

The kinetic gas universe

-

Lifting the Einstein Vlasov system to the tangent bundle

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University of Tartu

Joint work with N. Voicu and M. Hohmann

Virtual Conference of the Polish Society on Relativity 2020



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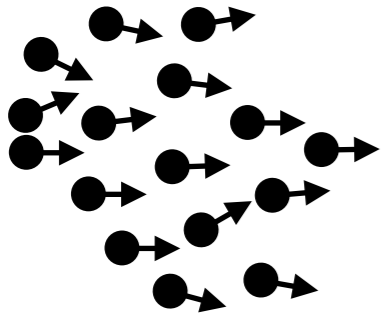
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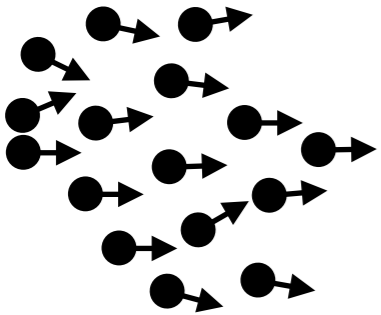
The gravitating kinetic gas

The kinetic gas

The kinetic gas



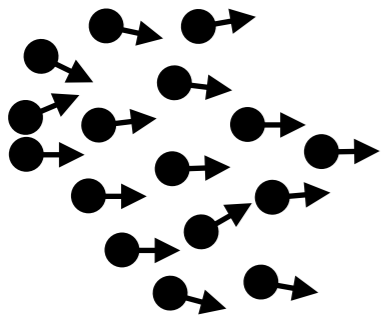
The kinetic gas



$$\phi(x, \dot{x})$$

$$\frac{\partial \phi}{\partial x^a} \dot{x}^a + \frac{\partial \phi}{\partial \dot{x}^a} \ddot{x}^a = C$$

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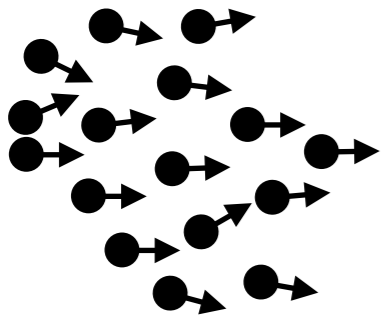
Einstein Vlasov equations

$$R^{ab} - \frac{1}{2} g^{ab} R = \frac{8\pi G}{c^4} T^{ab}$$
$$T^{ab} = \int_{V_x} \frac{\dot{x}^a \dot{x}^b}{g_{cd} \dot{x}^c \dot{x}^d} \phi(x, \dot{x}) \Sigma_x$$

[Andreasson 2011]

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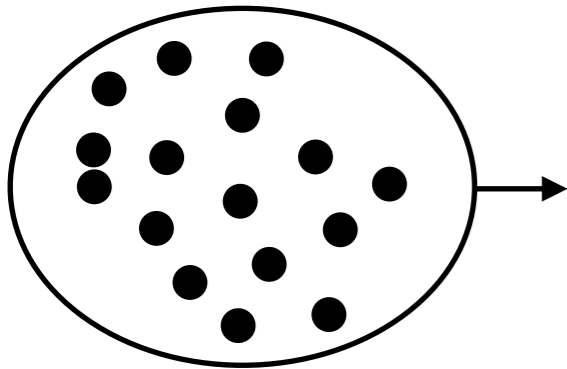
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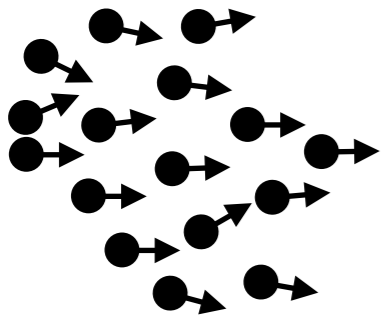
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The fluid



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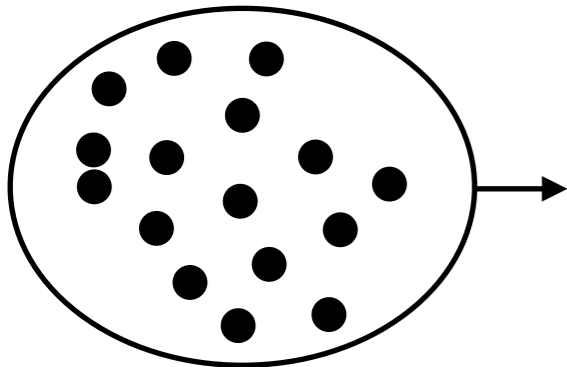
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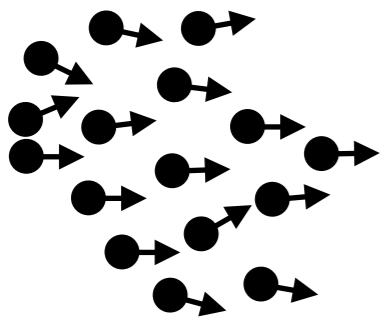
The fluid



$$T^{ab} = \alpha(x) U^a U^b + \beta(x) g^{ab}$$
$$\nabla_a T^{ab} = 0$$

The gravitating kinetic gas

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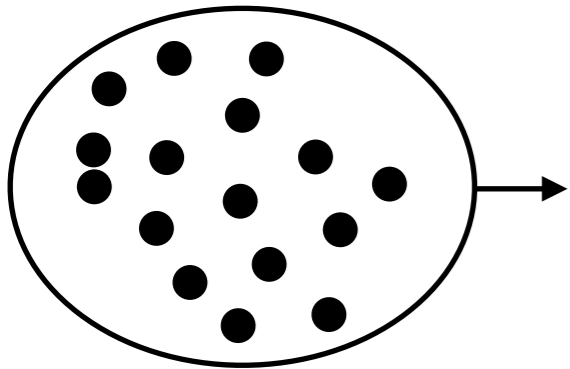
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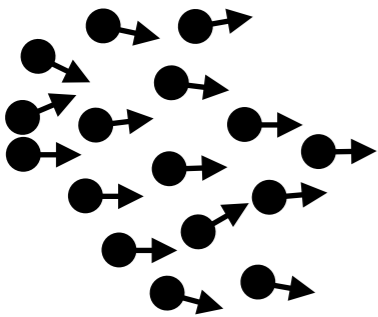
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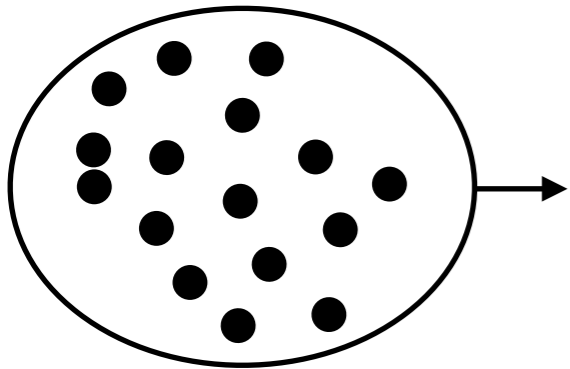
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the universe

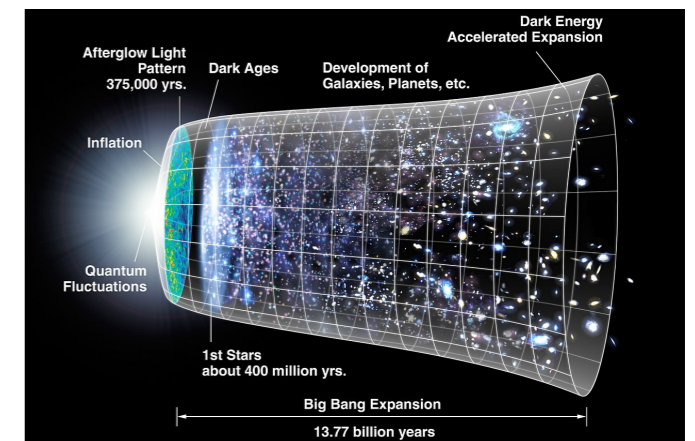
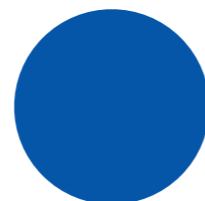
gas planets



accretion discs

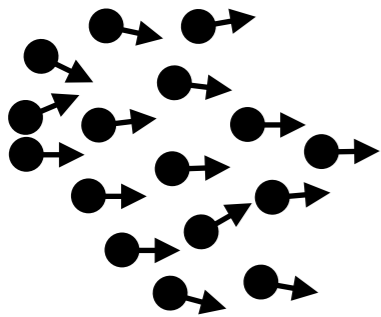


neutron stars



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Information is lost through velocity averaging!

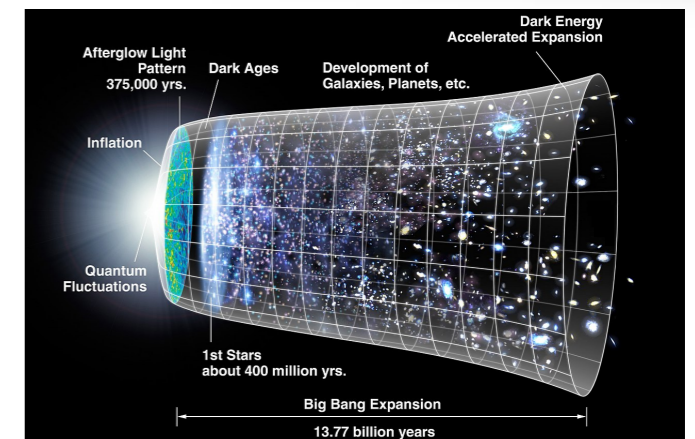
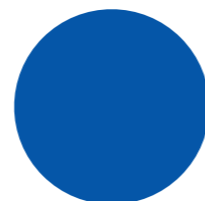
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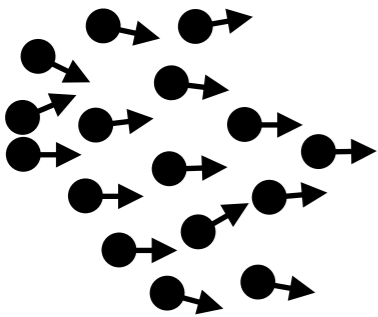


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Information is lost through velocity averaging!

This talk: The gravitational field of kinetic gases without velocity averaging

$$G(x, \dot{x}) = \hat{\kappa}^2 \phi(x, \dot{x})$$

Taking the velocity distribution of the gas particles into account!

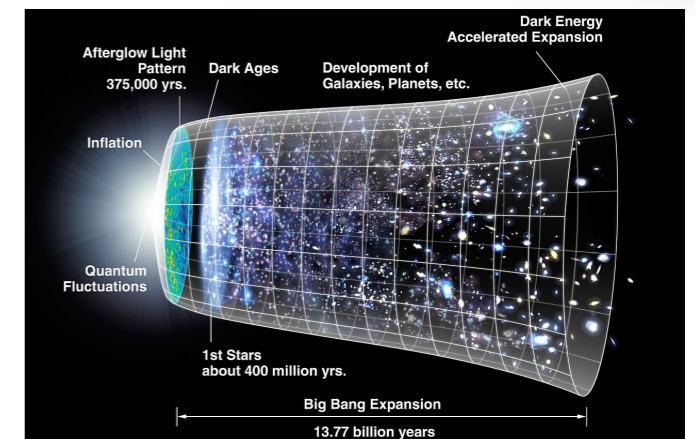
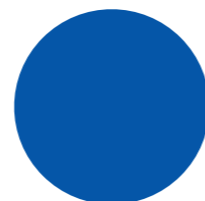
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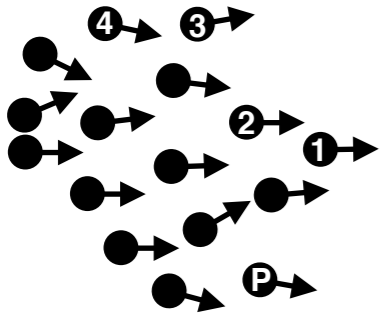


neutron stars



Point particles and Finsler geometry

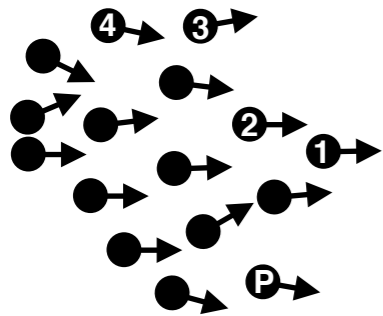
The relativistic gas



Individual gas particles

$$\begin{aligned} S[x_I] &= m \int_{s_1}^{s_2} ds_I \sqrt{L(x_I, \dot{x}_I)} \\ &= m \int_0^{t_I} d\tau_I = mt_I \end{aligned}$$

The relativistic gas



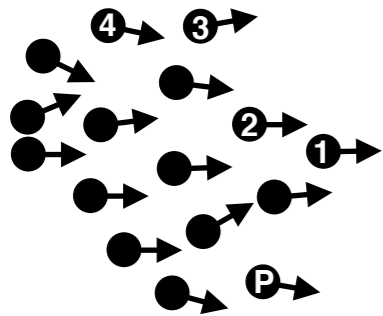
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Finsler geometry

Geometry defined by $L(x, \dot{x})$

The relativistic gas



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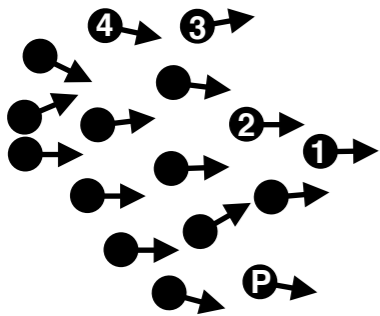
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⇒ parametrization invariant action

The relativistic gas



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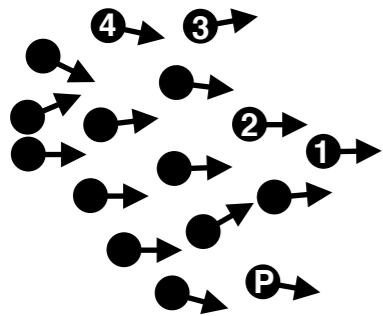
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The relativistic gas



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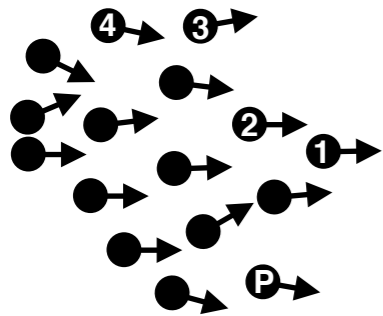
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The relativistic gas



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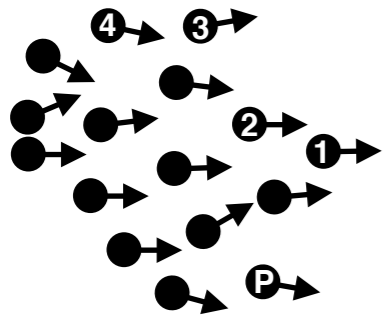
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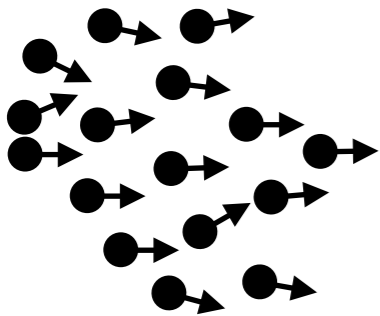
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The relativistic gas

Kinetic gas

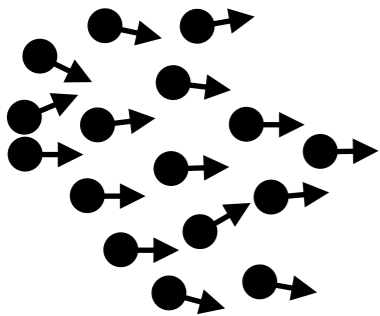


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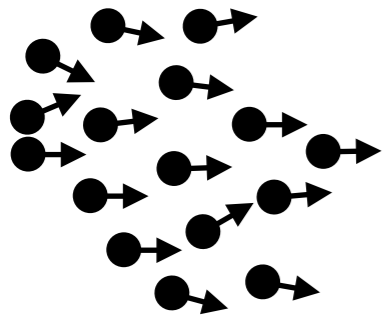
$$S_{gas} = m \sum_{i=1}^P \int_0^t ds_i \sqrt{L(x_I, \dot{x}_I)}$$

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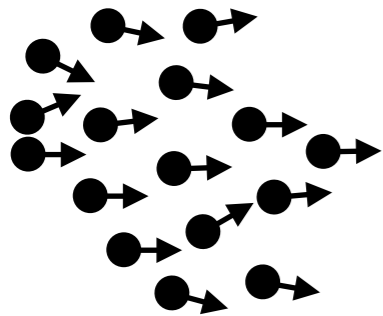
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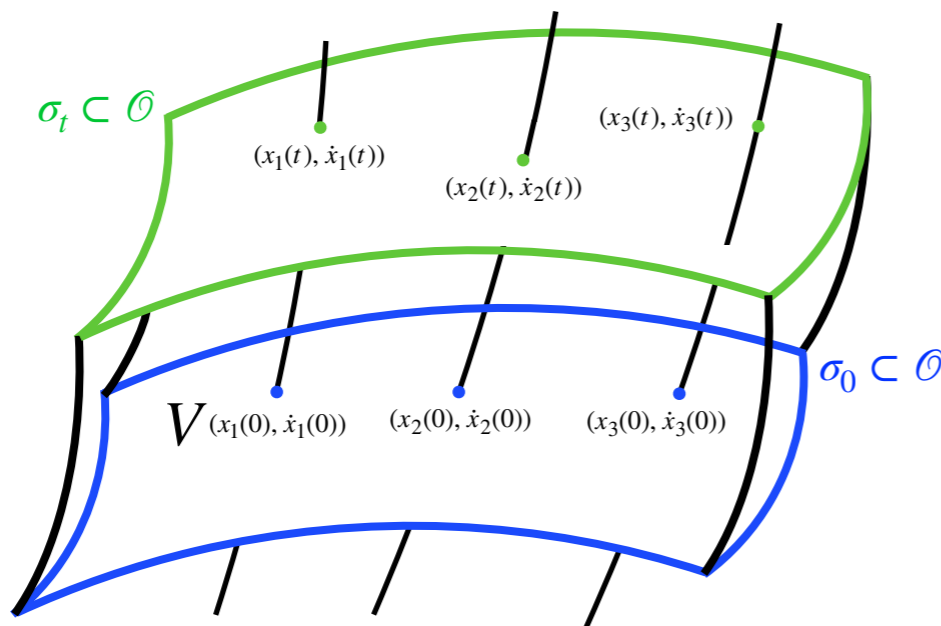
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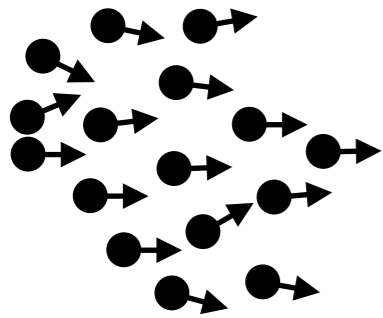


Finsler geometry

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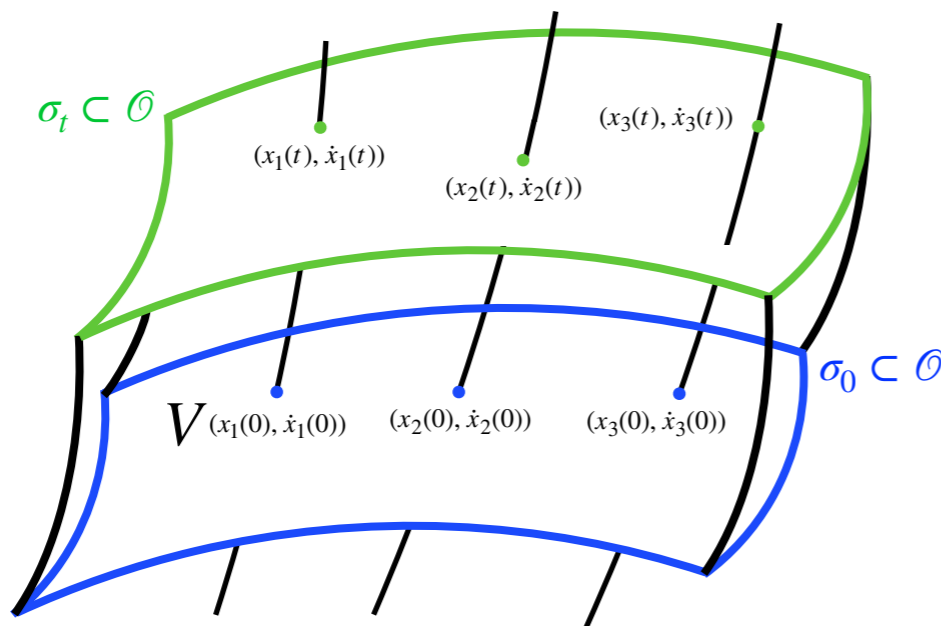
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 &= m \int_V \phi \Sigma
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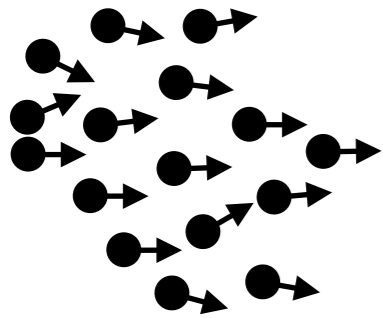


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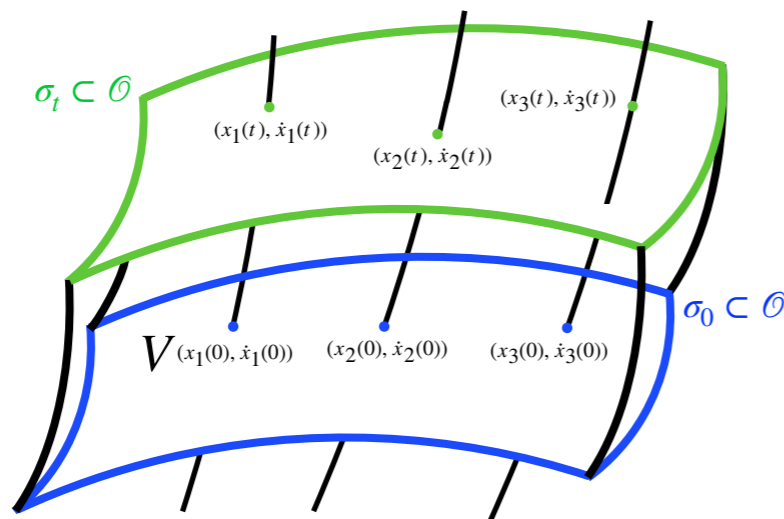
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$$S_{gas} = m \sum_{i=1}^P \int_0^t ds_i \sqrt{L(x_I, \dot{x}_I)}$$

$$= mPt = m \int_0^t P ds$$

$$= m \int_V \phi \Sigma$$

$$\Sigma \sim \det g^L d^4x \wedge d^3u$$



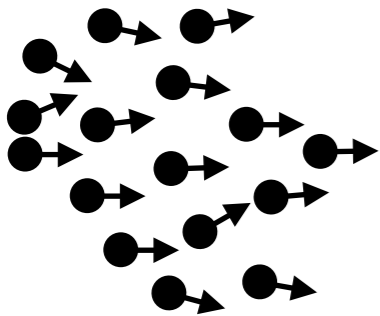
Finsler geometry

Geometry defined by $L(x, \dot{x})$

- $L(x, \lambda \dot{x}) = \lambda^2 L(x, \dot{x})$;
 \Rightarrow parametrization invariant action
- $g_{ab}^L = \frac{1}{2} \partial_a \partial_b L$ non-degenerate;
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The relativistic gas

Kinetic gas action



$$S_{gas} = m \int_V \phi \Sigma$$

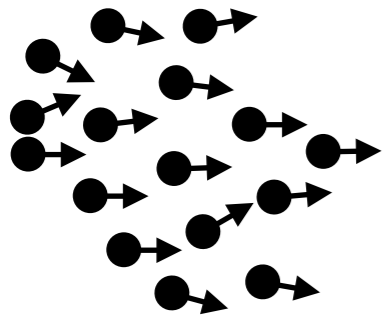
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Covariance under coordinate changes

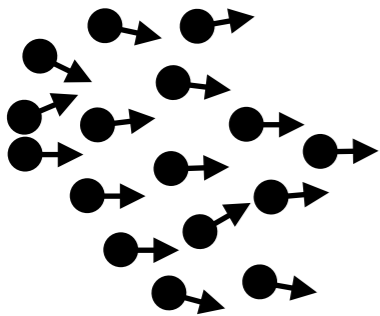
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$$\delta_\xi S = m \int_V \xi^b \nabla_{\delta_a} \left(\frac{\phi \dot{x}^a \dot{x}_b}{L} \right) \Sigma + m \int_{\partial V} J,$$

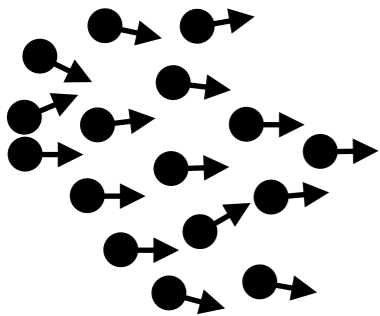
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Energy momentum distribution tensor

[Gotay, Marsden 1992]

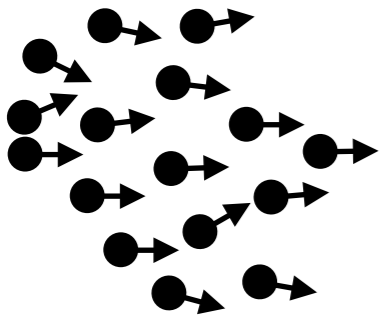
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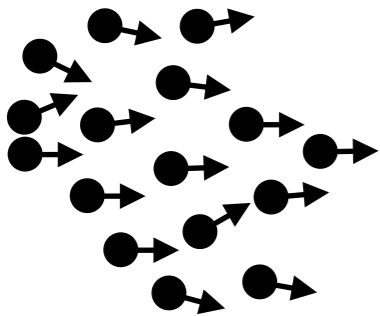
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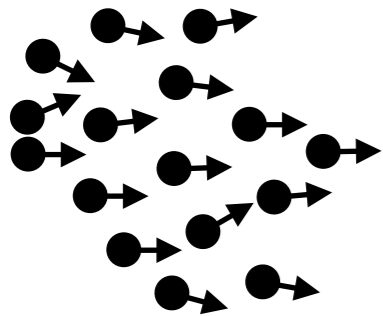
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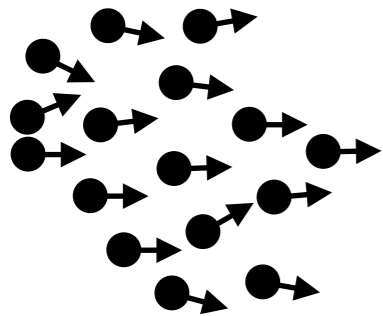
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Ge

$$L = g_{ab}(x) \dot{x}^a \dot{x}^b$$

• L
=

• g

Lorentzian Geometry

$\text{sign}(g^L) = (+, -, -, -)$ on $\mathcal{Y} \subset TM$

• gas particle trajectories satisfy

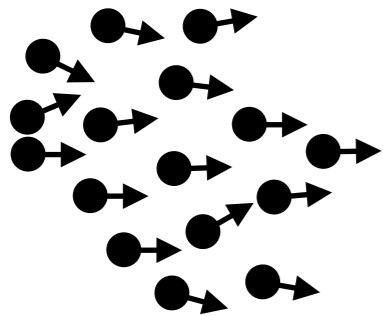
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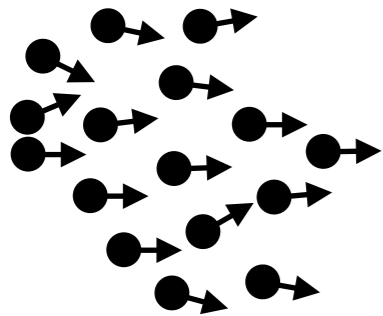
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The relativistic gas



Kinetic gas action

$$S_{gas} = m \int_V \phi \Sigma$$

Variational principle for the Einstein-Vlasov equations [Andersson, Korzynski 2019]

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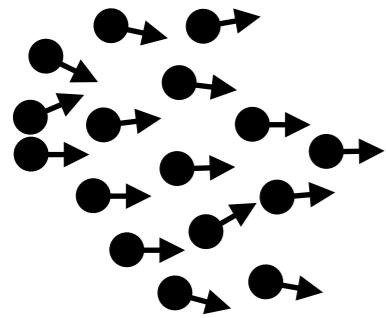
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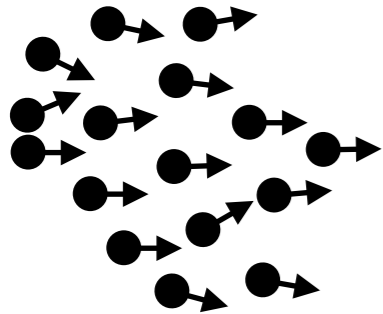
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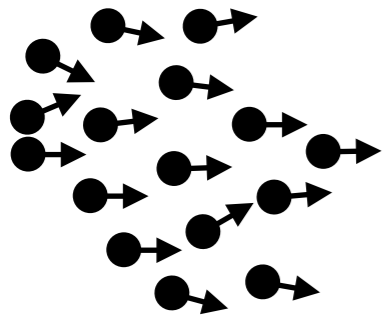
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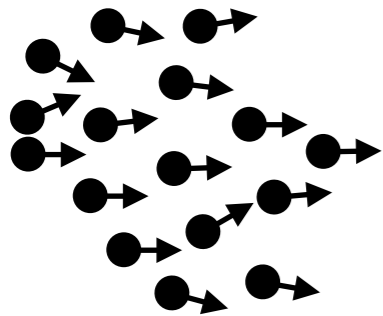
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$$S[L] = \tilde{\kappa} \int_V R_0 \Sigma + \int_V \phi \Sigma$$

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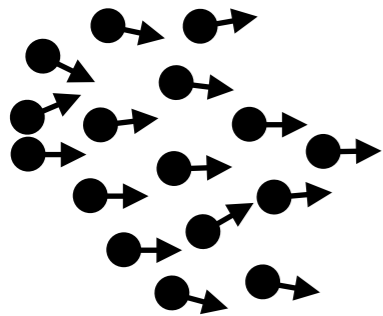
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The 1PDF of a relativistic kinetic gas as source of gravity

$$\frac{1}{2} g^{Lab} \dot{\partial}_a \dot{\partial}_b R - 3 \frac{R}{L} - g^{Lab} (\nabla_{\delta_a} P_b - P_a P_b + \dot{\partial}_a (\nabla P_b)) = -\frac{1}{2\tilde{\kappa}} \phi$$

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The averaged equation and the Einstein Vlasov system

$$\underbrace{\int_{V_x} \frac{\dot{x}^a \dot{x}_b}{L} G(x, \dot{x}) \Sigma_x}_{\mathfrak{G}^a_b(x)} = -\frac{1}{2\tilde{\kappa}} \int_{V_x} \frac{\dot{x}^a \dot{x}_b}{L} \phi \Sigma_x = T^a_b(x) \qquad R^a_b - \frac{1}{2} \delta_b^a R = \frac{8\pi G}{c^4} \int_{V_x} \frac{\dot{x}^a \dot{x}_b}{g_{cd}(x) \dot{x}^c \dot{x}^d} \phi \Sigma_x$$

- Under which conditions is $\mathfrak{G}^a_b(x) \sim R^a_b - \frac{1}{2} \delta_b^a R$?

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- Under which conditions is $\mathfrak{G}^a_b(x) \sim R^a_b - \frac{1}{2} \delta_b^a R$?

Homogeneous and Isotropic Ansatz

$$\Phi = \rho(x) e^{1 - \frac{i^2}{i^2 - a(t)^2 w^2}}$$