



universität
wien

Factorization of Inclusive Semileptonic Off-Shell and Boosted Top Quark Decays in the Endpoint region

Christoph Regner

with

André H. Hoang, Simon Plätzer and Ines Ruffa

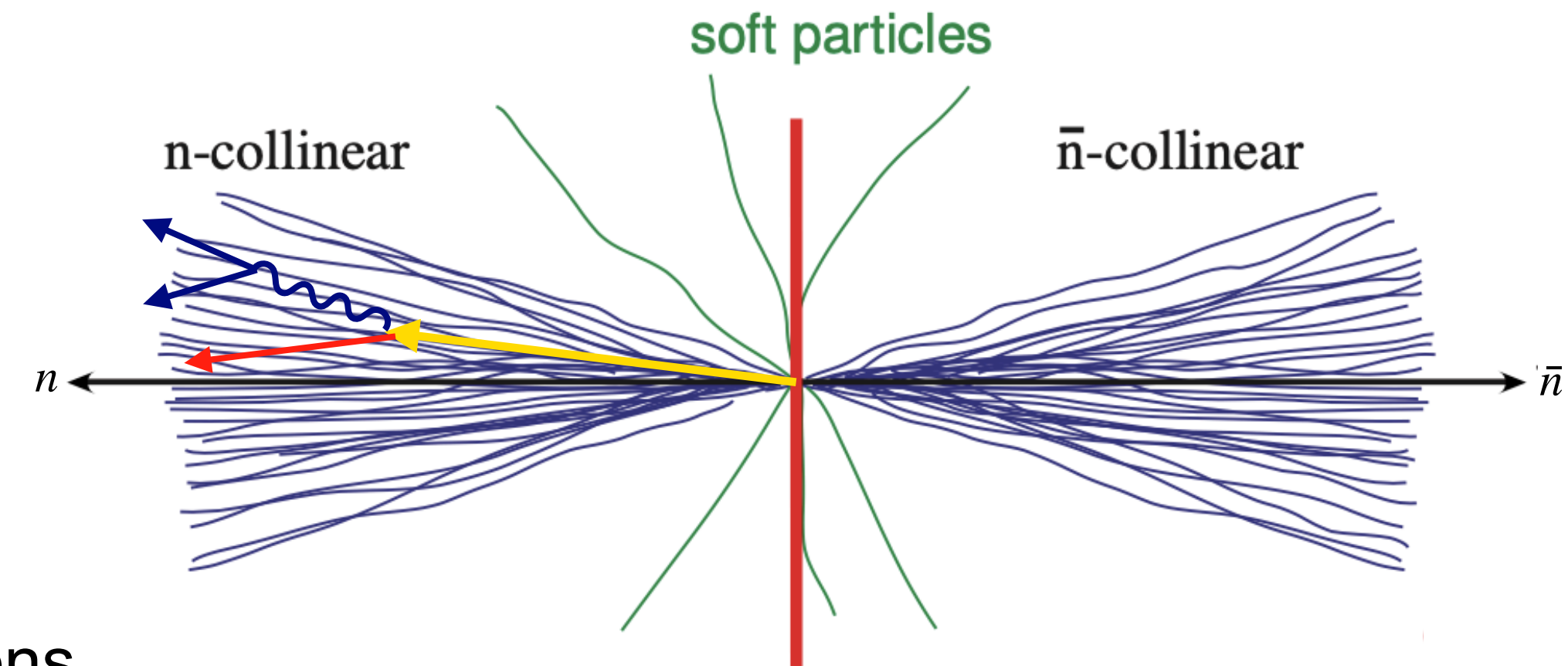
fdk Π Doktoratskolleg
Particles and Interactions

FWF
Der Wissenschaftsfonds.

VDS
VIENNA - DOCTORAL - SCHOOL - PHYSICS

Motivation — Top Quark Physics

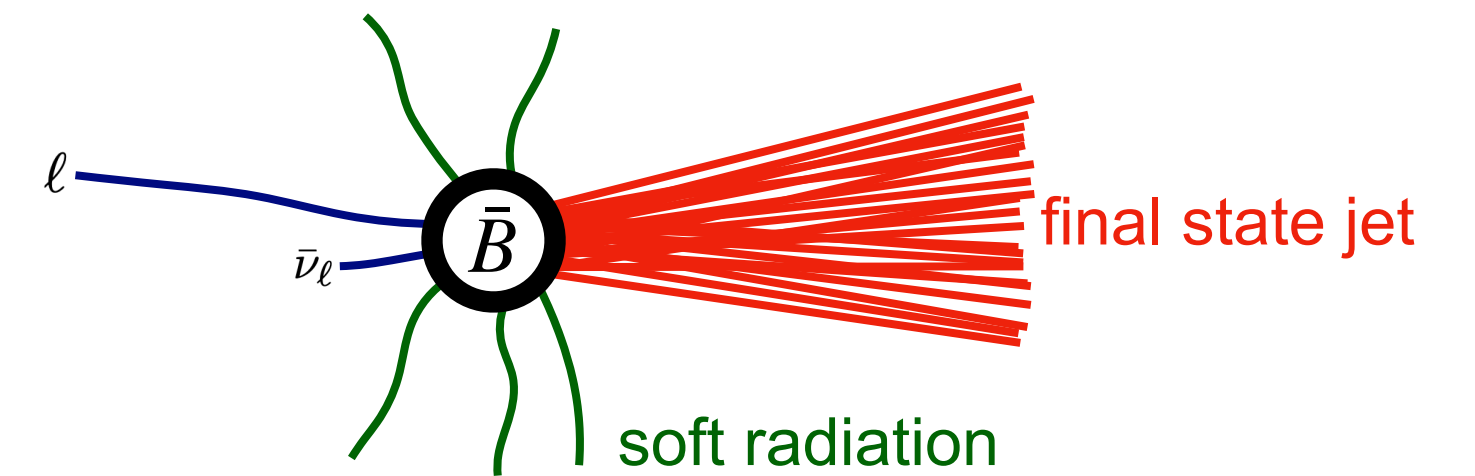
- The **top quark** is the heaviest known elementary particle.
- Due to its large mass the top quark plays a key role in consistency checks of the Standard Model and new-physics searches.
- Studies of top quark production and its decay commonly based on two approaches:
 - **Narrow-width (NW) limit:**
 - Top quark treated as on-shell particle.
 - Factorization of top production and decay dynamics.
 - **Off-shell fixed-order computation:**
 - Accounts for non-resonant, non-factorizable and finite lifetime effects.
 - Start-of-the-art: fixed-order NLO QCD.
- **New approach:** Combine properties of NW limit and off-shell computations.
 - **QCD factorization theorem** for off-shell boosted top quarks (\longrightarrow **SCET, bHQET**)
 - Merge factorization approaches for boosted top production and semileptonic B decays.
 - Top quark state based on a measurement (and not on the concept of a “top particle”).
 - Incorporate resummed QCD corrections for differential top decay observables.
- **Aim:** Analytic control of top mass dependent decay observables.



Soft-Collinear Effective Theory (SCET)

- **SCET**: Used to describe energetic QCD processes where the final state particles have large energies compared to their invariant mass.

- ▶ B meson decays to light particles in kinematic endpoint regions, e.g. $\bar{B} \rightarrow X_u \ell \bar{\nu}_\ell$ ($m_X \ll E_X \sim \frac{m_B}{2}$)
- ▶ Jet production in pp collisions and e^+e^- collisions. ($m_J \ll E_J$)
- ▶ ...



- Momentum modes:

collinear dofs :	$p_n \sim Q(\lambda^2, 1, \lambda)$
	$p_{\bar{n}} \sim Q(1, \lambda^2, \lambda)$
soft dofs :	$p_s \sim Q(\lambda^2, \lambda^2, \lambda^2)$
hard dofs :	$p_h \sim Q(1, 1, 1)$

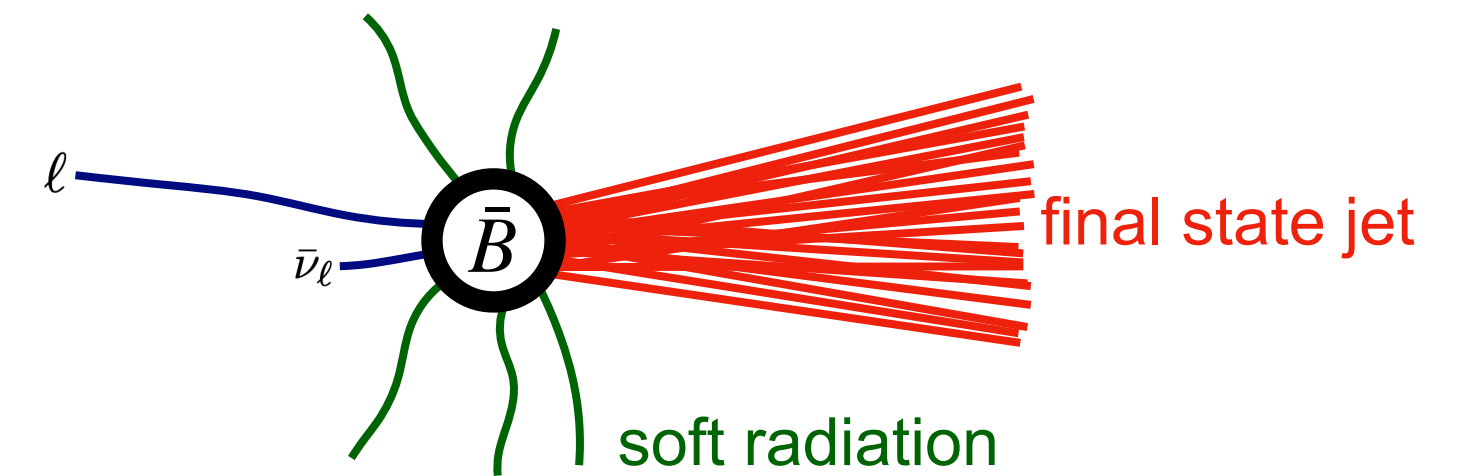
$$p^\mu = p^+ \frac{\bar{n}^\mu}{2} + p^- \frac{n^\mu}{2} + p_\perp^\mu \equiv (p^+, p^-, p_\perp)$$

$$p^2 = p^+ p^- + p_\perp^2$$

Soft-Collinear Effective Theory (SCET)

- **SCET:** Used to describe energetic QCD processes where the final state particles have large energies compared to their invariant mass.

- ▶ B meson decays to light particles in kinematic endpoint regions, e.g. $\bar{B} \rightarrow X_u \ell \bar{\nu}_\ell$ ($m_X \ll E_X \sim \frac{m_B}{2}$)
- ▶ Jet production in pp collisions and e^+e^- collisions. ($m_J \ll E_J$)
- ▶ ...

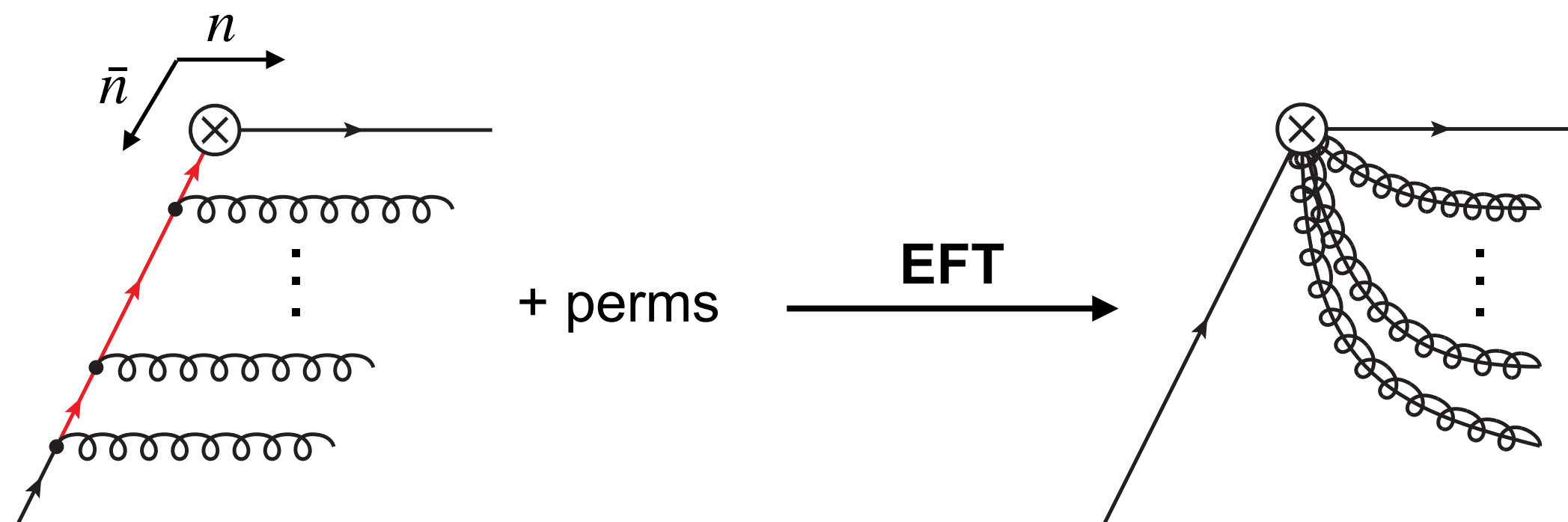


- Leading order collinear quark Lagrangian:

$$\mathcal{L}_n = \bar{\xi}_n \left[i \not{n} \cdot D_s + g \not{n} \cdot A_n + g \not{n} \cdot A_s + i \not{D}_n^\perp W_n^\dagger \frac{1}{\not{n} \cdot \mathcal{P}} W_n i \not{D}_n^\perp \right] \frac{\not{n}}{2} \xi_n$$

collinear Wilson line

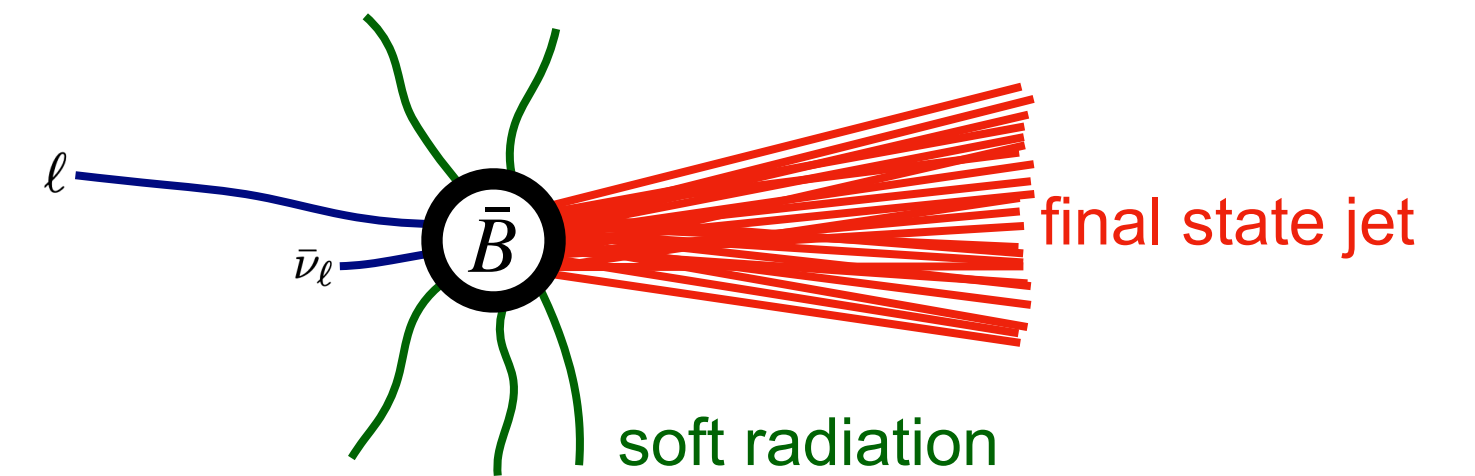
$$W_n(x) = P \exp \left(-ig_s \int_{-\infty}^0 ds \not{n} \cdot A_n(s\bar{n} + x) \right)$$



Soft-Collinear Effective Theory (SCET)

- **SCET:** Used to describe energetic QCD processes where the final state particles have large energies compared to their invariant mass.

- ▶ B meson decays to light particles in kinematic endpoint regions, e.g. $\bar{B} \rightarrow X_u \ell \bar{\nu}_\ell$ ($m_X \ll E_X \sim \frac{m_B}{2}$)
- ▶ Jet production in pp collisions and e^+e^- collisions. ($m_J \ll E_J$)
- ▶ ...



- Leading order collinear quark Lagrangian:

$$\mathcal{L}_n = \bar{\xi}_n \left[\underbrace{in \cdot D_s}_{\text{collinear-soft coupling}} + gn \cdot A_n + gn \cdot A_s + i\not{D}_n^\perp W_n^\dagger \frac{1}{\bar{n} \cdot \mathcal{P}} W_n i\not{D}_n^\perp \right] \frac{\not{n}}{2} \xi_n$$

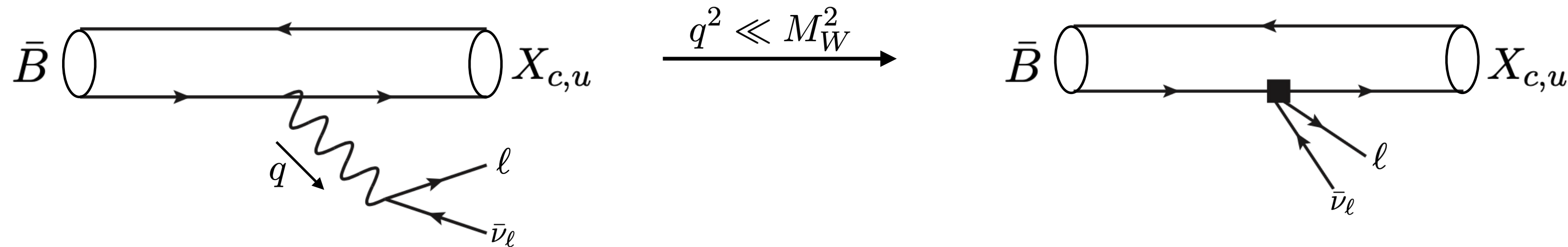
collinear-soft coupling $iD_s^\mu = i\partial^\mu + gA_s^\mu$

$$\begin{aligned} \xi_n &\rightarrow Y_n \xi_n, & W_n &\rightarrow Y_n W_n Y_n^\dagger \\ Y_n(x) &= \bar{P} \exp \left(-ig_s \int_0^\infty ds n \cdot A_s(ns + x) \right) \end{aligned}$$

$$\mathcal{L}_n = \bar{\xi}_n in \cdot \partial_s \frac{\not{n}}{2} \xi_n + \dots$$

Recap: Semileptonic B decays

- Inclusive semileptonic decays $\bar{B} \rightarrow X_c \ell \bar{\nu}_\ell$ and $\bar{B} \rightarrow X_u \ell \bar{\nu}_\ell$ allow extraction of $|V_{cb}|$ and $|V_{ub}|$ from measurements of the decay spectra.



$$\frac{d^3\Gamma}{dE_\ell dq^2 dh^2} \sim G_F^2 L_{\mu\nu} W^{\mu\nu}$$

$$W^{\mu\nu} = \frac{1}{2m_B} \frac{1}{\pi} \text{Im} \left[\langle \bar{B} | i \int d^4x e^{-iq \cdot x} T \{ J^{\dagger\mu}(0) J^\nu(x) \} | \bar{B} \rangle \right], \quad J^\nu = (\bar{c} \gamma^\nu P_L b)$$

- $\bar{B} \rightarrow X_c \ell \bar{\nu}_\ell$: $W^{\mu\nu}$ can be studied by using a **local OPE** within HQET.
 - Non-pert. physics encoded in matrix elements of local operators.
- $\bar{B} \rightarrow X_u \ell \bar{\nu}_\ell$: Cuts on E_ℓ or h^2 needed to eliminate $b \rightarrow c$ background events.
 - Restriction to phase space region of energetic jets with small invariant mass \longrightarrow **SCET**
 - OPE not applicable in this region \longrightarrow need to rely on **factorization tools**.

Factorization for Semileptonic B decays

- Factorized form of differential decay rate in the endpoint region:

Bauer, Fleming, Pirjol, Stewart '01
Bosch, Lange, Neubert, Paz '04

$$\frac{d^3\Gamma}{dE_\ell dq^2 dh^2} \sim \Gamma_0 H(E_\ell, \mu) \int dk^+ J(k^+, \mu) S_{\text{shape}}(m_B - m_b - k^+, \mu)$$

hard function

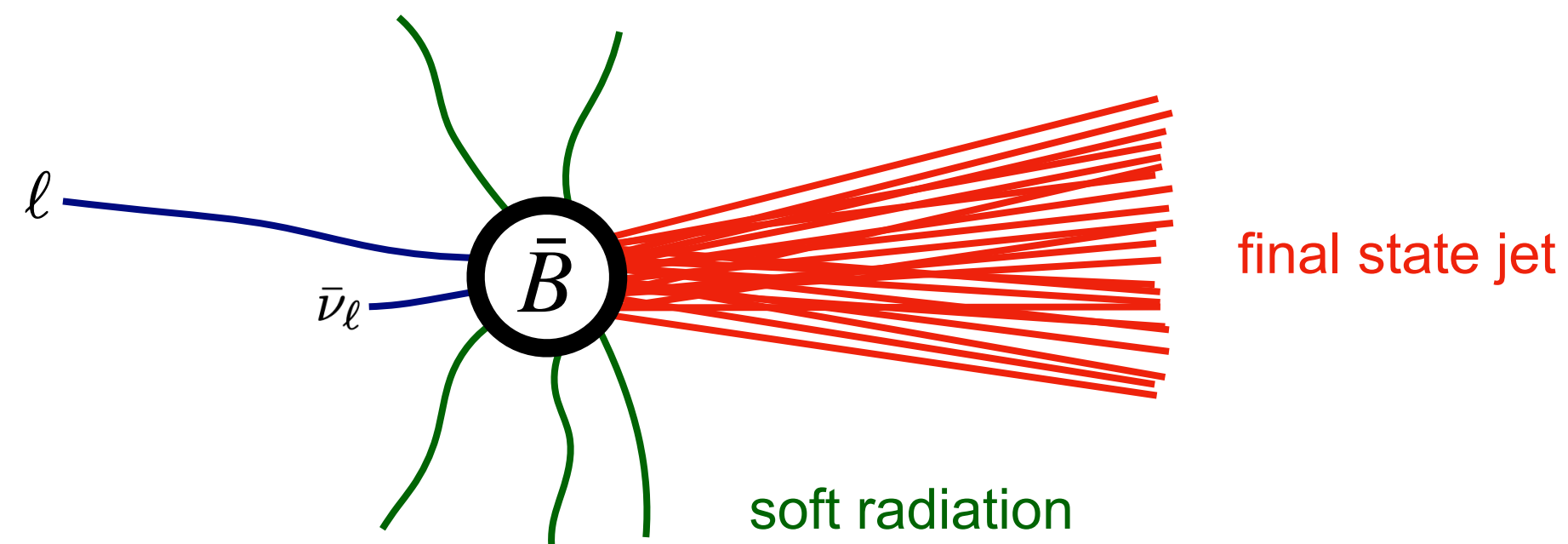
$$H \sim C^2$$

soft function

$$S_{\text{shape}}(k^+) = \langle \bar{B} | \bar{h}_v Y_n \delta(k^+ - i n \cdot \partial) Y_n^\dagger h_v | \bar{B} \rangle$$

jet function

$$J(k) \sim \text{Im} \left[\int d^4x e^{ik \cdot x} \langle 0 | T[\bar{\xi}_n W_n(0) \not{n} W_n^\dagger \xi_n(x)] | 0 \rangle \right]$$



Factorization for Semileptonic B decays

- Factorized form of differential decay rate in the endpoint region:

Bauer, Fleming, Pirjol, Stewart '01
Bosch, Lange, Neubert, Paz '04

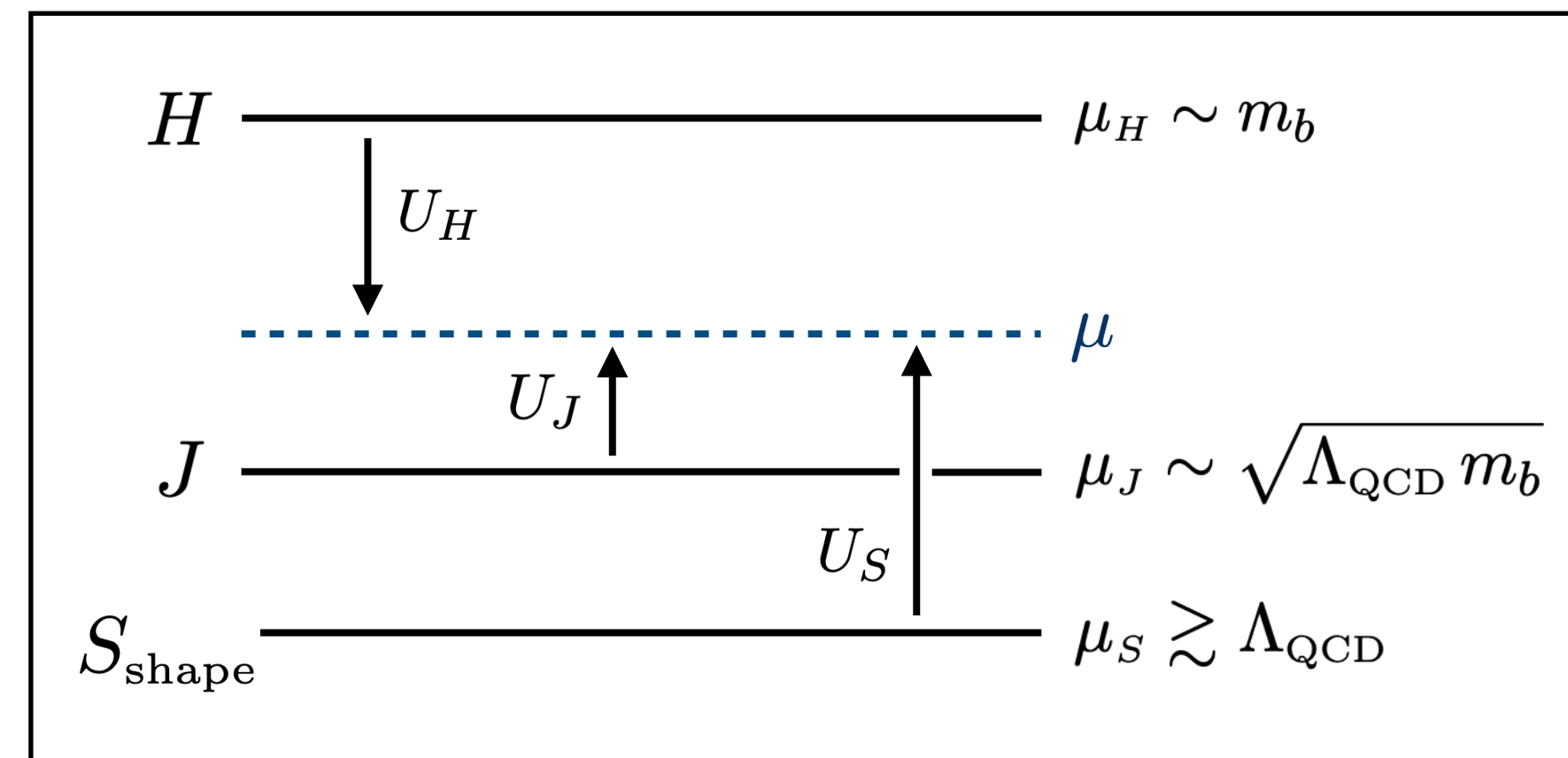
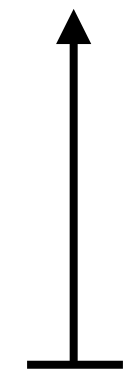
$$\frac{d^3\Gamma}{dE_\ell dq^2 dh^2} \sim \Gamma_0 H(E_\ell, \mu) \int dk^+ J(k^+, \mu) S_{\text{shape}}(m_B - m_b - k^+, \mu)$$

↑ hard function $H \sim C^2$
↑ jet function
 ↑ soft function

$$S_{\text{shape}}(k^+) = \langle \bar{B} | \bar{h}_v Y_n \delta(k^+ - i n \cdot \partial) Y_n^\dagger h_v | \bar{B} \rangle$$

$$J(k) \sim \text{Im} \left[\int d^4x e^{ik \cdot x} \langle 0 | T[\bar{\xi}_n W_n(0) \not{n} W_n^\dagger \xi_n(x)] | 0 \rangle \right]$$

hard function: $\mu_H \sim m_b$
jet function: $\mu_J \sim \sqrt{\Lambda_{\text{QCD}} m_b}$
soft function: $\mu_S \gtrsim \Lambda_{\text{QCD}}$

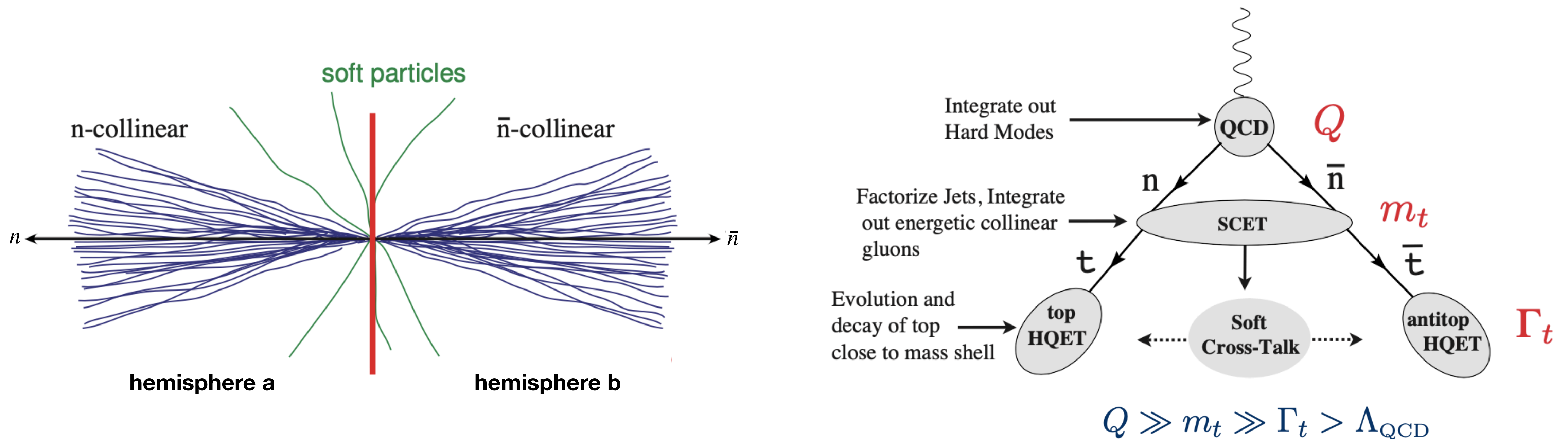


- Each sector depends only on a single physical scale \longrightarrow Large logarithms can be avoided.
- RGEs can be used to evolve the distinct sectors to a common scale.

Boosted top pair production in e^+e^- collisions

- Factorization approach for boosted top jet production in $e^+e^- \rightarrow t\bar{t}$ at c.m. energies $Q \gg m_t$.
 - Dijet region for factorization characterized by $s_{a,b} = M_{a,b}^2 - m_t^2 \ll m_t^2$.
 - Top state defined by measurements of $M_{a,b}$.

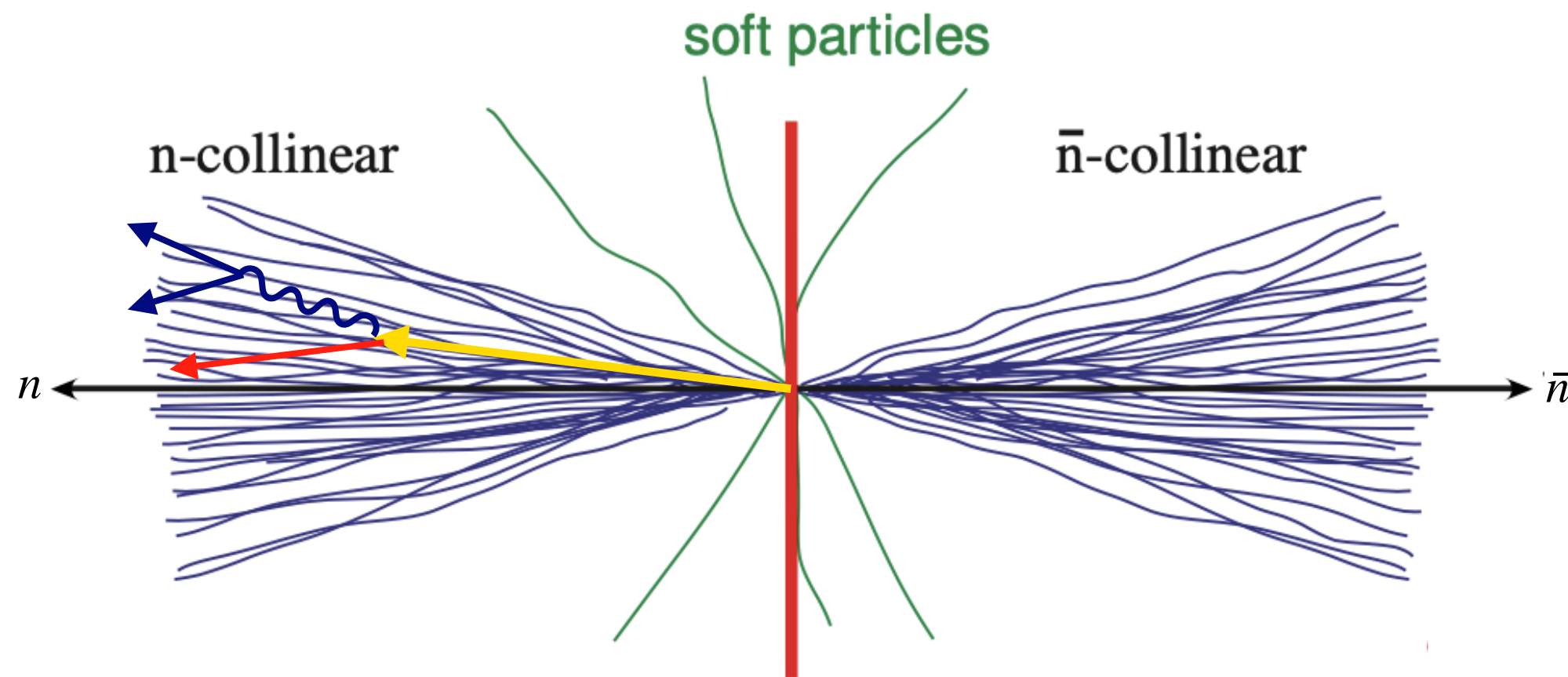
$$\frac{d^2\sigma}{dM_a^2 dM_b^2} \sim \sigma_0 H_Q(Q, \mu) \int dl dl' J_t(s_a - Ql, \mu) J_{\bar{t}}(s_b - Ql', \mu) S(l, l', \mu)$$



[Fleming, Hoang, Mantry, Stewart: Phys.Rev.D77:074010 (2008)]

Boosted top pair production in e^+e^- collisions

$$\frac{d^2\sigma}{dM_a^2 dM_b^2} \sim \sigma_0 H_Q(Q, \mu) \int dl dl' J_t(s_a - Ql, \mu) J_{\bar{t}}(s_b - Ql', \mu) S(l, l', \mu)$$



$$\frac{d^2\sigma}{dM_a^2 dM_b^2} \sim \sigma_0 H_Q [J_t \otimes J_{\bar{t}} \otimes S]$$

$J_t \sim H_m [J_b \otimes S_{ucs}]$

inclusive
 \downarrow
 differential

- Combination of factorization approaches for top quark production in e^+e^- collision with factorization methods known from heavy quark decays.
 - Examination of decay sensitive observables.
 - Treatment of finite lifetime effects and of the dynamics of the top quark decay products.
 - Study of effects at kinematic endpoint regions.
- Important application: Gauge invariant off-shell top quark decay.

Factorization approach for top jet function

- Starting point:
$$J_t \propto \frac{1}{N_c m_t} \sum_X (2\pi)^3 \delta^{(4)}(p - P_X) \text{Tr} \left[\langle 0 | \frac{\not{p}}{4} (W_n^\dagger h_v)(0) | X \rangle \langle X | (\bar{h}_v W_n)(0) | 0 \rangle \right]$$

Insert decay operator
into top jet function.

$$\mathcal{O} = (\bar{b} \gamma^\mu P_L t) (\bar{\nu}_\ell \gamma_\mu P_L \ell)$$

Match QCD current onto
heavy-to-light EFT current.

$$\mathcal{J}^\mu \sim \sum_i C_i (\bar{\chi}_{n'} \Gamma_i^\mu Y_{n'}^\dagger h_v)$$

Factorize momentum modes
into distinct sectors.

$(s_{a,b} = M_{a,b}^2 - m_t^2 \ll m_t^2)$

$$\frac{d^3 J_t(\hat{s}_a)}{dE_\ell dq^2 dh^2} \sim H(E_\ell, \mu) \int dr^+ J_b(r^+, \mu) S_{ucs}(\hat{s}_a - r^+, \mu)$$

$\hat{s}_a \equiv s_a/m_t$

Factorization approach for top jet function

• Starting point:
$$J_t \propto \frac{1}{N_c m_t} \sum_X (2\pi)^3 \delta^{(4)}(p - P_X) \text{Tr} \left[\langle 0 | \frac{\not{p}}{4} (W_n^\dagger h_v)(0) | X \rangle \langle X | (\bar{h}_v W_n)(0) | 0 \rangle \right]$$

Insert decay operator
into top jet function.

$$\mathcal{O} = (\bar{b} \gamma^\mu P_L t) (\bar{\nu}_\ell \gamma_\mu P_L \ell)$$

Match QCD current onto
heavy-to-light EFT current.

$$\mathcal{J}^\mu \sim \sum_i C_i (\bar{\chi}_{n'} \Gamma_i^\mu Y_{n'}^\dagger h_v)$$

Factorize momentum modes
into distinct sectors.

$$(s_{a,b} = M_{a,b}^2 - m_t^2 \ll m_t^2)$$

$$\frac{d^3 J_t(\hat{s}_a)}{dE_\ell dq^2 dh^2} \sim H(E_\ell, \mu) \int dr^+ J_b(r^+, \mu) S_{ucs}(\hat{s}_a - r^+, \mu)$$

$\hat{s}_a \equiv s_a/m_t$

Known from semi-leptonic B
decays

Factorization approach for top jet function

• Starting point:
$$J_t \propto \frac{1}{N_c m_t} \sum_X (2\pi)^3 \delta^{(4)}(p - P_X) \text{Tr} \left[\langle 0 | \frac{\not{p}}{4} (W_n^\dagger h_v)(0) | X \rangle \langle X | (\bar{h}_v W_n)(0) | 0 \rangle \right]$$

Insert decay operator into top jet function.

$$\mathcal{O} = (\bar{b} \gamma^\mu P_L t) (\bar{\nu}_\ell \gamma_\mu P_L \ell)$$

Match QCD current onto heavy-to-light EFT current.

$$\mathcal{J}^\mu \sim \sum_i C_i (\bar{\chi}_{n'} \Gamma_i^\mu Y_{n'}^\dagger h_v)$$

Factorize momentum modes into distinct sectors.

$(s_{a,b} = M_{a,b}^2 - m_t^2 \ll m_t^2)$

$$\frac{d^3 J_t(\hat{s}_a)}{dE_\ell dq^2 dh^2} \sim H(E_\ell, \mu) \int dr^+ J_b(r^+, \mu) S_{ucs}(\hat{s}_a - r^+, \mu)$$

$\hat{s}_a \equiv s_a/m_t$

“bHQET ultracollinear-soft function”

$$S_{ucs}(r^+) \propto \int d^4 y_1 \int d^4 y_2 \sum_{X_{uc}} \delta^{(4)}(k_{uc}^+ - r^+) \times \text{Tr} \left[\langle 0 | (\bar{h}_v Y_{n'}) (0) (W_n^\dagger h_v)(y_2) | X_{uc} \rangle \langle X_{uc} | (\bar{h}_v W_n)(y_1) (Y_{n'}^\dagger h_v)(0) | 0 \rangle \right]$$

- New ingredient
- Can be computed perturbatively.
- Top width acts as infrared cut-off.
- Describes Fermi motion of the decaying top in the measured state with mass M_a .

Factorization approach for top jet function

• Starting point:
$$J_t \propto \frac{1}{N_c m_t} \sum_X (2\pi)^3 \delta^{(4)}(p - P_X) \text{Tr} \left[\langle 0 | \frac{\not{p}}{4} (W_n^\dagger h_v)(0) | X \rangle \langle X | (\bar{h}_v W_n)(0) | 0 \rangle \right]$$

Insert decay operator into top jet function.

$$\mathcal{O} = (\bar{b} \gamma^\mu P_L t) (\bar{\nu}_\ell \gamma_\mu P_L \ell)$$

Match QCD current onto heavy-to-light EFT current.

$$\mathcal{J}^\mu \sim \sum_i C_i (\bar{\chi}_{n'} \Gamma_i^\mu Y_{n'}^\dagger h_v)$$

Factorize momentum modes into distinct sectors.

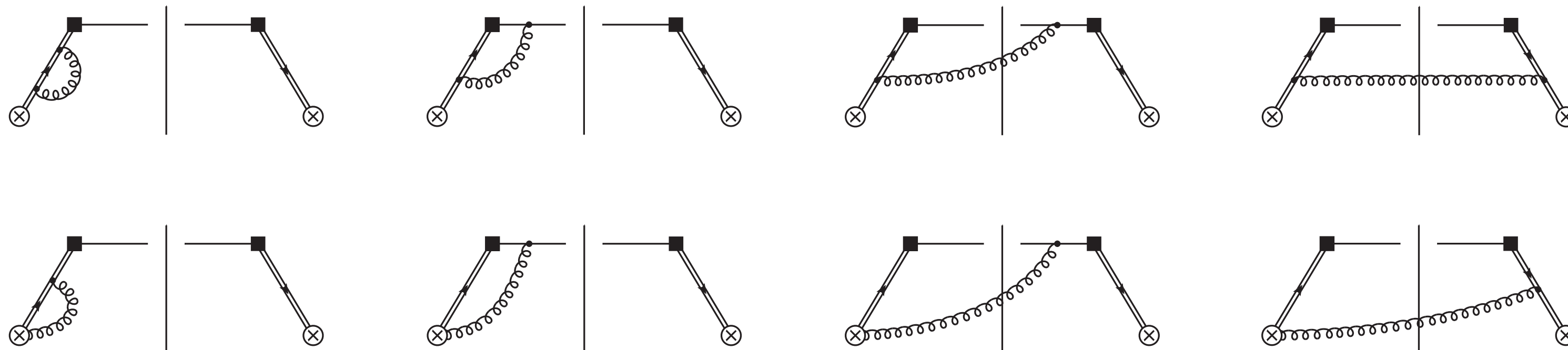
$(s_{a,b} = M_{a,b}^2 - m_t^2 \ll m_t^2)$

$$\frac{d^3 J_t(\hat{s}_a)}{dE_\ell dq^2 dh^2} \sim H(E_\ell, \mu) \int dr^+ J_b(r^+, \mu) S_{ucs}(\hat{s}_a - r^+, \mu)$$

$\hat{s}_a \equiv s_a/m_t$

“bHQET ultracollinear-soft function”

- New ingredient
- Can be computed perturbatively.
- Top width acts as infrared cut-off.
- Describes Fermi motion of the decaying top in the measured state with mass M_a .



Summary & Outlook

Summary:

- Factorization is an important theoretical tool for calculations of processes involving hadrons.
 - Allows to factorize cross sections and decay rates into different parts that can either be calculated perturbatively or determined from data.
 - Sums large logarithms arising in perturbative calculations by means of RG equations.
- We aim to combine properties of the NW limit and off-shell computations for decaying top quark studies.
 - Merge existing factorization theorems for off-shell top production in e^+e^- collisions and for heavy quark decays.
 - Top state defined by measurements (and not by NW limit).
 - Leads to a gauge-invariant jet function for boosted top quarks including off-shell effects (up to leading order in m_t/Q).
 - Allows to study decay sensitive observables beyond the commonly employed NW limit.

Outlook:

- We want to study the NLO QCD corrections for off-shell and boosted top decays in the endpoint region.
 - Analysis of our factorized approach with NLL resummation.
- Comparison with predictions from Monte-Carlo event generators.