### TT tensors in flat spaces of any dimension

#### J. Tafel

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Condition

$$\nabla_i T^{ij} = 0$$

with symmetric  $T^{ij}$  appears in GR as

- the momentum constraint in the initial value problem
- gauge condition for the metric tensor
- "conservation law" of the energy-momentum.

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$$\alpha^{i_1\dots i_p}_{,i_p} = \mathbf{0} \quad \Rightarrow \quad \alpha^{i_1\dots i_p} = \beta^{i_1\dots i_{p+1}}_{,i_{p+1}} \cdot$$

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External indices can be added.

## T tensors in flat spaces

#### Proposition

All tensors satisfying  $T^{ij}_{\ ,j}=0$  have the form

$${\cal T}^{\it ij}={\cal R}^{\it ikjp}_{,\it kp}$$

where R<sup>ikjp</sup> is any tensor with algebraic symmetries of the Riemann tensor.

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**Proof.** From Poincare lemma

$$\mathcal{T}^{ij}=\mathcal{S}^{ijk}_{,k}\ ,\ \ \mathcal{S}^{ikj}=-\mathcal{S}^{ijk}\ .$$

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Symmetry  $T^{ij} = T^{ji}$  yields

$$S^{[ij]k}_{,k} = 0$$

$$\mathcal{S}^{[ij]k} = \mathcal{V}^{ijkp}_{,p}$$

and

$$V^{(ij)kp} = 0 = V^{ij(kp)} .$$

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Let  $R^{ijkp} = -V^{ikjp} - V^{jpik} + 2V^{[ikjp]}$ 

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D=2

$$T^{ij} = \epsilon^{ik} \epsilon^{jp} R_{,kp}$$

D=3

$$T^{ij} = -\epsilon^{ikl}\epsilon^{jps}G_{ls,kp}$$

where *R* and  $G_{ij}$  are, respectively, "the Ricci scalar" and "the Einstein tensor" corresponding to  $R^{ijkp}$ .

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### Gauge transformations

In dimension D=2 the potential *R* is defined up to translation by  $a_i x^i + b$ .

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$$(\xi^{ikjpr} + \xi^{jpikr} - 2\xi^{[ikjpr]})_{,r}$$

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Can the Weyl part or the Ricci part in the decomposition

be eliminated?

T tensors in D = 3 and all analytic T tensors in  $D \ge 4$  are given by

$$T_{ij} = a igtriangleq R_{ij} - 2aR^k_{(i,j)k} + bR_{,ij} + (aR^{k
ho}_{\ ,k
ho} - bigtriangleq R)g_{ij} \quad (*)$$

where R<sub>ij</sub> is a symmetric tensor undergoing gauge transformation

$$R_{ij} \longrightarrow R_{ij} + \xi_{(i,j)} - \xi_{,k}^{k} g_{ij}$$

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- Equations (\*) are linearized Einstein equations for the first corrections R<sub>ij</sub> to the flat metric g<sub>ij</sub>.

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- Equations (\*) are linearized Einstein equations for the first corrections R<sub>ij</sub> to the flat metric g<sub>ij</sub>.
- Hence, the gauge freedom of *R<sub>ij</sub>* is deduced (up to functions of D-1 variables).

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#### Proposition

In  $D \ge 3$  TT tensors are given by

$${\cal T}^{ij}={\cal R}^{ikjp}_{\ ,kp} \ ,$$

iff the "Ricci tensor" satisfies

$${m R}^{ij} = {m S}^{(ij)k}_{\ ,r} \ , \ \ {m S}^{ijk} = - {m S}^{ikj} \ .$$

Gauge transformations of S<sup>ijk</sup>:

$${\cal S}^{ijk} \longrightarrow {\cal S}^{ijk} - 2 g^{i[j} \xi^{k]} + \chi^{ijkr}_{,r} + \eta^{ijk} \; ,$$

where

$$\chi^{ijkr} = \chi^{i[jkr]} , \quad \eta^{ijk} = \eta^{[ijk]}$$

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$$T^{ij} = \epsilon^{kl(i} (\triangle A^{j)}_{k} - A^{j)p}_{kp,})_{,l}$$

where A<sub>ij</sub> is a symmetric tensor defined up to the transformation

$$A_{ij} \longrightarrow A_{ij} + \chi_{(i,j)} + \eta g_{ij}$$

with arbitrary functions  $\chi_i$  and  $\eta$ .

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## Gauge conditions for TT tensors

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- Is  $R^{ij} = 0$  possible?

#### Proposition

In dimension  $D \ge 4$  every analytic TT tensor can be put into the form

$${\cal T}^{ij}={\cal C}^{ikj
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where C<sup>ikjp</sup> is a tensor with all symmetries of the Weyl tensor.

- The base of the

J. Tafel (University of Warsaw) TT tensors in flat spaces of any dimension

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- Covariant description of TT tensors in conformally flat spaces? Not yet.

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- Description of symmetric TT tensors in flat spaces? Yes, but in D=3 the approach of Conboye and Ó Murchadha is better.
- Covariant description of TT tensors in conformally flat spaces? Not yet.
- Generalization to curved spaces? Only for spaces of constant curvature.