

Universal electromagnetic fields¹

Marcello Ortaggio

Institute of Mathematics, Academy of Sciences of the Czech Republic

Kazimierz Dolny – September 27th, 2017

¹Joint work with V. Pravda, arXiv:1708.08017

Contents

- 1 Modified electrodynamics
- 2 Schrödinger's "universal" solutions
- 3 Universal solutions
- 4 Gravity analog: universal spacetimes

Modified electrodynamics

Maxwell's theory (in vacuum)

- $L = I_1$ ($I_1 \equiv F_{ab}F^{ab}$, $I_2 \equiv F_{ab}*F^{ab}$)
- $d\mathbf{F} = 0$, $*d*\mathbf{F} = 0$

Modified electrodynamics

Maxwell's theory (in vacuum)

- $L = I_1$ ($I_1 \equiv F_{ab}F^{ab}$, $I_2 \equiv F_{ab}{}^*F^{ab}$)
- $d\mathbf{F} = 0$, $*d*\mathbf{F} = 0$

"Corrections"

- 1 Classical theories (to cure infinite self-energies)
 - Mie 1912
 - Born-Infeld 1933, 1934:
 $L = \sqrt{b^4 + 2b^2 I_1 + I_2^2} - b^2$ ($\rightarrow I_1$ for $b \rightarrow \infty$)
 - general NLE: $L = L(I_1, I_2)$ (e.g. [Plebański'70])
 - Bopp 1940, Podolsky 1942: $L = I_1 + a^2(\nabla\mathbf{F})^2$

Modified electrodynamics

Maxwell's theory (in vacuum)

- $L = I_1$ ($I_1 \equiv F_{ab}F^{ab}$, $I_2 \equiv F_{ab}{}^*F^{ab}$)
- $d\mathbf{F} = 0$, $*d*\mathbf{F} = 0$

"Corrections"

- 1 Classical theories (to cure infinite self-energies)
 - Mie 1912
 - Born-Infeld 1933, 1934:
$$L = \sqrt{b^4 + 2b^2 I_1 + I_2^2} - b^2 \quad (\rightarrow I_1 \text{ for } b \rightarrow \infty)$$
 - general NLE: $L = L(I_1, I_2)$ (e.g. [Plebański'70])
 - Bopp 1940, Podolsky 1942: $L = I_1 + a^2(\nabla\mathbf{F})^2$
- 2 Vacuum polarization in QED: Heisenberg-Euler (NLE)
[Heisenberg-Euler'36, Weisskopf'36, Schwinger'51]
- 3 String theory: Born-Infeld as a low-energy Lagrangian
[Fradkin-Tseytlin'85, Bergshoeff-Sezgin-Pope-Townsend'87, ...]

Schrödinger's "universal" solutions

Field equations of NLE:

$$d\mathbf{F} = 0, \quad *d * \tilde{\mathbf{F}} = 0,$$

$$\text{where } \tilde{\mathbf{F}} \equiv \alpha\mathbf{F} + \beta * \mathbf{F}, \quad \alpha = \alpha(I_1, I_2), \quad \beta = \beta(I_1, I_2).$$

Difficult to find exact solutions. But:

Schrödinger's "universal" solutions

Field equations of NLE:

$$d\mathbf{F} = 0, \quad *d*\tilde{\mathbf{F}} = 0,$$

$$\text{where } \tilde{\mathbf{F}} \equiv \alpha\mathbf{F} + \beta*\mathbf{F}, \quad \alpha = \alpha(I_1, I_2), \quad \beta = \beta(I_1, I_2).$$

Difficult to find exact solutions. But:

Null fields

- $I_1 = 0 = I_2 \quad (\Leftrightarrow E^2 - B^2 = 0 = \vec{E} \cdot \vec{B})$

Schrödinger's "universal" solutions

Field equations of NLE:

$$d\mathbf{F} = 0, \quad *d*\tilde{\mathbf{F}} = 0,$$

$$\text{where } \tilde{\mathbf{F}} \equiv \alpha\mathbf{F} + \beta*\mathbf{F}, \quad \alpha = \alpha(I_1, I_2), \quad \beta = \beta(I_1, I_2).$$

Difficult to find exact solutions. But:

Null fields

- $I_1 = 0 = I_2$ ($\Leftrightarrow E^2 - B^2 = 0 = \vec{E} \cdot \vec{B}$)
- $\Rightarrow \alpha$ and β are constant
- $d\mathbf{F} = 0, *d*\mathbf{F} = 0 \Rightarrow *d*\tilde{\mathbf{F}} = 0$

Schrödinger's "universal" solutions

Field equations of NLE:

$$d\mathbf{F} = 0, \quad *d*\tilde{\mathbf{F}} = 0,$$

$$\text{where } \tilde{\mathbf{F}} \equiv \alpha\mathbf{F} + \beta*\mathbf{F}, \quad \alpha = \alpha(I_1, I_2), \quad \beta = \beta(I_1, I_2).$$

Difficult to find exact solutions. But:

Null fields

- $I_1 = 0 = I_2$ ($\Leftrightarrow E^2 - B^2 = 0 = \vec{E} \cdot \vec{B}$)
- $\Rightarrow \alpha$ and β are constant
- $d\mathbf{F} = 0, *d*\mathbf{F} = 0 \Rightarrow *d*\tilde{\mathbf{F}} = 0$

all null Maxwell fields satisfy any NLE! [Schrödinger'35,'43]

E.g. Born-Infeld:

$$F_{[ab,c]} = 0, \quad \left(\frac{F^{ab} - \mathcal{G}^* F^{ab}}{\sqrt{1 + \mathcal{F} - \mathcal{G}^2}} \right)_{;b} = 0 \quad (\mathcal{F} \equiv \frac{1}{2}I_1, \mathcal{G} \equiv \frac{1}{4}I_2).$$

E.g. Born-Infeld:

$$F_{[ab,c]} = 0, \quad \left(\frac{F^{ab} - \mathcal{G}^* F^{ab}}{\sqrt{1 + \mathcal{F} - \mathcal{G}^2}} \right)_{;b} = 0 \quad (\mathcal{F} \equiv \frac{1}{2}I_1, \quad \mathcal{G} \equiv \frac{1}{4}I_2).$$

Null fields:

- $\mathbf{F} = \ell \wedge \omega$, with $\ell_a \ell^a = 0 = \ell^a \omega_a$
- plane waves [Schwinger'51, Synge'55]
- asymptotic behavior of radiative systems [Penrose-Rindler'86]
- approximate field of high-energy sources [Bergmann'42, Synge'55, Robinson-Rózga'84]

E.g. Born-Infeld:

$$F_{[ab,c]} = 0, \quad \left(\frac{F^{ab} - \mathcal{G}^* F^{ab}}{\sqrt{1 + \mathcal{F} - \mathcal{G}^2}} \right)_{;b} = 0 \quad (\mathcal{F} \equiv \frac{1}{2} I_1, \quad \mathcal{G} \equiv \frac{1}{4} I_2).$$

Null fields:

- $\mathbf{F} = \ell \wedge \omega$, with $\ell_a \ell^a = 0 = \ell^a \omega_a$
- plane waves [Schwinger'51, Synge'55]
- asymptotic behavior of radiative systems [Penrose-Rindler'86]
- approximate field of high-energy sources [Bergmann'42, Synge'55, Robinson-Rózga'84]

Solve also theories $L = L(\mathbf{F}, \mathbf{F}^2, \dots, \nabla^{(k)} \mathbf{F}, \dots)$ with higher-derivatives? (QED, string theory, ...)

E.g. Born-Infeld:

$$F_{[ab,c]} = 0, \quad \left(\frac{F^{ab} - \mathcal{G}^* F^{ab}}{\sqrt{1 + \mathcal{F} - \mathcal{G}^2}} \right)_{;b} = 0 \quad (\mathcal{F} \equiv \frac{1}{2} I_1, \quad \mathcal{G} \equiv \frac{1}{4} I_2).$$

Null fields:

- $\mathbf{F} = \ell \wedge \omega$, with $\ell_a \ell^a = 0 = \ell^a \omega_a$
- plane waves [Schwinger'51, Synge'55]
- asymptotic behavior of radiative systems [Penrose-Rindler'86]
- approximate field of high-energy sources [Bergmann'42, Synge'55, Robinson-Róžga'84]

Solve also theories $L = L(\mathbf{F}, \mathbf{F}^2, \dots, \nabla^{(k)} \mathbf{F}, \dots)$ with higher-derivatives? (QED, string theory, ...)

The previous argument breaks down when $\tilde{\mathbf{F}}$ contains $\nabla^{(k)} \mathbf{F} \dots$

Universal solutions

Definition: F universal

$$dF = 0, \quad *d*\tilde{F} = 0,$$

\tilde{F} = any 2-form “polynomial” in $F, F^2, \dots, \nabla^{(k)}F, (\nabla^{(k)}F)^2, \dots$

Universal solutions

Definition: F universal

$$dF = 0, \quad *d*\tilde{F} = 0,$$

\tilde{F} = any 2-form “polynomial” in $F, F^2, \dots, \nabla^{(k)}F, (\nabla^{(k)}F)^2, \dots$

e.g. $\tilde{F}_{ab} = c_1 F_{ab} + c_2 F_{c[a} F_{b]}{}^c + c_3 F_{ab;c}{}^c + c_4 F_{ab;cd} F^{cd} + \dots$

Universal solutions

Definition: F universal

$$dF = 0, \quad *d*\tilde{F} = 0,$$

\tilde{F} = any 2-form “polynomial” in $F, F^2, \dots, \nabla^{(k)}F, (\nabla^{(k)}F)^2, \dots$

e.g. $\tilde{F}_{ab} = c_1 F_{ab} + c_2 F_{c[a} F_{b]}^c + c_3 F_{ab;c}{}^c + c_4 F_{ab;cd} F^{cd} + \dots$

Theorem

F is universal provided

- ① F is null: $F = \ell \wedge \omega$
- ②
 - (a) ℓ is a *Kundt* vector field ($\kappa = \sigma = \omega = \theta = 0$)
 - (b) the spacetime is of Petrov type *III* (w.r.t. ℓ)
 - (c) $S_{ab}\ell^b = 0 = S_{[ab}\ell_{c]}$, where $S_{ab} \equiv R_{ab} - \frac{R}{4}g_{ab}$ (traceless-Ricci).

Universal solutions

Definition: F universal

$$dF = 0, \quad *d*\tilde{F} = 0,$$

\tilde{F} = any 2-form “polynomial” in $F, F^2, \dots, \nabla^{(k)}F, (\nabla^{(k)}F)^2, \dots$

e.g. $\tilde{F}_{ab} = c_1 F_{ab} + c_2 F_{c[a} F_{b]}^c + c_3 F_{ab;c}{}^c + c_4 F_{ab;cd} F^{cd} + \dots$

Theorem

F is universal provided

- ① F is null: $F = \ell \wedge \omega$
- ②
 - (a) ℓ is a *Kundt* vector field ($\kappa = \sigma = \omega = \theta = 0$)
 - (b) the spacetime is of Petrov type *III* (w.r.t. ℓ)
 - (c) $S_{ab}\ell^b = 0 = S_{[ab}\ell_{c]}$, where $S_{ab} \equiv R_{ab} - \frac{R}{4}g_{ab}$ (traceless-Ricci).

2(b,c) simply mean: $R_{abcd} = \frac{R}{6}g_{a[c}g_{d]b} + (\text{terms with “many” } \ell)$

Explicitly:

$$\begin{aligned}
 ds^2 &= 2P^{-2}d\zeta d\bar{\zeta} - 2du (dr + Wd\zeta + \bar{W}d\bar{\zeta} + Hdu), \\
 \mathbf{F} &= du \wedge [f(u, \zeta)d\zeta + \bar{f}(u, \bar{\zeta})d\bar{\zeta}] \quad (\ell_a dx^a = du),
 \end{aligned}$$

where (some constraints omitted ...)

$$\begin{aligned}
 P &= P(u, \zeta, \bar{\zeta}), & W &= rg^{(1)}(u, \zeta, \bar{\zeta}) + g^{(0)}(u, \zeta, \bar{\zeta}), \\
 H &= r^2 H^{(2)}(u, \zeta, \bar{\zeta}) + rH^{(1)}(u, \zeta, \bar{\zeta}) + H^{(0)}(u, \zeta, \bar{\zeta}),
 \end{aligned}$$

Explicitly:

$$ds^2 = 2P^{-2}d\zeta d\bar{\zeta} - 2du (dr + Wd\zeta + \bar{W}d\bar{\zeta} + Hdu),$$

$$\mathbf{F} = du \wedge [f(u, \zeta)d\zeta + \bar{f}(u, \bar{\zeta})d\bar{\zeta}] \quad (\ell_a dx^a = du),$$

where (some constraints omitted ...)

$$P = P(u, \zeta, \bar{\zeta}), \quad W = rg^{(1)}(u, \zeta, \bar{\zeta}) + g^{(0)}(u, \zeta, \bar{\zeta}),$$

$$H = r^2 H^{(2)}(u, \zeta, \bar{\zeta}) + rH^{(1)}(u, \zeta, \bar{\zeta}) + H^{(0)}(u, \zeta, \bar{\zeta}),$$

- all the invariants of \mathbf{F} vanish (including derivatives)

[M.O.-Pravda'16]

- also p -forms in n -dimensions

Explicitly:

$$ds^2 = 2P^{-2}d\zeta d\bar{\zeta} - 2du (dr + Wd\zeta + \bar{W}d\bar{\zeta} + Hdu),$$

$$\mathbf{F} = du \wedge [f(u, \zeta)d\zeta + \bar{f}(u, \bar{\zeta})d\bar{\zeta}] \quad (\ell_a dx^a = du),$$

where (some constraints omitted ...)

$$P = P(u, \zeta, \bar{\zeta}), \quad W = rg^{(1)}(u, \zeta, \bar{\zeta}) + g^{(0)}(u, \zeta, \bar{\zeta}),$$

$$H = r^2 H^{(2)}(u, \zeta, \bar{\zeta}) + rH^{(1)}(u, \zeta, \bar{\zeta}) + H^{(0)}(u, \zeta, \bar{\zeta}),$$

- all the invariants of \mathbf{F} vanish (including derivatives)
- [M.O.-Pravda'16]
- also p -forms in n -dimensions
- holds in Minkowski, (A)dS (plus possible waves)
- test fields (no backreaction)
- coupling to gravity also possible: further restrictions

cf. also [Güven'87, Horowitz-Steif'90, Coley'02]

NLE in [Kichenassamy'59, Kremer-Kichenassamy'60, Peres'60]

A simple example:

$$ds^2 = 2du \left[dr + \frac{1}{2}x(r - e^x) du \right] + e^x(dx^2 + e^{2u}dy^2)$$

$$\mathbf{F} = e^{x/2}c(u)du \wedge \left(-\cos \frac{ye^u}{2}dx + e^u \sin \frac{ye^u}{2}dy \right)$$

- $\ell = \partial_r$ is Kundt and *recurrent* ($\ell_{a;b} \sim \ell_a \ell_b$)
- Petrov type III, Ricci-flat [Petrov'62]
- \mathbf{F} is universal

Sketchy proof

Recall: $\tilde{\mathbf{F}}$ = “polynomial” in $\mathbf{F}, \mathbf{F}^2, \dots, \nabla^{(k)} \mathbf{F}, (\nabla^{(k)} \mathbf{F})^2, \dots$

Sketchy proof

Recall: $\tilde{\mathbf{F}}$ = “polynomial” in $\mathbf{F}, \mathbf{F}^2, \dots, \nabla^{(k)}\mathbf{F}, (\nabla^{(k)}\mathbf{F})^2, \dots$

- assumptions $\Rightarrow \mathbf{F}$ is VSI (all invariants are zero) [M.O.-Pravda'16]

Sketchy proof

Recall: \tilde{F} = “polynomial” in $F, F^2, \dots, \nabla^{(k)} F, (\nabla^{(k)} F)^2, \dots$

- assumptions $\Rightarrow F$ is VSI (all invariants are zero) [M.O.-Pravda'16]
- \Rightarrow all $\nabla^{(k)} F$ are VSI \Rightarrow contain “many” ℓ (type III) [Hervik'11]

Sketchy proof

Recall: $\tilde{\mathbf{F}}$ = “polynomial” in $\mathbf{F}, \mathbf{F}^2, \dots, \nabla^{(k)}\mathbf{F}, (\nabla^{(k)}\mathbf{F})^2, \dots$

- assumptions $\Rightarrow \mathbf{F}$ is VSI (all invariants are zero) [M.O.-Pravda'16]
- \Rightarrow all $\nabla^{(k)}\mathbf{F}$ are VSI \Rightarrow contain “many” ℓ (type III) [Hervik'11]
- any quadratic 2-tensor $\sim \ell_a \ell_b$, but antisymmetry $\Rightarrow 0!$

Sketchy proof

Recall: \tilde{F} = “polynomial” in $F, F^2, \dots, \nabla^{(k)} F, (\nabla^{(k)} F)^2, \dots$

- assumptions $\Rightarrow F$ is VSI (all invariants are zero) [M.O.-Pravda'16]
- \Rightarrow all $\nabla^{(k)} F$ are VSI \Rightarrow contain “many” ℓ (type III) [Hervik'11]
- any quadratic 2-tensor $\sim \ell_a \ell_b$, but antisymmetry $\Rightarrow 0!$
- $\Rightarrow \tilde{F}$ linear in F (and $*F$) and $\nabla^{(k)} F$

$$\text{e.g. } \tilde{F}_{ab} = c_1 F_{ab} + c_2 F_{ab;c}{}^c + c_3 \underbrace{F_{[a|c;{}^c|b]}}_{=0!} + c_4 F_{ab;cd}{}^{cd} + \dots$$

Sketchy proof

Recall: \tilde{F} = “polynomial” in $F, F^2, \dots, \nabla^{(k)} F, (\nabla^{(k)} F)^2, \dots$

- assumptions $\Rightarrow F$ is VSI (all invariants are zero) [M.O.-Pravda'16]
- \Rightarrow all $\nabla^{(k)} F$ are VSI \Rightarrow contain “many” ℓ (type III) [Hervik'11]
- any quadratic 2-tensor $\sim \ell_a \ell_b$, but antisymmetry $\Rightarrow 0!$

- $\Rightarrow \tilde{F}$ linear in F (and $*F$) and $\nabla^{(k)} F$

e.g. $\tilde{F}_{ab} = c_1 F_{ab} + c_2 F_{ab;c}{}^c + c_3 \underbrace{F_{[a|c;{}^c|b]}}_{=0!} + c_4 F_{ab;cd}{}^{cd} + \dots$

- use the Ricci id. “[∇, ∇] $F = F \cdot \text{Riem}$ ” to isolate terms

$F_{[a|c;{}^{c\dots}|b] \dots} = 0$ and $F_{ab;c\dots}{}^{c\dots}$

Sketchy proof

Recall: \tilde{F} = “polynomial” in $F, F^2, \dots, \nabla^{(k)} F, (\nabla^{(k)} F)^2, \dots$

- assumptions $\Rightarrow F$ is VSI (all invariants are zero) [M.O.-Pravda'16]
- \Rightarrow all $\nabla^{(k)} F$ are VSI \Rightarrow contain “many” ℓ (type III) [Hervik'11]
- any quadratic 2-tensor $\sim \ell_a \ell_b$, but antisymmetry $\Rightarrow 0!$
- $\Rightarrow \tilde{F}$ linear in F (and $*F$) and $\nabla^{(k)} F$
 e.g. $\tilde{F}_{ab} = c_1 F_{ab} + c_2 F_{ab;c}{}^c + c_3 \underbrace{F_{[a|c;{}^c|b]}}_{=0!} + c_4 F_{ab;cd}{}^{cd} + \dots$
- use the Ricci id. “[∇, ∇] $F = F \cdot \text{Riem}$ ” to isolate terms
 $F_{[a|c;{}^{c\dots}|b] \dots} = 0$ and $F_{ab;c}{}^{c\dots}$
- Weitzenböck id. with $\Delta F = 0$ and $R_{abcd} = \frac{R}{6} g_{a[c} g_{d]b} + \dots$
 $\Rightarrow F_{ab;c}{}^c \sim R F_{ab}$

Sketchy proof

Recall: \tilde{F} = “polynomial” in $F, F^2, \dots, \nabla^{(k)}F, (\nabla^{(k)}F)^2, \dots$

- assumptions $\Rightarrow F$ is VSI (all invariants are zero) [M.O.-Pravda'16]
- \Rightarrow all $\nabla^{(k)}F$ are VSI \Rightarrow contain “many” ℓ (type III) [Hervik'11]
- any quadratic 2-tensor $\sim \ell_a \ell_b$, but antisymmetry $\Rightarrow 0!$
- $\Rightarrow \tilde{F}$ linear in F (and $*F$) and $\nabla^{(k)}F$
 e.g. $\tilde{F}_{ab} = c_1 F_{ab} + c_2 F_{ab;c}{}^c + c_3 \underbrace{F_{[a|c;{}^c|b]}}_{=0!} + c_4 F_{ab;cd}{}^{cd} + \dots$
- use the Ricci id. “[∇, ∇]F = F · Riem” to isolate terms
 $F_{[a|c;{}^{c\dots}|b] \dots} = 0$ and $F_{ab;c}{}^{c\dots}$
- Weitzenböck id. with $\Delta F = 0$ and $R_{abcd} = \frac{R}{6} g_{a[c} g_{d]b} + \dots$
 $\Rightarrow F_{ab;c}{}^c \sim R F_{ab}$
- $\Rightarrow \tilde{F}_{ac;{}^c} = 0$, Q.E.D.

Gravity analog: universal spacetimes

- Certain type N vacuum pp -waves solve *any* $L(\text{Riem}, \nabla\text{Riem}, \dots)$
e.g. quadratic gravity, Lovelock gravity, string corrections, etc.
[Deser'75, Güven'87, Amati-Klimčik'89, Horowitz-Steif'90]

Gravity analog: universal spacetimes

- Certain type N vacuum pp -waves solve *any* $L(\text{Riem}, \nabla\text{Riem}, \dots)$
e.g. quadratic gravity, Lovelock gravity, string corrections, etc.
[Deser'75, Güven'87, Amati-Klimčík'89, Horowitz-Steif'90]
- formal definition of “universal” metrics and more examples in 4D [Coley-Gibbons-Hervik-Pope'08]

Gravity analog: universal spacetimes

- Certain type N vacuum *pp*-waves solve *any* $L(\text{Riem}, \nabla\text{Riem}, \dots)$
e.g. quadratic gravity, Lovelock gravity, string corrections, etc.
[Deser'75, Güven'87, Amati-Klimčík'89, Horowitz-Steif'90]
- formal definition of “universal” metrics and more examples in 4D [Coley-Gibbons-Hervik-Pope'08]
- results for type III/N and II in 4D and HD
for example: *all Kundt Einstein spacetimes of Weyl type N are universal*
[Hervik-Pravda-Pravdová'14, Hervik-Málek-Pravda-Pravdová'15, Hervik-Pravda-Pravdová'17]

Gravity analog: universal spacetimes

- Certain type N vacuum pp -waves solve *any* $L(\text{Riem}, \nabla \text{Riem}, \dots)$
e.g. quadratic gravity, Lovelock gravity, string corrections, etc.
[Deser'75, Güven'87, Amati-Klimčík'89, Horowitz-Steif'90]
- formal definition of “universal” metrics and more examples in 4D [Coley-Gibbons-Hervik-Pope'08]
- results for type III/N and II in 4D and HD
for example: *all Kundt Einstein spacetimes of Weyl type N are universal*
[Hervik-Pravda-Pravdová'14, Hervik-Málek-Pravda-Pravdová'15, Hervik-Pravda-Pravdová'17]
- slightly different approach [Gürses-Şişman-Tekin'14,'17]