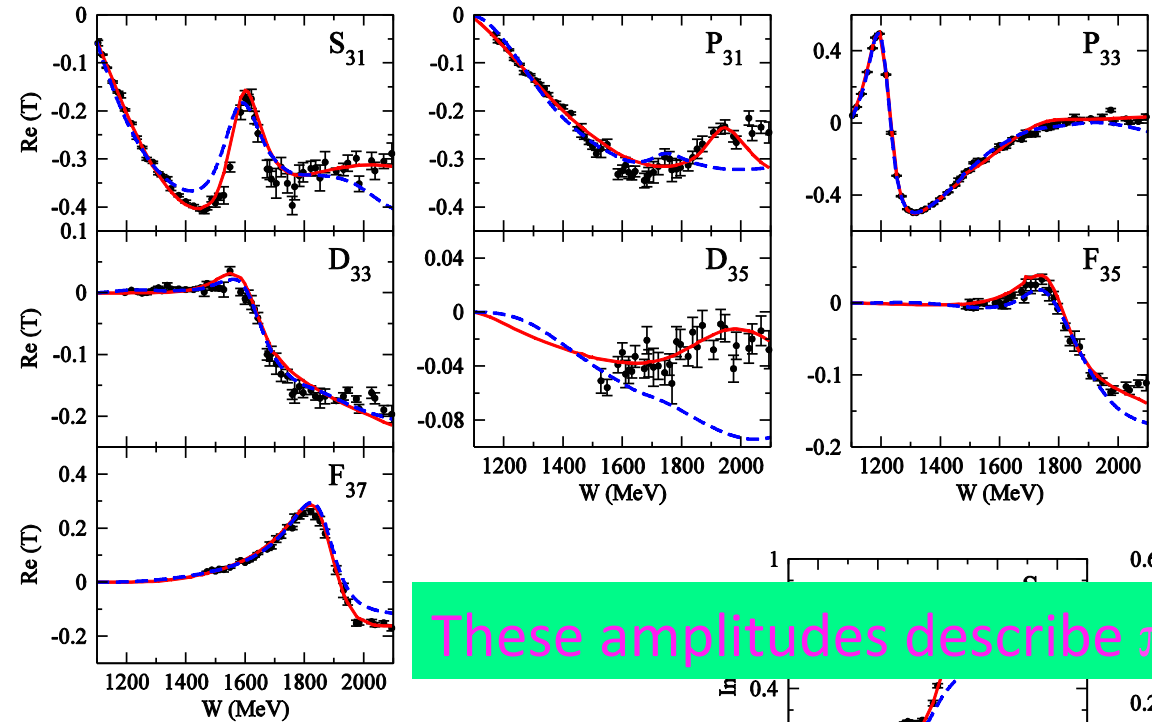
Recent MAMI data included
PRL 112, 142001 (2014)

Elementary amplitudes obtained by this fit
go into our deuteron reaction model

 $\gamma n \rightarrow \pi^+ p$ 

Recent JLab data included
PRC 96, 035204 (2017)

Partial wave amplitudes of πN scattering



Real part

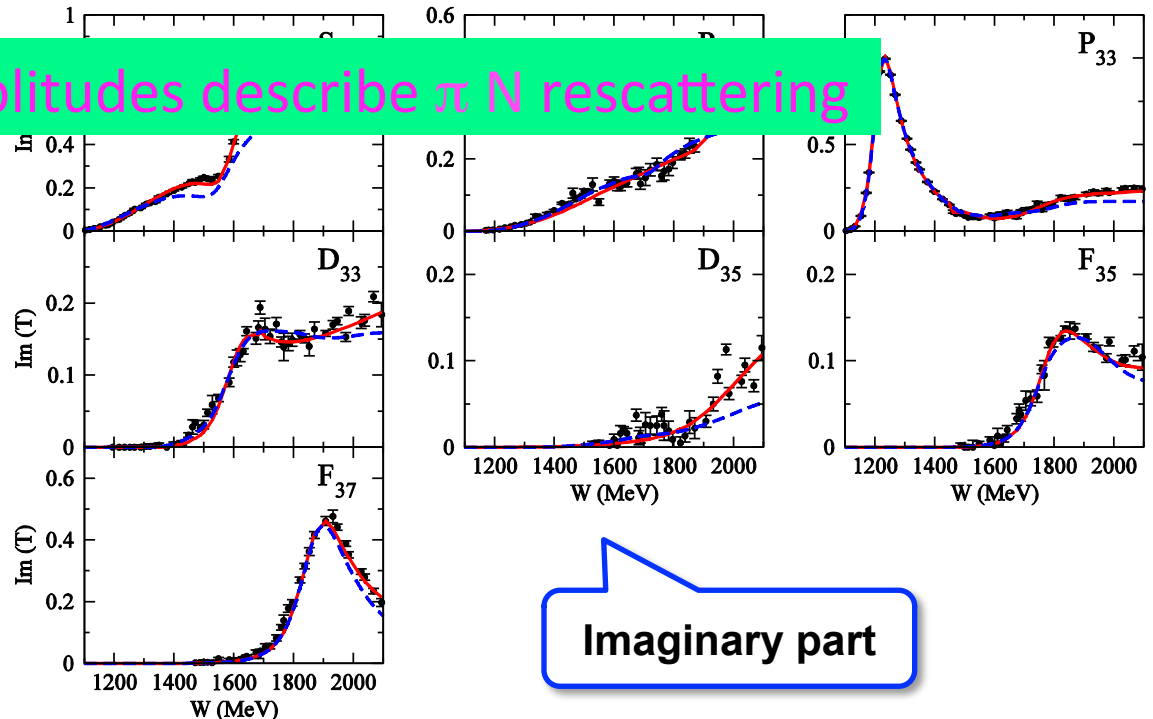
$$I = \frac{3}{2}$$

These amplitudes describe πN rescattering

— Kamano, Nakamura, Lee, Sato,
PRC 88 (2013)

- - - Previous model
(fitted to $\pi N \rightarrow \pi N$ data only)
[PRC76 065201 (2007)]

Data: SAID πN amplitude

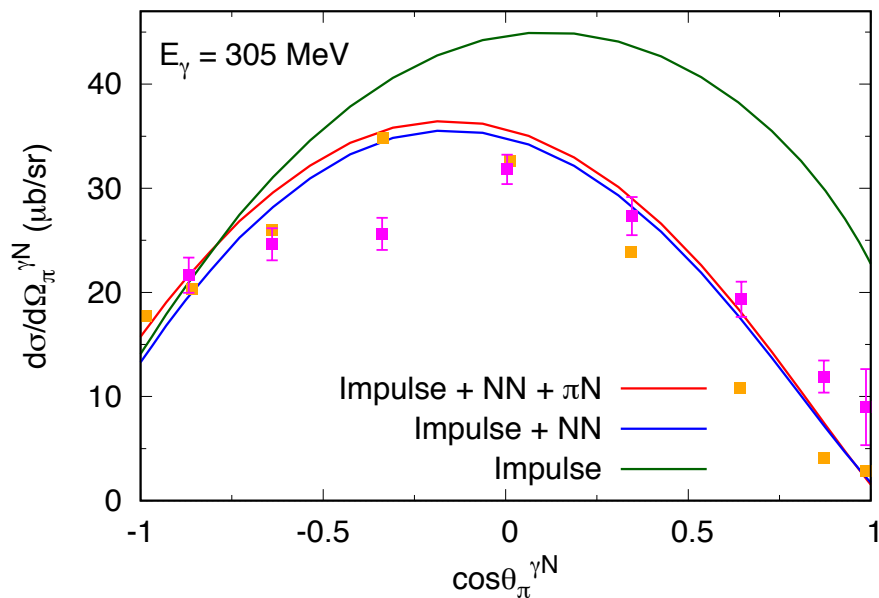


Imaginary part

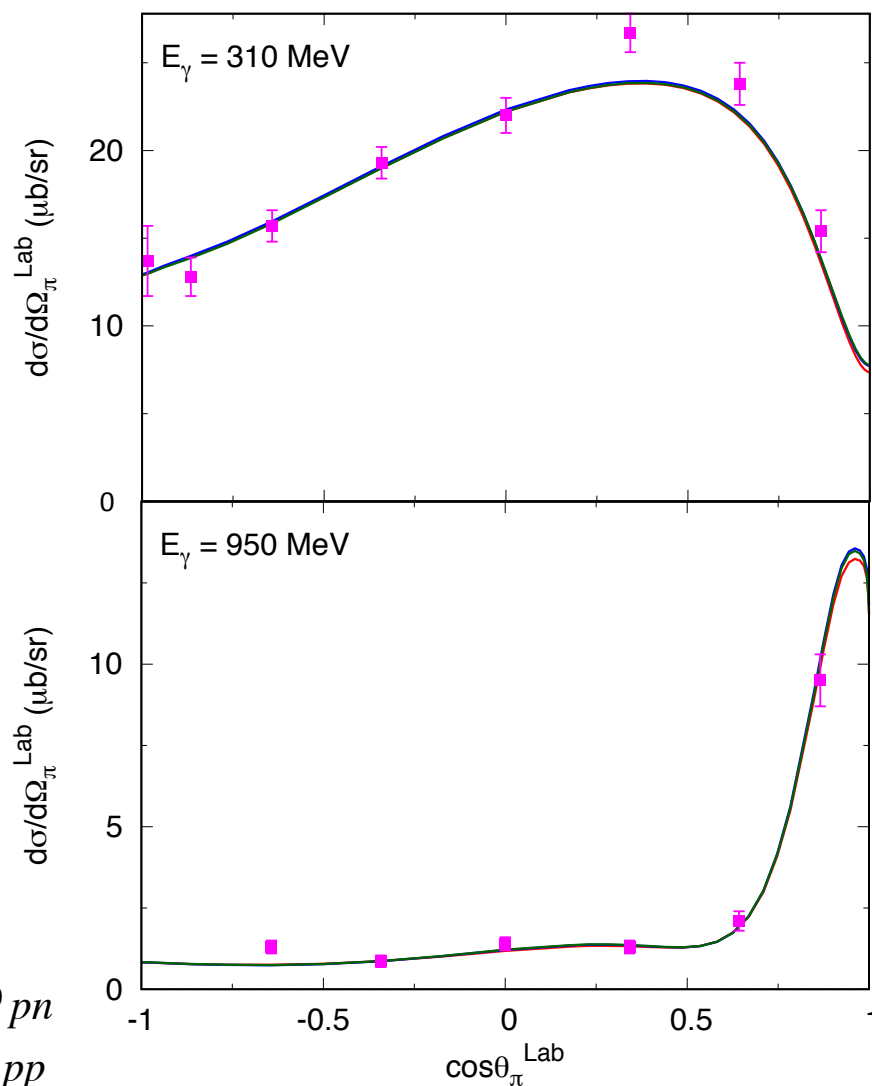
Results

$\gamma d \rightarrow \pi NN$: model predictions and data

$\gamma d \rightarrow \pi^0 pn$



$\gamma d \rightarrow \pi^- pp$



- Large NN FSI effect for π^0 productions
 \leftarrow NN and deuteron wave fn. are orthogonal
- FSI effects are small for π^- productions
- Reasonable agreement with data

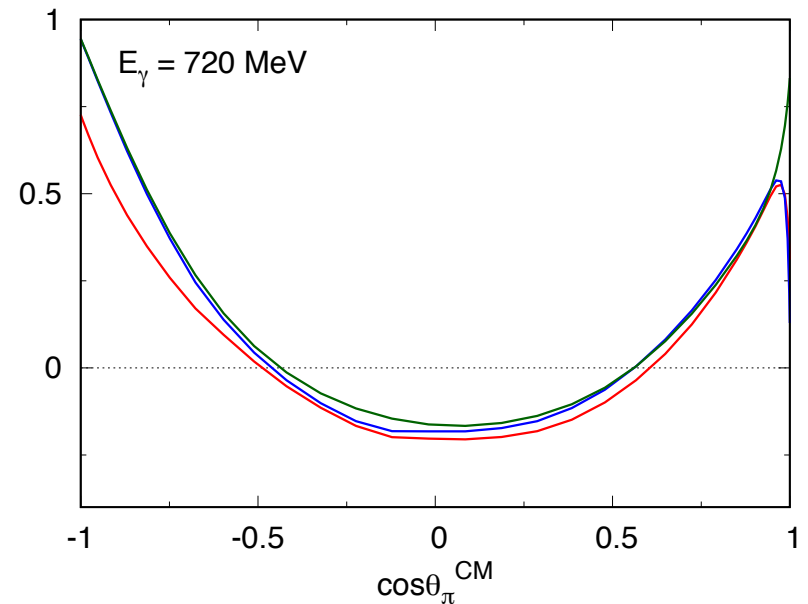
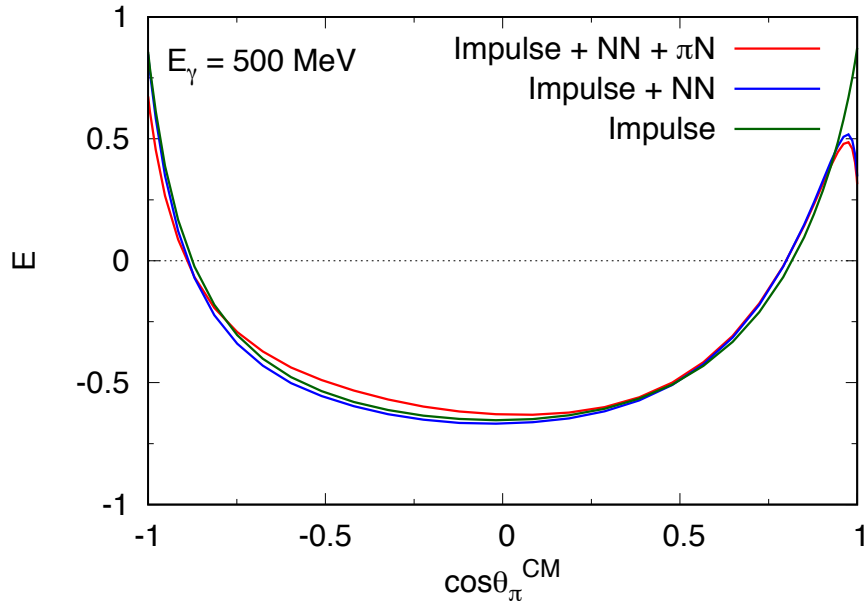
Data: EPJA 6, 309 (1999); 10, 365 (2001) for $\gamma d \rightarrow \pi^0 pn$

NPB 65, 158 (1973) for $\gamma d \rightarrow \pi^- pp$

Model predictions for E of $\gamma d \rightarrow \pi N N$

Polarization observable E :
$$E \equiv \frac{\sigma_{+-} - \sigma_{++}}{\sigma_{+-} + \sigma_{++}} \quad \sigma_{\pm\pm} \equiv \sigma(s_\gamma^z = +1, s_d^z = \pm 1)$$

$\gamma d \rightarrow \pi^0 pn$

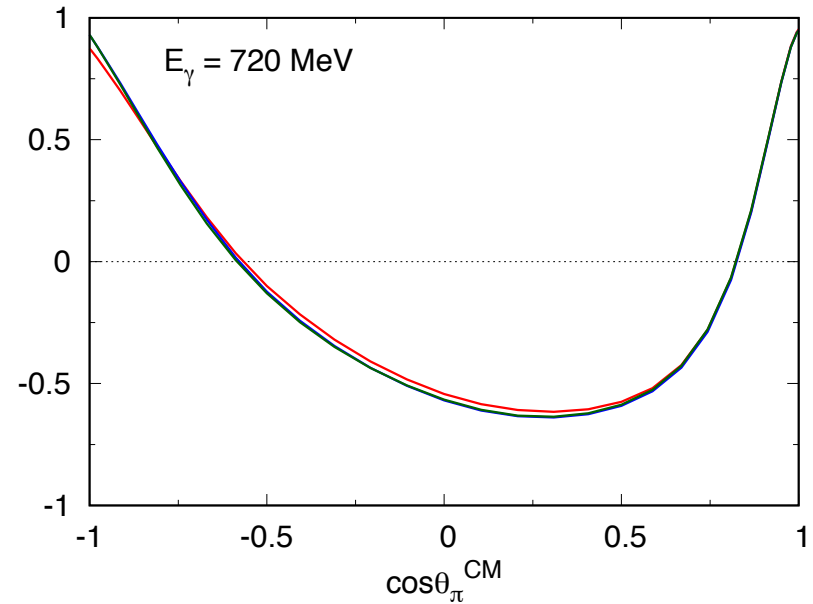
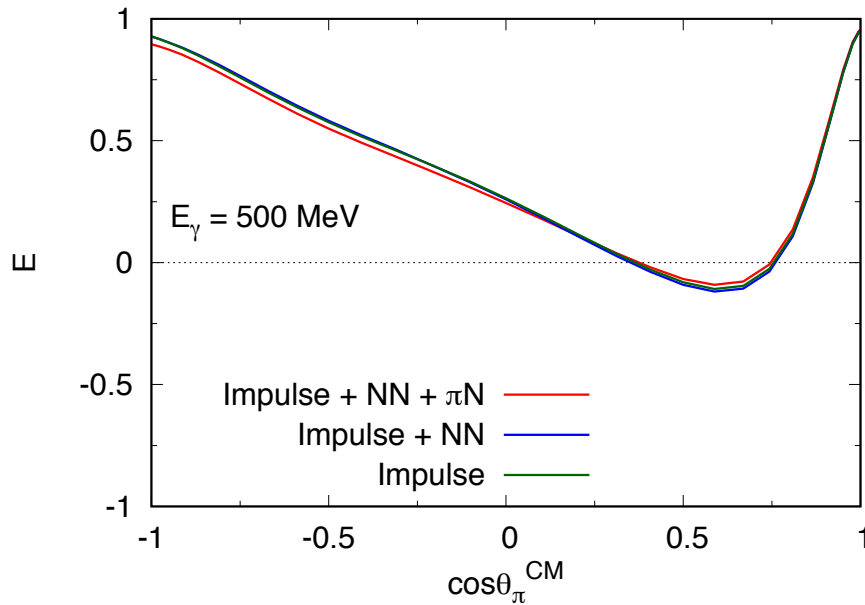


- FSI effects are smaller than $d\sigma/d\Omega_\pi$, but still visible
 → not completely cancelled out in the ratio

Model predictions for E of $\gamma d \rightarrow \pi N N$

Polarization observable E :
$$E \equiv \frac{\sigma_{+-} - \sigma_{++}}{\sigma_{+-} + \sigma_{++}} \quad \sigma_{\pm\pm} \equiv \sigma(s_\gamma^z = +1, s_d^z = \pm 1)$$

$\gamma d \rightarrow \pi^- pp$



- FSI effects are small as the same for $d\sigma/d\Omega_\pi$

Quasi-free $\gamma'n'$ cross sections
extracted from γd cross sections

How to extract $\sigma(\gamma' n' \rightarrow \pi^0 n)$ from $\sigma(\gamma d \rightarrow \pi^0 np)$

First we need to establish a formula to relate $\sigma(\gamma' n' \rightarrow \pi^0 n)$ and $\sigma(\gamma d \rightarrow \pi^0 np)$

Ideal situation : only this mechanism contributes \rightarrow

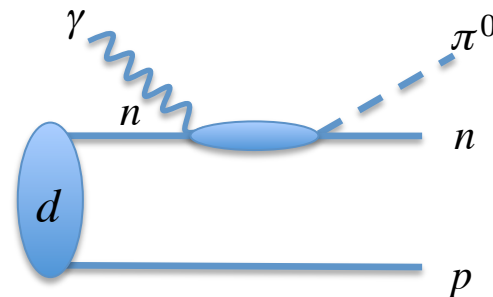
No FSI, No exchange terms, No deuteron D-wave

“quasi-free mechanism”



$$\frac{d^2\sigma(\gamma d \rightarrow \pi^0 np; E_\gamma)}{dW d\cos\theta_\pi^*} = \phi(W) \frac{\tilde{E}_\gamma}{E_\gamma} \frac{d\sigma(\gamma n \rightarrow \pi^0 n; W)}{d\cos\theta_\pi^*}$$

E_γ (\tilde{E}_γ) : γ - d (γ - n') Lab frame



W : $\pi^0 n$ invariant mass $\cos\theta_\pi^* \equiv \hat{q}_\gamma \cdot \hat{k}_\pi$ in π^0 - n CM frame

$$\phi(W) \equiv \int dp_p^3 \delta(W - W(\vec{p}_p, E_\gamma)) |\psi_d(\vec{p}_p)|^2, \quad W(\vec{p}_p, E_\gamma) \equiv \sqrt{(E_\gamma + m_d - E_p(\vec{p}_p))^2 - (\vec{q}_\gamma - \vec{p}_p)^2}$$

$\phi(W)$: Probability of a photon of E_γ interacting with the nucleon with W

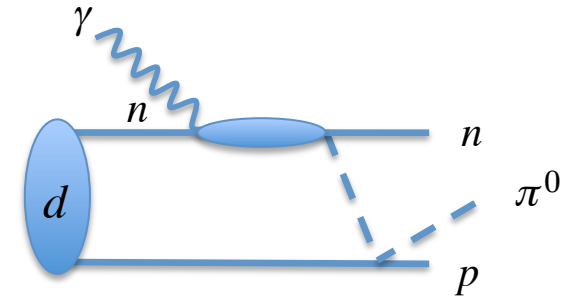
How to extract $\sigma(\gamma n \rightarrow \pi^0 n)$ from $\sigma(\gamma d \rightarrow \pi^0 np)$

In reality, photon hits the other nucleon and FSI contribute

→ kinematical cuts are applied to remove them

Counterpart from our model :

$$\left. \frac{d^2 \sigma(\gamma d \rightarrow \pi^0 np)}{dW d \cos \theta_\pi^*} \right|_{\text{cuts}}$$



Assuming that this is solely from quasi-free contribution integrated over the same phase-space

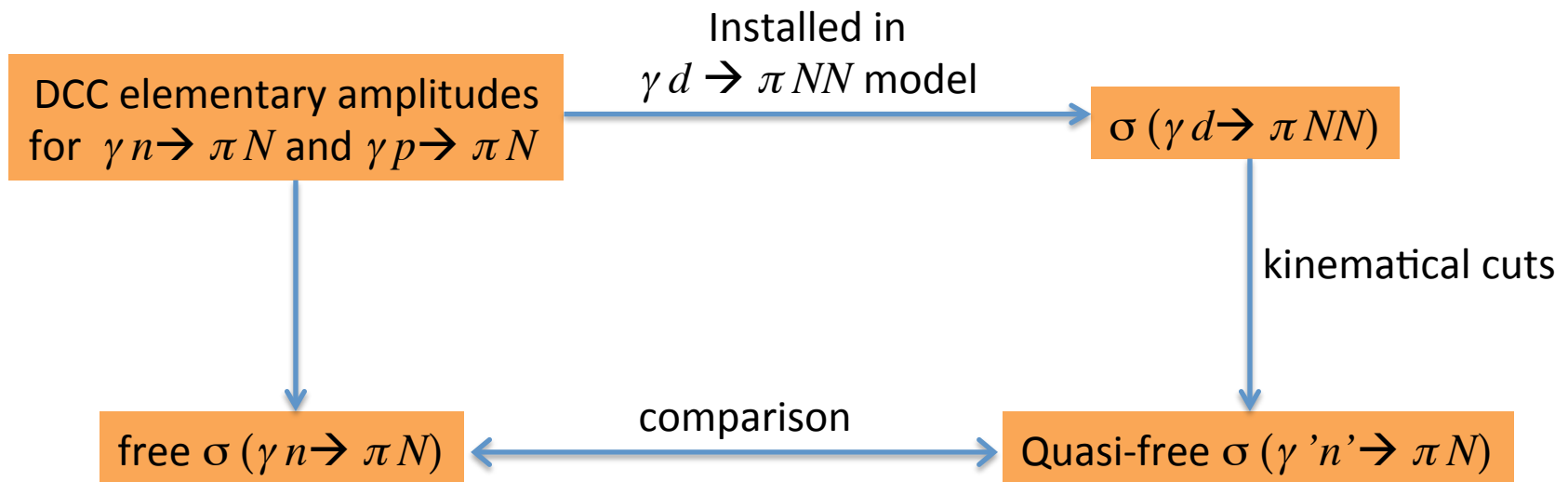
$$\left. \frac{d^2 \sigma(\gamma d \rightarrow \pi^0 np)}{dW d \cos \theta_\pi^*} \right|_{\text{cuts}} = \phi_{\text{mod}}(W) \frac{\tilde{E}_\gamma}{E_\gamma} \frac{d\sigma(\gamma n \rightarrow \pi^0 n)}{d \cos \theta_\pi^*}, \quad \phi_{\text{mod}}(W) \equiv \int_{\text{cuts}} dp_p^3 \delta(W - W(\vec{p}_p, E_\gamma)) |\psi_d(\vec{p}_p)|^2$$



Formula to extract $\sigma(\gamma n \rightarrow \pi^0 n)$ from $\sigma(\gamma d \rightarrow \pi^0 np)$ from either experiment or model

Examine within our model how much $\sigma(\gamma n \rightarrow \pi^0 n)$ deviates from 'free' $\sigma(\gamma n \rightarrow \pi^0 n)$ due to FSI

Q: How FSI and kinematical cuts could distort $\gamma'n'$ observables ($d\sigma/d\Omega, \Sigma, E, G$) ?



Kinematical cuts

For extracting $\gamma' n' \rightarrow \pi^- p$ from $\gamma d \rightarrow \pi^- pp$, we use :

[1] CLAS@JLab, PRC 86 (2012)
 [2] CLAS@JLab, PRC 96 (2017)
 [3] CLAS@JLab, PRL 118 (2017)

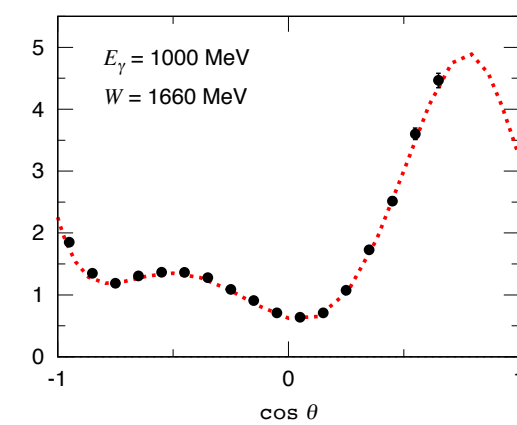
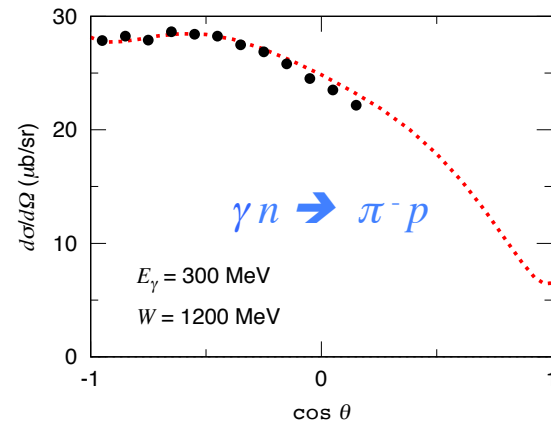
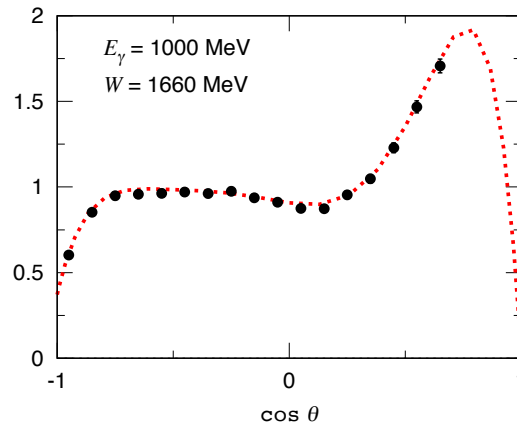
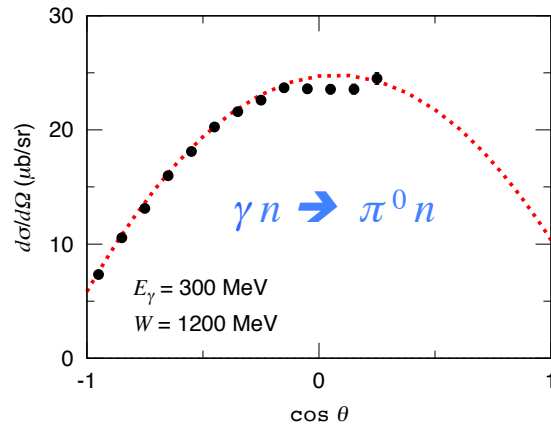
Ref.	[1] ($d\sigma/d\Omega$)	[2] ($d\sigma/d\Omega$)	[3] (E)
E_γ (MeV)	301 – 455	445 – 2510	700 - 2400
π^- momentum (MeV)	> 80	> 100	> 400
Faster proton momentum (MeV)	> 270	> 360	> 400
Slower proton momentum (MeV)	< 270	< 200	< 100
$\Delta\phi_\pi = \phi_\pi - \phi_{\text{faster proton}} $	–	–	$160^\circ < \Delta\phi < 200^\circ$

For extracting $\gamma' n' \rightarrow \pi^0 n$ from $\gamma d \rightarrow \pi^0 pn$, the same cuts are used after modification
 “ π^- ” \rightarrow “ π^0 ”, “Faster proton” \rightarrow “neutron”, “Slower proton” \rightarrow “proton”

NOTE: Different kinematical cuts were used in A2@MAMI analysis on $\gamma n \rightarrow \pi^0 n$

Check the extraction method in ideal case

Q: With quasi-free mechanism only, examine if the formula $\frac{d^2\sigma(\gamma d \rightarrow \pi^0 np)}{dW d\cos\theta_\pi} \Big|_{\text{cuts}} = \phi_{\text{mod}}(W) \frac{\tilde{E}_\gamma}{E_\gamma} \frac{d\sigma(\gamma n \rightarrow \pi^0 n)}{d\cos\theta_\pi}$ returns $\sigma(\gamma 'n' \rightarrow \pi N)$ that agree with 'free' $\sigma(\gamma n \rightarrow \pi N)$



- : extracted $\sigma(\gamma 'n' \rightarrow \pi N)$
- ⋯ : free $\sigma(\gamma n \rightarrow \pi N)$

Note:

- * Phase-space integrals for $\sigma(\gamma d)$ are done with Monte-Carlo method to easily implement any cuts; statistical error bars are very small
- * Forward pion productions are invisible due to the kinematical cuts

Free $\sigma(\gamma n \rightarrow \pi N)$ are well reproduced as expected for the ideal case

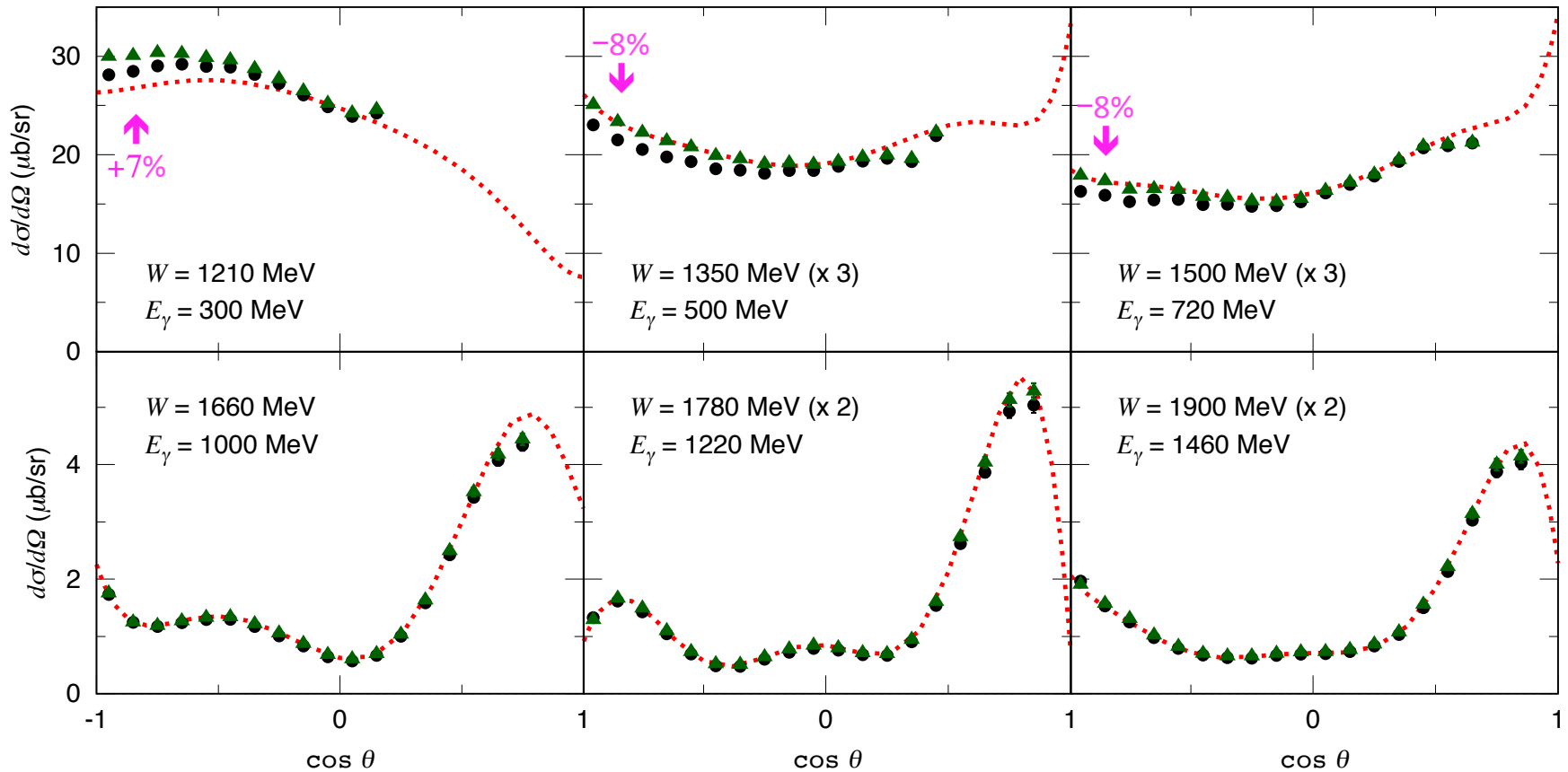
Now go on to realistic case !

$d\sigma/d\Omega_\pi$ for $\gamma' n' \rightarrow \pi^- p$

▲ : impulse + NN FSI

● : impulse + NN FSI + πN FSI

⋯ : free $\sigma(\gamma n \rightarrow \pi N)$



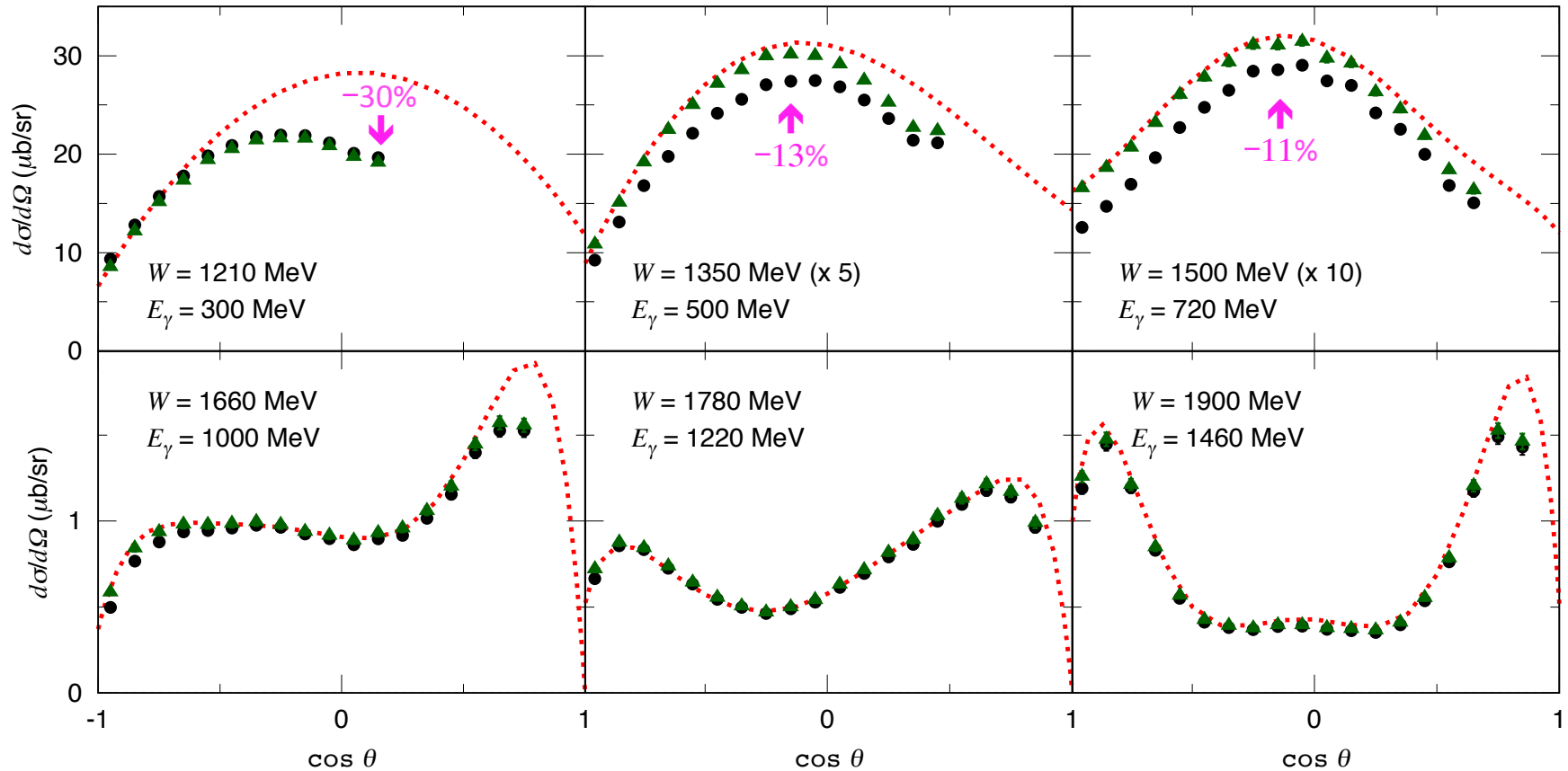
- πN FSI reduce $d\sigma/d\Omega_\pi$ of backward pion
- kinematical cuts cannot remove the FSI effect
- Larger FSI effects for smaller E_γ

$d\sigma/d\Omega_\pi$ for $\gamma 'n' \rightarrow \pi^0 n$

▲ : impulse + NN FSI

● : impulse + NN FSI + πN FSI

..... : free $\sigma(\gamma n \rightarrow \pi N)$



- Significant FSI effects reduce $d\sigma/d\Omega_\pi$; πN and NN FSI are comparably important
- kinematical cuts cannot remove FSI effect
- Larger FSI effects for smaller E_γ

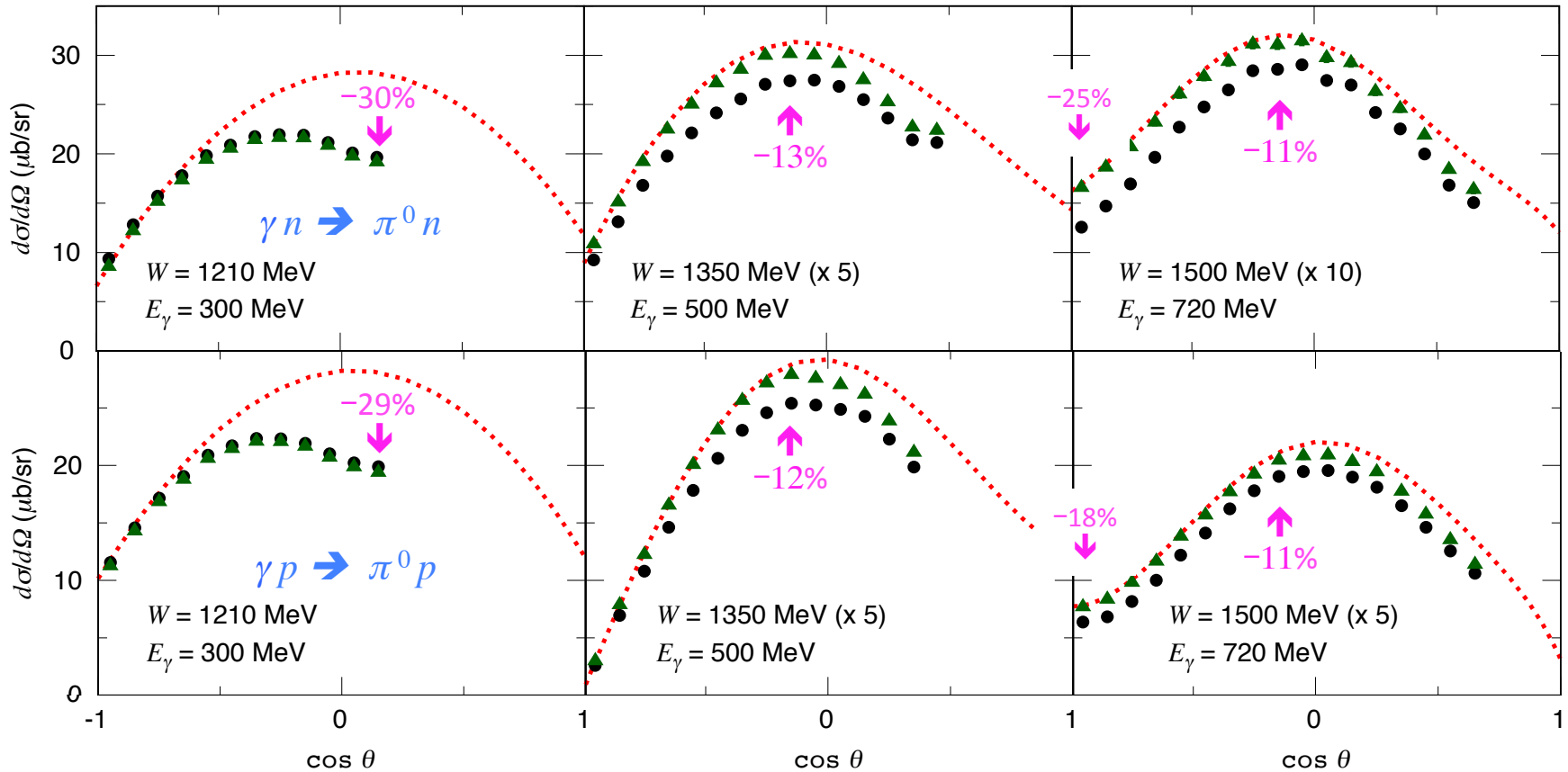
Comparison



▲ : impulse + NN FSI

● : impulse + NN FSI + π N FSI

..... : free $\sigma(\gamma n \rightarrow \pi N)$



- FSI effects on $\gamma' n' \rightarrow \pi^0 n$ and $\gamma' p' \rightarrow \pi^0 p$ are generally similar (a few % difference)
 - ↔ Same FSI effects are assumed in A2@MAMI analysis [PRL 112 (2014)]
- But sometimes more different

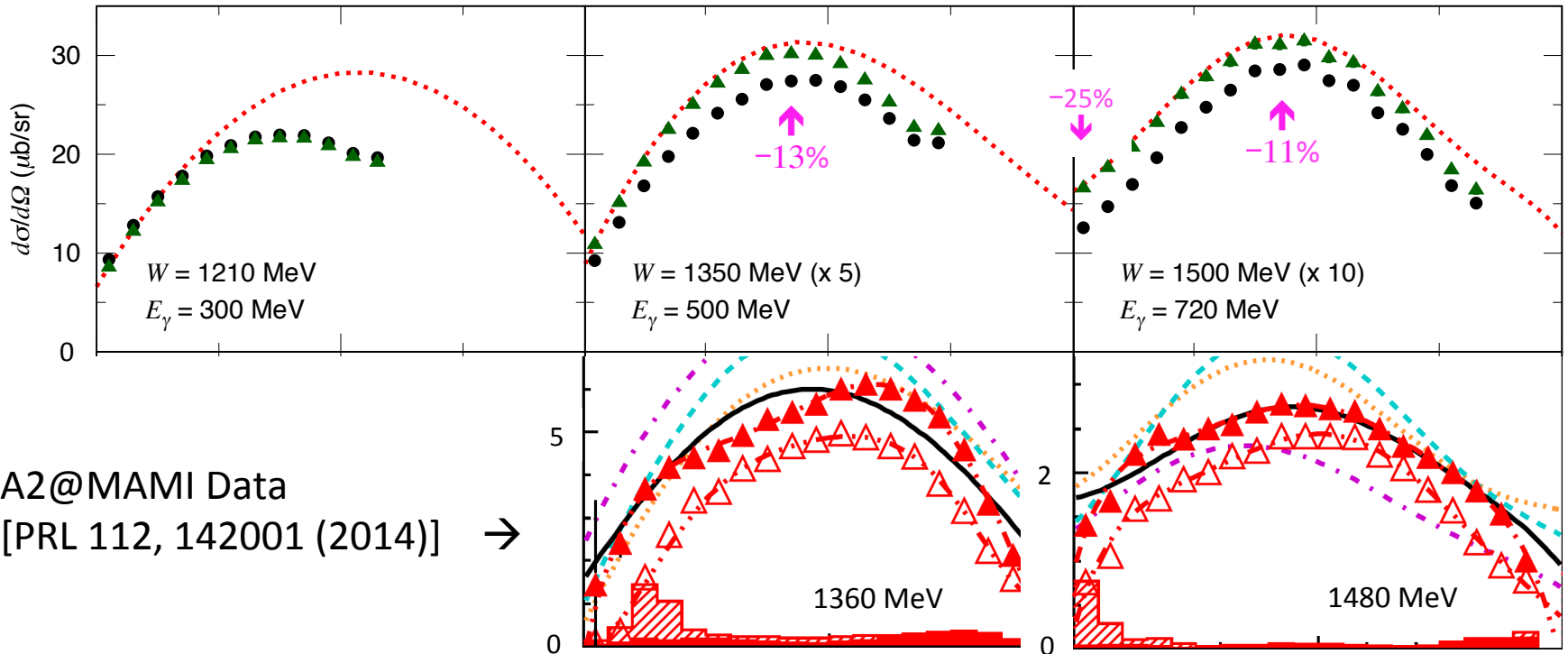
Comparison with Data : FSI effect



▲ : impulse + NN FSI

● : impulse + NN FSI + πN FSI

⋯ : free $\sigma(\gamma n \rightarrow \pi N)$



- Fairly good agreement on the FSI effects estimated in A2@MAMI analysis

- A2@MAMI analysis assumed $\frac{\sigma(\gamma'n' \rightarrow \pi^0 n)}{\text{free } \sigma(\gamma n \rightarrow \pi^0 n)} = \frac{\sigma(\gamma'p' \rightarrow \pi^0 p)}{\text{free } \sigma(\gamma p \rightarrow \pi^0 p)}$ no theoretical estimate

- FSI effects can depend on the cuts; MAMI analysis uses different cuts

Extraction of polarization observables for $\gamma 'n' \rightarrow \pi^0 n$ from $\gamma d \rightarrow \pi^0 np$

A possible complication : relativistic nucleon spin rotation

← Lorentz boost from γn CM frame to γd Lab frame

We confirmed this effect is very small for Σ, E, G to be discussed here → ignored

Polarized cross sections can be extracted using formula similar to the unpolarized one

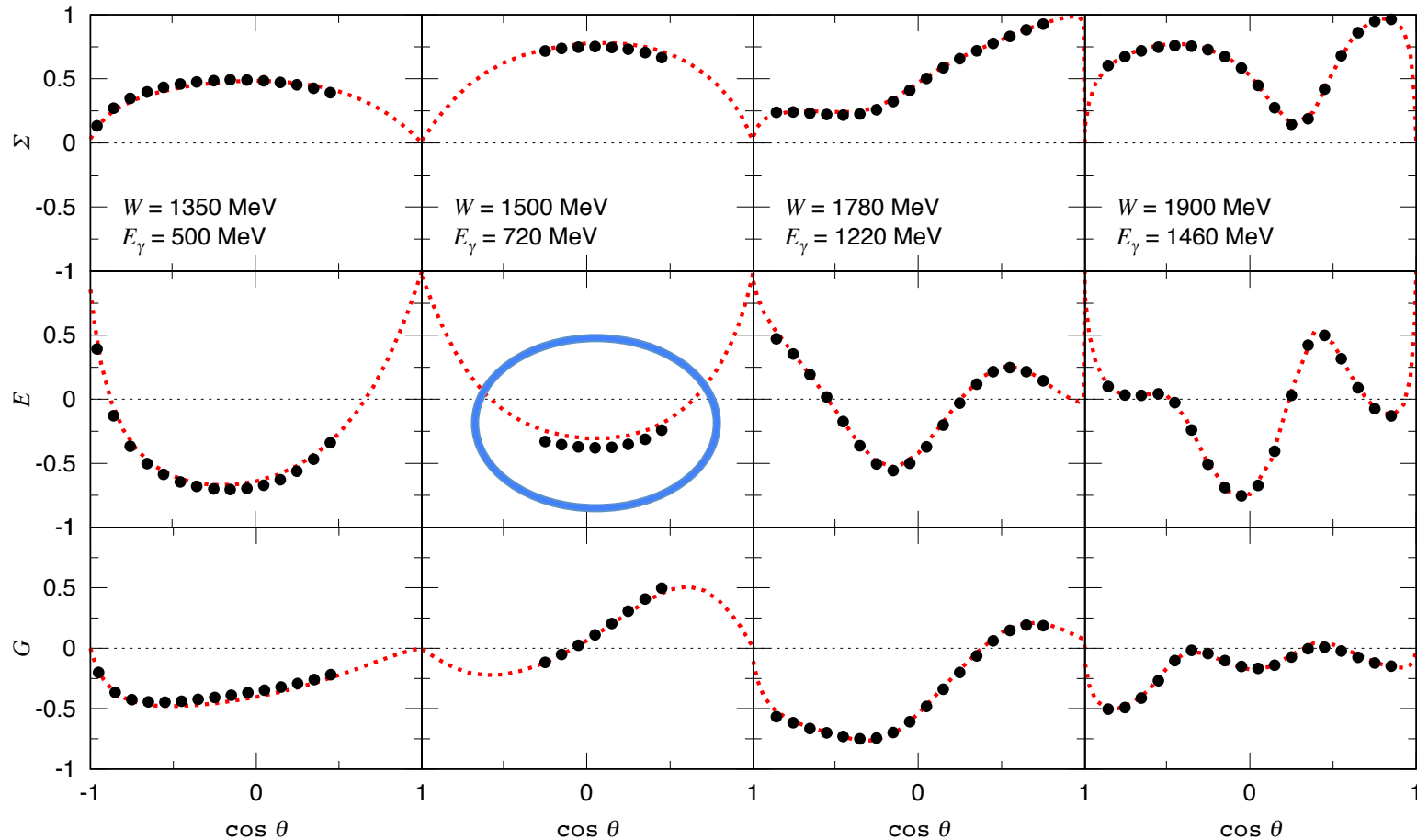
$$\left. \frac{d^2\sigma(\gamma d \rightarrow \pi^0 np; \lambda, s_d = +1)}{dW d\cos\theta_\pi^*} \right|_{\text{cuts}} = \phi_{\text{mod}}(W) \frac{\tilde{E}_\gamma}{E_\gamma} \frac{d\sigma(\gamma n \rightarrow \pi^0 n; \lambda, s_N = +1/2)}{d\cos\theta_\pi^*}$$

$$E = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{d\sigma(\lambda = -1, s_N = +1/2) - d\sigma(\lambda = +1, s_N = +1/2)}{d\sigma(\lambda = -1, s_N = +1/2) + d\sigma(\lambda = +1, s_N = +1/2)}$$

Σ, E, G for $\gamma n \rightarrow \pi^0 n$

● : impulse + NN FSI + πN FSI

⋯ : free $\sigma(\gamma n \rightarrow \pi N)$



- Free $\gamma n \rightarrow \pi N$ polarization asymmetries are well reproduced
- FSI effects are very small ; canceled by taking the ratios; One exception
- No FSI effects are assumed in A2@MAMI analysis for E [PLB 770 (2017)]

(More) common extraction method

Extraction formula used so far in this presentation

$$\left. \frac{d^2\sigma(\gamma d \rightarrow \pi^0 np; E_\gamma)}{dW d\cos\theta_\pi^*} \right|_{\text{cuts}} = \phi_{\text{mod}}(W) \frac{\tilde{E}_\gamma}{E_\gamma} \frac{d\sigma(\gamma n \rightarrow \pi^0 n; W)}{d\cos\theta_\pi^*}, \quad \phi_{\text{mod}}(W) \equiv \int^{\text{cuts}} dp_p^3 \delta(W - W(\vec{p}_p, E_\gamma)) |\psi_d(\vec{p}_p)|^2$$

Integrate both sides with respect to W , assuming dominant contribution from nucleon-at-rest in deuteron

(and/or fairly weak W -dependence)

(More) commonly used extraction formula

$$\left. \frac{d^2\sigma(\gamma d \rightarrow \pi^0 np; E_\gamma)}{d\cos\theta_\pi^*} \right|_{\text{cuts}} = \frac{d\sigma(\gamma n \rightarrow \pi^0 n; \bar{W})}{d\cos\theta_\pi^*} \int^{\text{cuts}} dp_p^3 \frac{\tilde{E}_\gamma}{E_\gamma} |\psi_d(\vec{p}_p)|^2, \quad \bar{W}^2 = (E_\gamma + m_N)^2 - E_\gamma^2$$

Used in recent CLAS@JLab analyses: PRC 96 (2017) for $d\sigma/d\Omega_\pi$

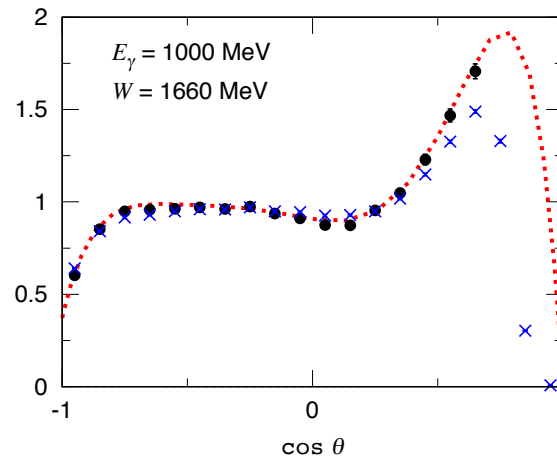
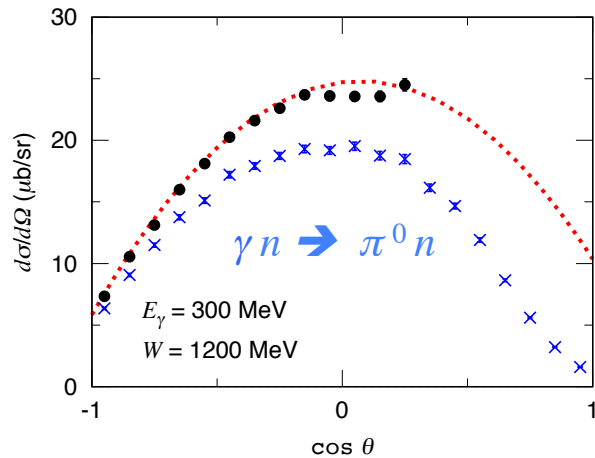
PRL 118 (2017) for E

also in theoretical analysis

Tarasov et al., PRC 84, 035203 (2011)

Examine the validity of this formula without **W -cut**

(More) common extraction method



Quasi-free mechanism only;
no FSI, no exchange terms

- : extracted with W -cut
- × : without W -cut
- ⋯ : free $\sigma(\gamma n \rightarrow \pi N)$

- $E_\gamma = 300 \text{ MeV}$: Significant difference in $\Delta(1232)$ -region

1. Sharp peak of $\sigma_{\gamma n}$ at $W \sim 1.2 \text{ GeV}$

2. $\sigma_{\gamma d}$ is an average of $\sigma_{\gamma n}$ over $W \sim 1.18\text{--}1.21 \text{ GeV}$ ($W < 1.2 \text{ GeV}$ in $\cos \theta > 0.3$)
because of Fermi motion

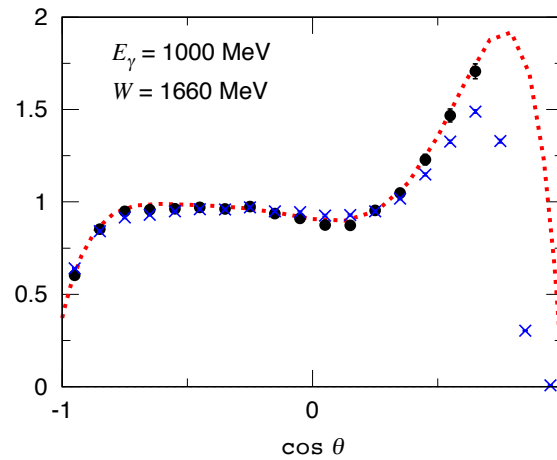
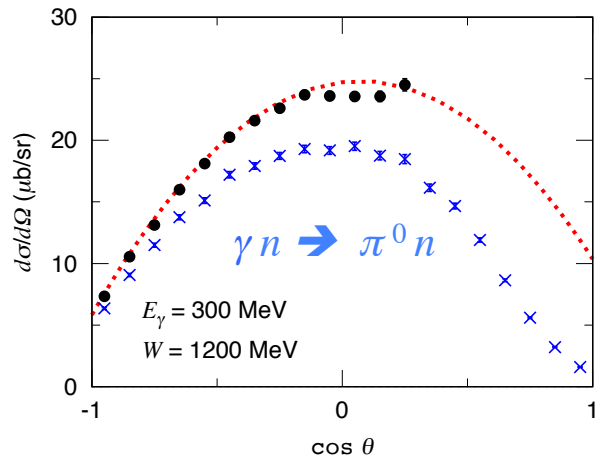
→ extracted $\sigma_{\gamma n}$ without W -cut is necessarily smaller than $\sigma_{\gamma n}$ at $W = 1.2 \text{ GeV}$

- $E_\gamma = 1 \text{ GeV}$: Good agreement between with and without W -cut in $\cos \theta < 0.3$

1. Mild and monotonic W -dependence of $\sigma_{\gamma n}$ for $W \sim 1.6\text{--}1.7 \text{ GeV}$

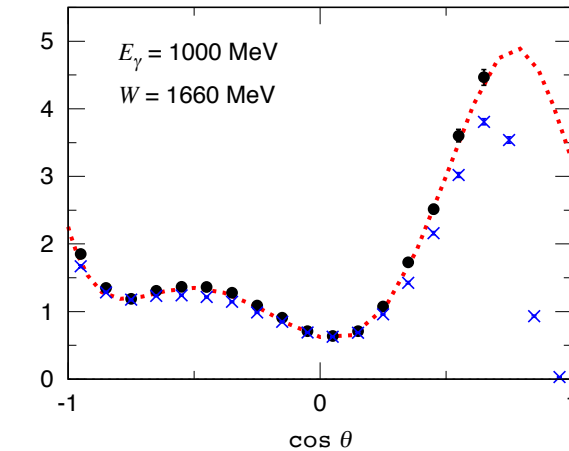
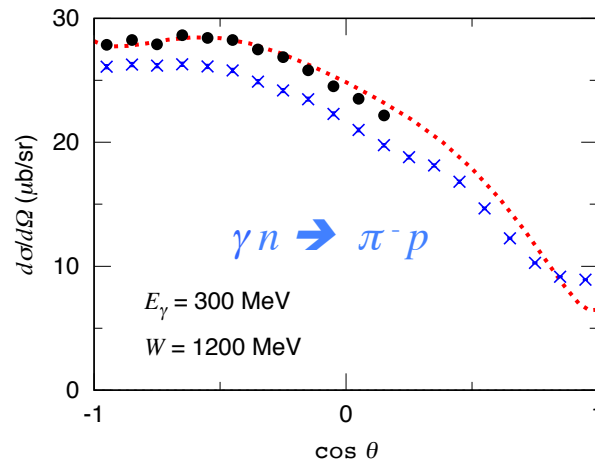
→ W -dependence is canceled in $\sigma_{\gamma d}$ by W -average; cancel is incomplete in $\cos \theta > 0.3$

(More) common extraction method



Quasi-free mechanism only;
no FSI, no exchange terms

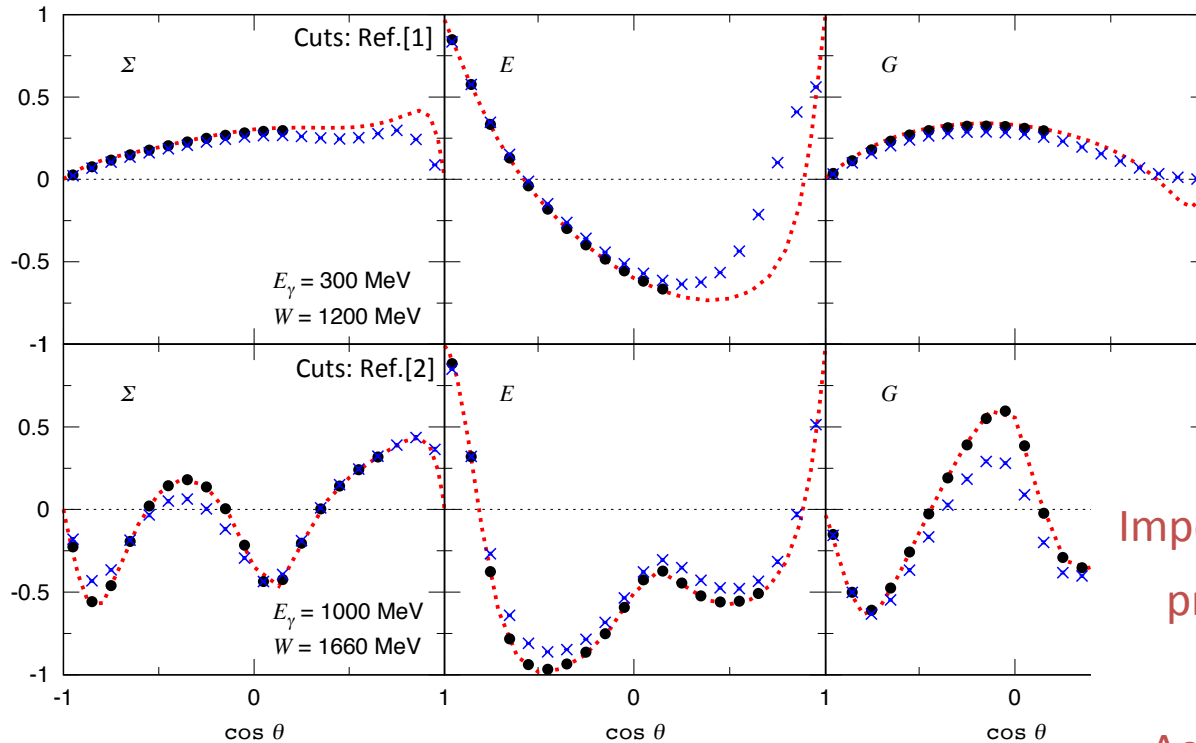
- : extracted with W -cut
- × : without W -cut
- ⋯ : free $\sigma(\gamma n \rightarrow \pi N)$



Similar result for $\gamma n \rightarrow \pi^- p$

(More) common extraction method

Polarization observables for $\gamma n \rightarrow \pi^- p$



Quasi-free mechanism only;
no FSI, no exchange terms

- : extracted with W -cut
- × : without W -cut
- ⋯ : free $\sigma(\gamma n \rightarrow \pi N)$

Important to apply W -cut to suppress
problematic Fermi motion effect



Accuracy of extracted observables

Reasons for deviations from free ones

- E at $E_\gamma=300\text{MeV}$, $\cos \theta > 0.3$: nucleon-at-rest kinematics is not allowed by cuts;
contributions are from different W
- G at $E_\gamma=1\text{GeV}$: Average non-monotonic W -dependence \rightarrow W -dependence not cancelled

Conclusion

Conclusions

- $d\sigma/d\Omega_\pi$ and Σ , E , G for $\gamma' n' \rightarrow \pi^0 n$, $\pi^- p$ are extracted from $\gamma d \rightarrow \pi NN$ using kinematical cuts of CLAS@JLab analyses
- $d\sigma/d\Omega_\pi$ for $\gamma' n' \rightarrow \pi^0 n$ are significantly distorted by FSI effects;
FSI effects on $\gamma' n' \rightarrow \pi^0 n$ and $\gamma' p' \rightarrow \pi^0 p$ are sometimes visibly different
FSI effects remain even after the kinematical cuts being applied
- Σ , E , G for $\gamma' n' \rightarrow \pi N$ are mostly unaffected by FSI effects
FSI effects are mostly cancelled in the ratios
- Without applying the W -cut, extracted observables can be rather largely distorted by Fermi motion

Future plan

- Similar study for electroproductions
- Similar study for $\gamma' n' \rightarrow \eta n, K\Lambda, K\Sigma$
- More polarization observables including recoil polarization

Thank you very much
for your attention

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