

# **Link soliton in model with $U(1)_{B-L}$ and $U(1)_{PQ}$ symmetries**

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**Yu Hamada (KEK)**

**arXiv: 2303.XXXXX (work in progress)**

**w/ M. Eto (Yamagata U.) and M. Nitta (Keio U.)**



22nd Feb. 2023 @ Kagoshima workshop

$U(1)_{global} \times U(1)_{gauge}$  を自発的に破る模型では  
変態的なソリトンができます

$$\begin{array}{cc} U(1)_{PQ} & U(1)_{B-L} \\ \uparrow & \uparrow \end{array}$$

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# Introduction

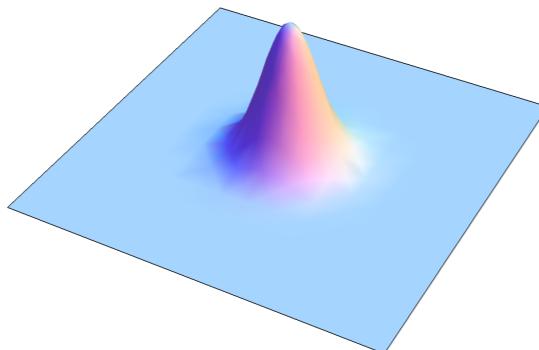
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# Soliton

- (素)粒子：真空まわりの場のゆらぎ



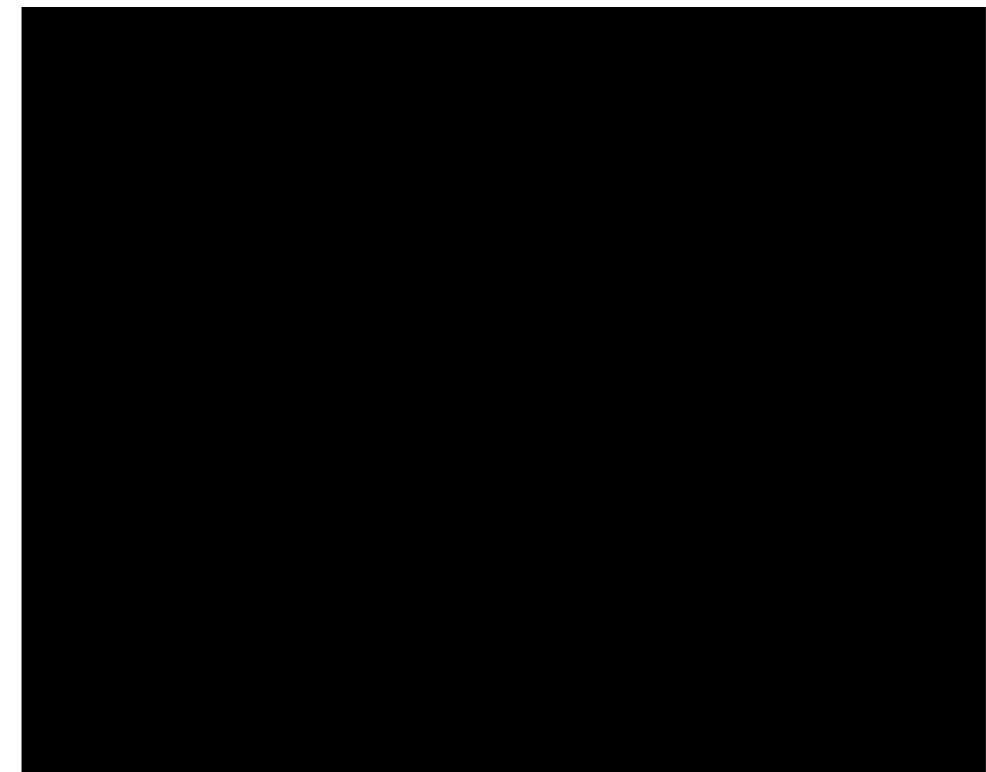
- ソリトン：素粒子でない古典的励起（エネルギーの“カタマリ”）



津波



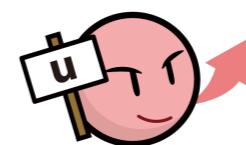
wikipedia  
“神奈川沖浪裏”



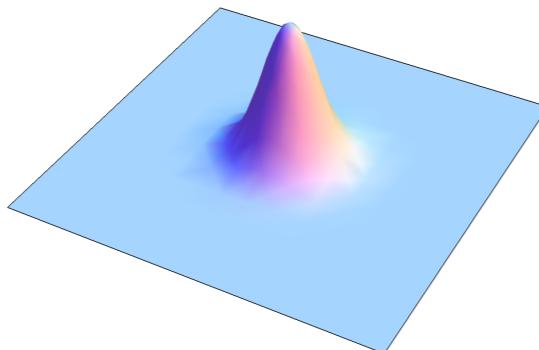
“Collision of KdV solitons”  
(from YouTube)

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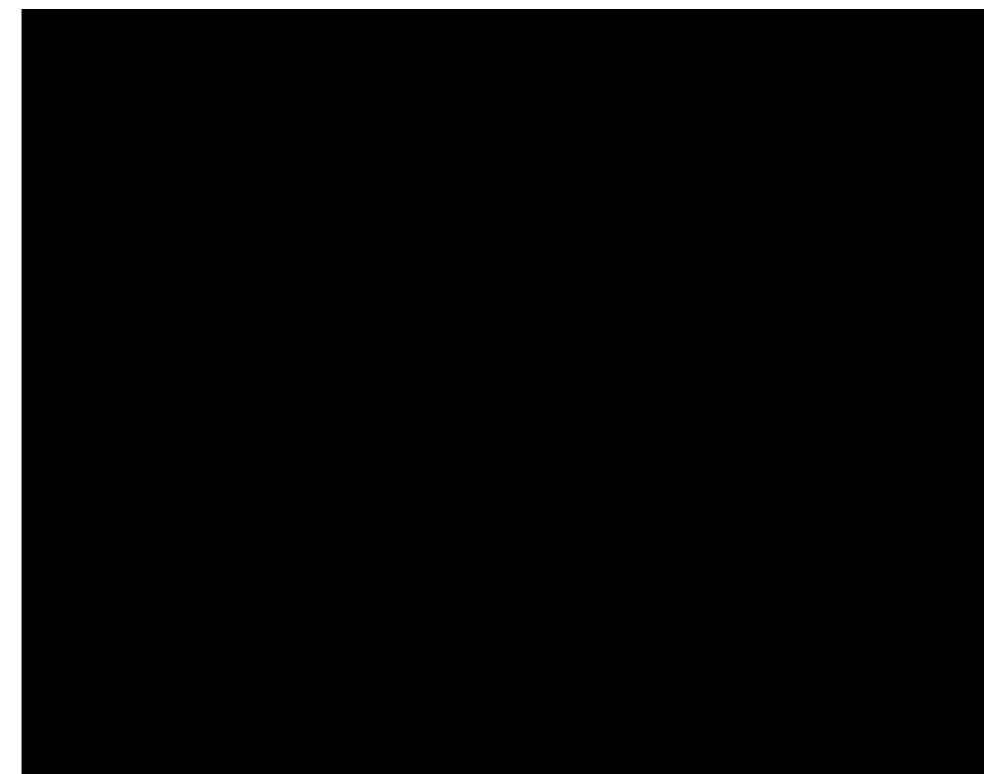
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# Eg.) Abrikosov-Nielsen-Olesen string

[Abrikosov '58]

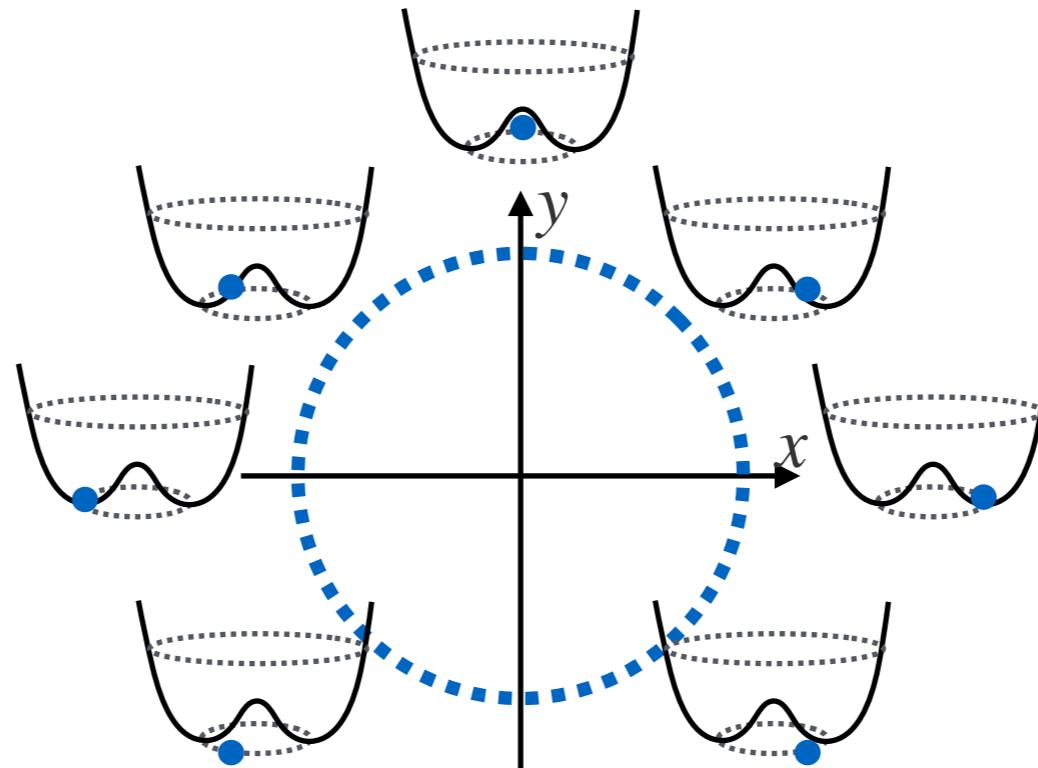
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- 3+1 D Abelian-Higgs model

$$\langle \phi \rangle = v \rightarrow \cancel{U(1)}$$

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + |D_\mu\phi|^2 + m^2|\phi|^2 - \lambda|\phi|^4$$

- z軸方向に一様な配位を仮定
- xy平面上で、各点で異なる真空の位相を選んでも良い



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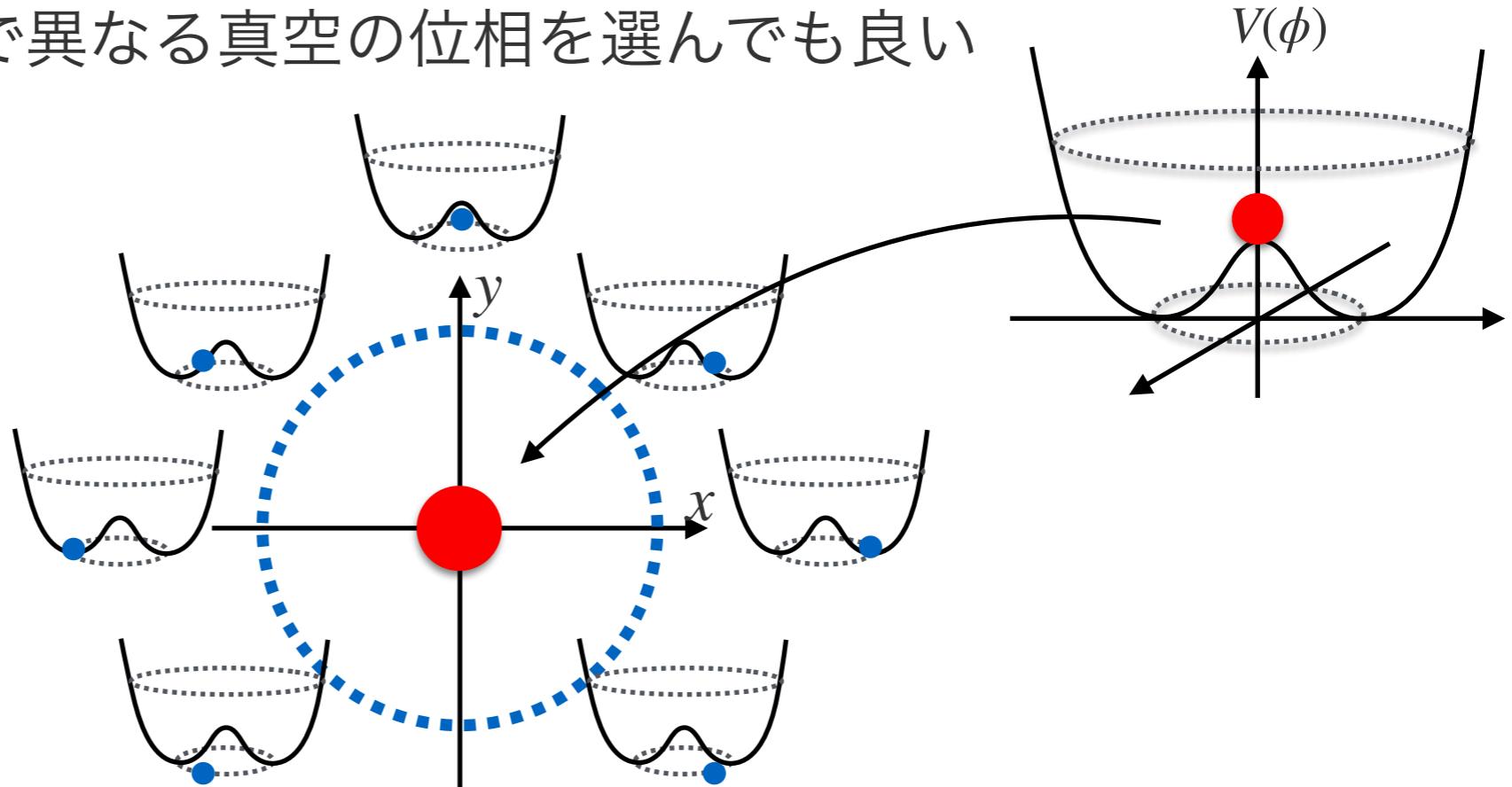
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→ 中心は必ず potential の頂点にいる → excitation (soliton)

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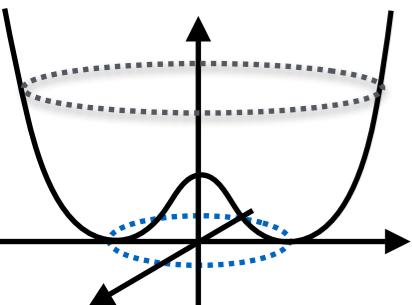
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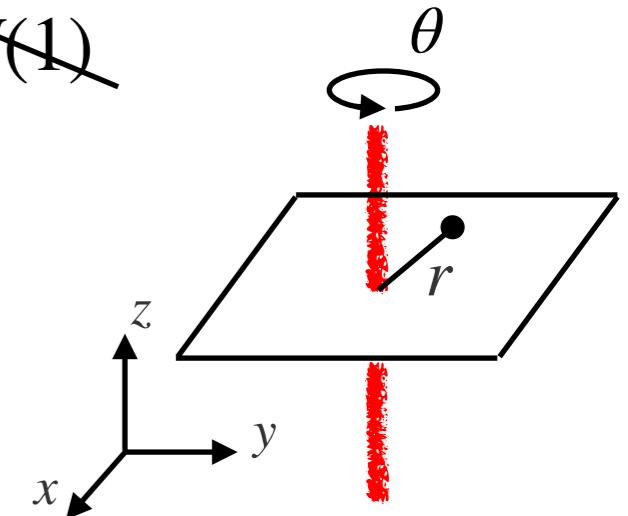
- **Field configuration:**



$$\phi(x) = v f(r) \underline{e^{i\theta}}$$

$$\vec{A}(x) = g^{-1} a(r) \vec{e}_\theta$$

$\phi$ 's phase has winding # = 1



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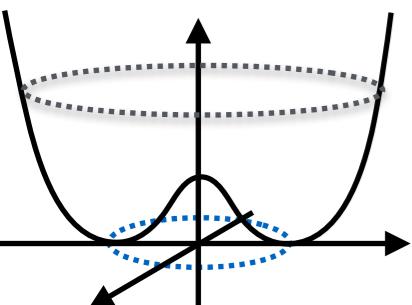
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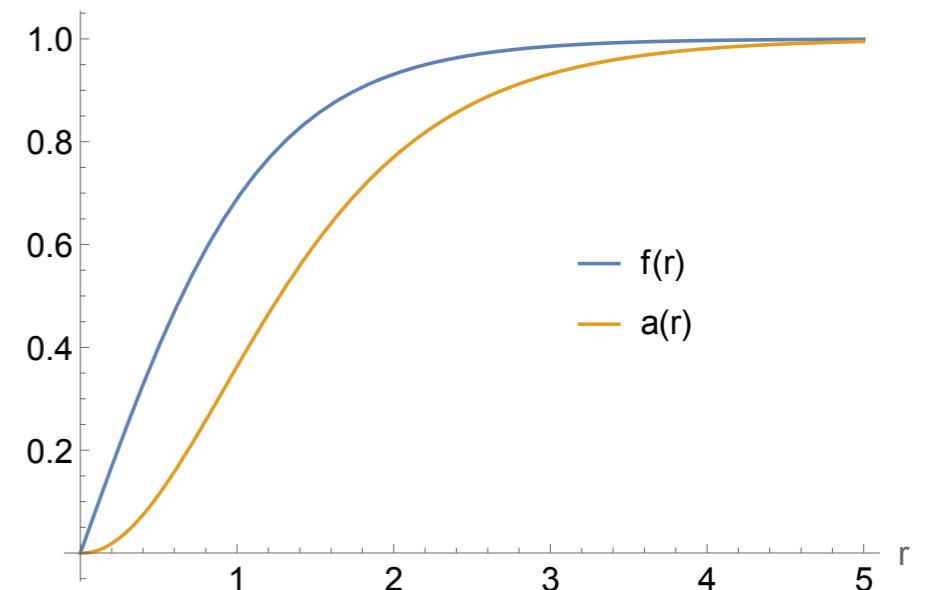
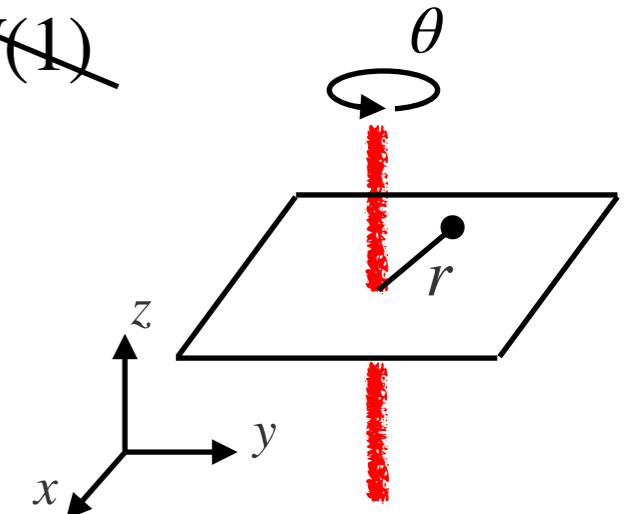
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- solving classical EOMs for  $f(r)$  and  $a(r)$ :

$$f'' + \frac{1}{r}f' - \frac{(1-a)^2}{r^2}f - \frac{1}{2}\frac{\partial V}{\partial f} = 0$$

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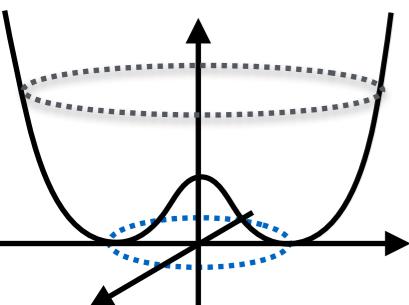
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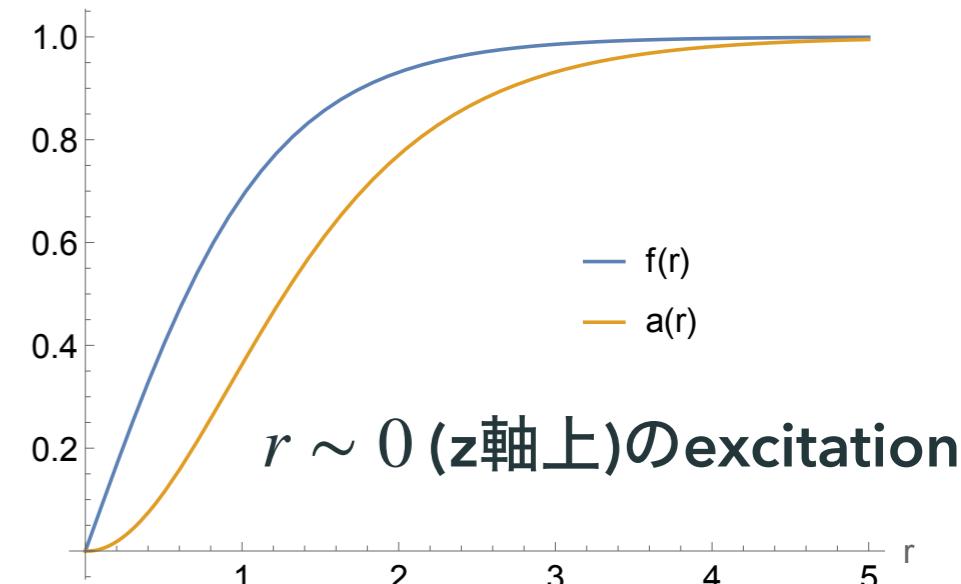
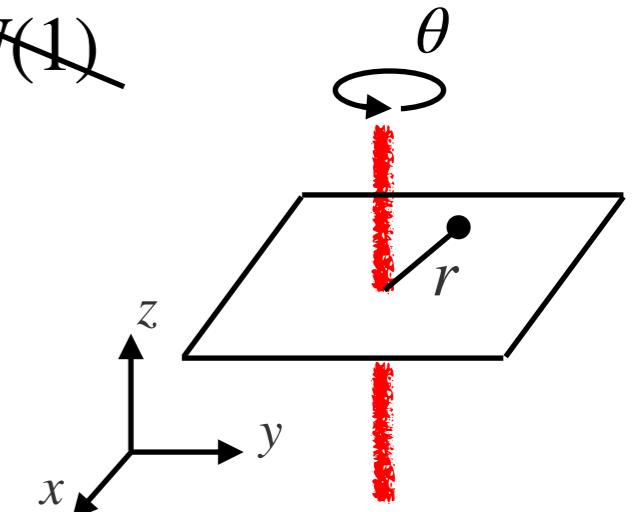
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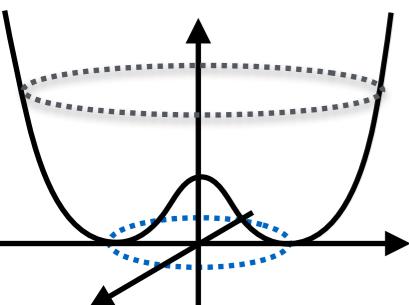
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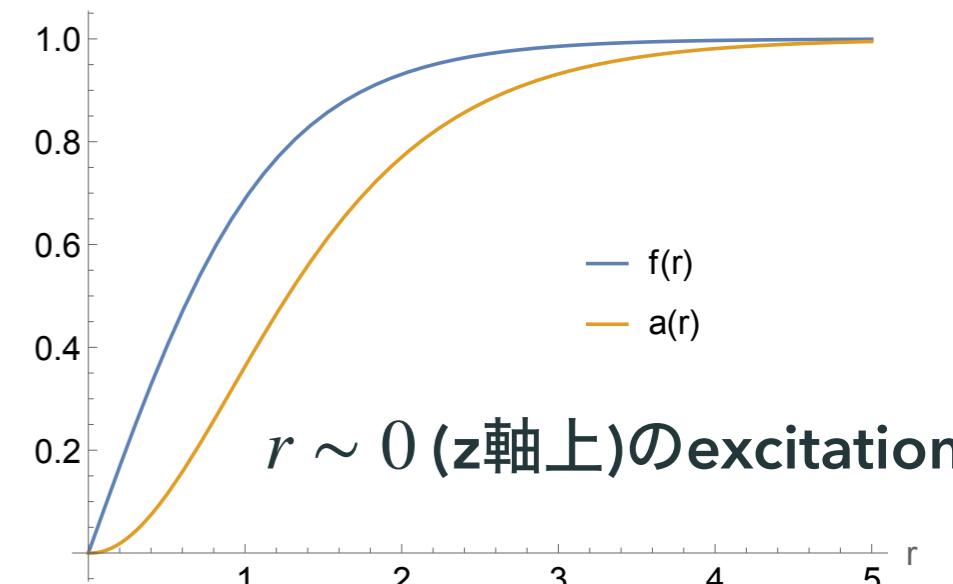
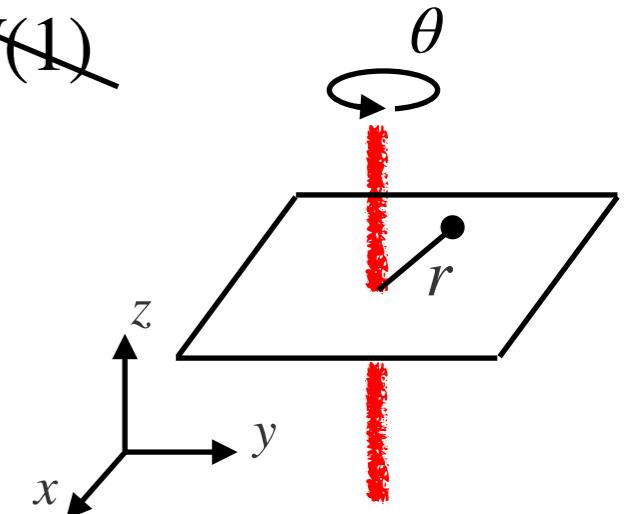
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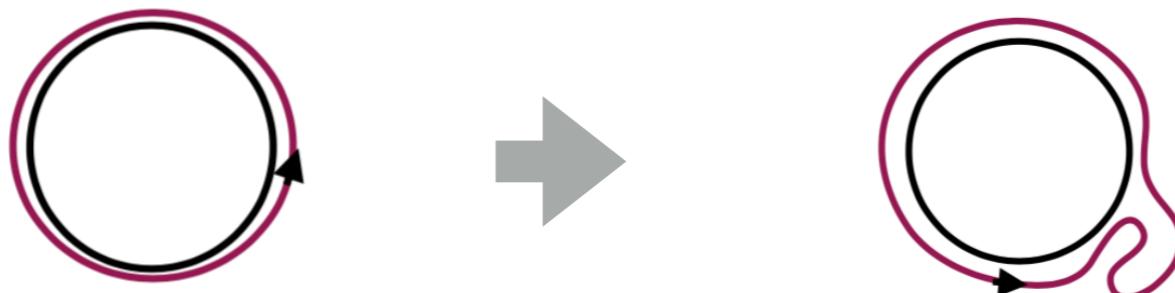
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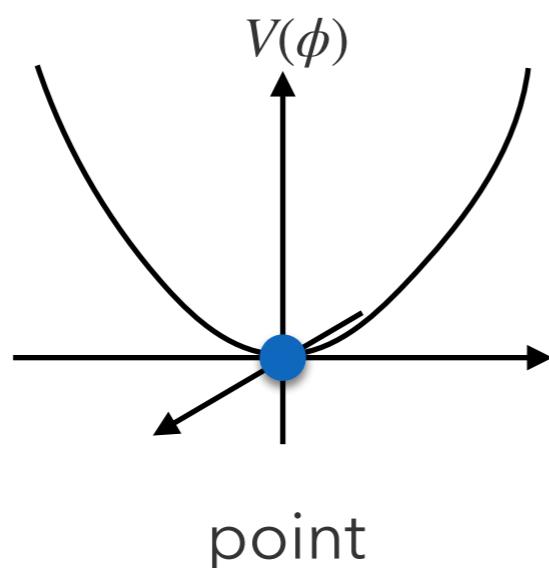
Quantized magnetic flux:  $\int d^2x B = \oint d\vec{l} \cdot \vec{A} = 2\pi/g$

# Topology of vacuum

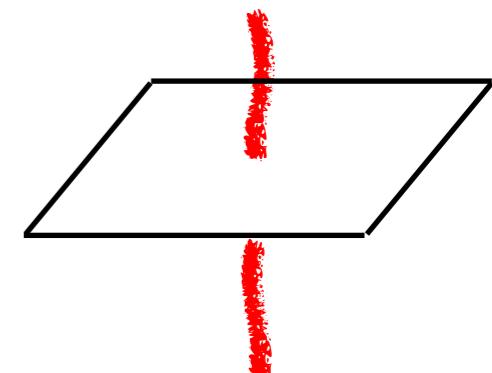
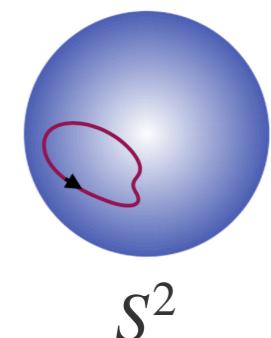
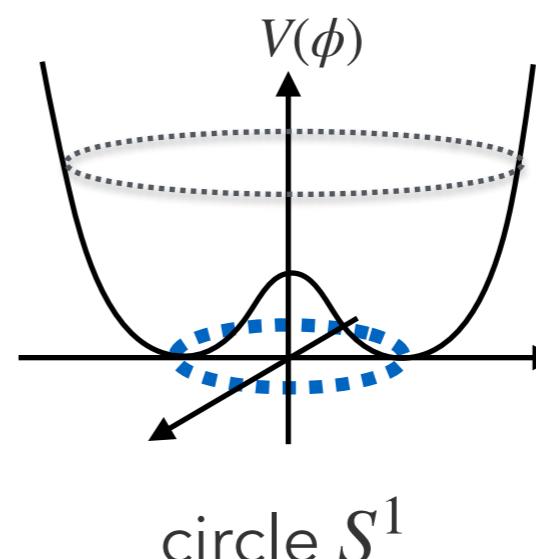
- このwindingはいかなる連續変形でも取り除けない → 安定



- 真空が“circle structure”であることが安定性を保証

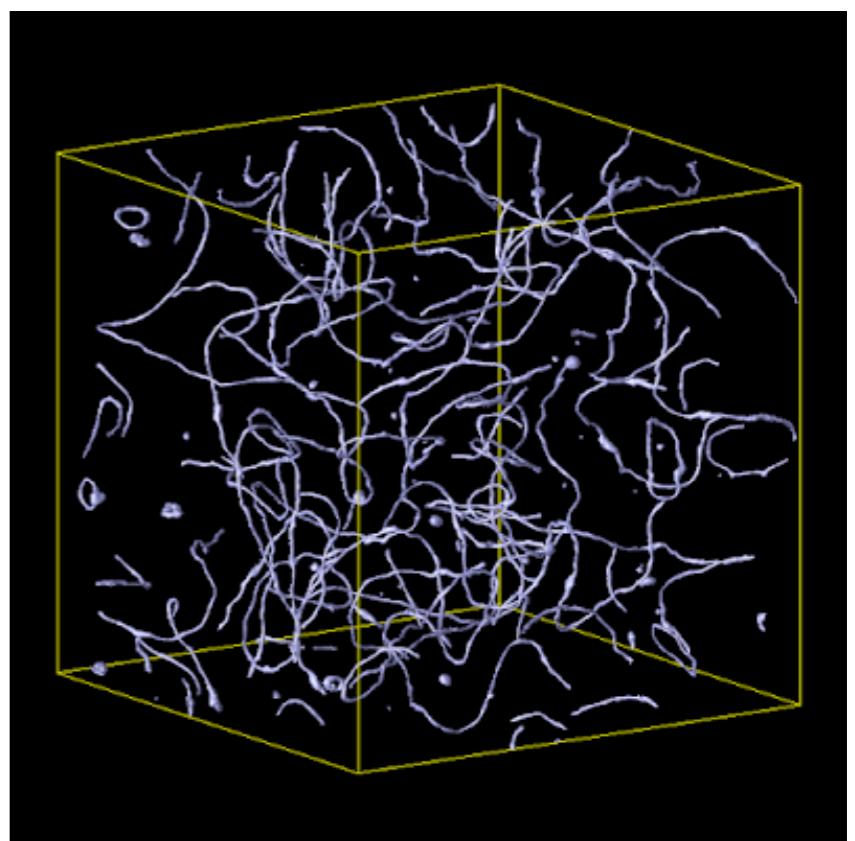
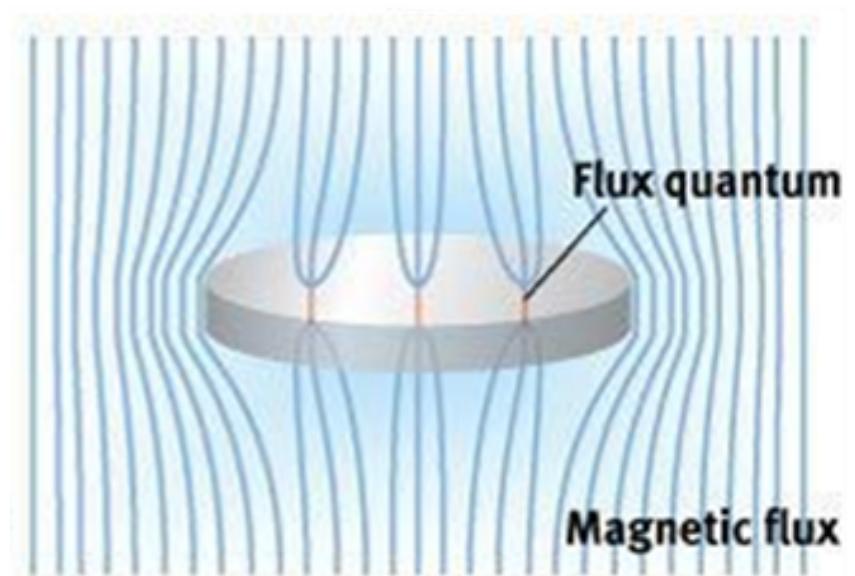


→このようなsolitonをvortex stringと呼ぶ



# Vortex string in many systems

- Magnetic flux tube in superconductor
  - characterize phases of supercond.
- Vortex string in the universe: **Cosmic string**
  - CMB observation
  - Gravitational wave
  - **strong evidence of new physics**,  
but haven't yet been discovered.
  - 昔からよく調べられている  
('90sに流行った->最近また流行ってる?)



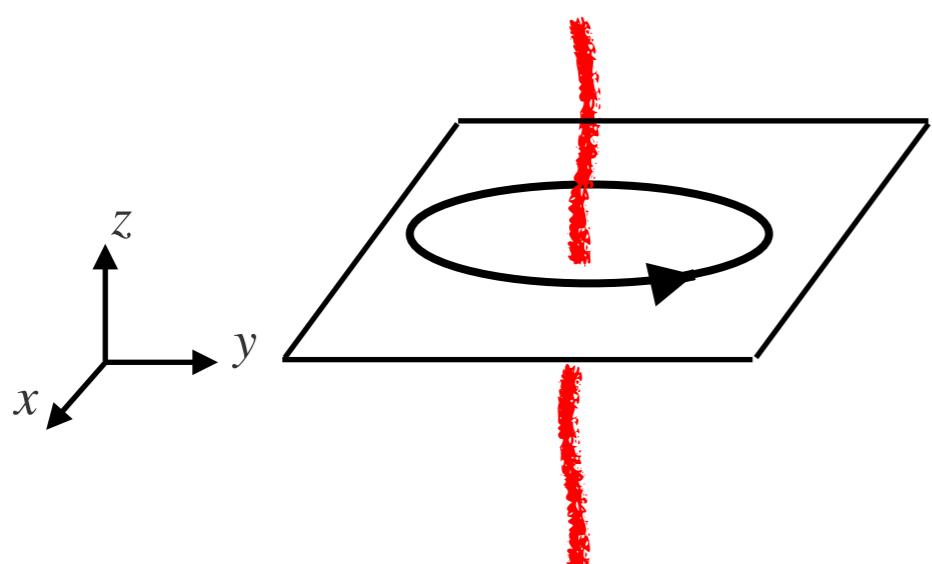
# Global vs Local strings

- SSB of **gauged**  $U(1)$  sym  $\rightarrow$  **local** vortex string

$\rightarrow$  **quantized magnetic flux:**  $\int d^2x B = 2\pi/g$

- SSB of **global**  $U(1)$  sym  $\rightarrow$  **global** vortex string

$\rightarrow$  **w/o magnetic flux**



$$\phi(x) = v f(r) e^{i\theta}$$

string周りでNG boson phaseが  
0から $2\pi$ に変化

$$\text{EOM: } f'' + \frac{1}{r}f' - \frac{(1-\cancel{\alpha})^2}{r^2}f - \frac{1}{2}\frac{\partial V}{\partial f} = 0$$

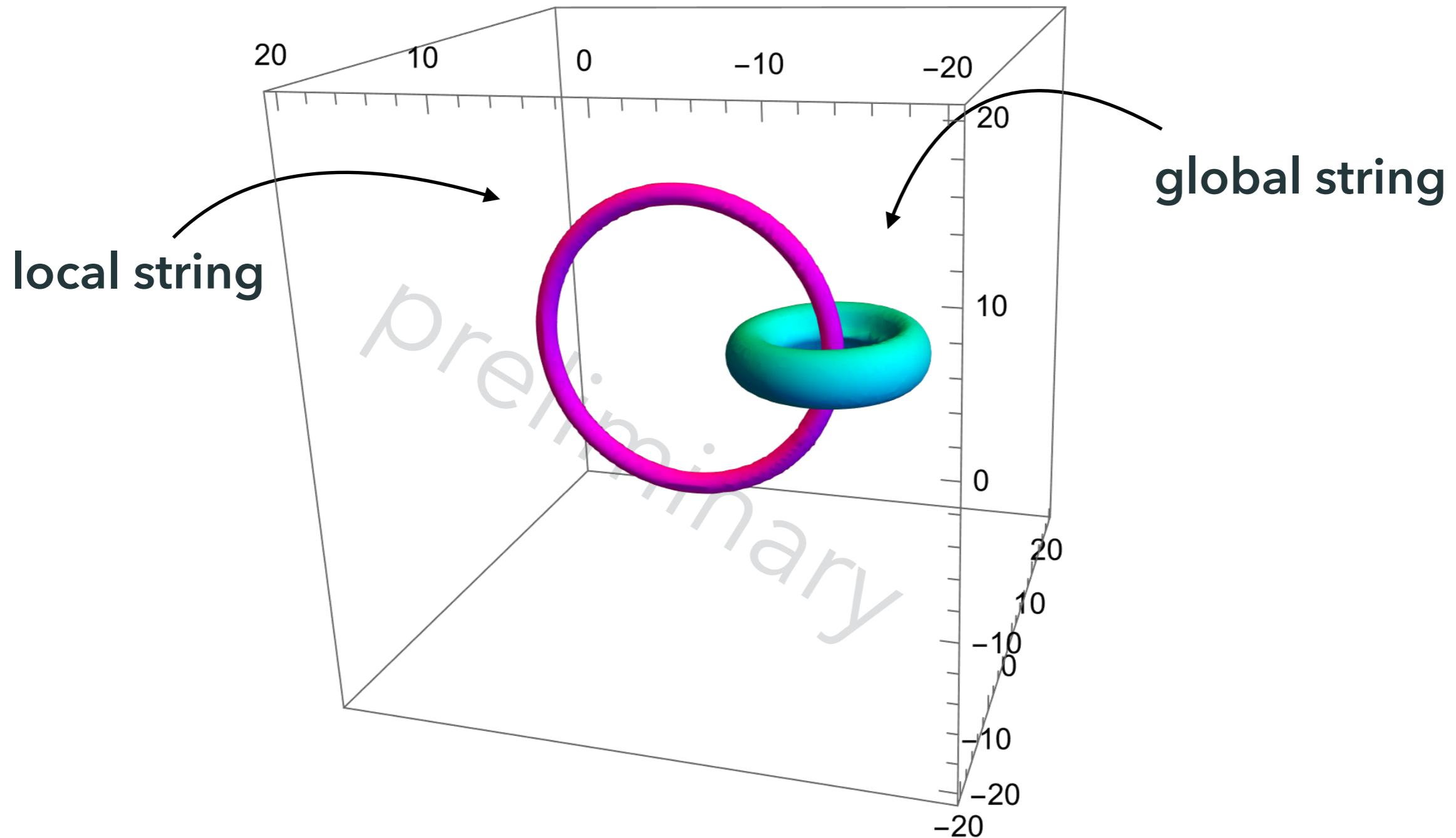
# Global vs Local strings



# Global vs Local strings



# Link soliton



link soliton made of local & global strings!

# Link soliton

- Massage of this talk:

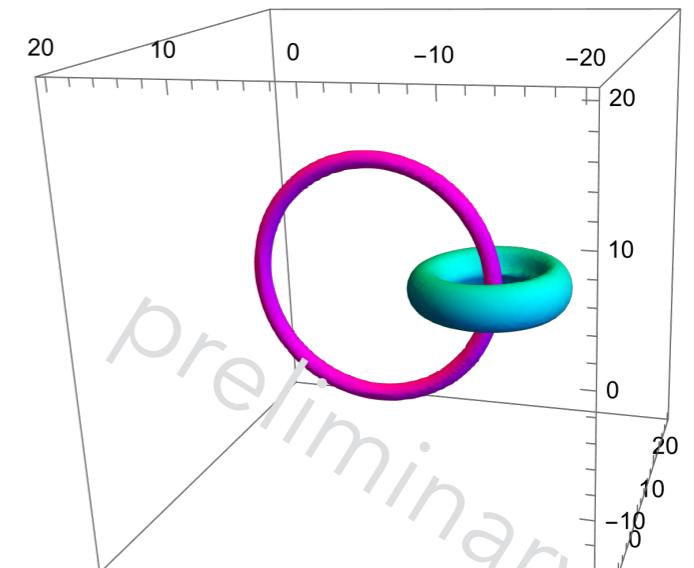
$U(1)_{global} \times U(1)_{gauge}$  を自発的に破る模型では link solitonが存在する！

- Key: Chern-Simons coupling  $\frac{c}{16\pi^2} \int d^4x aF\tilde{F}$
- motivativeなセットアップ:

$$\begin{cases} U(1)_{global} = U(1)_{PQ} & \text{QCD axion} \rightarrow \text{strong CP \& DM} \\ U(1)_{gauge} = U(1)_{B-L} & \text{RH}\nu \rightarrow \text{Type-I seesaw, GUT} \end{cases}$$

axion stringとB-L stringからなるlink

→バリオジェネシス、原子重力波での検出、  
GUTとの関係？ etc.



# Plan of talk

- Introduction
- Link soliton
- Application to Baryogenesis
- Summary

# Link soliton

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# The model

3+1D theory:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |\partial_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2)$$

$$D_\mu \phi_1 = (\partial_\mu - igA_\mu) \phi_1$$

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

- Symmetries:

$$U(1)_{gauge} : \phi_1 \rightarrow e^{i\theta_1} \phi_1 \quad U(1)_{global} : \phi_2 \rightarrow e^{i\theta_2} \phi_2$$

- For  $\kappa > 0$  &  $\lambda > 0$ , both symmetries are broken at the vacuum:

$$\langle \phi_1 \rangle = v_1, \langle \phi_2 \rangle = v_2$$

→ local string ( $\phi_1$  string) & global string ( $\phi_2$  string)

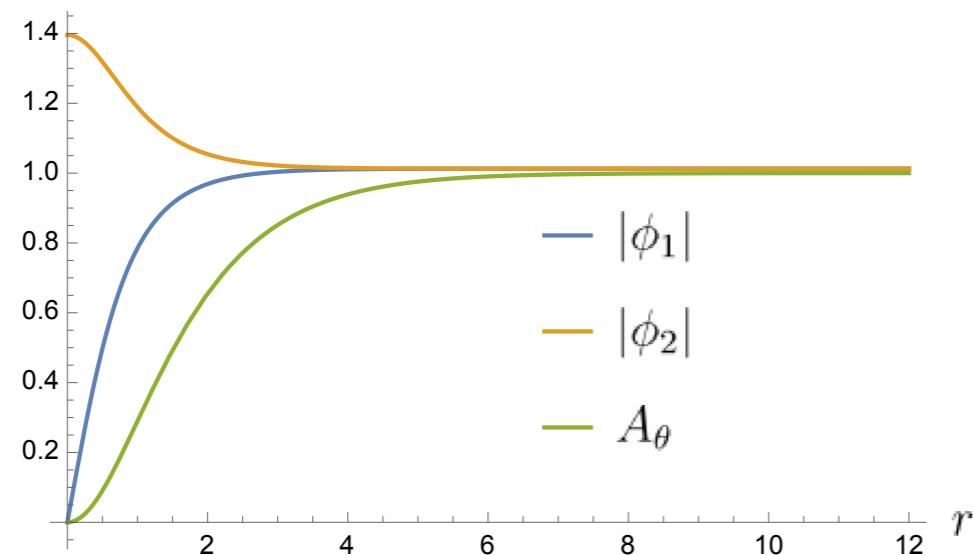
# $\phi_1$ string & $\phi_2$ string

- Field configuration for  $\phi_1$  string (local):

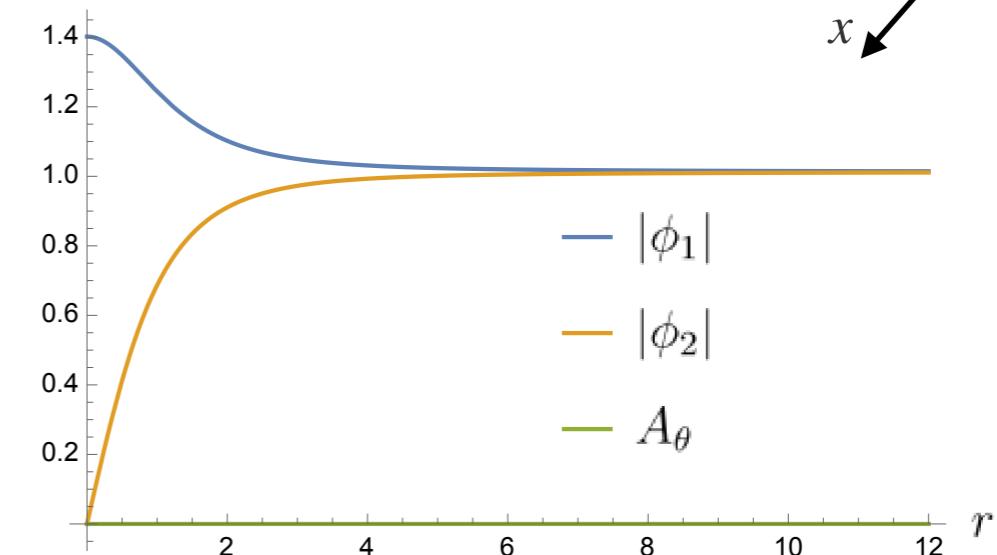
$$\phi_1(x) = v_1 f_1(r) \underline{e^{i\theta}} \quad \phi_2(x) = v_2 f_2(r) \quad \vec{A}(x) = g^{-1} a(r) \vec{e}_\theta$$

- Field configuration for  $\phi_2$  string (global):

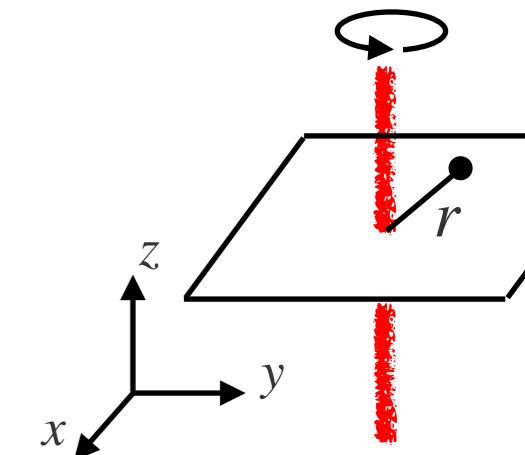
$$\phi_1(x) = v_1 h_1(r) \quad \phi_2(x) = v_2 h_2(r) \underline{e^{i\theta}} \quad A_\mu(x) = 0$$



$\phi_1$  string



$\phi_2$  string



# Chern-Simons coupling

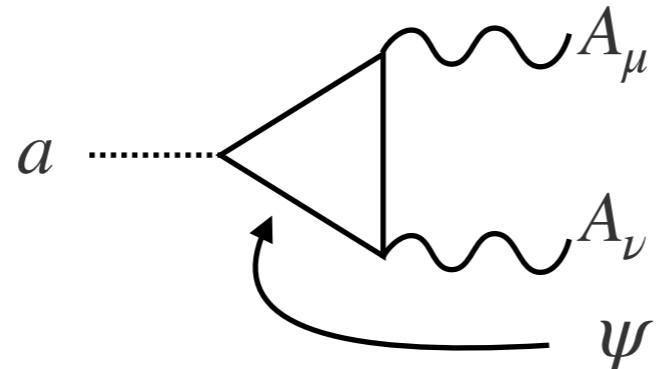
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- At the broken phase, CS coupling is induced by triangle anomaly.



dependent on matter sector

$$\Rightarrow c = \sum_f Q_{global}^f (Q_{gauge}^f)^2$$

taken as free parameter

Chern-Simons coupling

$$+\frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$a \equiv -i \arg(\phi_2)$$

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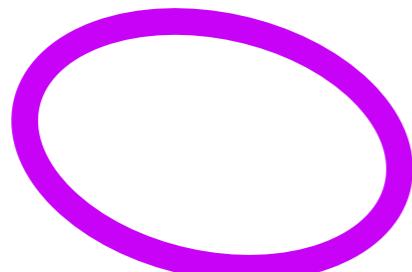
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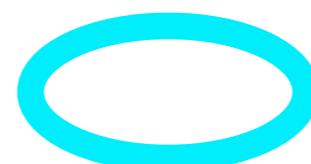
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- CS couplingは単独のstringには効かない

$\phi_1$  string



$\phi_2$  string



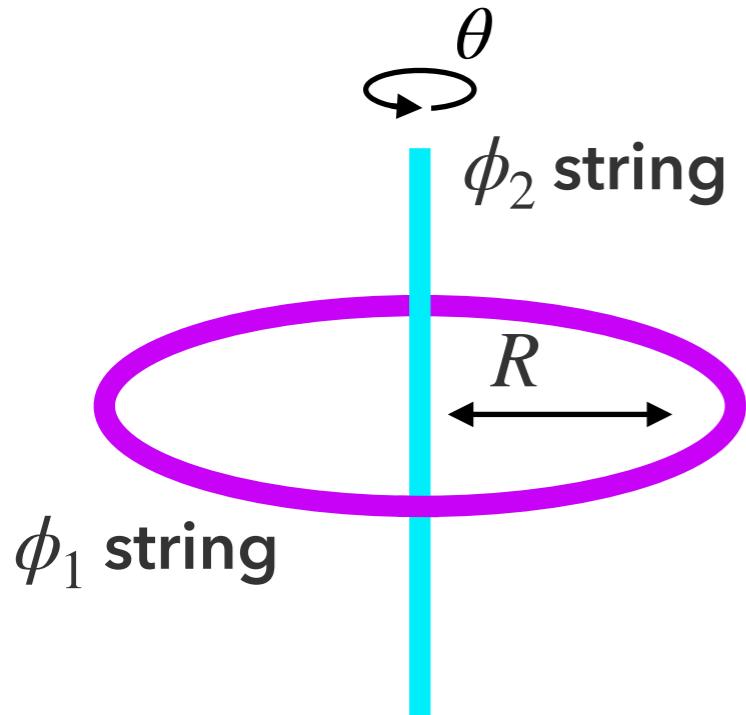
→ こういうループは縮んで消える

# Charged string

Rewriting CS coupling:

$$\frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu} \longrightarrow -\frac{c}{16\pi^2} (\partial_i a) A_0 B^i$$

linkしてるとき、 $\partial_i a$  と  $B_i$  は同じ向き  $\Rightarrow (\partial_i a) A_0 B^i = \frac{1}{R} A_0 |\vec{B}|$

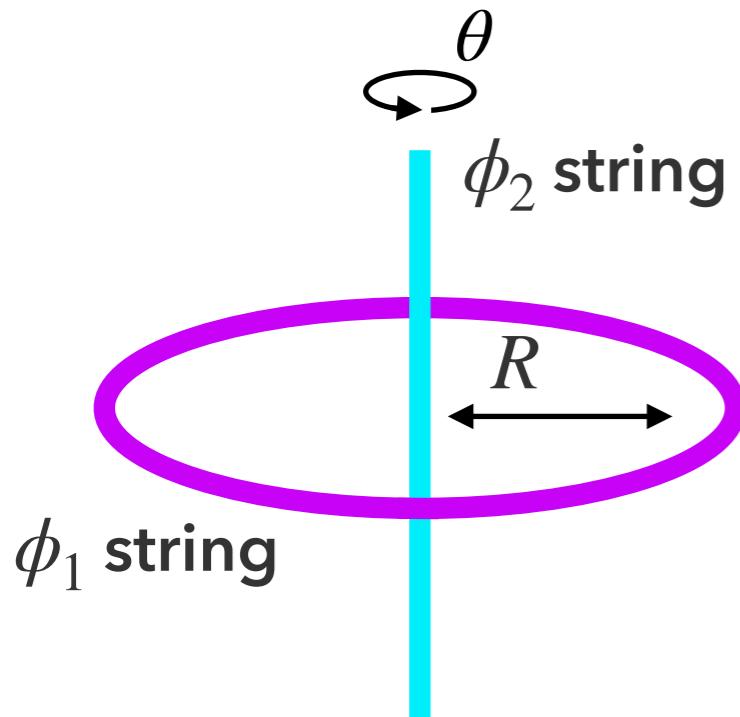


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Gauss law:

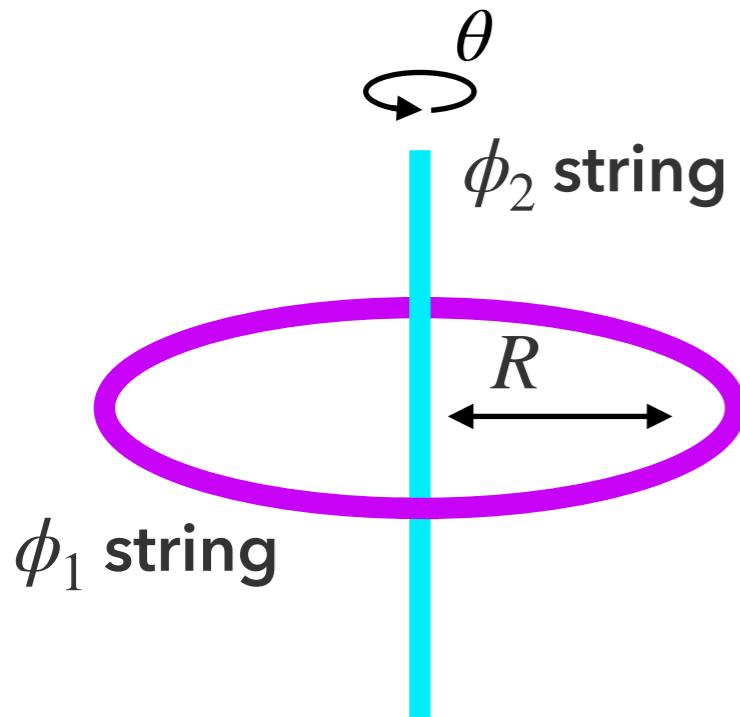
$$\frac{\delta \mathcal{L}}{\delta A_0} = \partial_i E_i - g^2 J^0 + \frac{g^2 c}{16\pi^2 R} |\vec{B}| = 0$$

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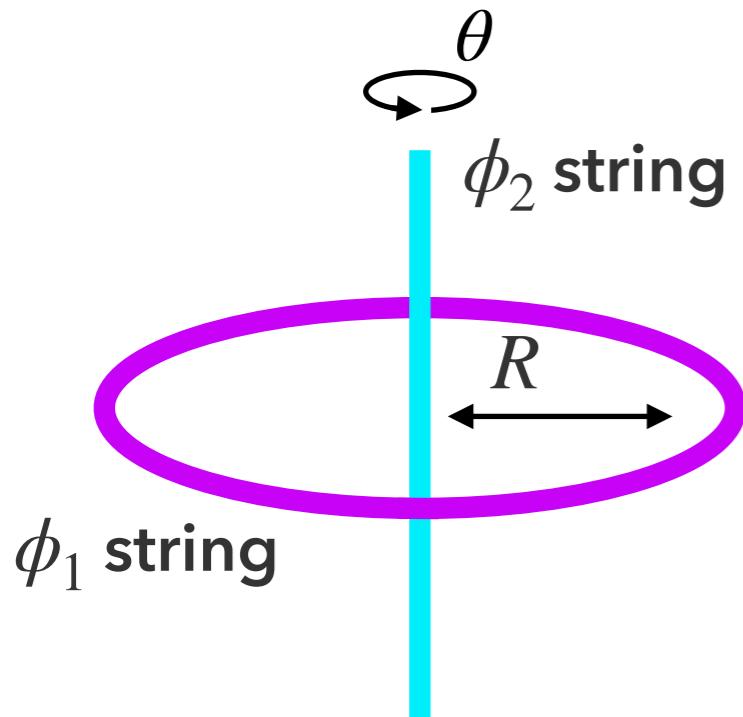
$$\Rightarrow \int d^3x J^0 = 2\pi R \int d^2x \frac{c}{16\pi^2 R} |\vec{B}| = \frac{c}{4g}$$

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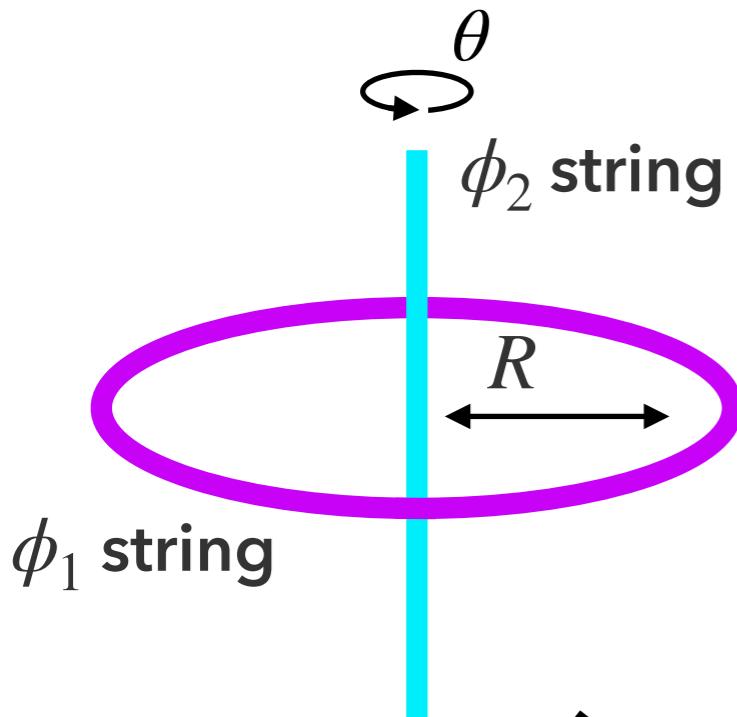
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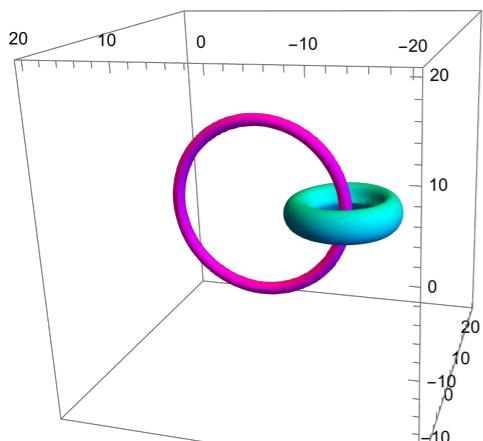
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$\phi_1$  stringが電荷を持つ  $\rightarrow$  電場による反発で縮まない

# Link stability

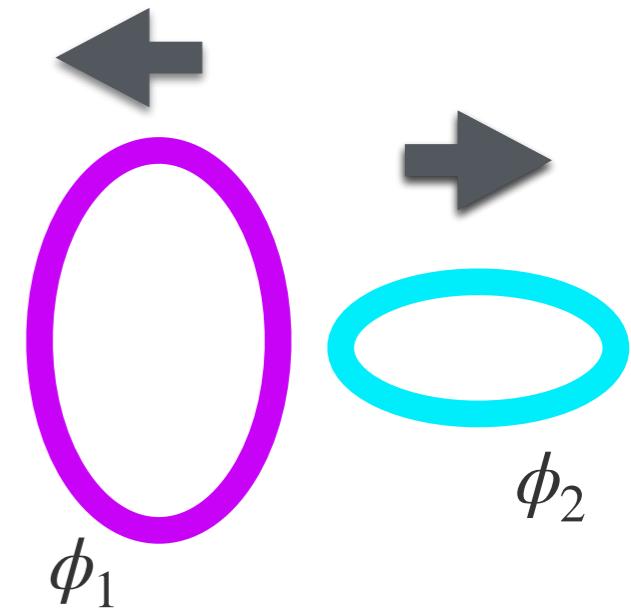
- linkを外してdecayできる？

→  $\lambda \gg g^2, \kappa, \chi$  と取っておけば外れない

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

→ non-linear  $\sigma$ -model w/  $O(4)$  sym. →  $O(3)$  sym.

**link = skyrmion** [Gudnason-Nitta '20]

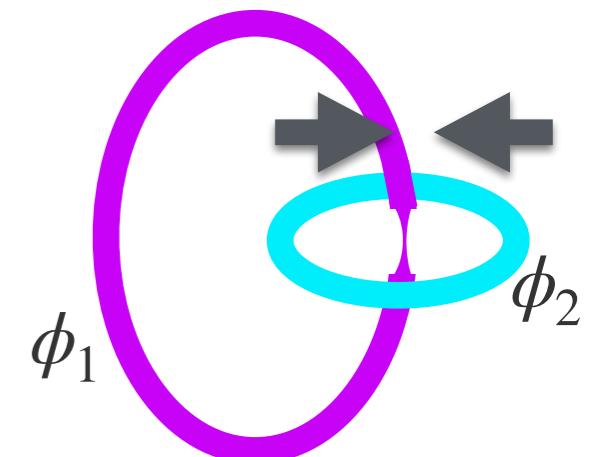


- $\phi_2$  stringは縮んでdecayできる？

→  $v_2/v_1 \ll 1$  と取っておけば縮まない

この2つの条件のもとで古典的に安定

ただし量子効果で崩壊できる（後述）



# Numerical calculation

Energy:

$$\begin{aligned}\mathcal{E} = & |D_i \phi_1|^2 + |\partial_i \phi_2|^2 + V(\phi_1, \phi_2) + \frac{1}{2g^2} (\partial_i A_j)^2 \\ & - g^2 |\phi_1|^2 A_0^2 - \frac{1}{2g^2} (\partial_i A_0)^2 - \frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}\end{aligned}$$

- Not positive definite  $\rightarrow$  remove  $A_0$  by solving Gauss law:

$$\frac{\delta \mathcal{L}}{\delta A_0} = \partial^2 A_0 - 2g^2 |\phi_1|^2 A_0 + \frac{g^2 c}{16\pi^2} (\vec{\nabla} a) \cdot \vec{B} = 0$$

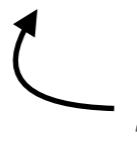
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$$\sim M_c^2 A_0 \sim \mathcal{O}(g^2 v_1^2 / c^2) A_0 \quad \text{for large } c$$

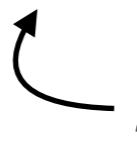
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 $\sim M_c^2 A_0 \sim \mathcal{O}(g^2 v_1^2 / c^2) A_0 \quad \text{for large } c$

$$\therefore A_0 \approx \frac{g^2 c}{16\pi^2} \frac{(\vec{\nabla} a) \cdot \vec{B}}{2g^2 |\phi_1|^2}$$

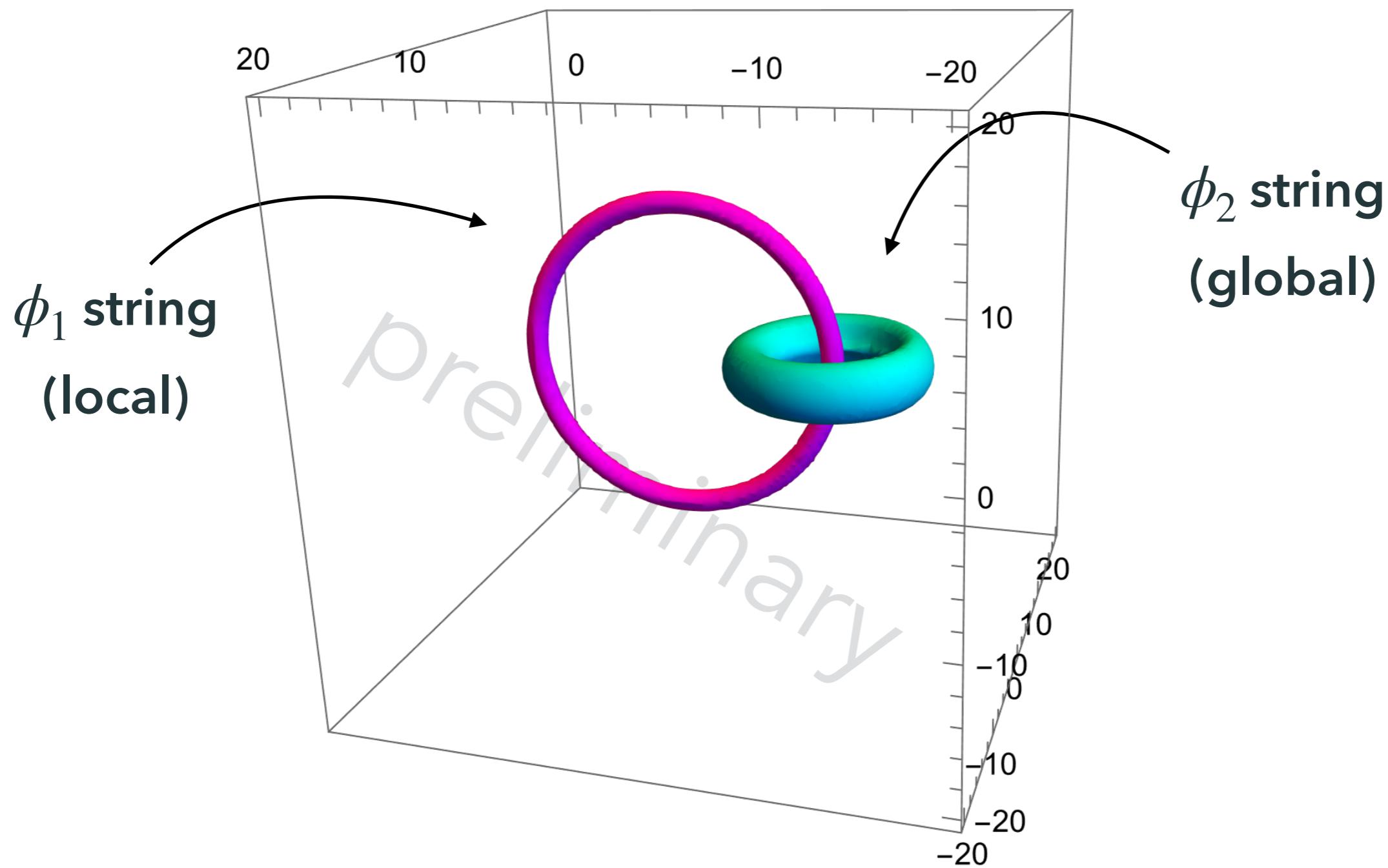
# Numerical calculation

Energy:

$$\begin{aligned}\mathcal{E} = & |D_i \phi_1|^2 + |\partial_i \phi_2|^2 + V(\phi_1, \phi_2) + \frac{1}{2g^2} (\partial_i A_j)^2 \\ & + \left( \frac{g^2 c}{16\pi^2} \right)^2 \frac{\left( (\vec{\nabla} a) \cdot \vec{B} \right)^2}{2g^2 |\phi_1|^2}\end{aligned}$$

- positive definite -> no obstacle
- Minimizing energy via non-linear conjugate gradient method
- CPU 3584-cores parallelizing on YITP computer cluster
- lattice spacing =  $0.2/gv_1$ ,  $N = 200^3$ , converged w/ O(1) days

# Numerical solution

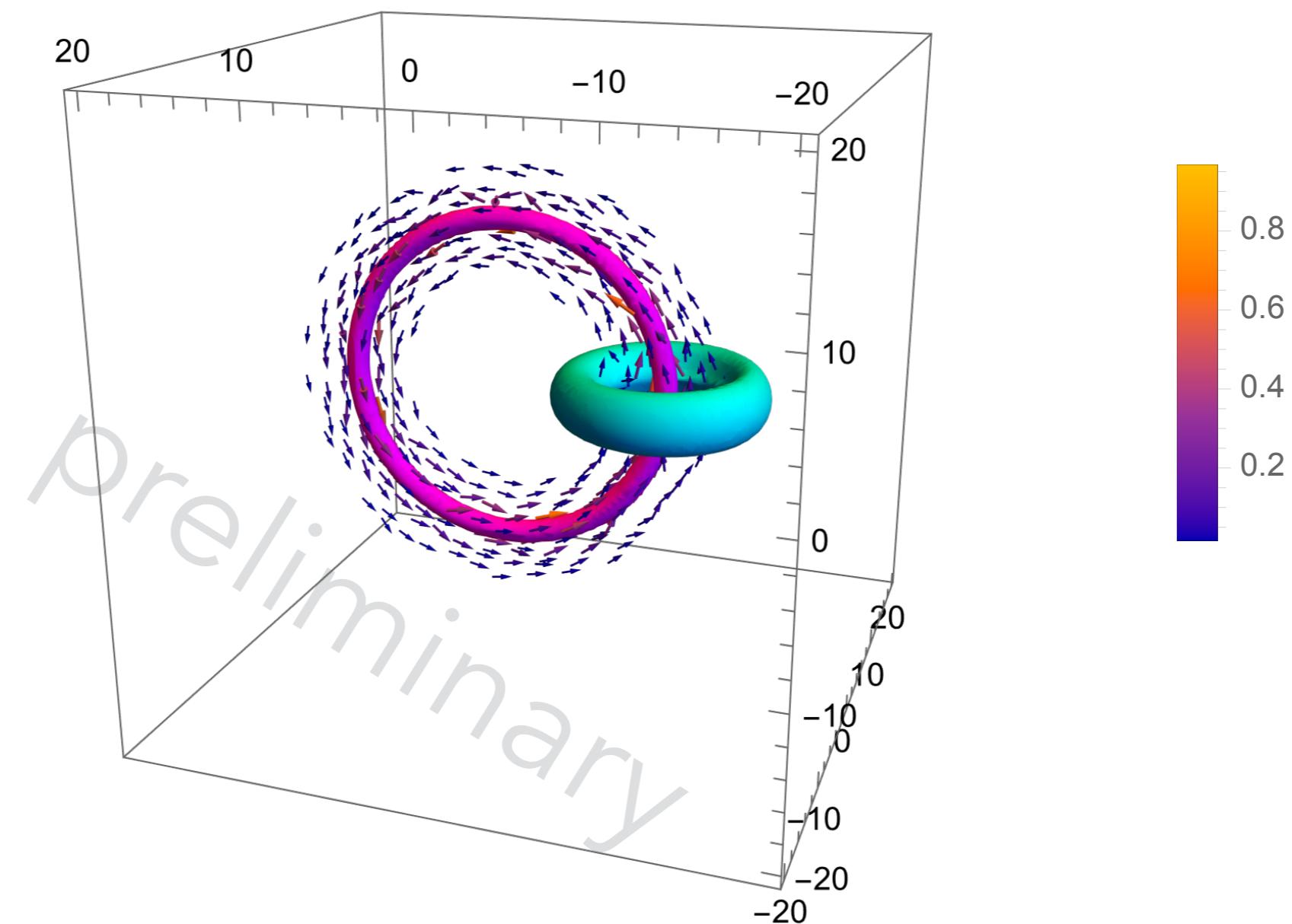


$$\lambda/g^2 = 100, \chi/g^2 = 19.5, \kappa/g^2 = 0.1, v_2/v_1 = 0.05$$

$$g^2 c/(16\pi^2) = 16$$

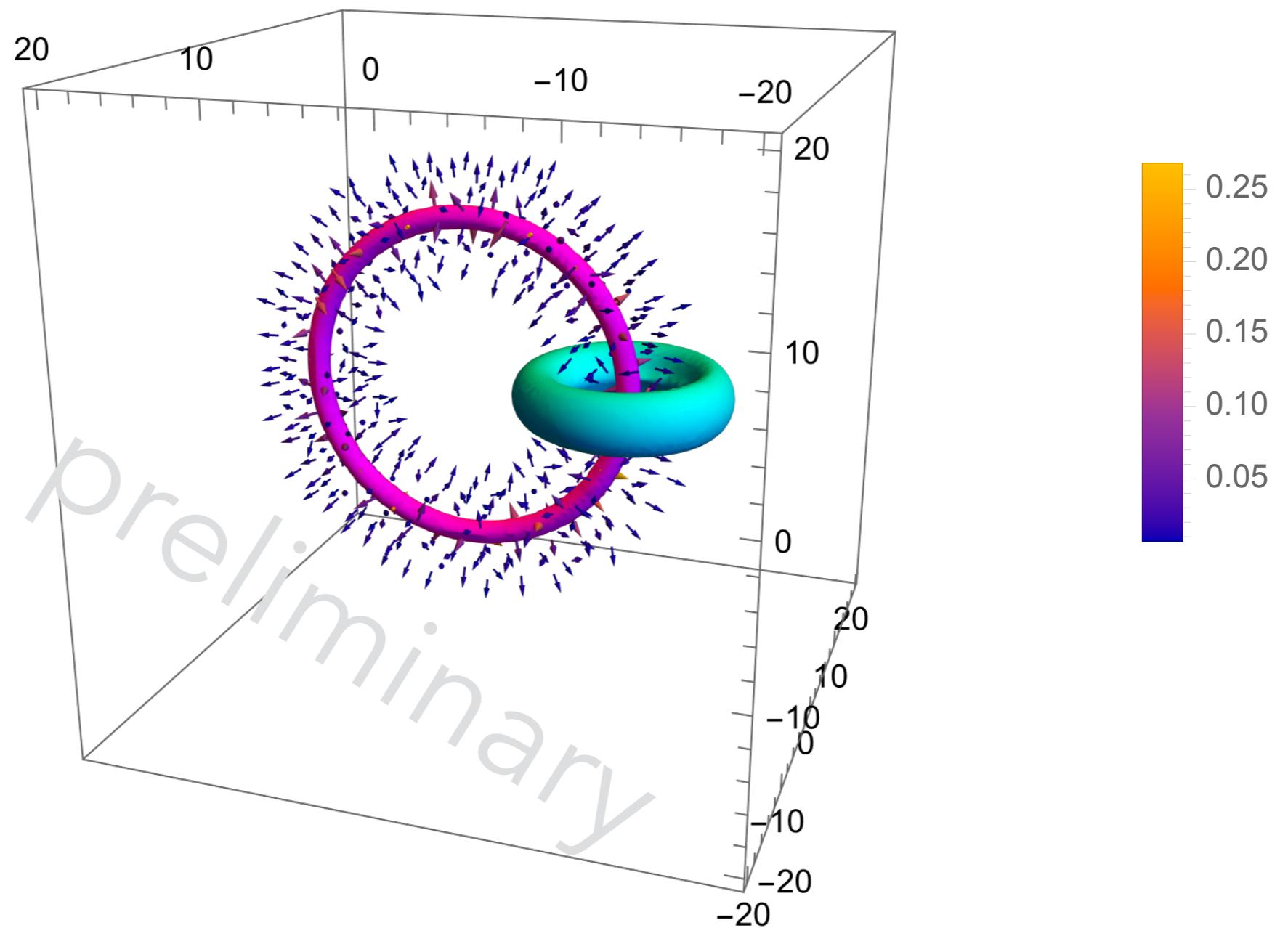
# Magnetic field

$$\vec{B}$$



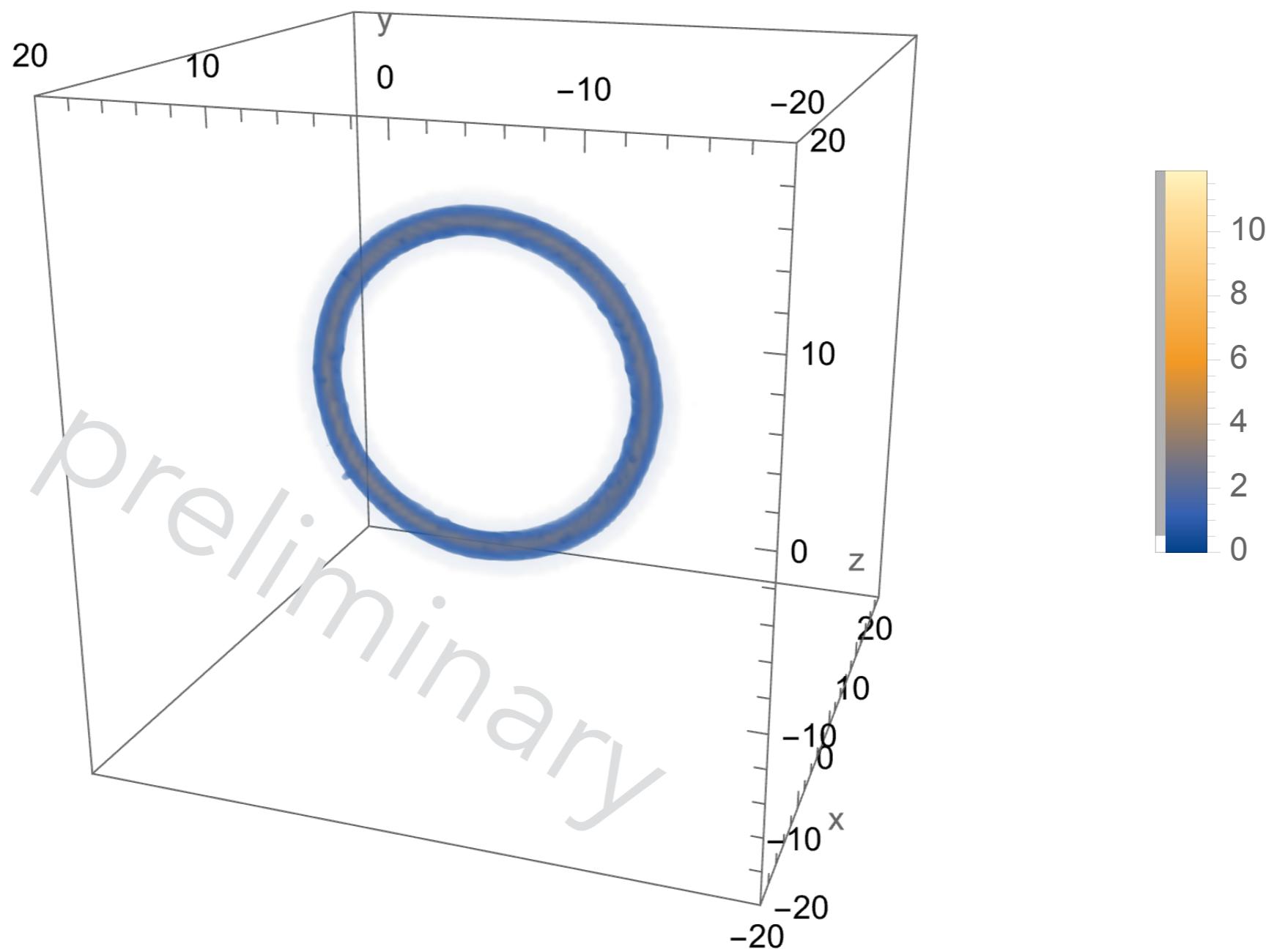
# Electric field

$$\vec{E} = \vec{\nabla} A_0$$



# Energy

Energy is dominated by  $\phi_1$  string



total energy:  $E \sim 504gv_1$

# Plan of talk

- Introduction
- Link soliton
- Application to Baryogenesis
- Summary

# Application to Baryogenesis

---

# General setup

More general charge assignment:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2)$$

$$D_\mu \phi_1 = (\partial_\mu - igq_1 A_\mu) \phi_1 \quad D_\mu \phi_2 = (\partial_\mu - igq_2 A_\mu) \phi_2 \quad q_2/q_1 \ll 1$$

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

- Symmetries:

$$U(1)_{gauge} : \begin{cases} \phi_1 \rightarrow e^{iq_1\theta_1} \phi_1 \\ \phi_2 \rightarrow e^{iq_2\theta_1} \phi_2 \end{cases}$$

$$U(1)_{global} : \begin{cases} \phi_1 \rightarrow e^{-iq_2\theta_1} \phi_1 \\ \phi_2 \rightarrow e^{iq_1\theta_1} \phi_2 \end{cases}$$

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$$U(1)_{gauge} : \begin{cases} \phi_1 \rightarrow e^{iq_1\theta_1} \phi_1 \\ \phi_2 \rightarrow e^{iq_2\theta_1} \phi_2 \end{cases} \quad U(1)_{global} : \begin{cases} \phi_1 \rightarrow e^{-iq_2\theta_1} \phi_1 \\ \phi_2 \rightarrow e^{iq_1\theta_1} \phi_2 \end{cases}$$

→ Also  $\phi_2$  string contains magnetic flux, but the solution is almost the same.

# General setup

More general charge assignment:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2) + \frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

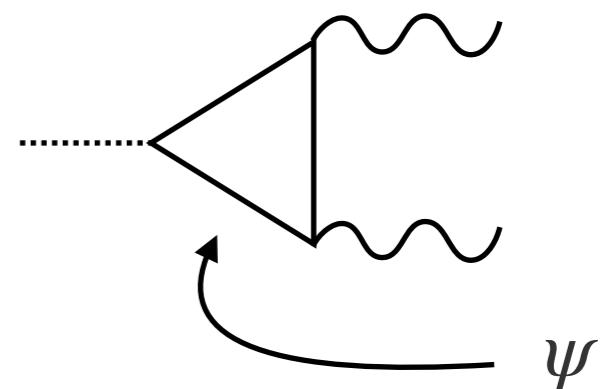
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$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

- Definition of "axion" is more complicated

$$a \equiv \frac{1}{i\sqrt{q_2^2 v_1^2 + q_1^2 v_2^2}} (-q_2 v_1 \arg(\phi_1) + q_1 v_2 \arg(\phi_2))$$

- Triangle anomaly is also complicated  
→  $c$  will be taken as free parameter.



# The model

$$\mathcal{L} = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2) + \frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

- Natural setup:  $U(1)_{gauge} = U(1)_{B-L}$  &  $U(1)_{global} = U(1)_{PQ}$

Type-I seesaw  $\rightarrow \nu$ -mass

QCD axion  $\rightarrow$  strong CP & Dark matter

$$\Rightarrow v_1 \sim v_2 \sim 10^{9-12} \text{ GeV}$$

- Axion quality problem can be avoided by gauged PQ mechanism.

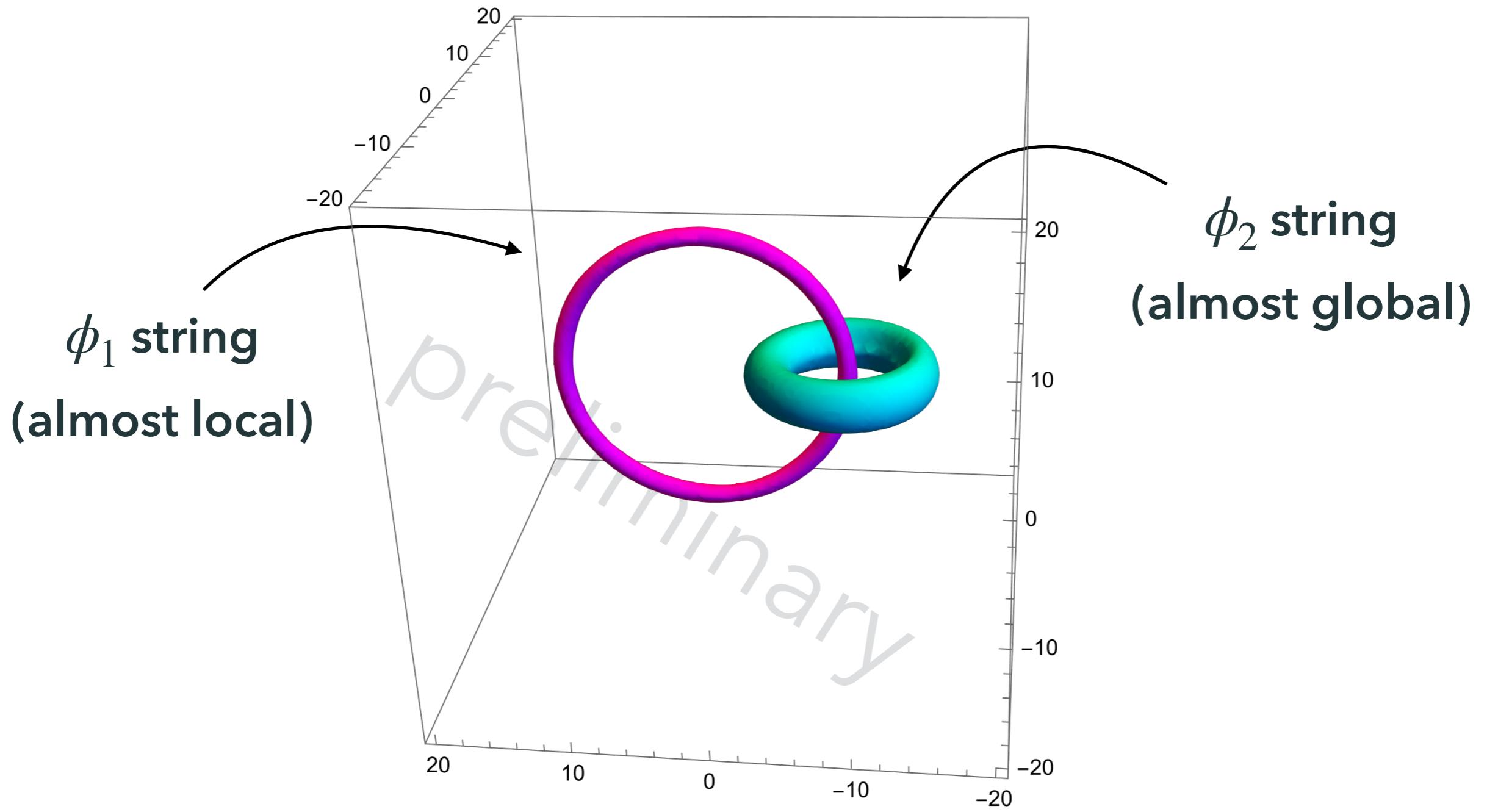
$$\Rightarrow q_1 = 1, q_2 = 0.1$$

[Fukuda-Ibe-Suzuki-Yanagida '17]

- Assume kinetic mixing with  $U(1)_Y$  in SM:  $\mathcal{L} \supset \frac{\epsilon}{2} Y_{\mu\nu} F^{\mu\nu}$

# Numerical solution

Solution is almost same as that w/  $q_2 = 0$ .



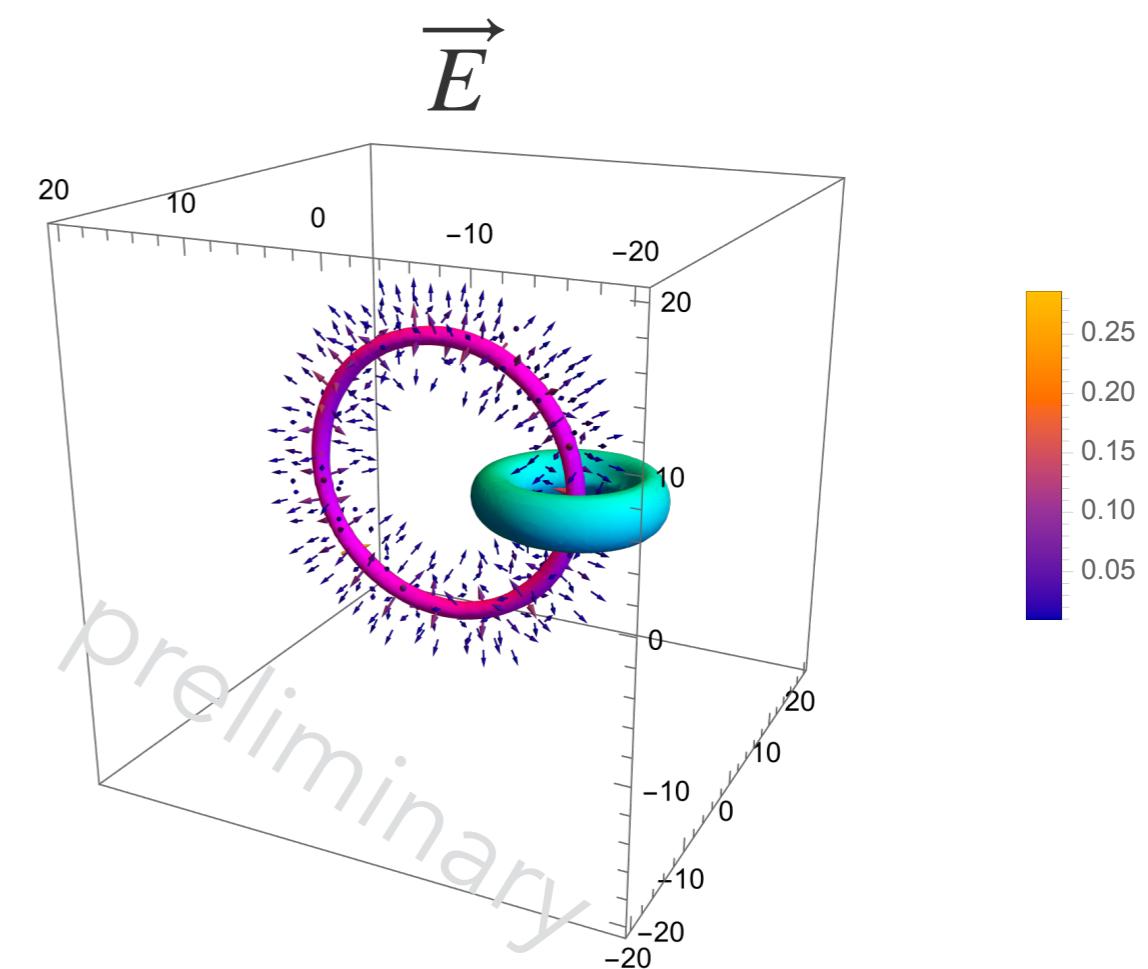
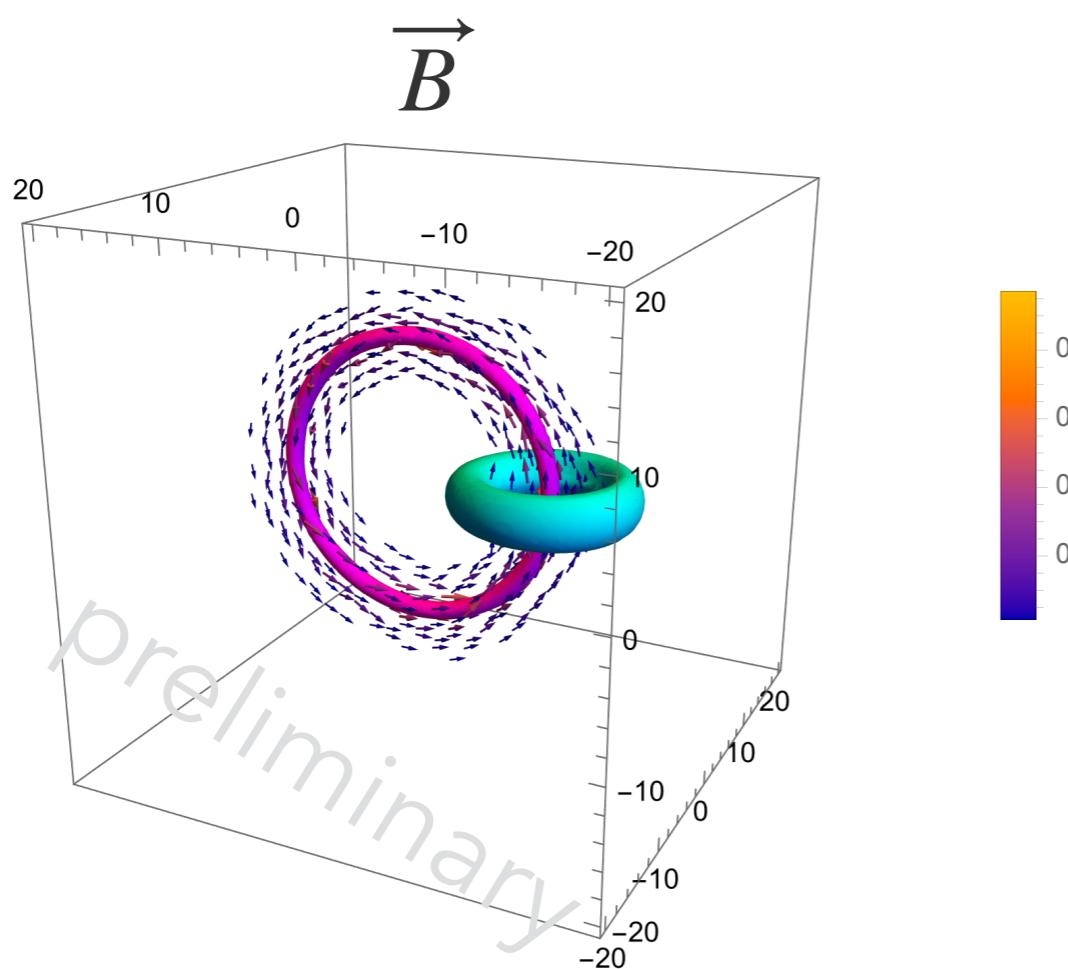
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$$g^2 c/(16\pi^2) = 16$$

$$q_2/q_1 = 0.1$$

# Magnetic & electric field

- Solution is almost same as that w/  $q_2 = 0$ .
- But note that  $\phi_2$  string also contains small  $B$  &  $E$  fluxes (cannot be seen).



# Helical magnetic field

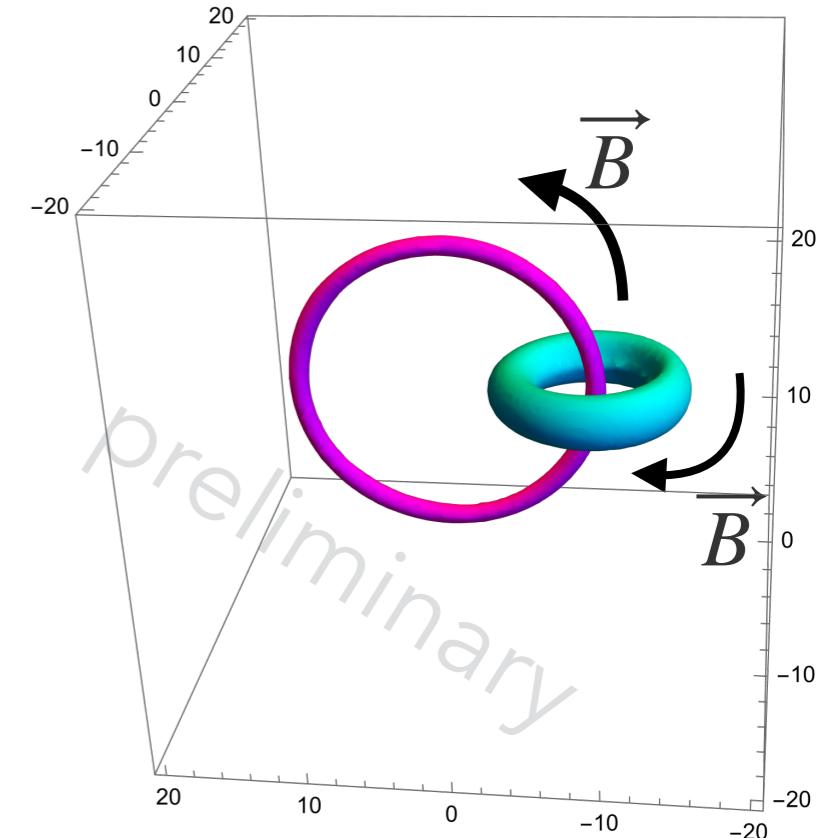
- Since the magnetic fluxes are linked, this soliton has finite helicity (Chern-Simons number):

$$N_{CS}[A] \equiv \frac{1}{16\pi^2} \int d^3x A dA = \frac{1}{16\pi^2} \int d^3x \vec{A} \cdot \vec{B}$$

For  $q_1 = 1, q_2 = 0.1, N_{CS}[A] \simeq 0.28$

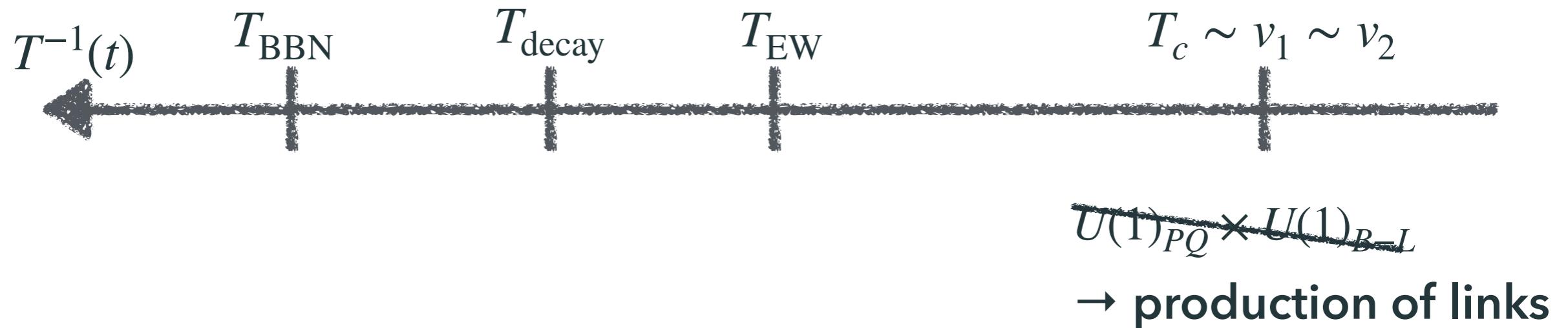
→ contains  $U(1)_Y$  helicity:  $N_{CS}[Y] \simeq \epsilon^2 N_{CS}[A]$

**can be used for baryogenesis!**

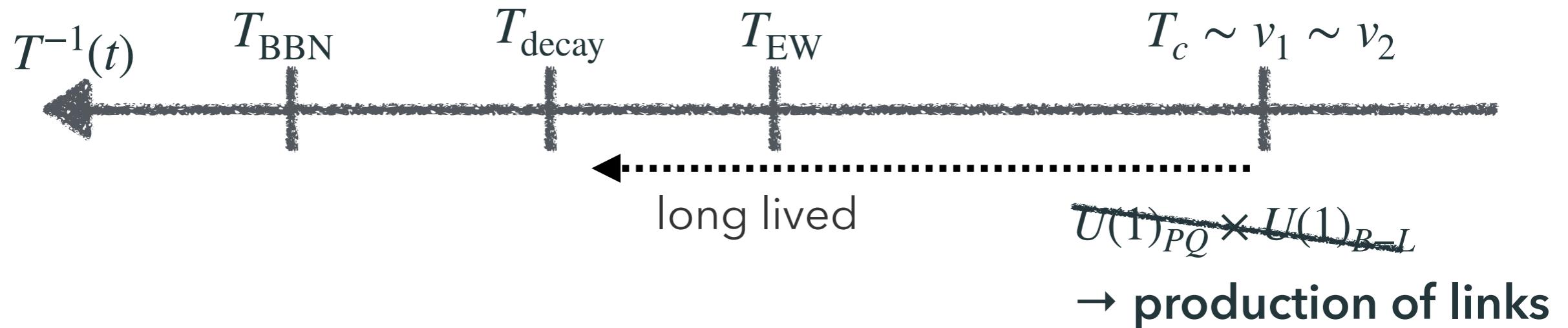


(cf: baryogenesis by helical  $U(1)_Y$  field) [Kamada-Long '16]

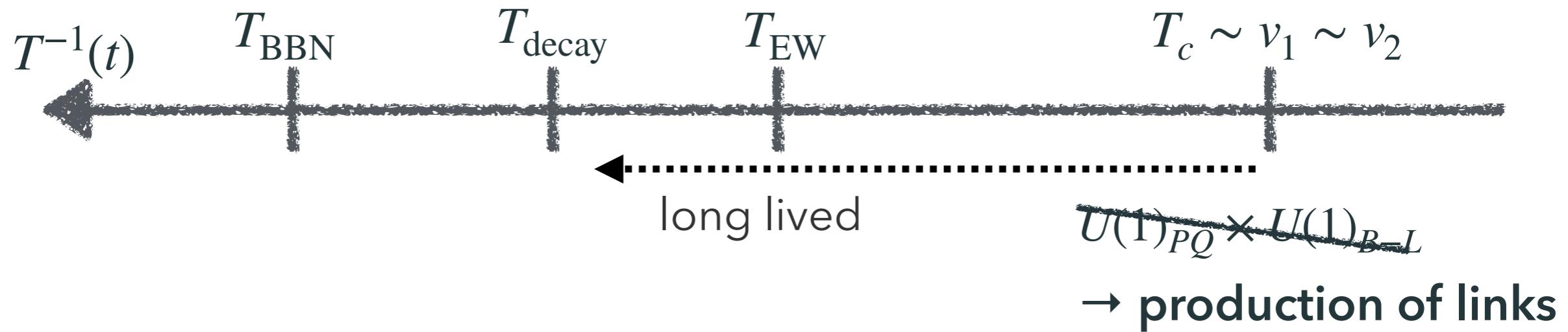
# Fate of link soliton



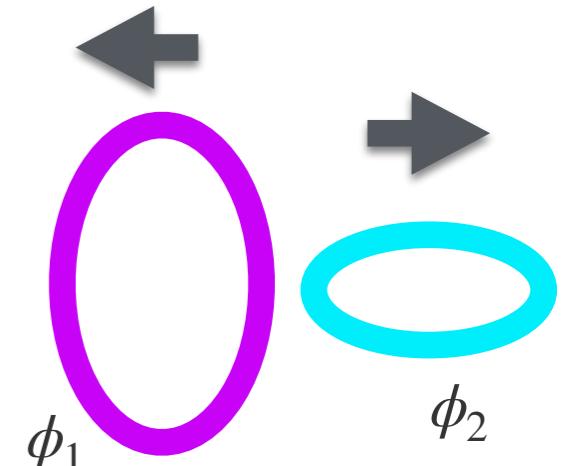
# Fate of link soliton



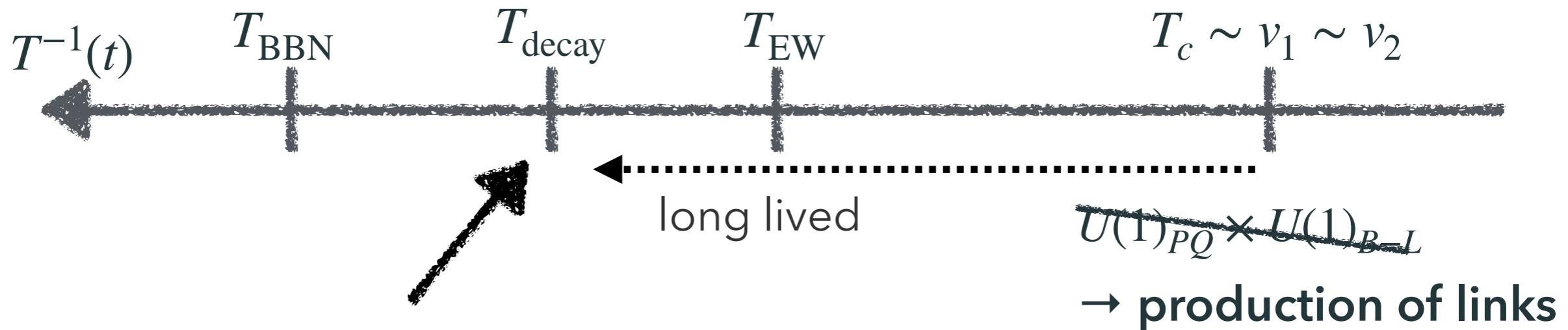
# Fate of link soliton



Links decay by tunneling effect after EW phase transition and before BBN.



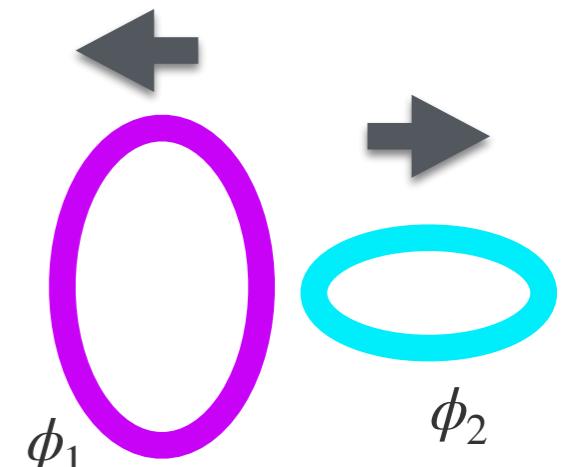
# Fate of link soliton



Links decay by tunneling effect after EW phase transition and before BBN.

→ change of helicity:  $\Delta N_{CS}[Y] \simeq \epsilon^2 N_{CS}[A]$

→ baryon # is generated through chiral anomaly:

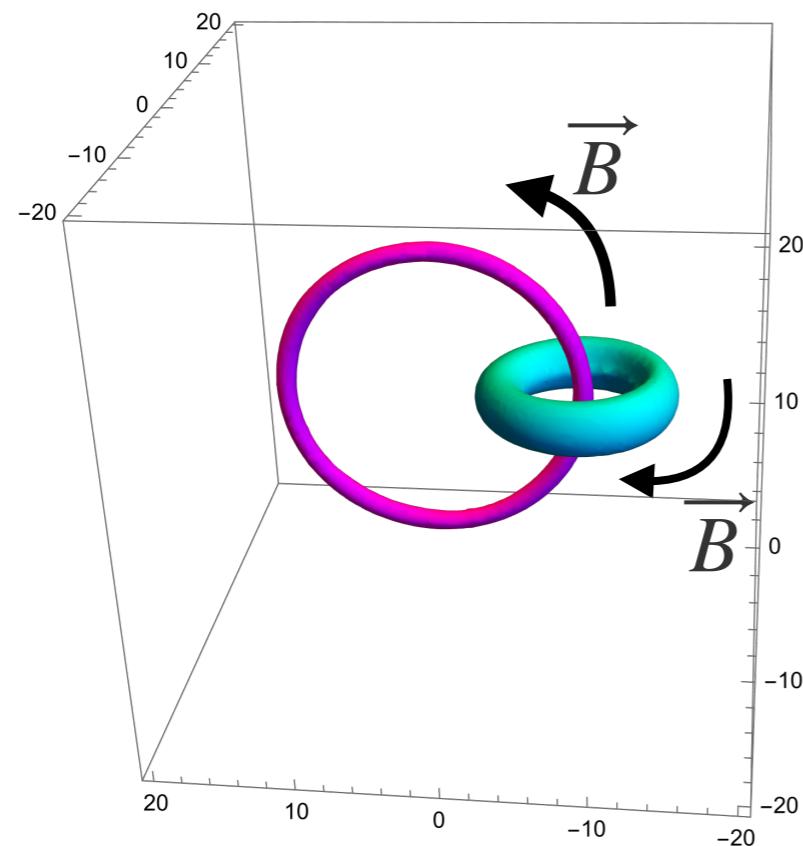


$\Delta B = \epsilon^2 N_{CS}[A] \text{ per link}$

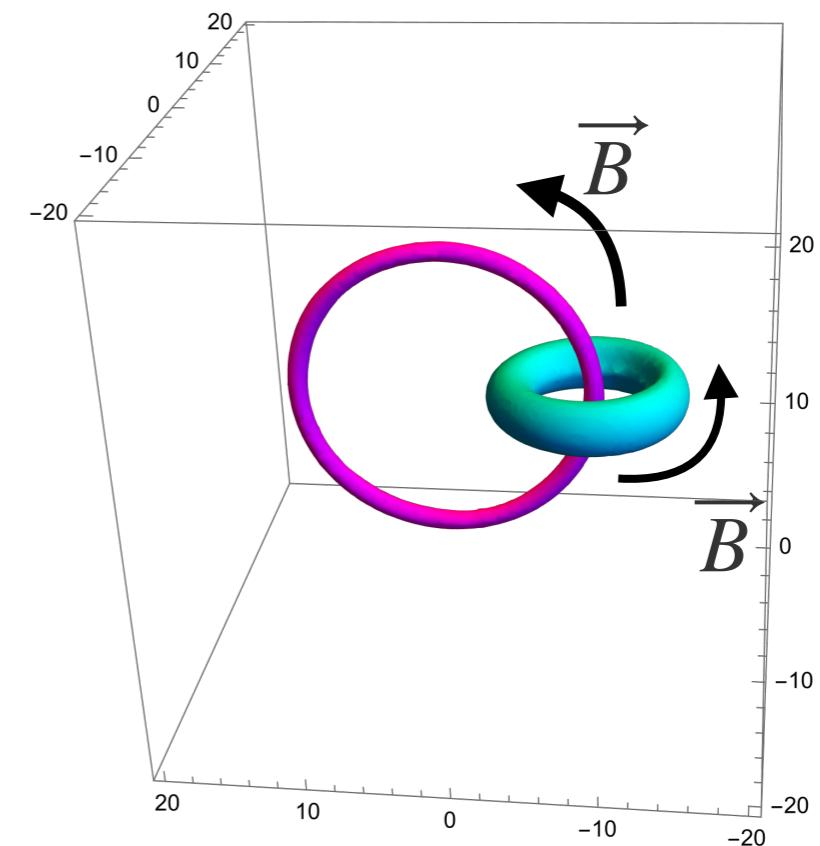
$$(\partial_\mu J_B^\mu \sim Y\tilde{Y} + \text{tr } W\tilde{W})$$

# Link and anti-link

- Anti-link solution also exists and its decay gives opposite baryon #.



link



anti-link

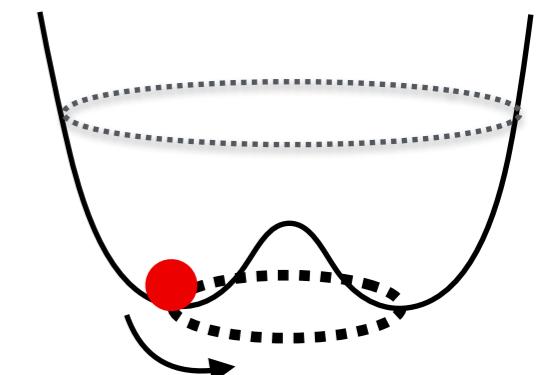
- Net baryon # is zero unless there is "chemical potential"  $\mu$   
btw link and anti-link at the production  $T \sim T_c$ .

# Origin of chemical potential?

- Choice (1): rotating pseudo scalar (axion-like particle)

$$\Delta \mathcal{L} = \frac{c}{16\pi^2} a' F_{\mu\nu} \tilde{F}^{\mu\nu} = -\frac{c}{16\pi^2} \underline{(\partial_0 a')} A dA \equiv \mu_{eff}$$

(cf: Affleck-Dine mechanism, axiogenesis [Co-Harigaya '19])



- Choice (2): chiral asymmetry  $Q_5$  from SO(10) GUT etc.

(cf: Chiral magnetic effect, Wash-in leptogenesis [Domcke+ '20])

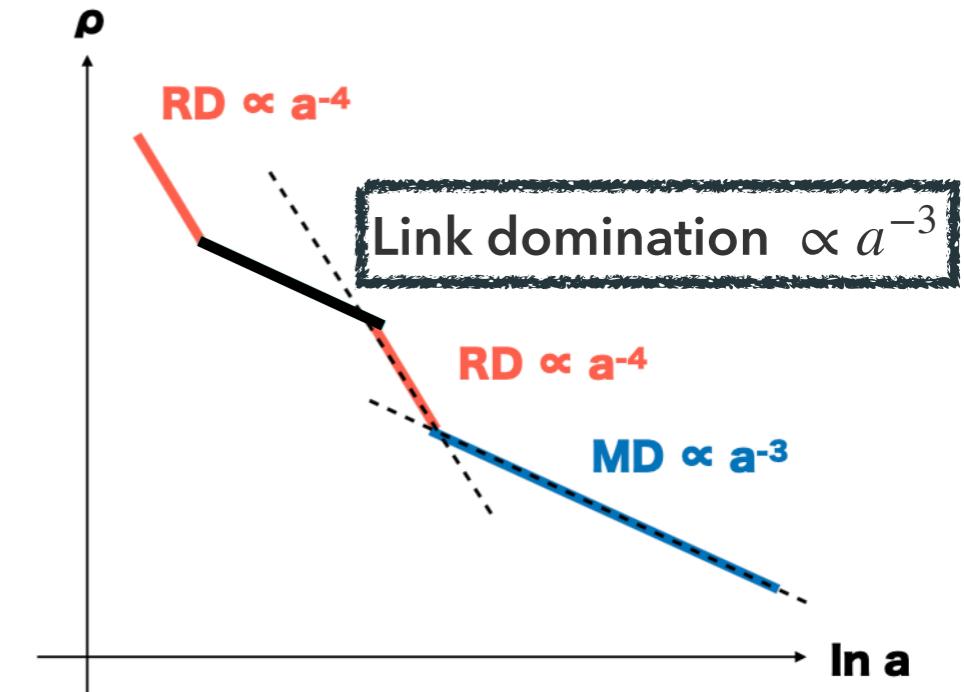
$$\eta_B = \frac{n_B}{s} \sim 10^{-10} \left( \frac{\epsilon}{0.1} \right)^2 \left( \frac{\mu_{eff}|_{T \sim v_1}}{0.1v_1} \right)$$

# Testability

- Before the links decay, they dominate the energy density of universe.

$$\frac{\rho_{link}}{\rho_\gamma} \Bigg|_{T \sim T_{EW}} \simeq \frac{M_{link} n_{link}}{g_* T^4} \Bigg|_{T \sim T_{EW}}$$

$$\simeq 10^{-4} \frac{v_1}{v_{EW}} \gg 1$$



- The entropy production due to decay cannot be ignored.
  - distorts spectrum of primordial gravitational wave
  - probed by primordial gravitational wave?

# Summary

- Massage of this talk:

$U(1)_{global} \times U(1)_{gauge}$  を自発的に破る模型では link solitonが存在する！

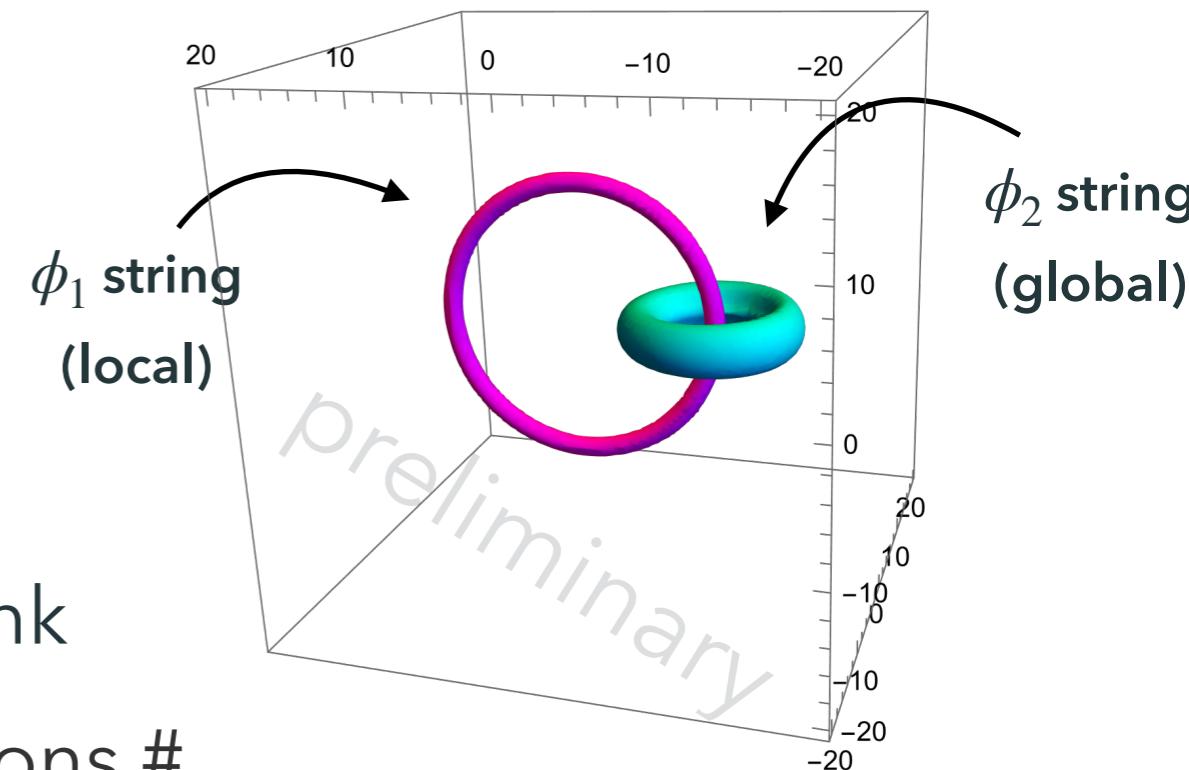
- Key: Chern-Simons coupling  $\frac{c}{16\pi^2} \int d^4x aF\tilde{F}$
- motivativeなセットアップ:

$$\begin{cases} U(1)_{global} = U(1)_{PQ} \\ U(1)_{gauge} = U(1)_{B-L} \end{cases}$$

axion stringとB-L stringからなるlink

- Links carry non-zero Chern-Simons #
- Baryon can be generated by decay of links

**link = "origin of baryon"**



# Backup

# Abelian-Higgs w/ Chern-Simons

Let's start from 2+1D Abelian-Higgs w/ CS term:

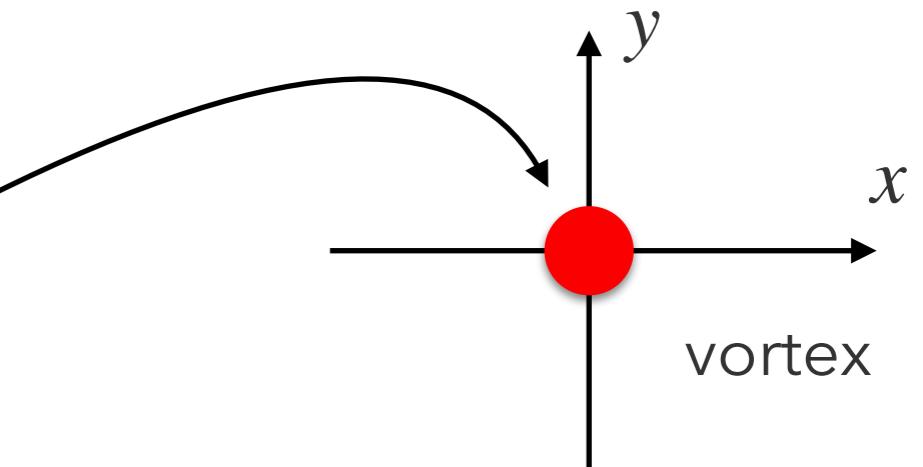
$$\mathcal{L} = |D_\mu \phi|^2 - \frac{1}{4} F_{\mu\nu}^2 - V(\phi) + g^2 c \epsilon^{\mu\nu\lambda} A_\mu \partial_\nu A_\lambda$$

$$V(\phi) = -m^2 |\phi|^2 + \lambda |\phi|^4$$

- For  $c = 0$ ,  $A_0$  is decoupled from static configurations.

→ static solution:  $\begin{cases} \phi = v f(r) e^{i\theta} \\ A_\theta = a(r)/g \\ A_0 = A_r = 0 \end{cases} \quad \begin{cases} f(0) = 0, f(\infty) = 1 \\ a(0) = 0, a(\infty) = 1 \end{cases}$

→ quantized magnetic flux  $\int d^2x B = 2\pi/g$



# Chern-Simons vortex

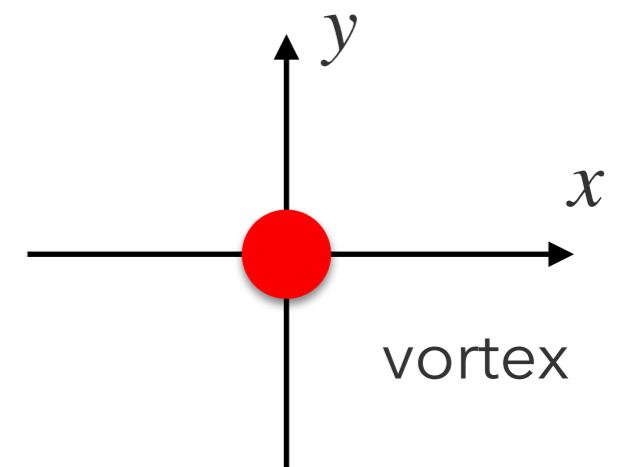
Let's start from 2+1D Abelian-Higgs w/ CS term:

$$\mathcal{L} = |D_\mu \phi|^2 - \frac{1}{4} F_{\mu\nu}^2 - V(\phi) + g^2 c \epsilon^{\mu\nu\lambda} A_\mu \partial_\nu A_\lambda$$

- For  $c \neq 0$ ,  $A_0$  is NOT, due to Gauss law constraint:

$$\frac{\delta \mathcal{L}}{\delta A_0} = \partial_i E_i - g^2 J^0 + g^2 c B = 0 \quad \begin{aligned} E_i &= \partial_i A_0 \\ J^0 &\equiv \phi^\dagger i D^0 \phi + (h.c.) \end{aligned}$$

→ magnetic flux sauces electric field!



# Chern-Simons vortex

Let's start from 2+1D Abelian-Higgs w/ CS term:

$$\mathcal{L} = |D_\mu \phi|^2 - \frac{1}{4} F_{\mu\nu}^2 - V(\phi) + g^2 c \epsilon^{\mu\nu\lambda} A_\mu \partial_\nu A_\lambda$$

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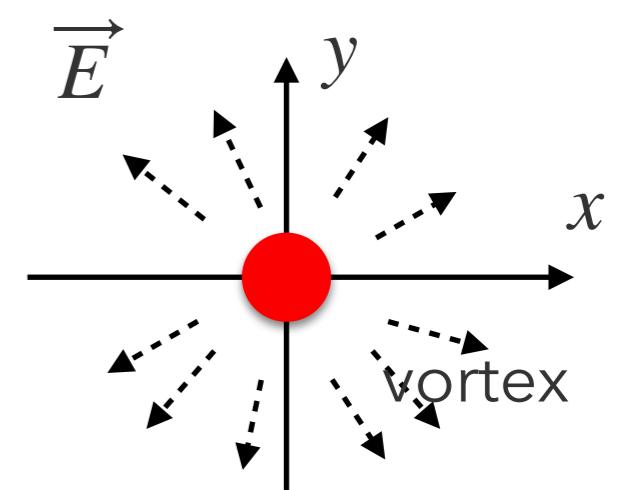
$$\frac{\delta \mathcal{L}}{\delta A_0} = \partial_i E_i - g^2 J^0 + g^2 c B = 0$$

$$E_i = \partial_i A_0$$
$$J^0 \equiv \phi^\dagger i D^0 \phi + (h.c.)$$

→ magnetic flux sauces electric field!

- quantized magnetic flux & electric charge

$$\int d^2x B = 2\pi/g \quad \int d^2x J^0 = 2\pi c/g$$



# Chern-Simons vortex

Let's start from 2+1D Abelian-Higgs w/ CS term:

$$\mathcal{L} = |D_\mu \phi|^2 - \frac{1}{4} F_{\mu\nu}^2 - V(\phi) + g^2 c \epsilon^{\mu\nu\lambda} A_\mu \partial_\nu A_\lambda$$

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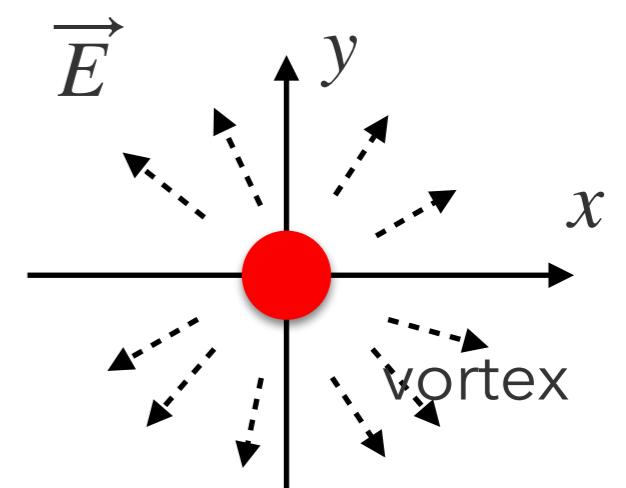
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→ magnetic flux sauces electric field!

- quantized magnetic flux & electric charge

$$\int d^2x B = 2\pi/g \quad \int d^2x J^0 = 2\pi c/g$$

called Chern-Simons vortex



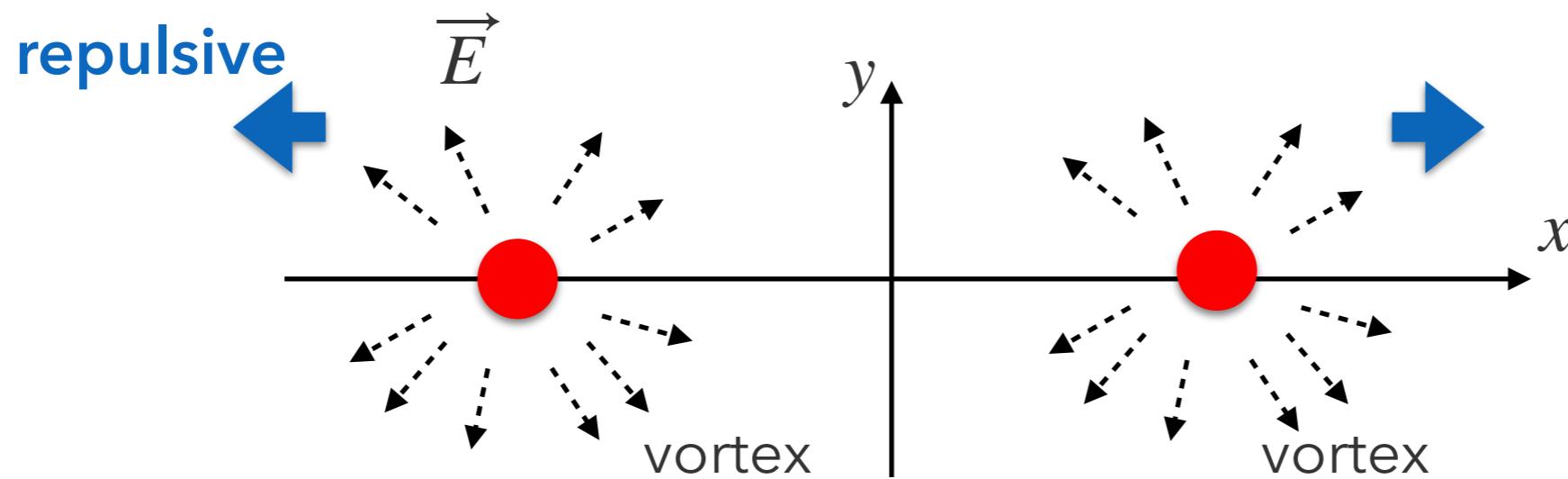
# Interaction of Chern-Simons vortices

- typical length scale of  $\vec{E}$ :  $l_E$

$$\text{w/ } l_E^{-1} \equiv gv \left( \frac{1}{2}\sqrt{4+c^2} - \frac{c}{2} \right) \simeq \mathcal{O}(gv/c) \text{ for } c \rightarrow \infty$$

# Interaction of Chern-Simons vortices

- Interaction btw two CS vortices

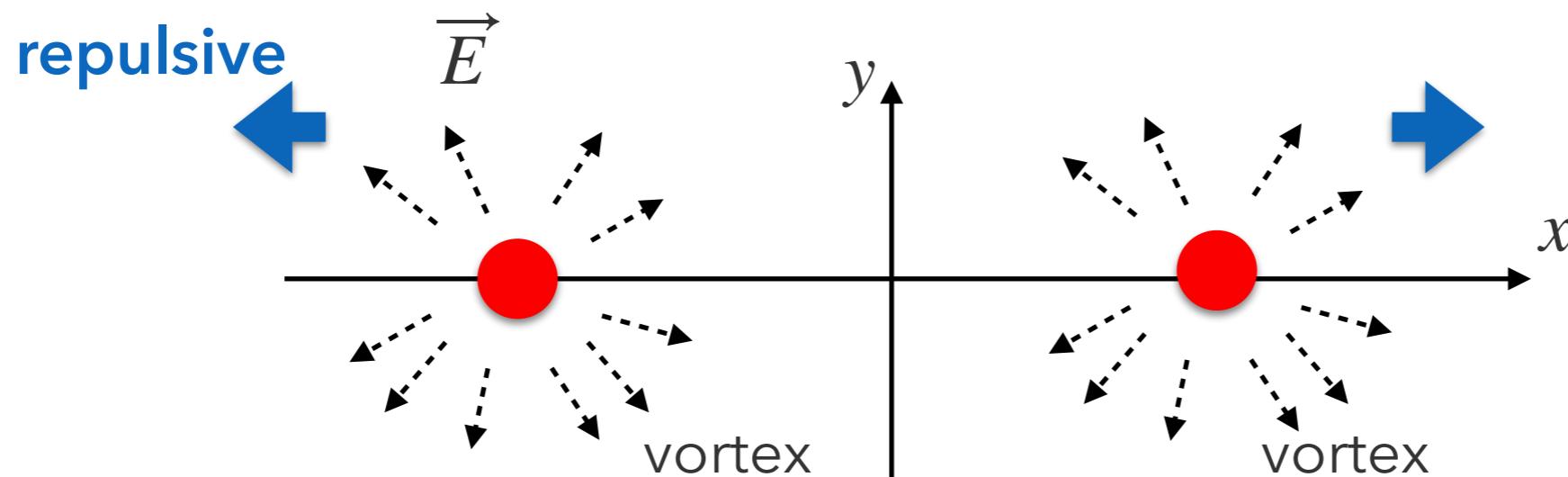


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- Interaction btw two CS vortices



- typical length scale of  $\vec{E}$ :  $l_E$

$$\text{w/ } l_E^{-1} \equiv gv \left( \frac{1}{2}\sqrt{4 + c^2} - \frac{c}{2} \right) \simeq \mathcal{O}(gv/c) \text{ for } c \rightarrow \infty$$

- For large  $c$ , **long-range repulsive force!**

# Relation to Skyrmion

For  $\lambda \gg g^2, \kappa, \chi$ ,

$$\begin{aligned} V(\phi) &= \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4 \\ &\rightarrow \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 \end{aligned}$$

→ non-linear sigma model w/  $O(4)$  symmetry,  
which breaks into  $O(3)$

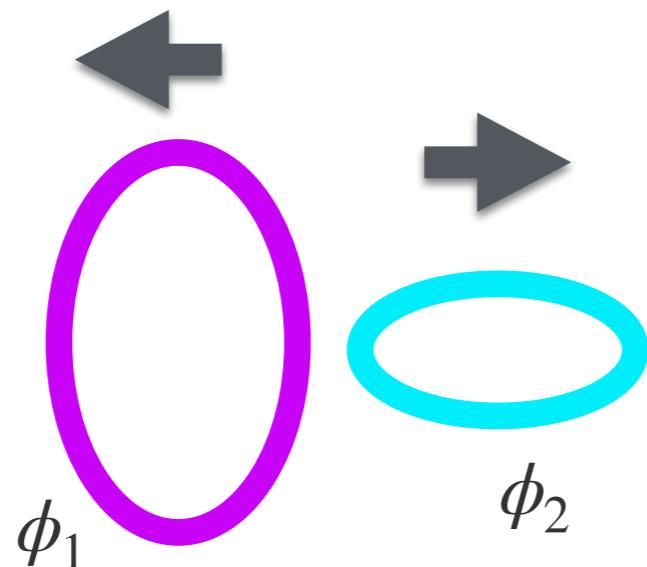
There exists Skyrmion defined by winding number:

$$N_{sk} = \int d^3x \epsilon^{ijk} \text{Tr} \left[ U^\dagger \partial_i U U^\dagger \partial_j U U^\dagger \partial_k U \right] \quad U = \begin{pmatrix} \text{Re } \phi_1 & \text{Im } \phi_2 \\ -\text{Im } \phi_1 & \text{Re } \phi_2 \end{pmatrix}$$

The link is nothing but the Skyrmion!

[Gudnason-Nitta '20]

# Decay of link soliton



$$\tau^{-1} \sim \Gamma \sim g\nu_1 \exp \left[ -\frac{4}{3} \sqrt{\frac{\lambda\nu_1}{g\nu_2}} \right]$$

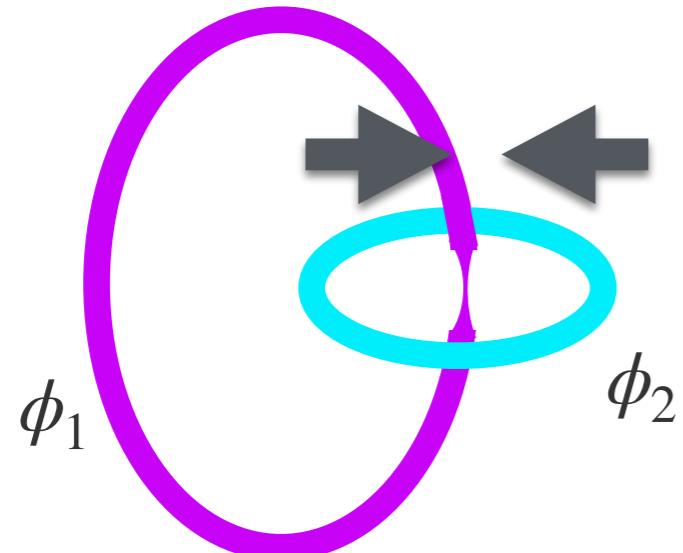
$$\frac{M_{\text{pl}}}{\nu_{EW}^2} < \tau < \frac{M_{\text{pl}}}{(1\text{MeV})^2}$$

$$\begin{aligned}\lambda &< 4\pi \\ g &< \sqrt{24\pi}\end{aligned}$$

$$\Leftrightarrow \log \frac{g\nu_1}{10^{11}\text{GeV}} + 60 \lesssim \frac{4}{3} \sqrt{\frac{\lambda\nu_1}{g\nu_2}} \lesssim \log \frac{g\nu_1}{10^{11}\text{GeV}} + 82$$

$$\therefore 100 \lesssim \frac{\lambda}{g} \lesssim 190$$

# Decay of link soliton



$$\tau^{-1} \sim \Gamma \sim g v_1 \exp \left[ -\frac{1}{g^2} \frac{v_1}{v_2} \# \right]$$

$$\frac{M_{\text{pl}}}{v_{EW}^2} < \tau < \frac{M_{\text{pl}}}{(1\text{MeV})^2}$$

$$\begin{aligned}\lambda &< 4\pi \\ g &< \sqrt{24\pi}\end{aligned}$$

$$\Leftrightarrow \log \frac{g v_1}{10^{11}\text{GeV}} + 60 \lesssim \frac{v_1}{g^2 v_2} \# \lesssim \log \frac{g v_1}{10^{11}\text{GeV}} + 82$$

our parameter choice:  $\frac{v_1}{v_2} = 20 \quad \rightarrow 0.49 \lesssim g \lesssim 0.58$

# Dimensionless unit

$$\begin{aligned}\mathcal{L} &= |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2) \\ V(\phi) &= \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4\end{aligned}$$

$$\mathcal{L} = \frac{1}{g^2} \left[ g^2 |D_\mu \phi_1|^2 + g^2 |D_\mu \phi_2|^2 - \frac{1}{4} F_{\mu\nu}^2 - g^2 V(\phi_1, \phi_2) \right]$$

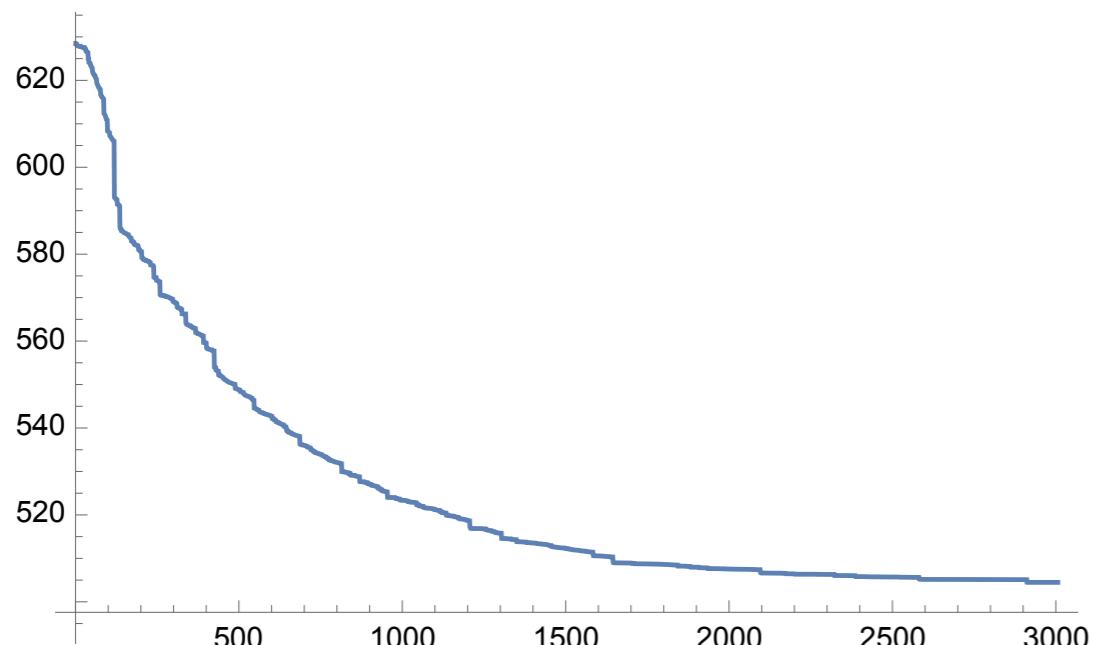
- field redefinition:  $g\phi_1 \rightarrow \phi$ ,  $g\phi_2 \rightarrow \phi_2$

$$\begin{aligned}\mathcal{L} &= \frac{1}{g^2} \left[ |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4} F_{\mu\nu}^2 - \tilde{V}(\phi_1, \phi_2) \right] \\ \tilde{V}(\phi) &= \frac{\lambda}{g^2} \left( |\phi_1|^2 + |\phi_2|^2 - g^2 \mu^2 \right)^2 - \frac{\kappa}{g^2} |\phi_1|^2 |\phi_2|^2 + \frac{\chi}{g^2} |\phi_2|^4\end{aligned}$$

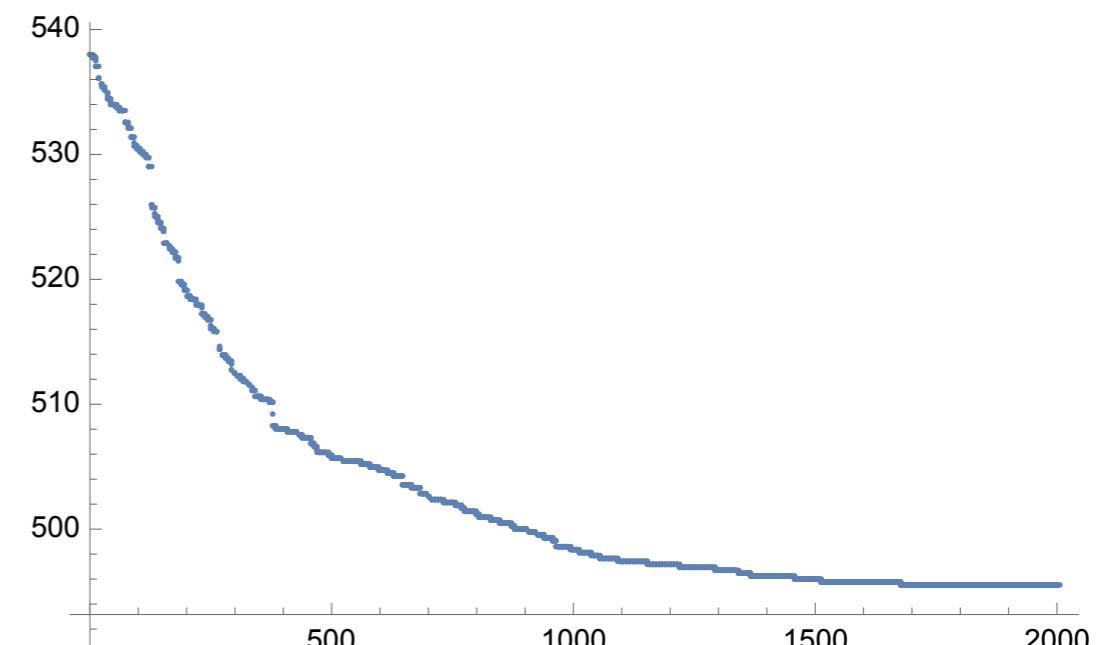
only dimensionful parameter  $g\mu$

# Energy convergence

$$q_2/q_1 = 0$$



$$q_2/q_1 = 0.1$$



# The model

3+1D theory:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |\partial_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2)$$

CS coupling

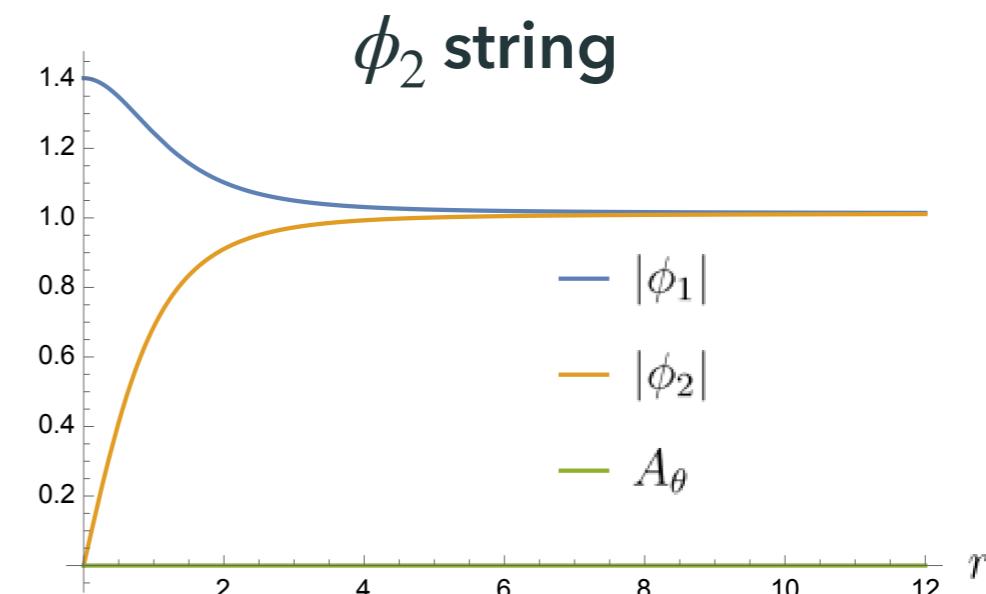
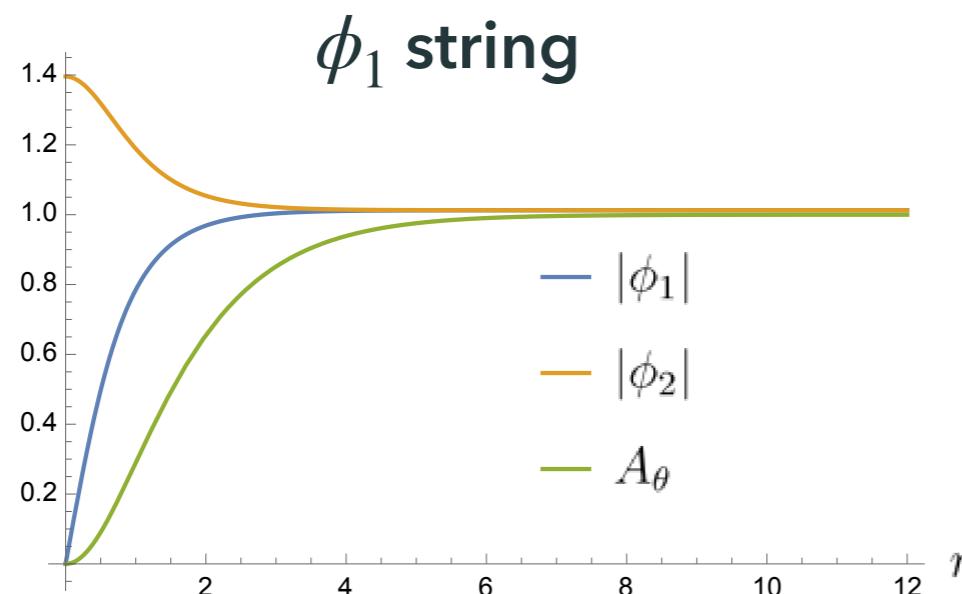
$$+\frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$D_\mu \phi_1 = (\partial_\mu - igA_\mu) \phi_1$$

$$a \equiv -i \arg(\phi_2)$$

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

- CS coupling does not affect single strings.



# Interaction of Chern-Simons vortices

- Single static solution:  $\begin{cases} \phi = v f(r) e^{i\theta} \\ A_\theta = a(r)/g \\ A_0 = b(r)/g \end{cases}$   $\begin{cases} f(0) = 0, f(\infty) = 1 \\ a(0) = 0, a(\infty) = 1 \\ b(0) = 0, b(\infty) = 0 \end{cases}$
- Asymptotic behavior at  $r \rightarrow \infty$ 
$$1 - f(r) \sim e^{-M_\phi r} \quad b(r) \sim 1 - a(r) \sim e^{-M_c r}$$
w/  $M_c \equiv gv \left( \frac{1}{2}\sqrt{4+c^2} - \frac{c}{2} \right) \simeq \mathcal{O}(gv/c)$  for  $c \rightarrow \infty$

# Interaction of Chern-Simons vortices

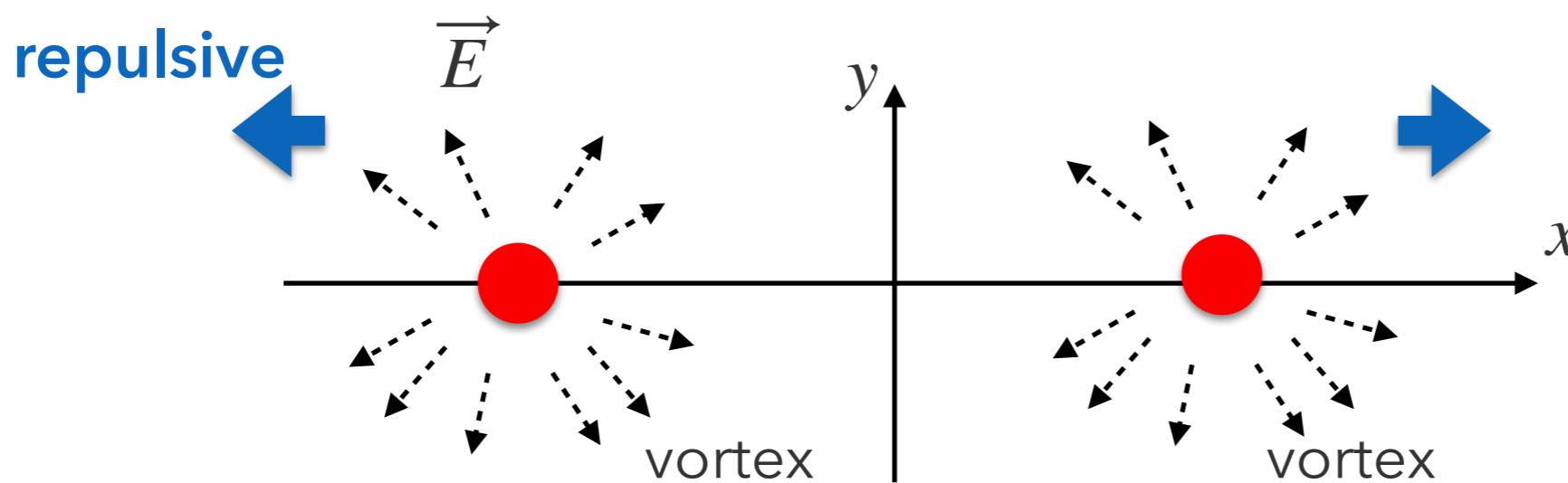
- Single static solution:  $\begin{cases} \phi = v f(r) e^{i\theta} \\ A_\theta = a(r)/g \\ A_0 = b(r)/g \end{cases}$   $\begin{cases} f(0) = 0, f(\infty) = 1 \\ a(0) = 0, a(\infty) = 1 \\ b(0) = 0, b(\infty) = 0 \end{cases}$

- Asymptotic behavior at  $r \rightarrow \infty$

$$1 - f(r) \sim e^{-M_\phi r} \quad b(r) \sim 1 - a(r) \sim e^{-M_c r}$$

$$\text{w/ } M_c \equiv gv \left( \frac{1}{2}\sqrt{4+c^2} - \frac{c}{2} \right) \simeq \mathcal{O}(gv/c) \text{ for } c \rightarrow \infty$$

- For large  $c$ , **long-range repulsive force!**



# Classical stability

- Delinking by passing through each other?

→ prevented by taking  $\lambda \gg g^2, \kappa, \chi$

**Overlap of strings ( $\phi_1 = \phi_2 = 0$ ) cost large energy**

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

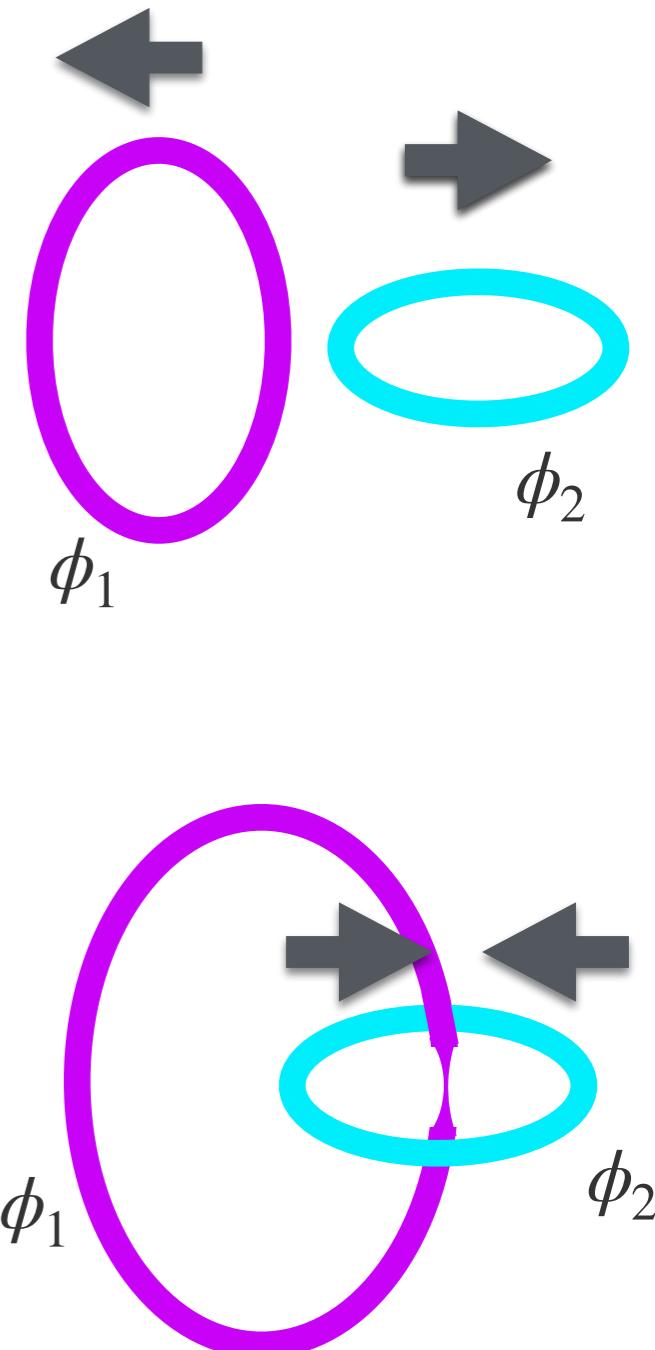
- $\phi_2$  string is not charged and thus can shrink ?

→ prevented by taking  $v_2/v_1 \ll 1$

**$\phi_2$  string is too light to pinch  $\phi_1$  string**



**classically stable (but decay by tunneling effect)**



# The full Lagrangian

$$\mathcal{L} = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2) + \frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$+ \mathcal{L}_{mat} + \mathcal{L}_{SM} - V_{portal}(\phi_1, \phi_2, H) + \mathcal{L}_{kin.mix.}$$

- Natural setup:  $U(1)_{gauge} = U(1)_{B-L}$  &  $U(1)_{global} = U(1)_{PQ}$

Type-I seesaw  $\rightarrow \nu$ -mass

QCD axion  $\rightarrow$  strong CP & Dark matter

$$\Rightarrow v_1 \sim v_2 \sim 10^{9-12} \text{ GeV}$$

- Assume kinetic mixing with  $U(1)_Y$  in SM:  $\mathcal{L} \supset \frac{\epsilon}{2} Y_{\mu\nu} F^{\mu\nu}$

# Baryon # from link

- Naively, anti-link is also produced  $\rightarrow n_{link} - n_{\overline{link}} = 0$  ?
- need "chemical potential"  $\mu$  (discussed later) when produced:

$$\frac{n_{link} - n_{\overline{link}}}{s} \simeq \frac{\mu}{T} \frac{n_{link}}{s} \simeq \frac{\mu}{T} 10^{-6}$$

We have used  $n_{link} \sim 10^{-4} T^3$

[Vachaspati '84]

$\rightarrow$  generated total baryon # due to decay:

$$\eta_B = \frac{n_B}{s} \sim 10^{-10} \left( \frac{\epsilon}{0.1} \right)^2 \left( \frac{\mu/v_1}{0.1} \right)$$