

# Superconducting Circuits for Quantum Computing and Quantum Sensing

Robert McDermott, UW-Madison

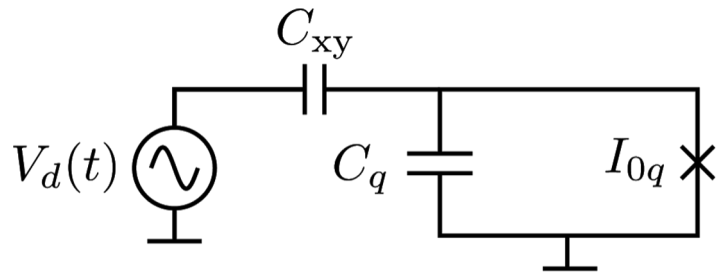


ICEPP Workshop  
March 9, 2022

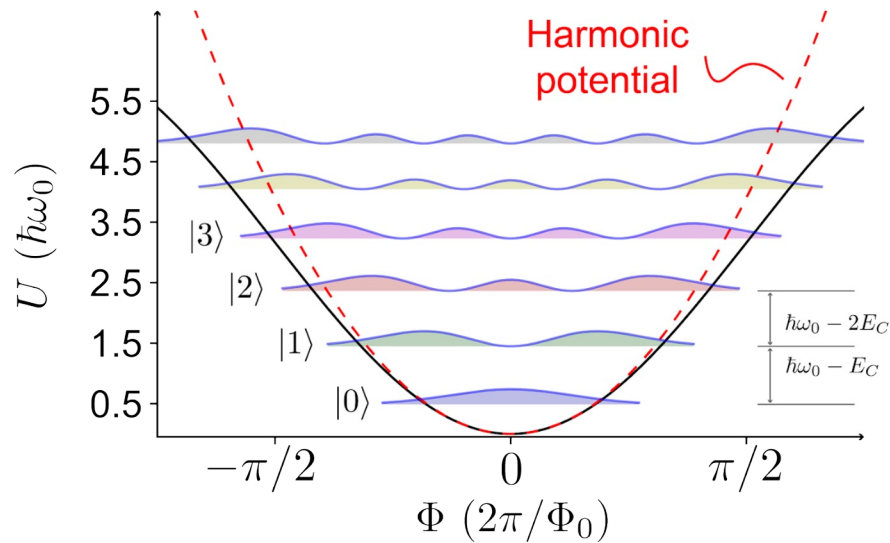
# Outline

- Introduction to SC qubits
- Requirements for fault tolerance
- Correlated errors from particle impacts; mitigation
- Photon-assisted QP poisoning
- Qubit-based sensors for dark matter detection

# Transmon qubit



$$\hat{H} = \frac{\hat{Q}^2}{2C_\Sigma} - E_J \cos\left(\frac{2\pi\hat{\Phi}}{\Phi_0}\right)$$

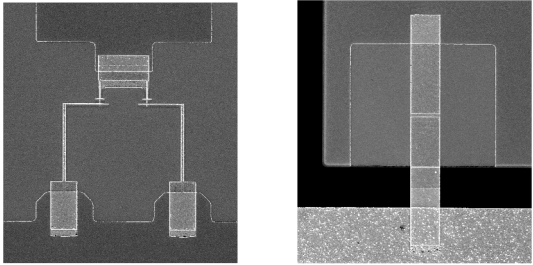
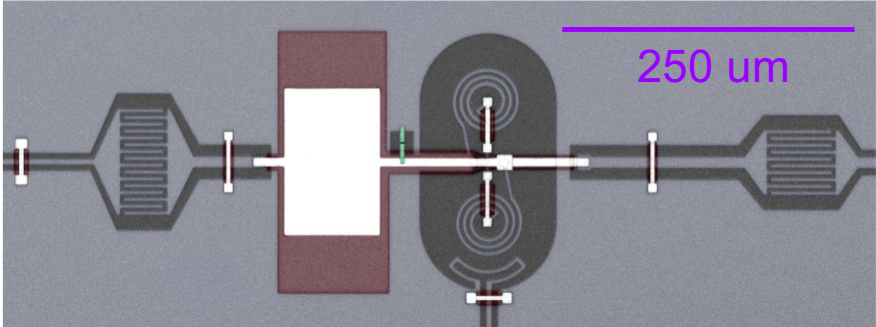
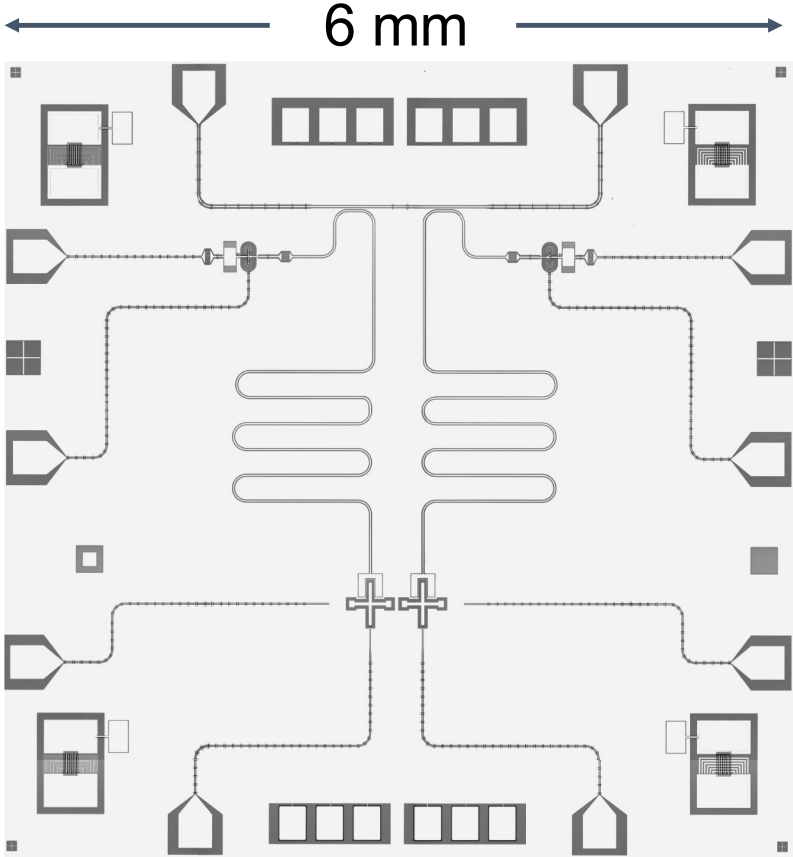


$$E_J = I_{0q}\Phi_0/2\pi$$

$$E_C = e^2/2C_\Sigma$$

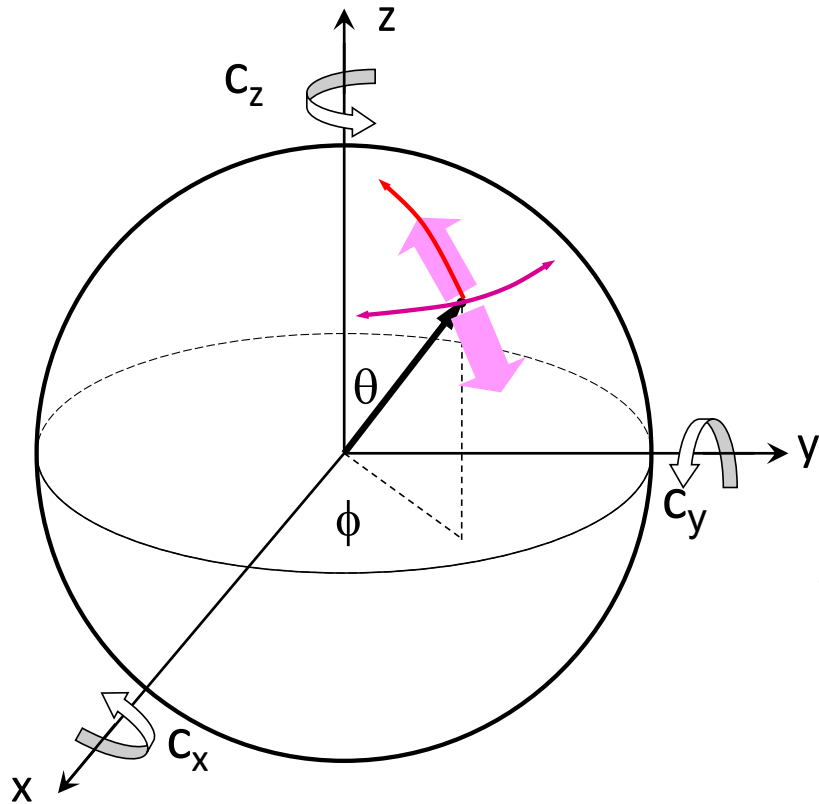
$$\omega_0 = \sqrt{8E_J E_C}/\hbar$$

# Chip layout



Opremcak *et al.*, PRX 11, 011027 (2021)

# Bloch sphere representation



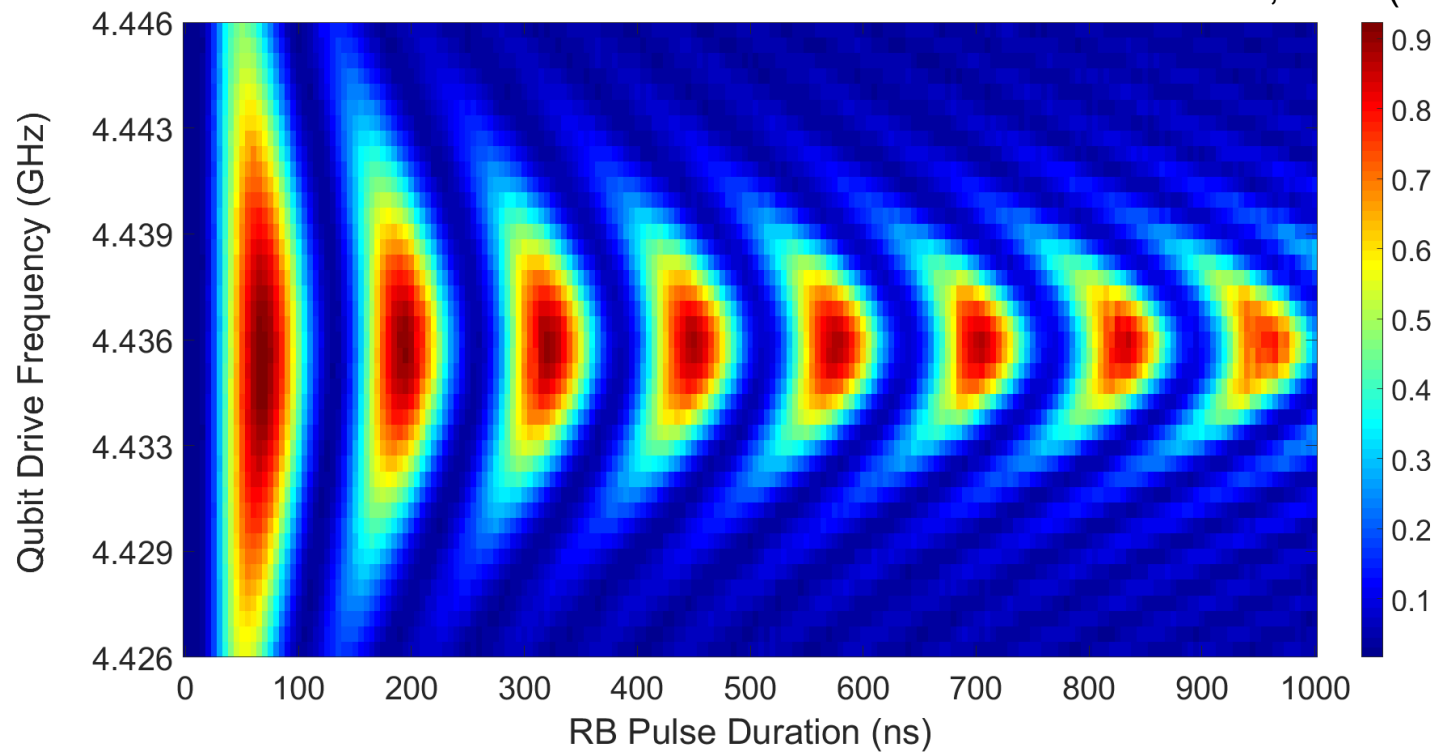
$$|\Psi\rangle = \cos \frac{\theta}{2} |0\rangle + e^{i\phi} \sin \frac{\theta}{2} |1\rangle$$

$$\hat{H} = \frac{\hat{Q}^2}{2C_\Sigma} - E_J \cos \left( \frac{2\pi\hat{\Phi}}{\Phi_0} \right) + \frac{C_{xy}}{C_\Sigma} V_d(t) \hat{Q}$$

controlled dynamically

# Qubit Rabi oscillations

Opremcak *et al.*,  
*Science* **361**, 1239 (2018)



# Towards fault-tolerant QC

The challenge:

- Pauli X and Z don't commute  
(impossible to simultaneously monitor bit and phase flips using naïve approach)

However:

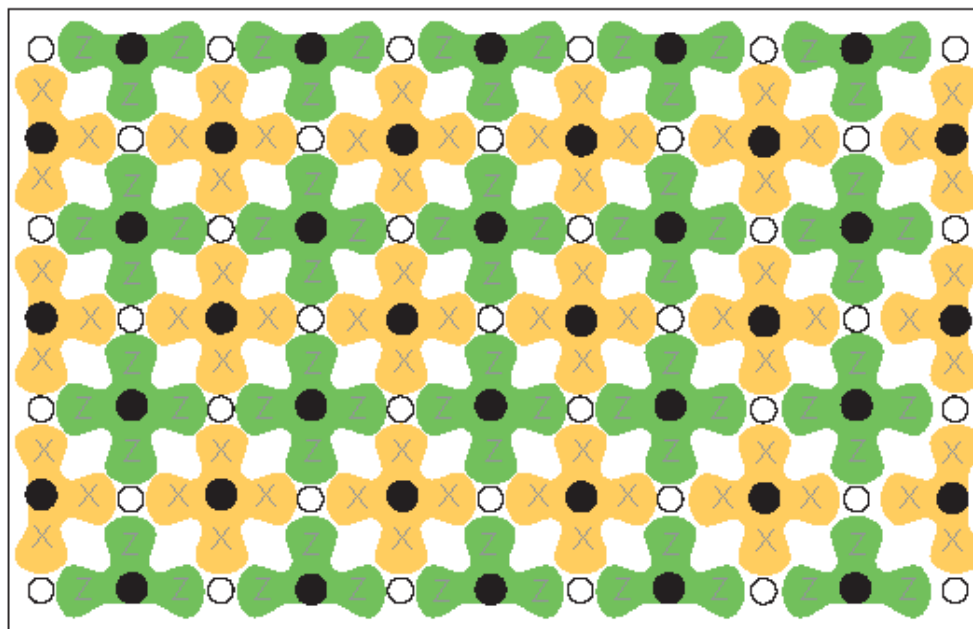
- Multiqubit parity operators (e.g., XXXX, ZZZZ) do commute!

possible in principle to uniquely identify errors in multiqubit arrays

# Towards fault-tolerant QC

Two-dimensional  
surface code

Fowler et al.,  
PRA 86, 032324  
(2012)



single QB gates: 99.9%  
entangling gates: 99%  
measurement: 90%

(generally target one order  
of magnitude lower infidelity)

Assumption: NO correlated errors



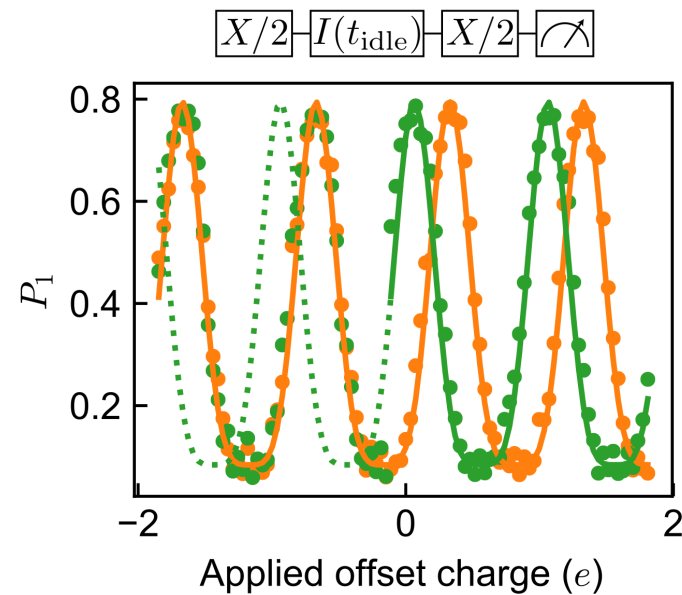
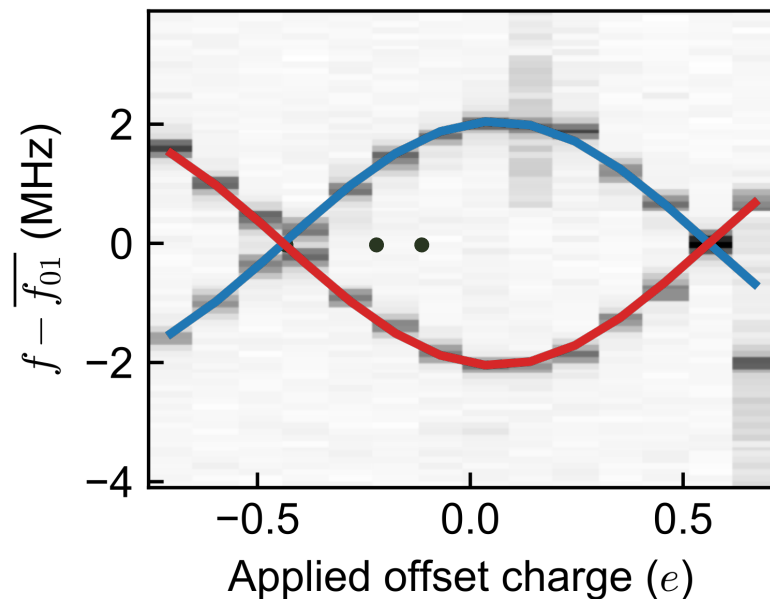
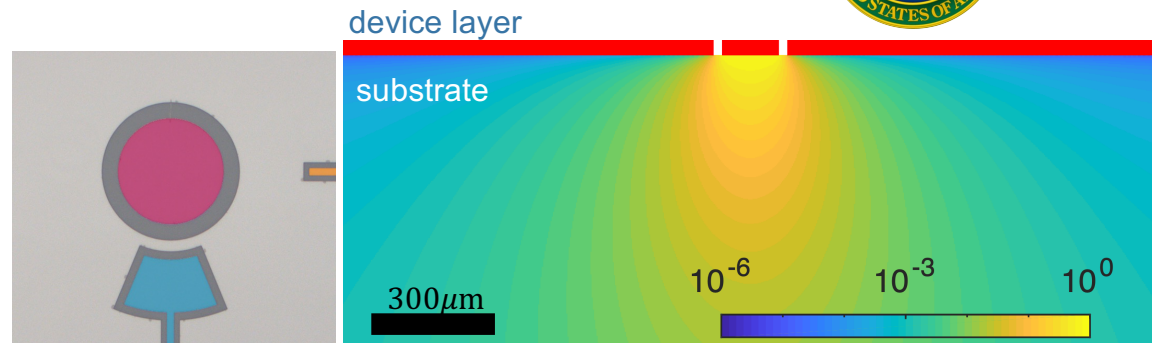
# Correlated errors in SC multiqubit circuits

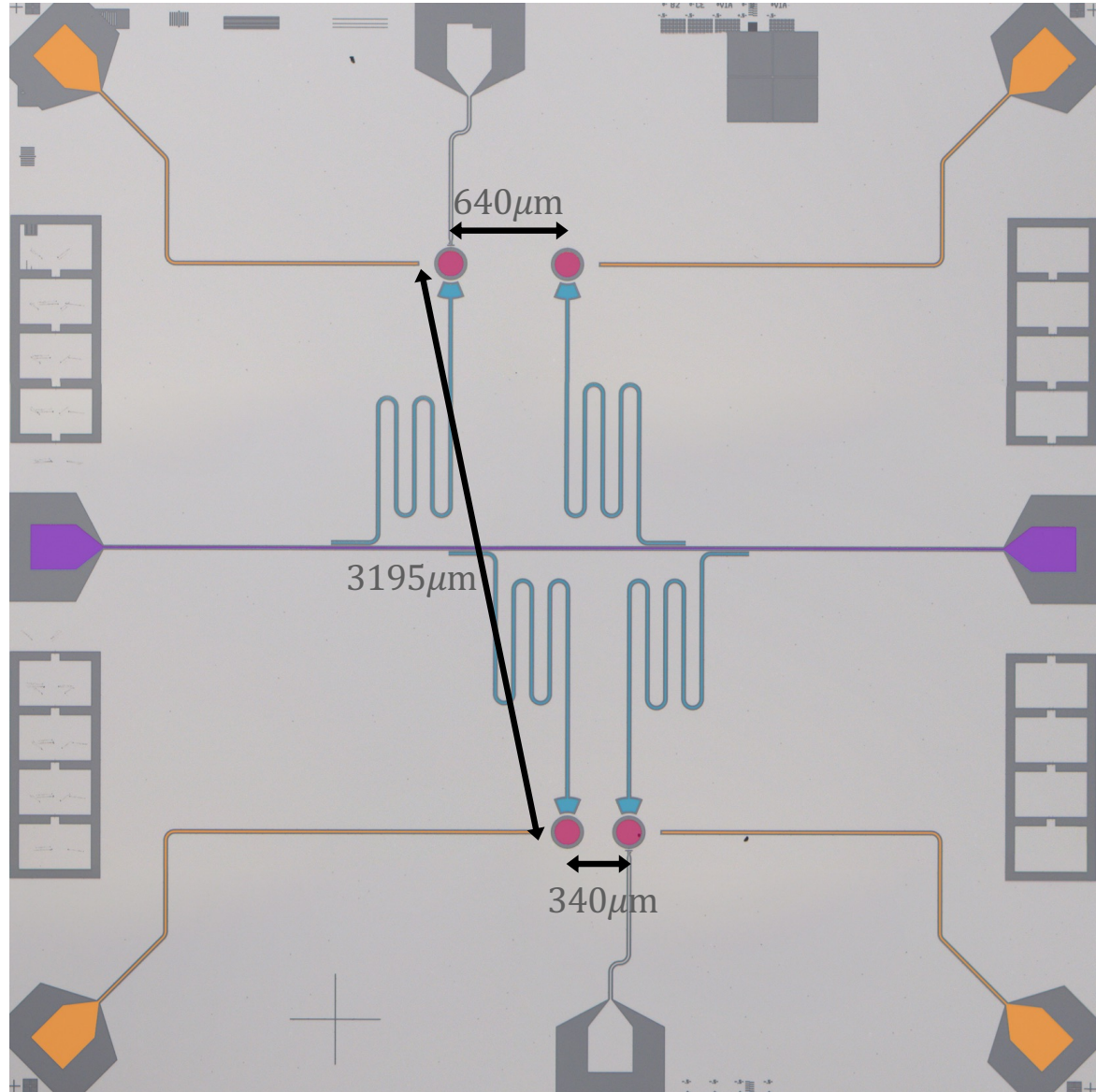
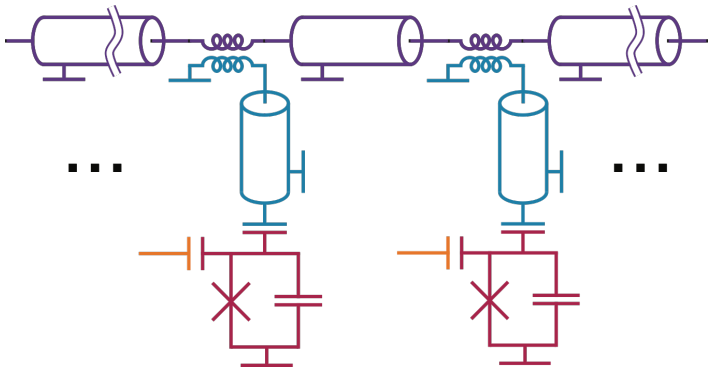
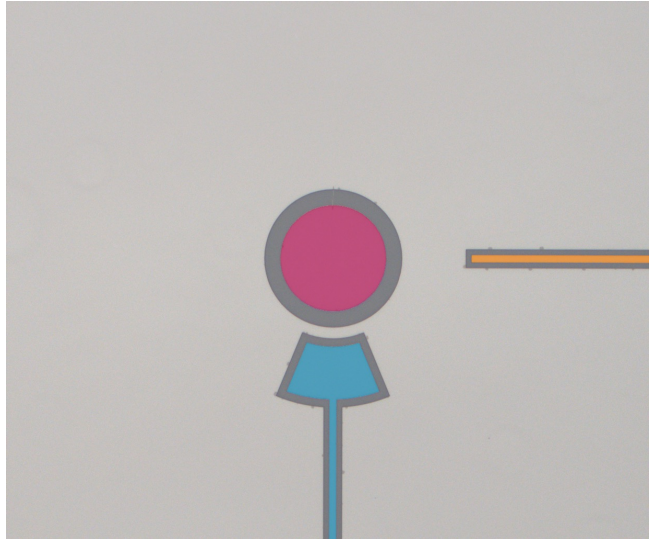
(prior DOE project)



Wilén *et al.*, Nature **594**, 7863 (2021)

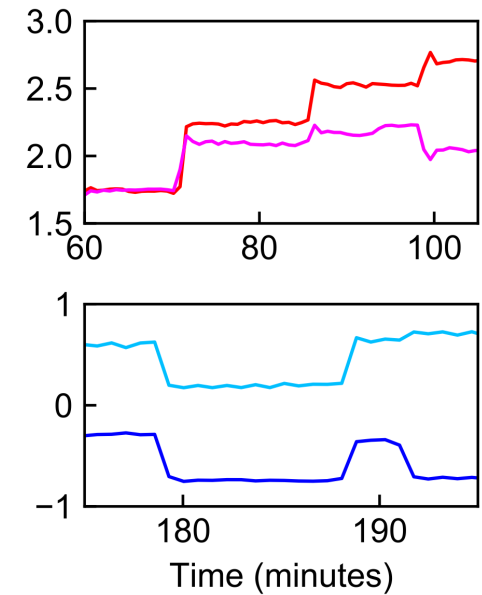
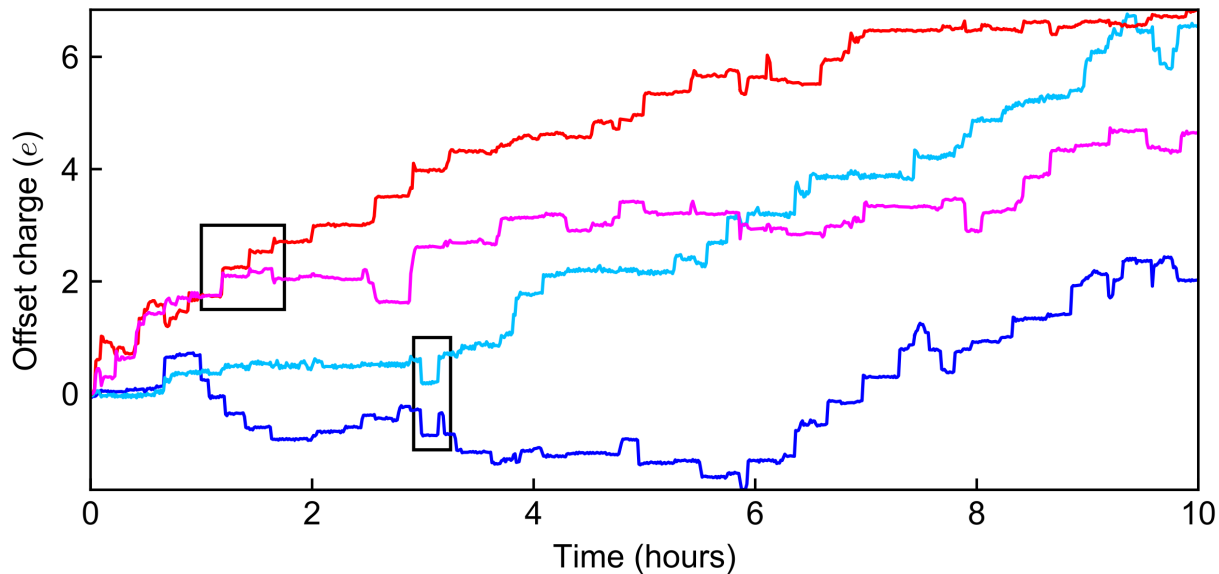
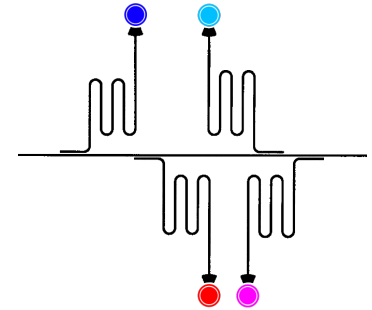
- Qubit as electrometer
- $E_J/E_C \sim 25$
- Ramsey-based charge tomography

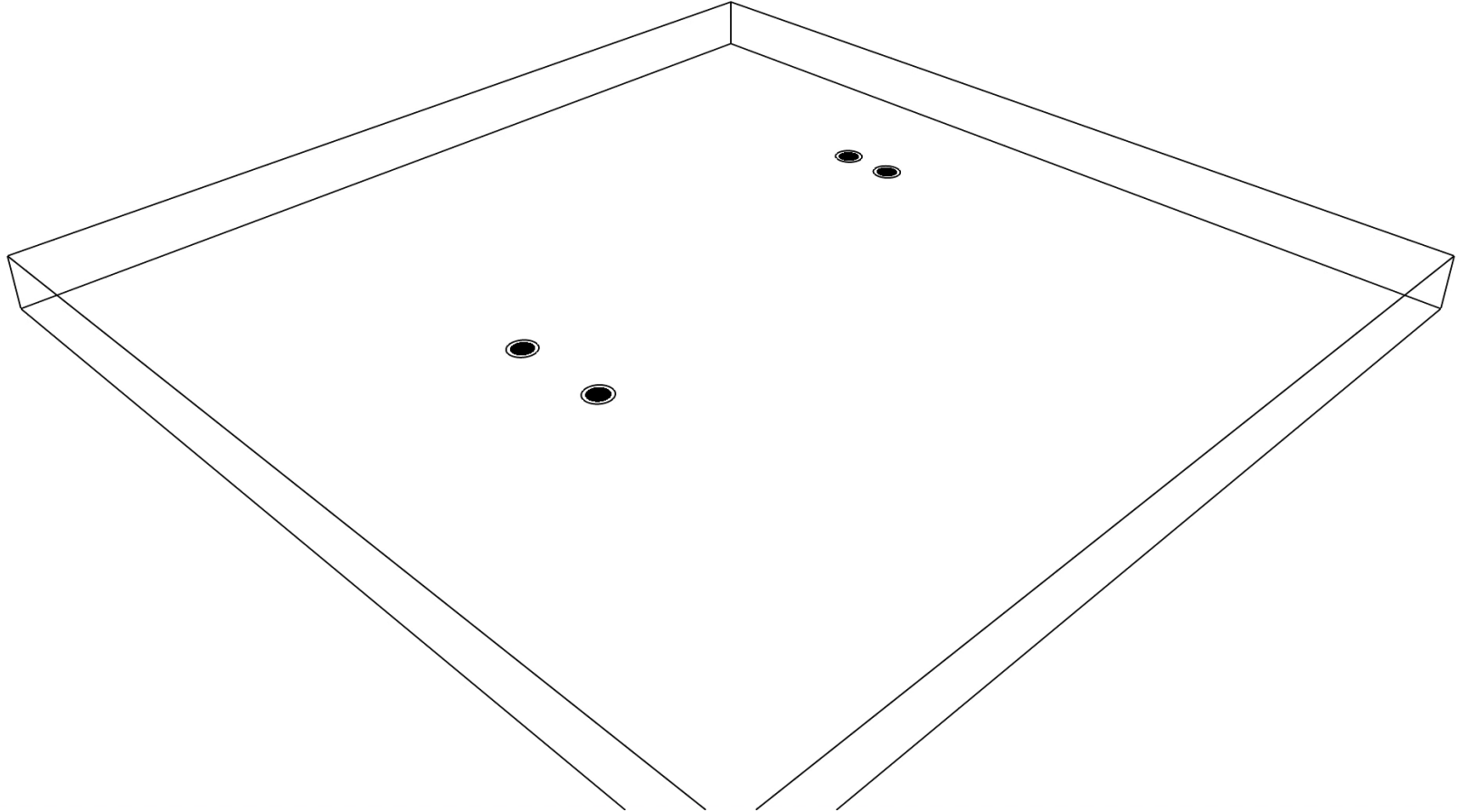




# Correlated charge fluctuations

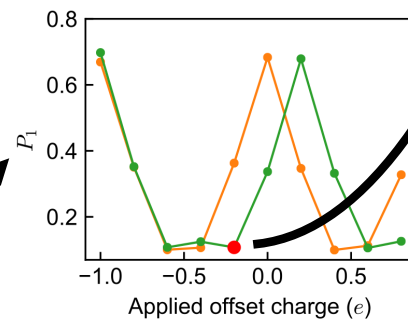
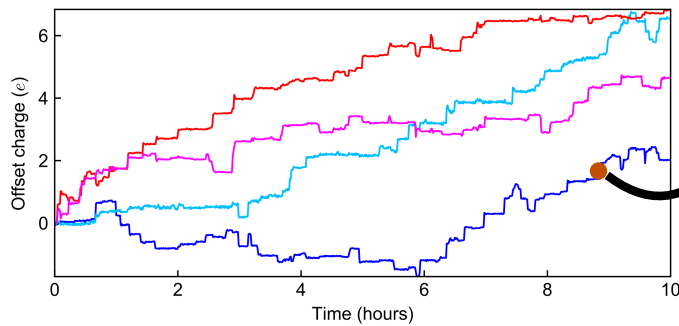
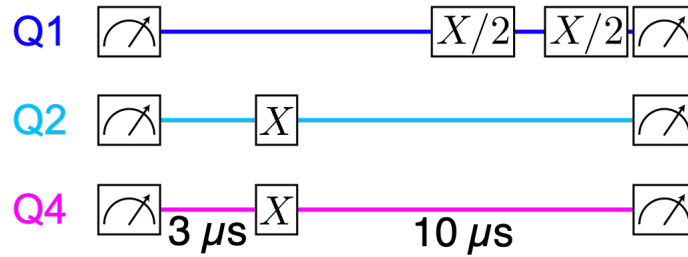
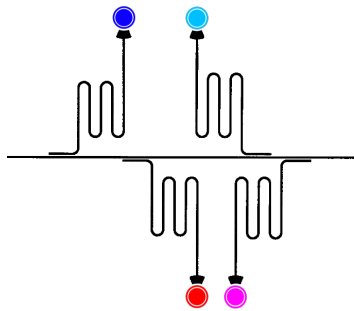
- Large discrete jumps  
Riste et al., Nat. Comm. **4**, 1913 (2013).  
Serniak et al., PRL **121**, 157701 (2018).  
Christensen et al., PRB **100**, 140503 (R)(2019).
- Correlated jumps in neighboring qubits





# Correlated relaxation errors

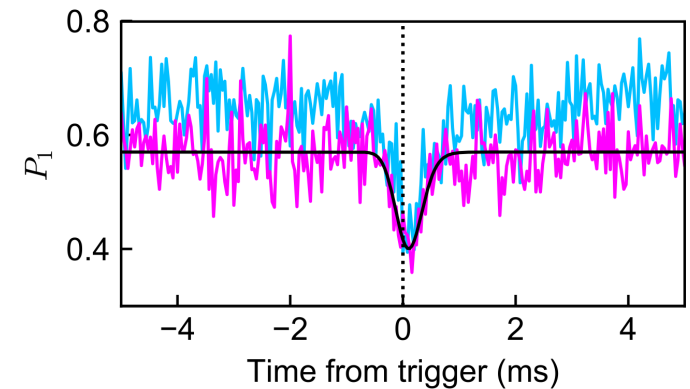
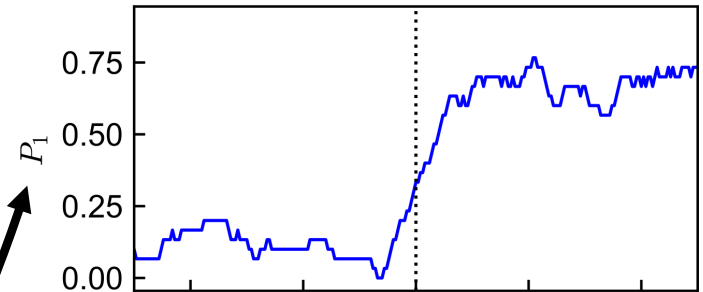
- Use jumps in Ramsey measurement as trigger
- Seen across distances > 3 mm



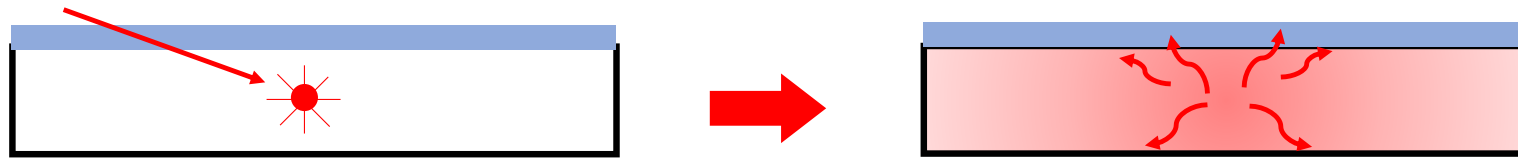
8k measurements

recovery timescale

$$\tau = 130 \pm 40 \mu s$$



# Phonon-mediated quasiparticle poisoning



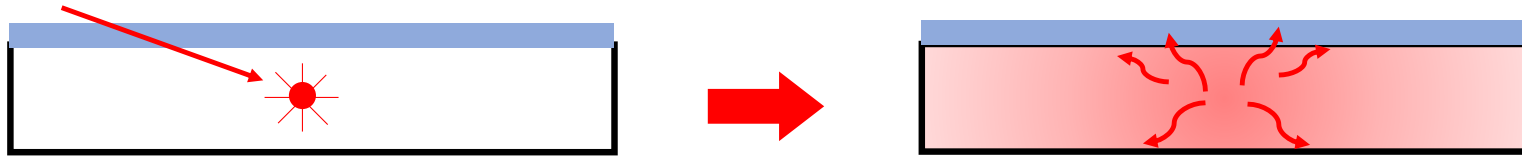
➤ 95% of deposited energy coupled to phonon reservoir  
[CDMS: e.g., Shutt et al., PRL 69, 3531 (1992)]

- Weak acoustic coupling to environment
- Few relaxation channels below SC gap edge
- Correlated pair-breaking across the chip

Possible mitigation strategies:

- Superconducting gap engineering
- Backside normal metallization
- Micromachining to suppress phonon transport
- Shielding to suppress absorption of HE particles

# Phonon-mediated quasiparticle poisoning



➤ 95% of deposited energy coupled to phonon reservoir  
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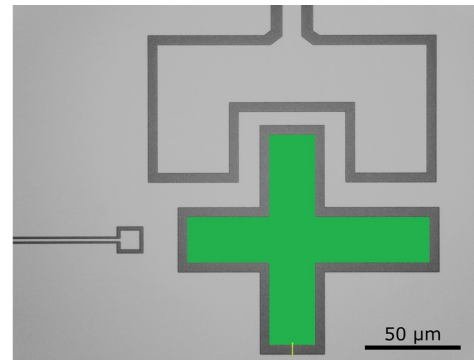
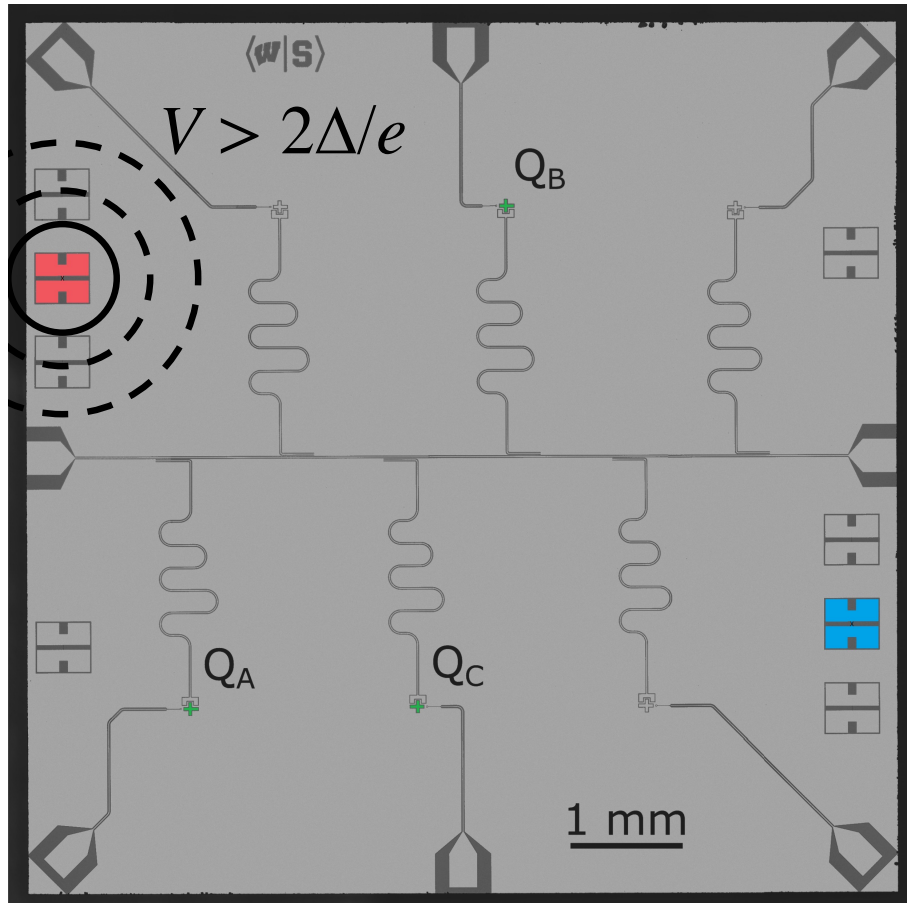
- Weak acoustic coupling to environment
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Possible mitigation strategies:

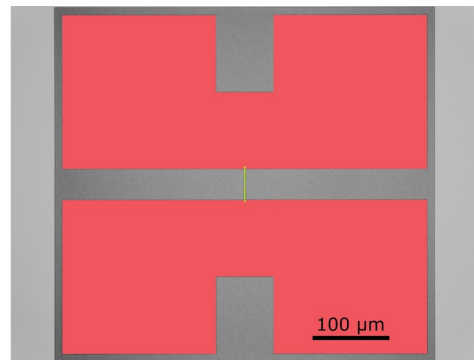
- Superconducting gap engineering
- **Backside normal metallization**
- Micromachining to suppress phonon transport
- Shielding to suppress absorption of HE particles

# Controlled study of phonon-mediated QP poisoning

(With B. Plourde group, Syracuse U.)



- Nb groundplane, 70 nm
- Al/AlOx/Al JJ, 360 x 150 nm<sup>2</sup>
- Cavities: 5.9 - 6.5 GHz
- Qubits: 4.9 - 5.1 GHz
- $E_J/E_C \sim 30$



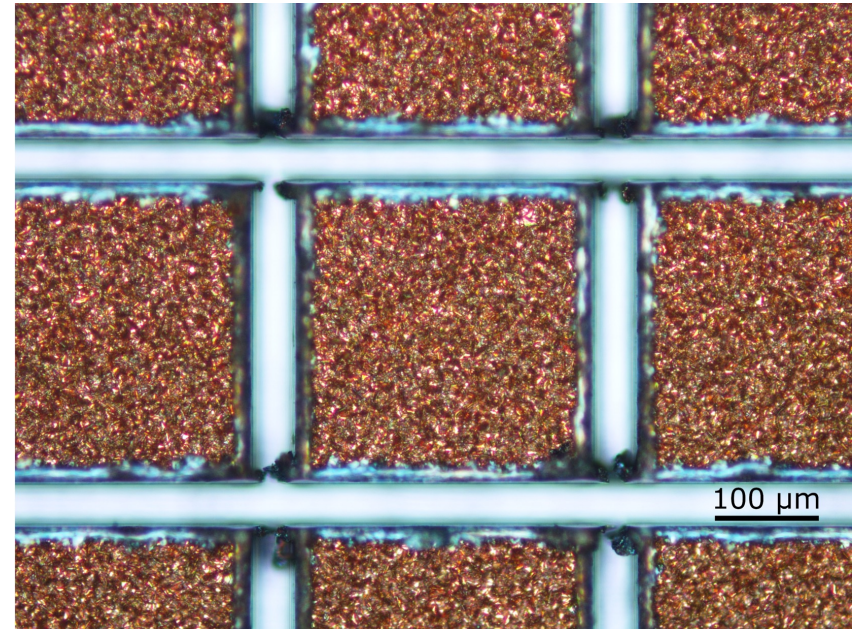
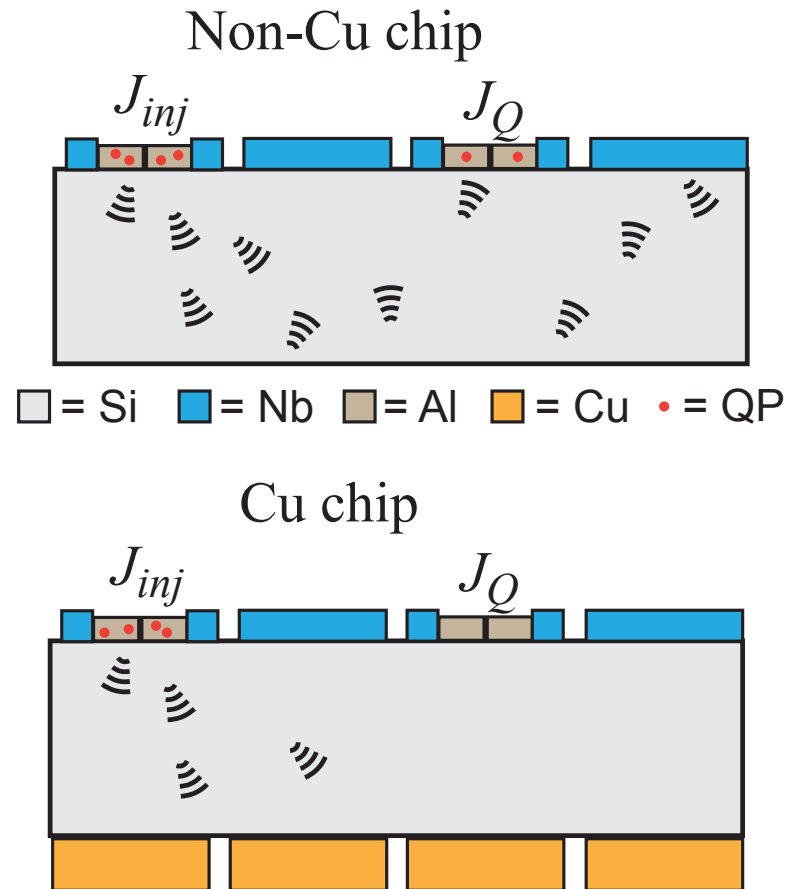
- Injector JJs: 1  $\mu\text{m}$  x 150 nm
- $I_c \sim 100$  nA
- No galvanic connection to groundplane



# Controlled study of phonon-mediated QP poisoning

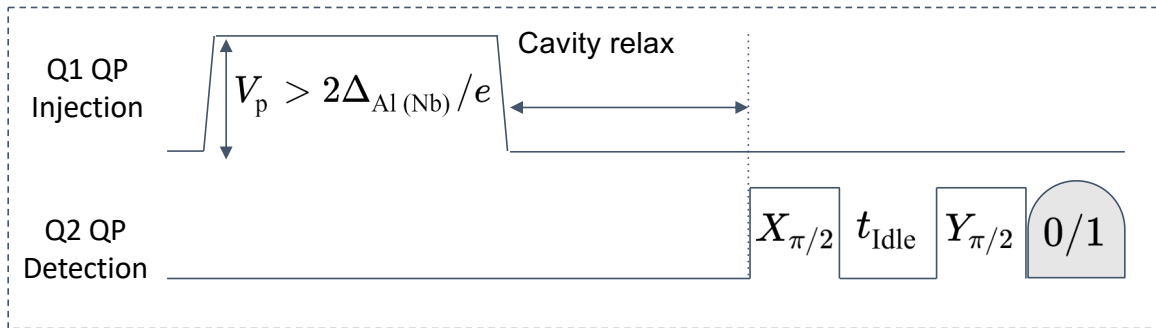
(With B. Plourde group, Syracuse U.)

Electroplated Cu  
10  $\mu\text{m}$  thick  
RRR  $\sim 42$

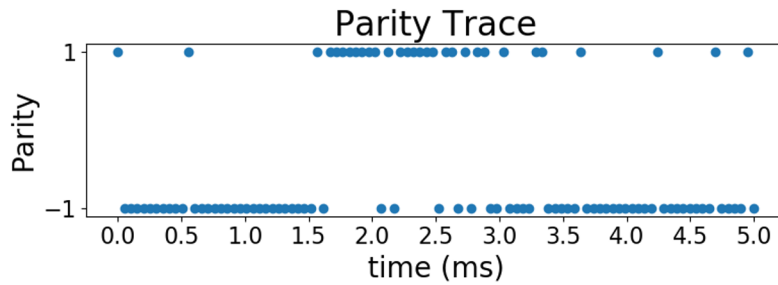


# Direct measurement of QP poisoning via charge parity

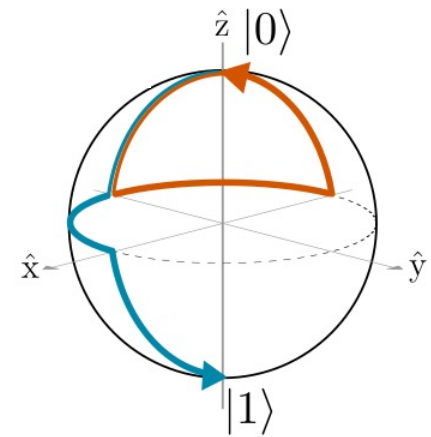
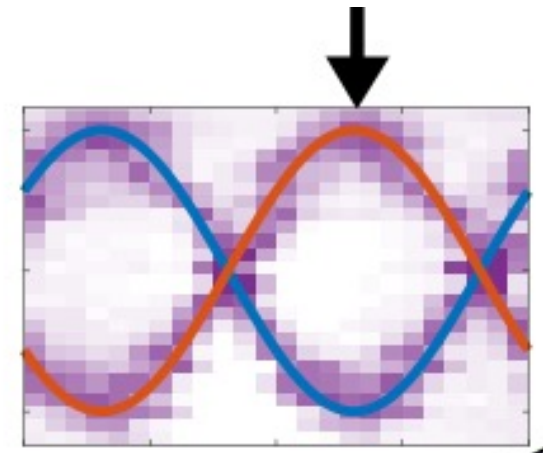
Riste et al., Nat. Comm. **4**, 1913 (2013)  
 Serniak et al., PRL **121**, 157701 (2018)



$\times 10000$

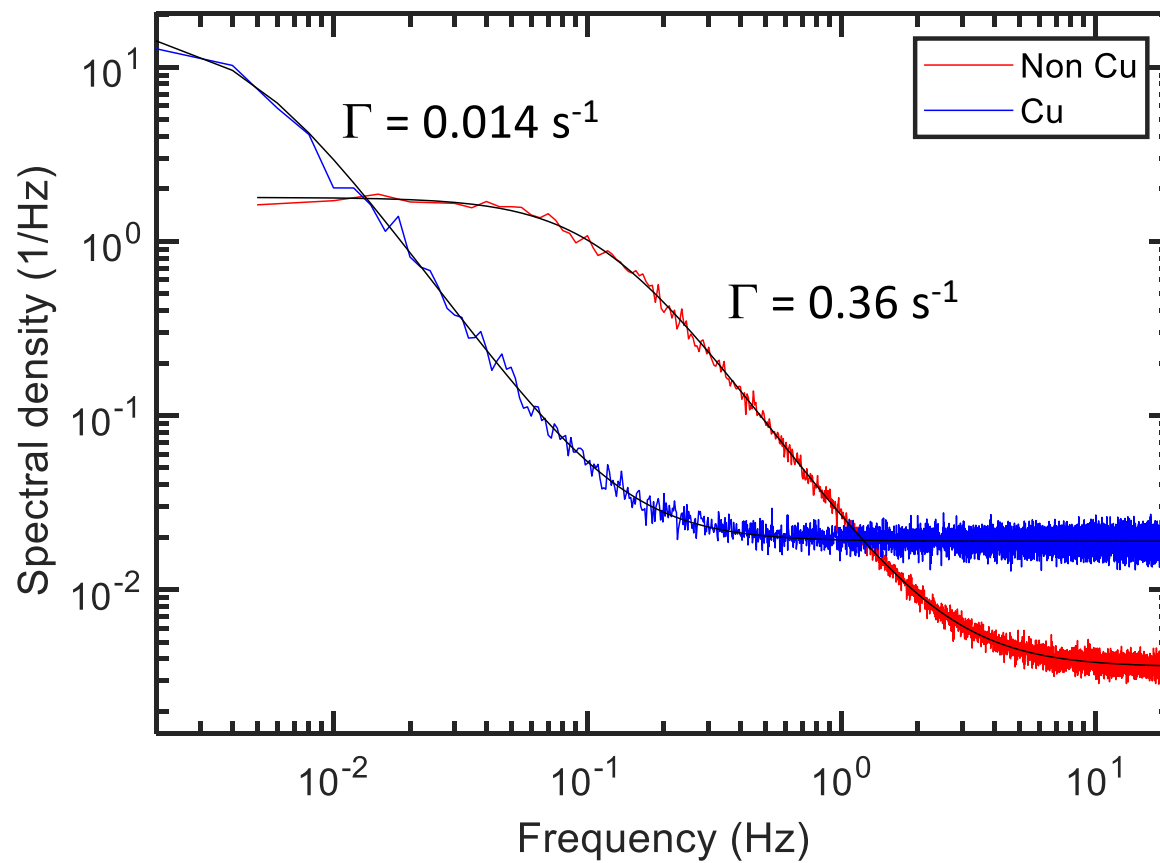


bias at maximum charge dispersion

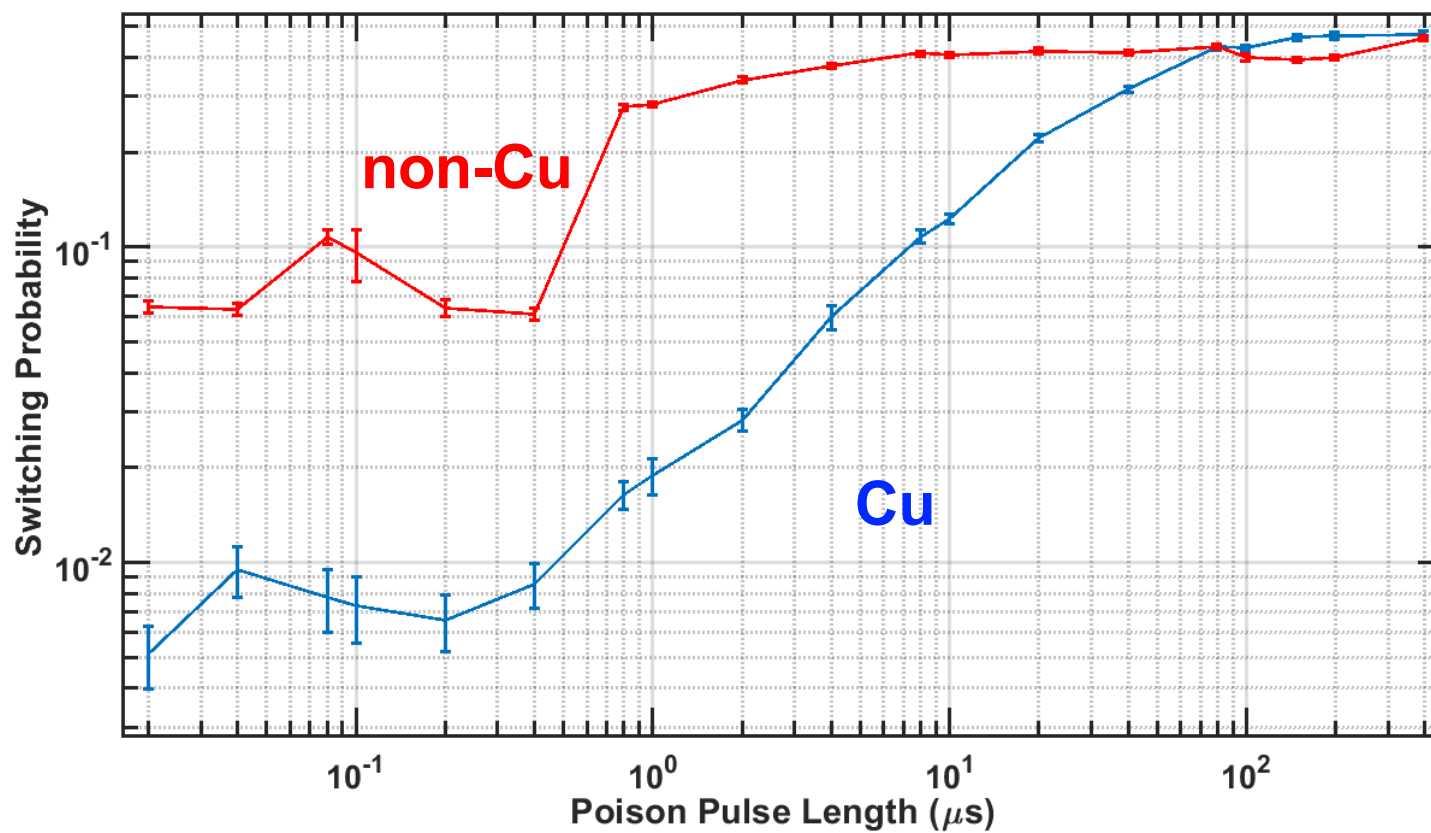


# Direct measurement of QP poisoning via charge parity

Baseline:  
no QP injection



# Simultaneous parity monitoring with phonon injection



## No injection: suppression of correlated error from impact events

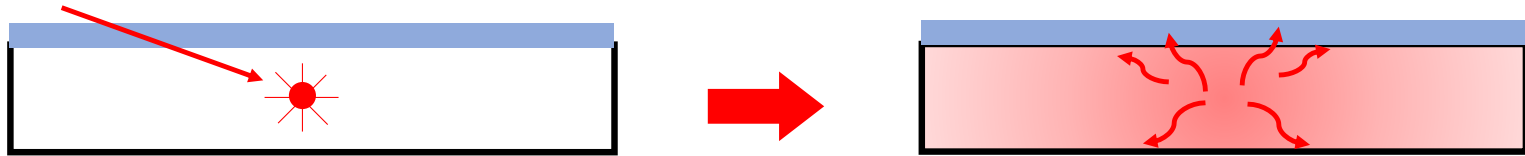
Observed parity switching rates ( $s^{-1}$ )							
Device	$Q_A$	$Q_B$	$Q_C$	$Q_A \wedge Q_B$	$Q_B \wedge Q_C$	$Q_A \wedge Q_C$	$Q_A \wedge Q_B \wedge Q_C$
Non-Cu	0.309(3)	0.313(4)	0.301(3)	0.073(2)	0.068(2)	0.064(2)	0.0134(9)
Cu	0.0250(4)	0.0344(5)	0.0545(6)	0.0014(1)	0.00357(2)	0.0062(3)	0.00012(6)
<i>spacing</i>	-	-	-	5.3 mm	4.5 mm	2.0 mm	-

Expected random coincident parity switching rates ( $s^{-1}$ )				
Device	$Q_A \wedge Q_B$	$Q_B \wedge Q_C$	$Q_A \wedge Q_C$	$Q_A \wedge Q_B \wedge Q_C$
Non-Cu	0.039(5)	0.038(5)	0.037(5)	0.005(6)
Cu	0.0003(6)	0.0007(8)	0.0005(7)	0.000007(852)

$$(r_A^{\text{obs}} \Delta t)(r_B^{\text{obs}} \Delta t)/\Delta t \quad (r_A^{\text{obs}} \Delta t)(r_B^{\text{obs}} \Delta t)(r_C^{\text{obs}} \Delta t)/\Delta t$$

Extracted poisoning rates ( $s^{-1}$ )							
Device	$Q_A$	$Q_B$	$Q_C$	$Q_A \wedge Q_B$	$Q_B \wedge Q_C$	$Q_A \wedge Q_C$	$Q_A \wedge Q_B \wedge Q_C$
Non-Cu	0.16(2)	0.16(2)	0.16(1)	0.21(1)	0.19(1)	0.18(1)	0.070(6)
Cu	0.021(2)	0.051(2)	0.072(2)	0.0049(7)	0.013(1)	0.024(1)	0.0002(5)
<i>spacing</i>	-	-	-	5.3 mm	4.5 mm	2.0 mm	-

# Are particle impacts the source of nonequilibrium QPs?



Density of QPs in SC devices far exceeds thermal equilibrium value expected from theory (typ.  $1 \mu\text{m}^{-3}$ )

- Poisoning of charge-sensitive devices [cf. Aumentado et al., PRL **92**, 066802 (2004)].
- Enhanced qubit relaxation [Corcoles et al., APL **99**, 181906 (2011)].
- Qubit initialization errors of order 1-3%

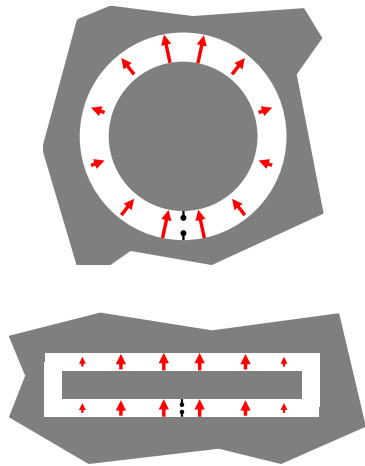
MIT/LL group: model for QP poisoning from gamma impacts [Vepsalainen et al., Nature **584**, 551 (2020)]

However, event rate too low and escape rate of athermal phonons too high to account for measured QP densities in most devices

# Antenna modes of the transmon qubit

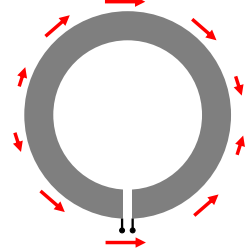
Rafferty *et al.*, arXiv:2103.06803 (2021)

Typ. SC  
qubit structures

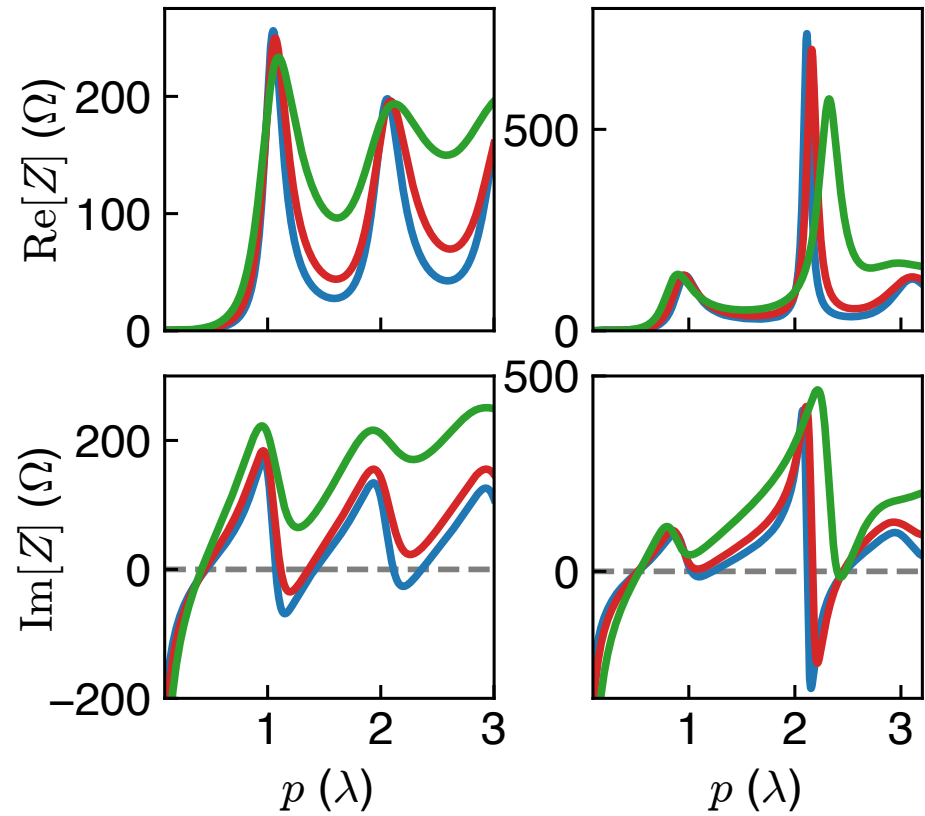
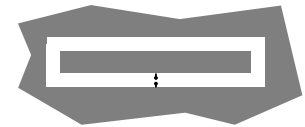
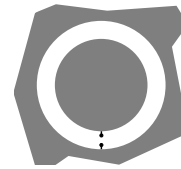
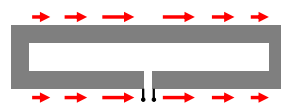


Wire antenna  
duals

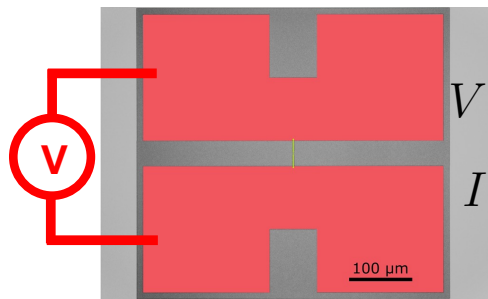
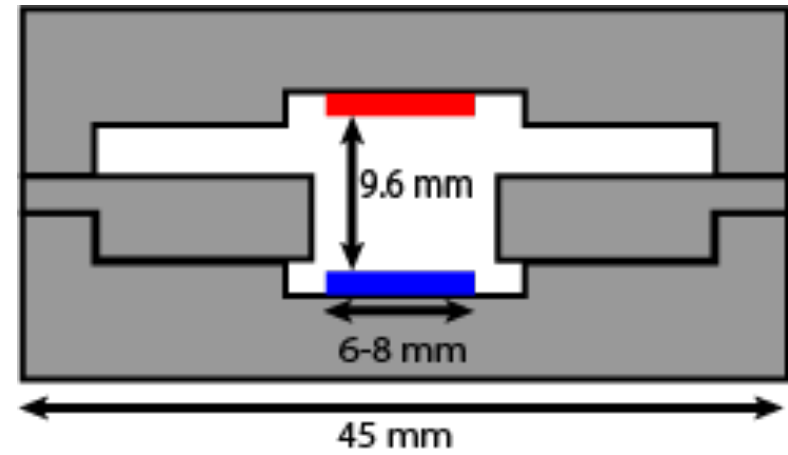
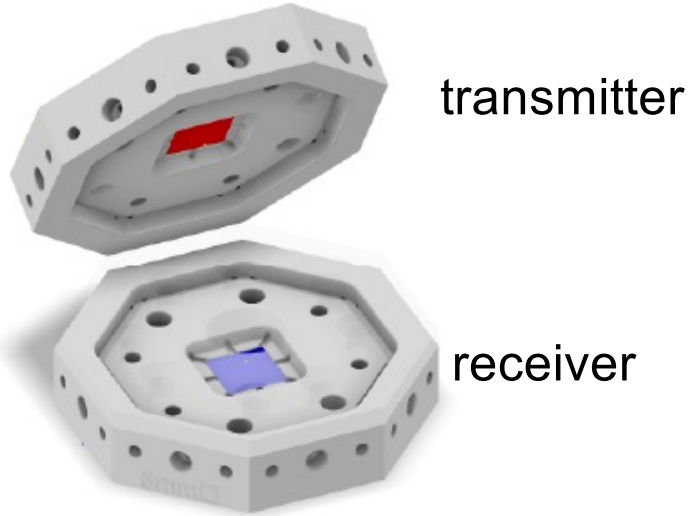
*resonant loop  
antenna*



*folded dipole  
antenna*

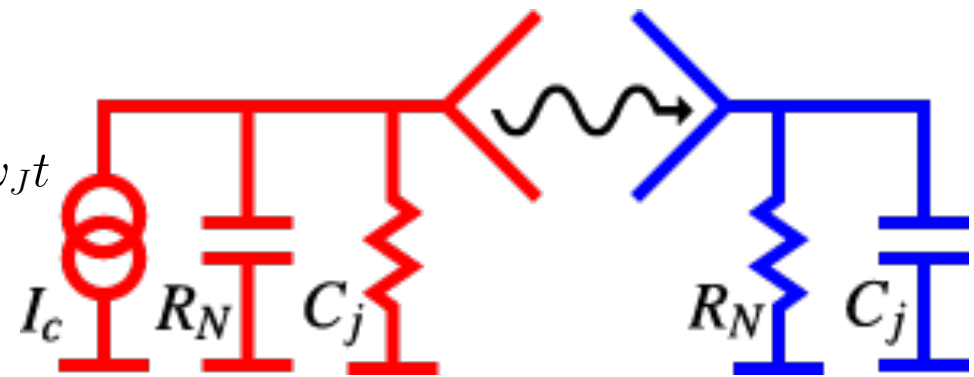


# Validation of antenna model



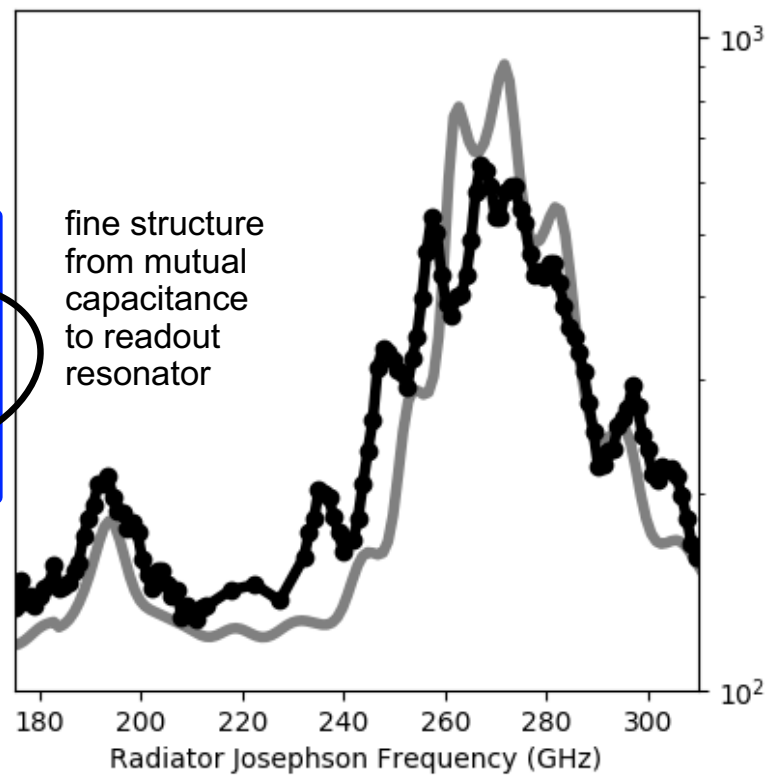
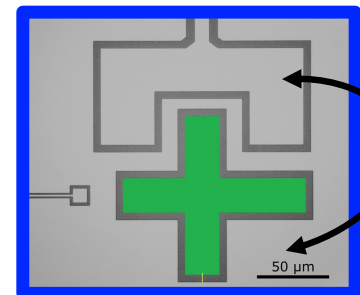
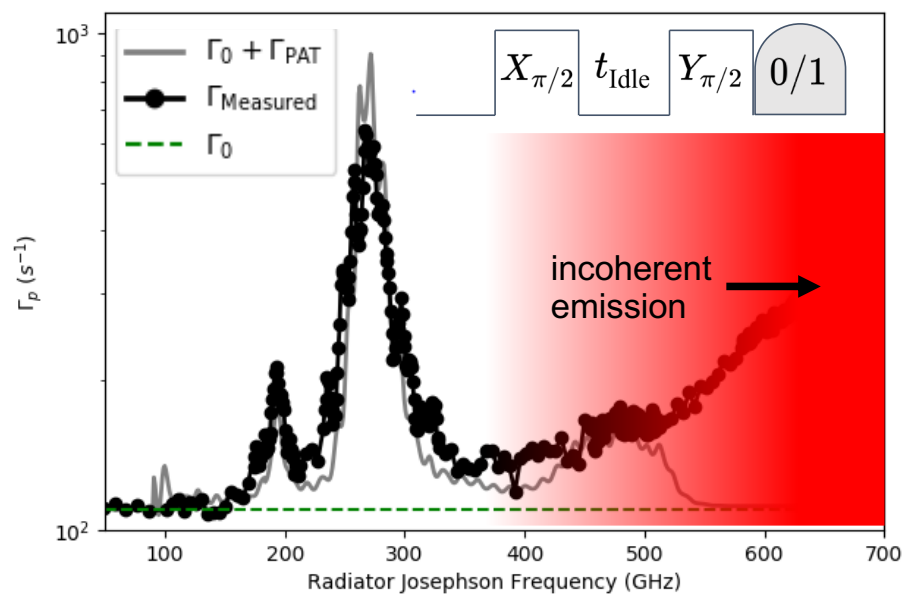
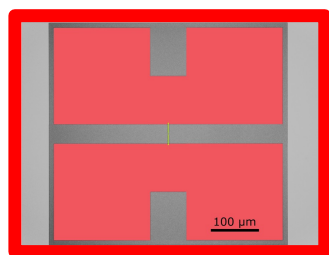
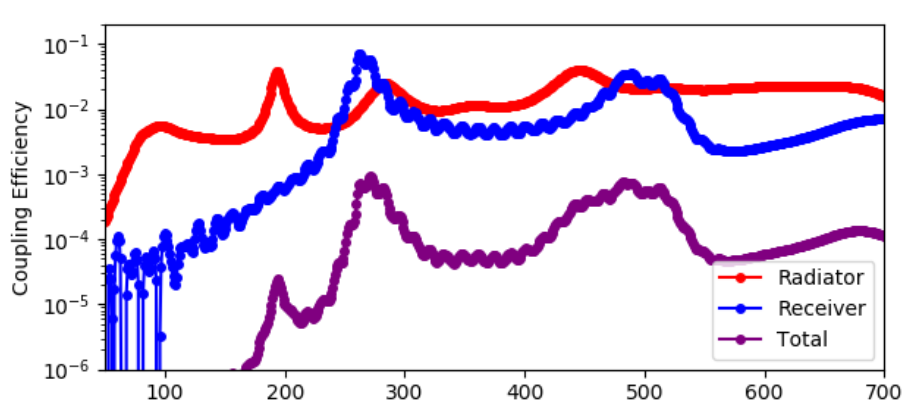
$$V = \frac{\Phi_0}{2\pi} \dot{\delta}$$

$$I = I_c \sin \delta \quad \Rightarrow \quad \begin{aligned} I &= I_c \sin \omega_J t \\ f_J &= V/\Phi_0 \end{aligned}$$

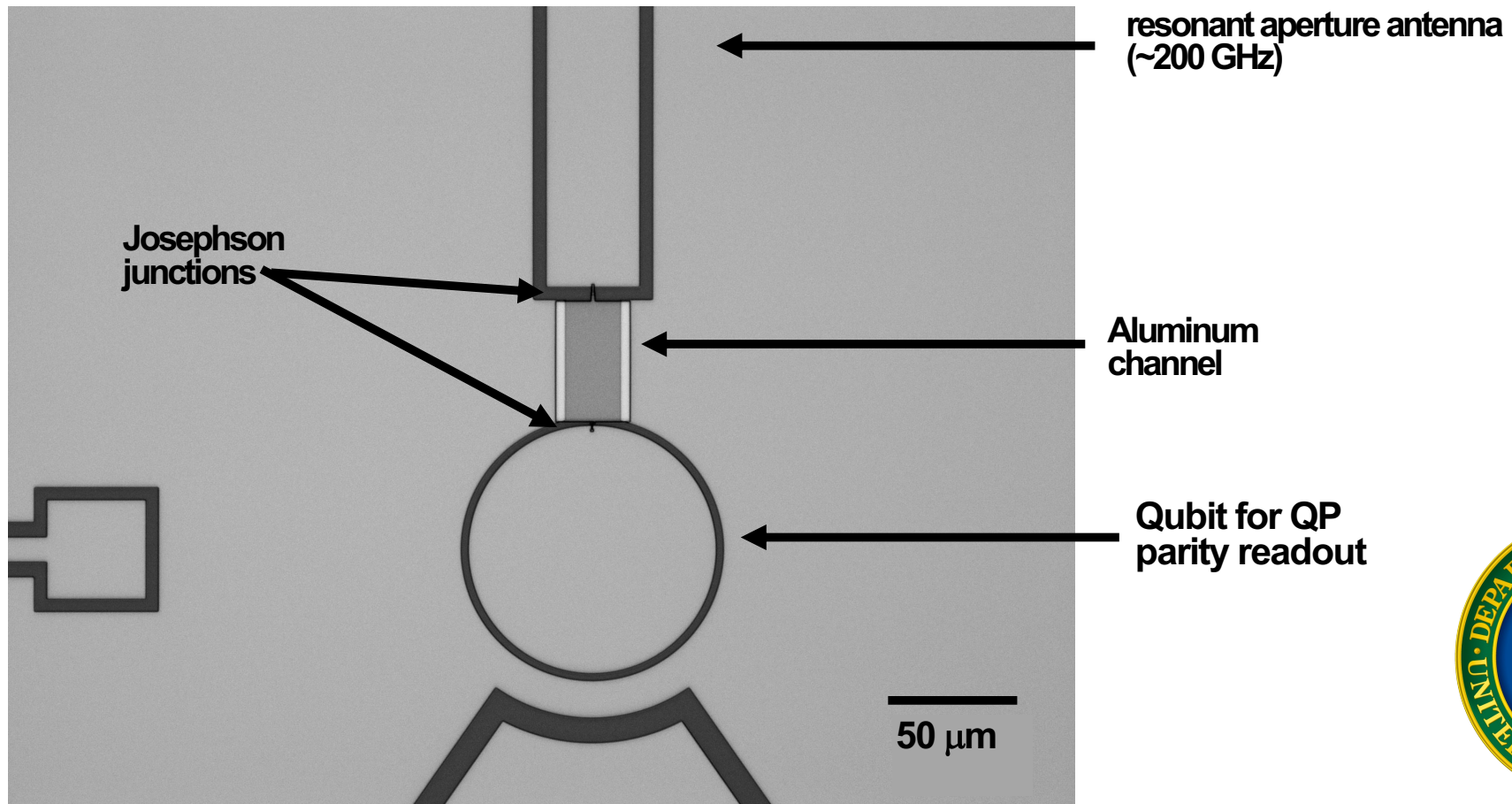




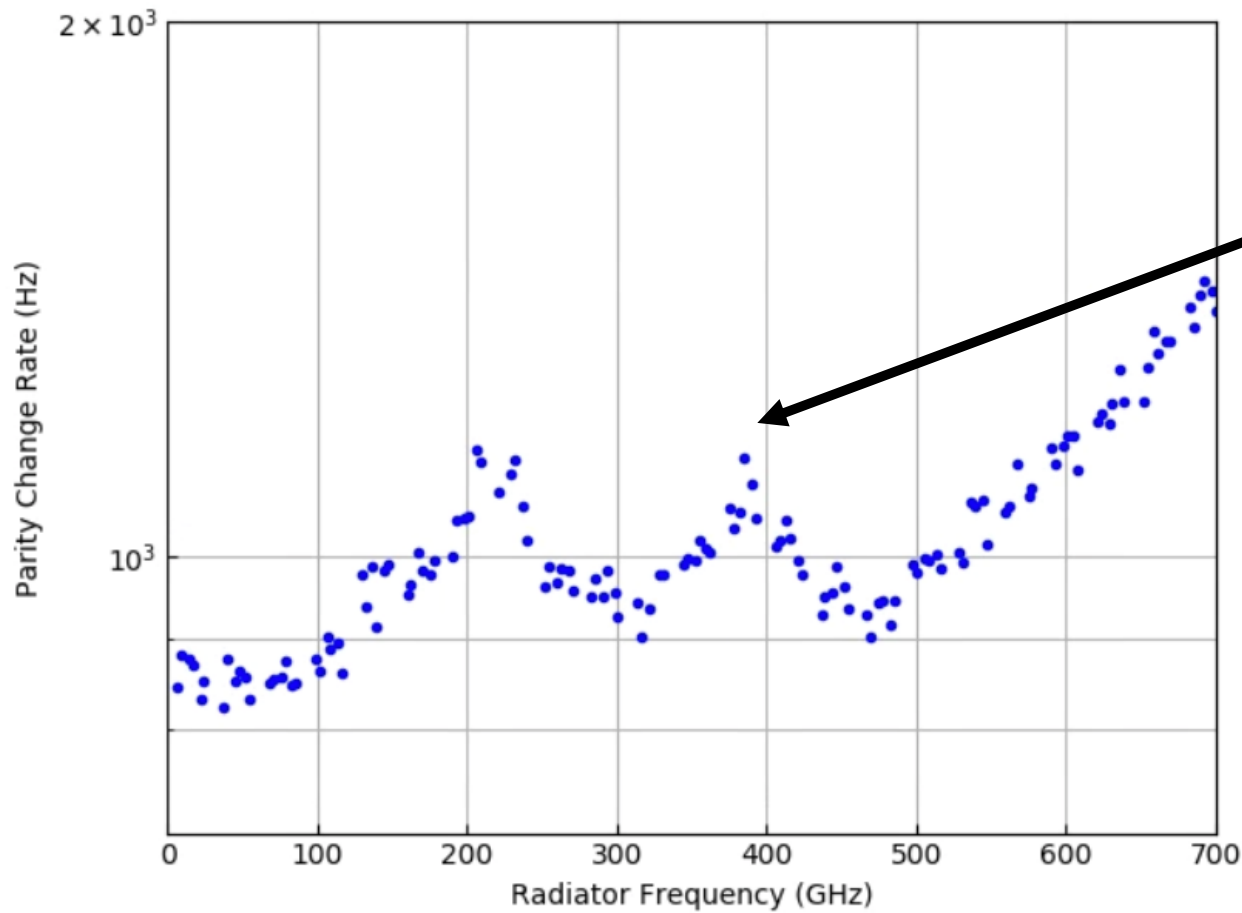
# Validation of antenna model



# Modular detector for axion search



# Modular detector for axion search



# Acknowledgments

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S. Abdullah (GS)

F. Schlenker (Res. Sci.)

C. D. Wilen (GS)

# Conclusion

- Particle impacts give rise to damaging correlated errors on length scales greater than 3 mm: relaxation errors must be mitigated
- Phonon downconversion yields dramatic suppression of correlated parity jumps
- Detailed understanding of photon-assisted QP poisoning
- New class of modular detectors for dark-matter searches

