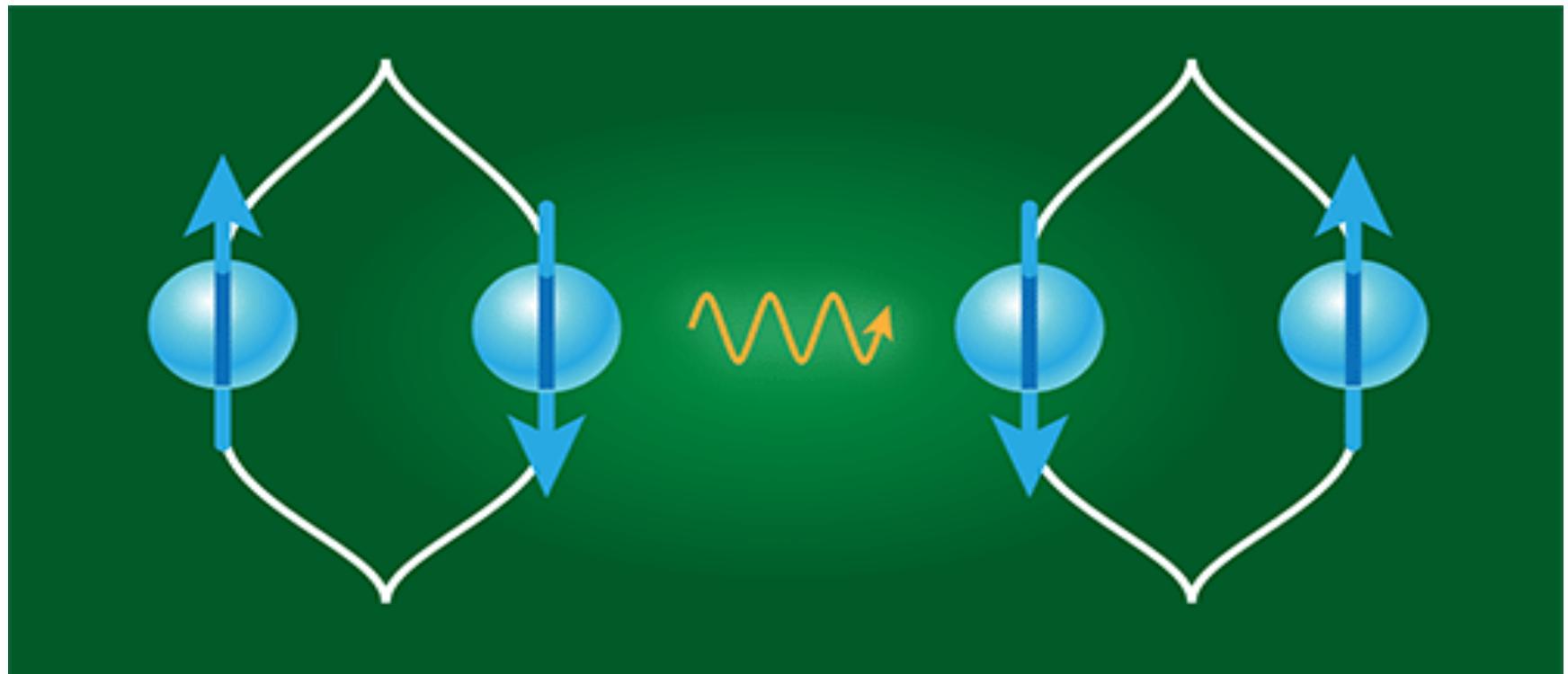


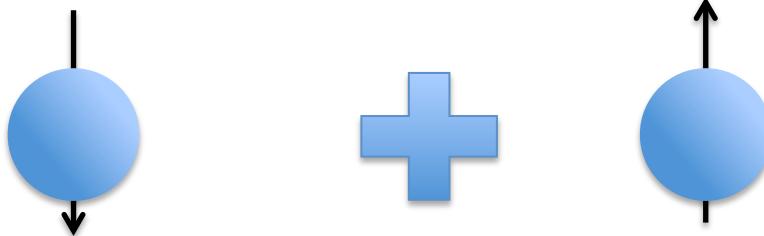
Testing the Quantum Coherent Behaviour of Gravity

Sougato Bose

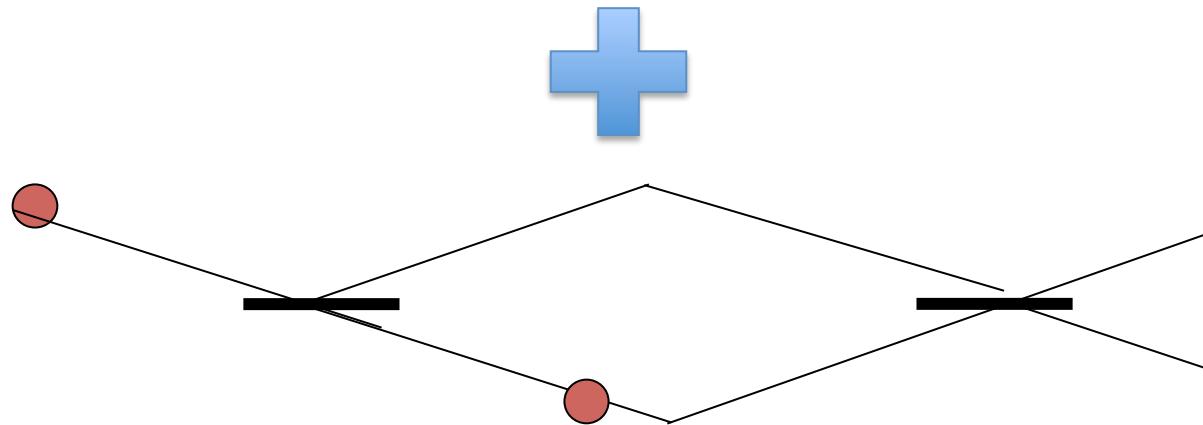
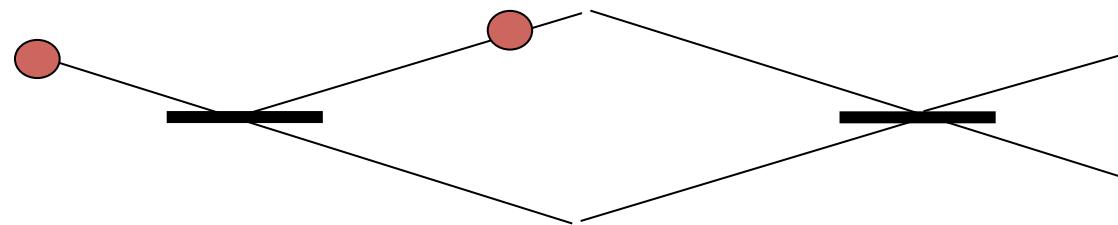
University College London



The Superposition Principle Underpins Quantum Mechanics

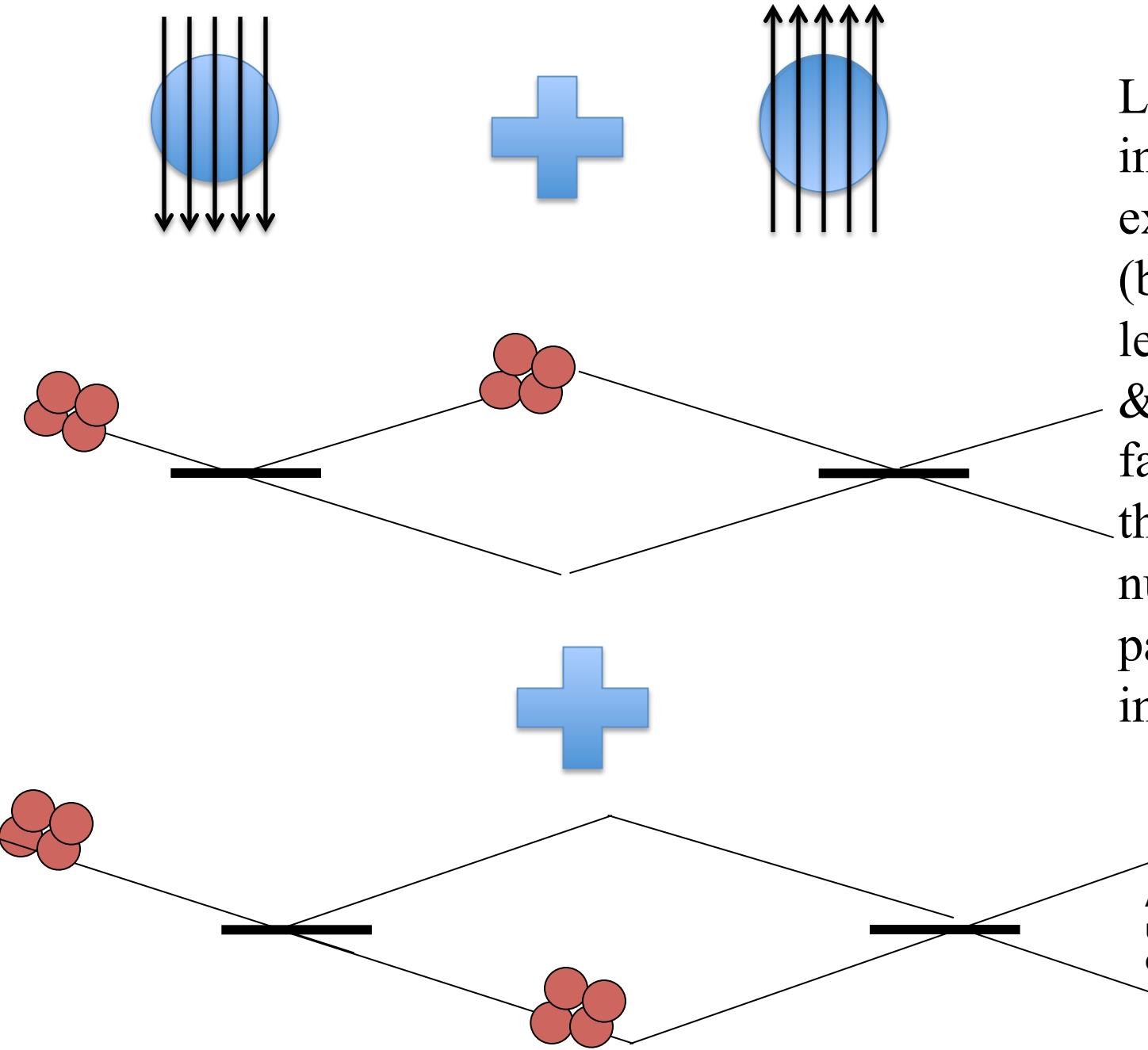


Very familiar
in experiments



Does gravity obey the quantum superposition principle?

To verify superpositions you need an interference experiment (to check for quantum coherence) — not clear in this case what to do.

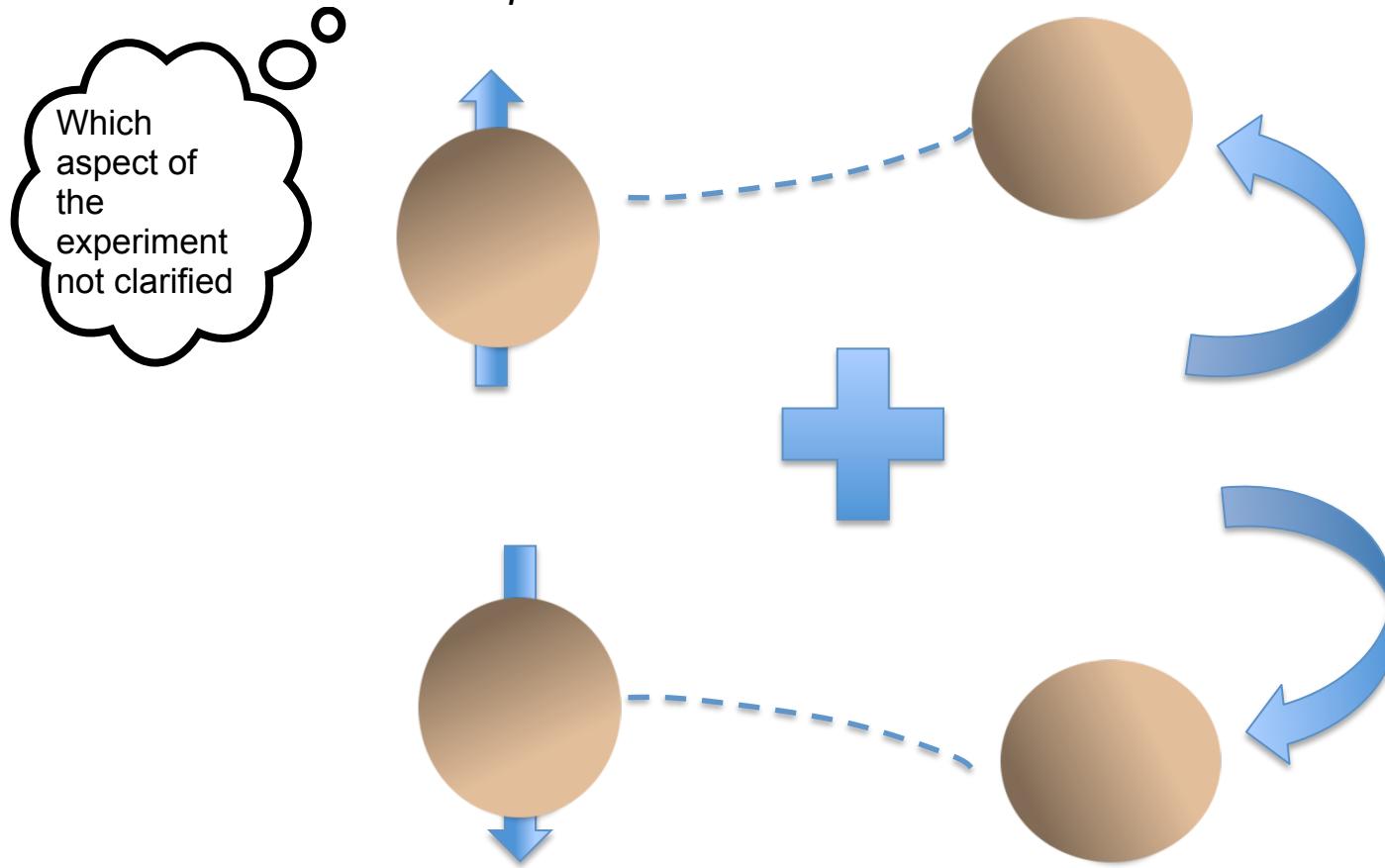


Less familiar
in
experiments
(becomes
less
& less
familiar as
the
number of
particles
increase)

Accomplished
up to 10^4 amu
objects (Arndt et. al.)

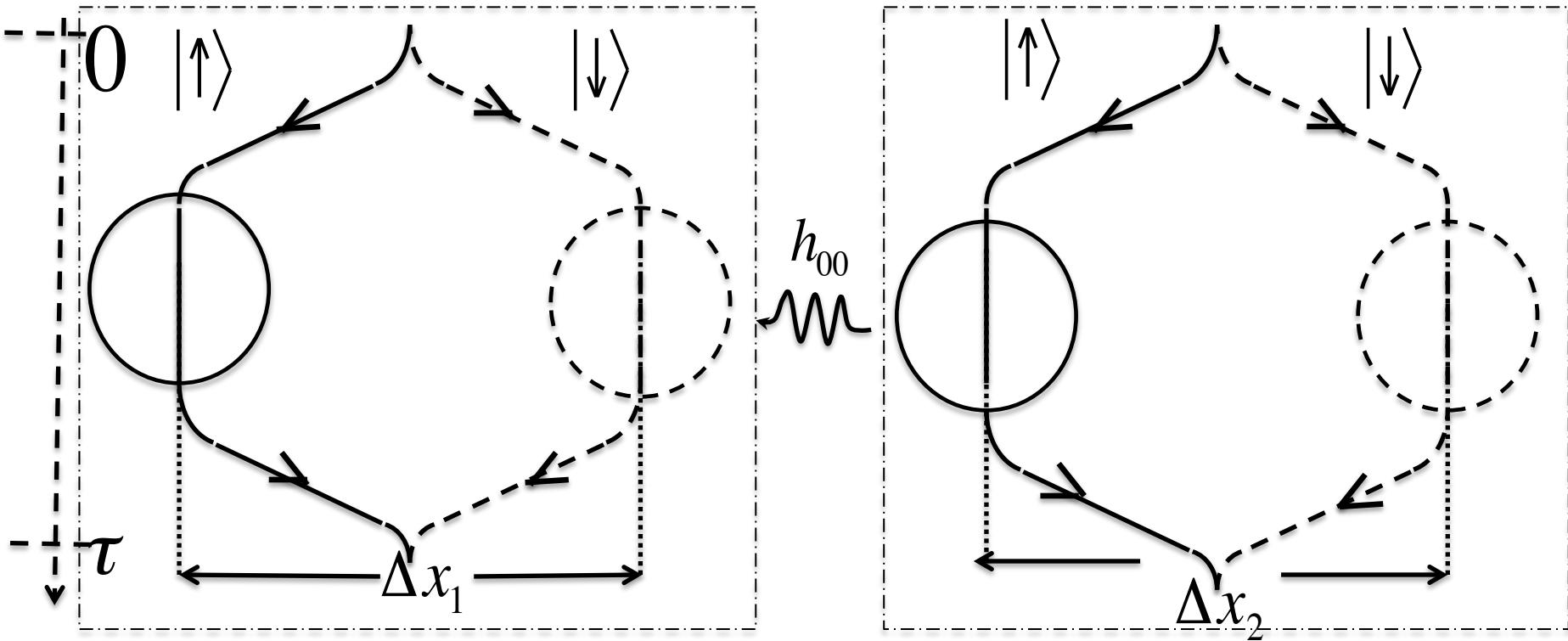
Such superpositions are also called **GHZ** states or NOON states or **Schroedinger** Cat States

Feynman, 1954: Motivation – to argue about the necessity to quantize gravity: “*if you believe in quantum mechanics up to any level then you have to believe in gravitational quantization in order to describe this experiment.*”

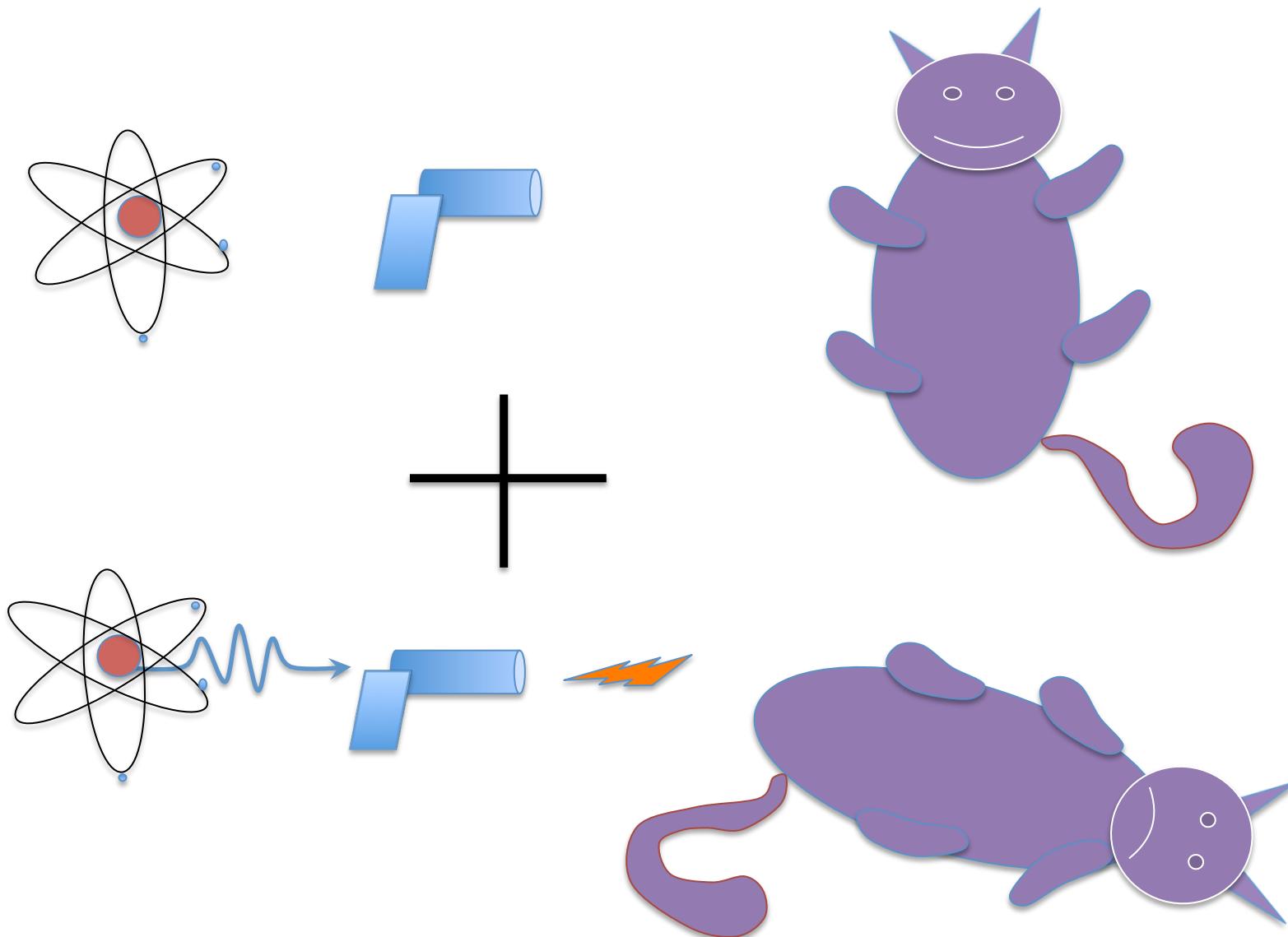


“*The only way to avoid quantization of gravity can in principle no longer play a role beyond a certain point in the chain, and you are not allowed to use quantum mechanics on such a large scale. But I would say that this is the only 'out' if you don't want to quantize gravity.*”

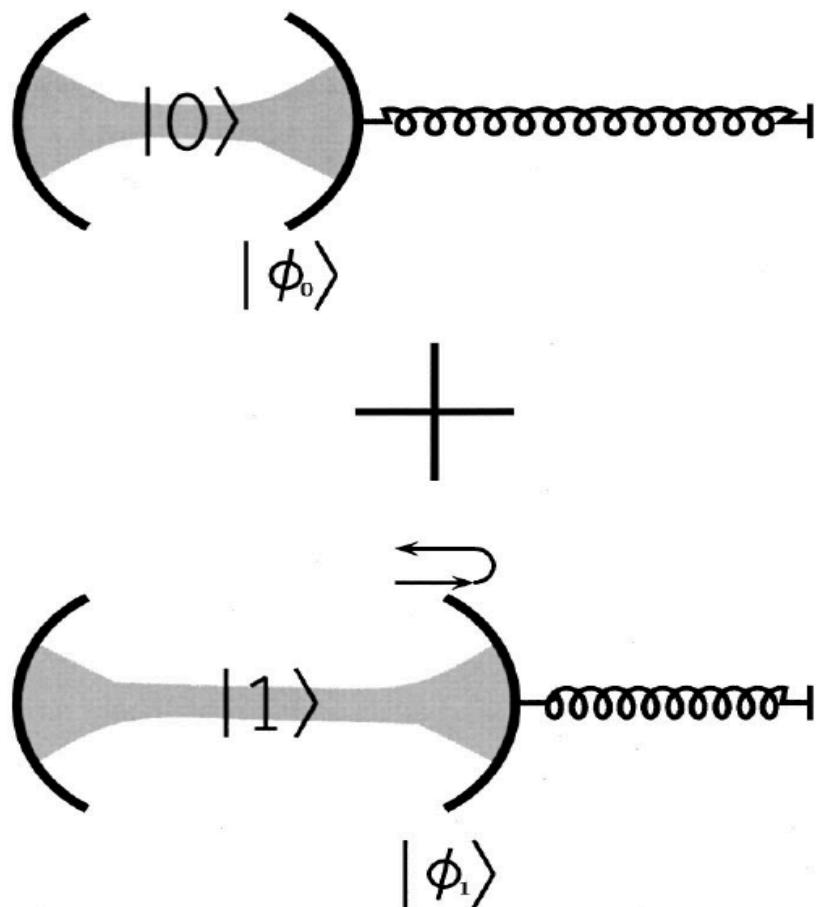
- Is Gravity a Quantum Entity?
-
- If Gravity Mediates Entanglement
It must be quantum entity



How to create the macroscopic superpositions (earliest idea is Schroedinger's Nucleo-Biological mechanism). **Coherent ancilla induced.**



Superpositions of States of a Macroscopic Object using an Ancillary Quantum System:



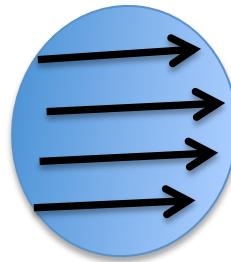
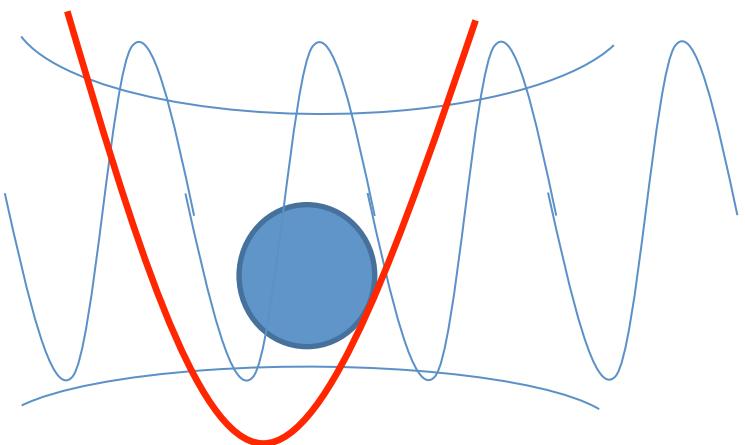
S. Bose, K. Jacobs, P. L. Knight,
Phys. Rev. A 59 (5), 3204
(1999). [arXiv: 1997].
Decoherence/partial coherence is used to certify superposition.

Armour, Blencowe, Schwab,
PRL 2002.
Marshall, Simon, Penrose,
Bouwmeester, PRL 2003.
Decoherence & Recoherence is used to certify superpositions

Bose, PRL 2006.

Ramsey Interferometry with a Levitated Thermal Mesoscopic Object

Diamond bead trapped in an optical trap. The bead contains a spin-1 NV center.

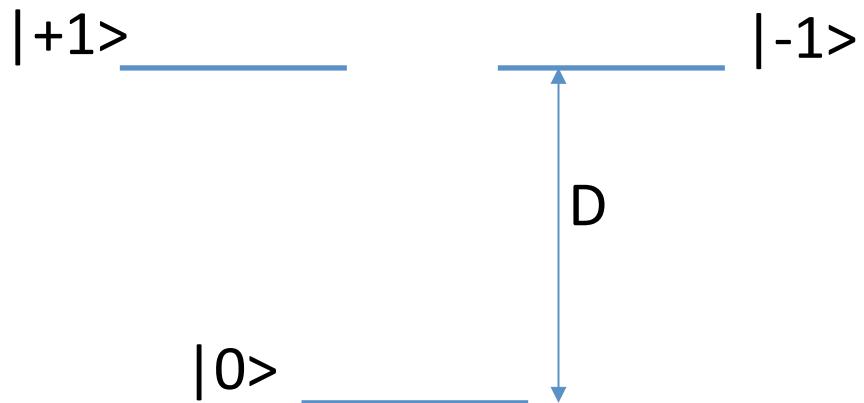


No cavity,
no cooling.

Exploits Spin-Motion
coupling mechanism
proposed by Rabl et.al.
2009.

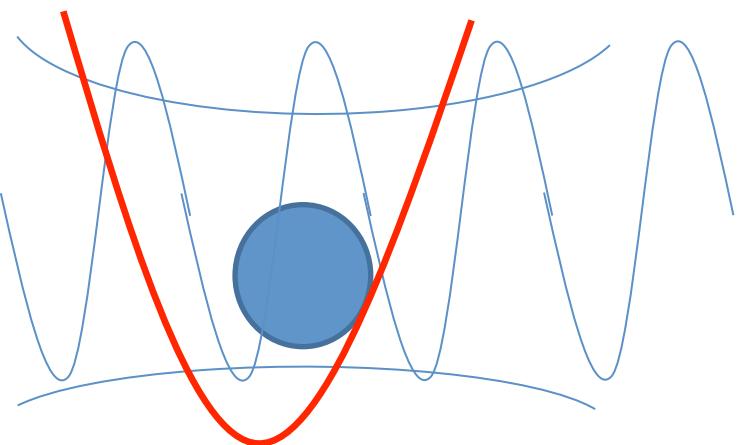
Initial State:

$$|\beta\rangle|0\rangle$$

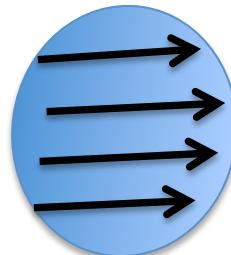


Ramsey Interferometry with a Levitated Thermal Mesoscopic Object

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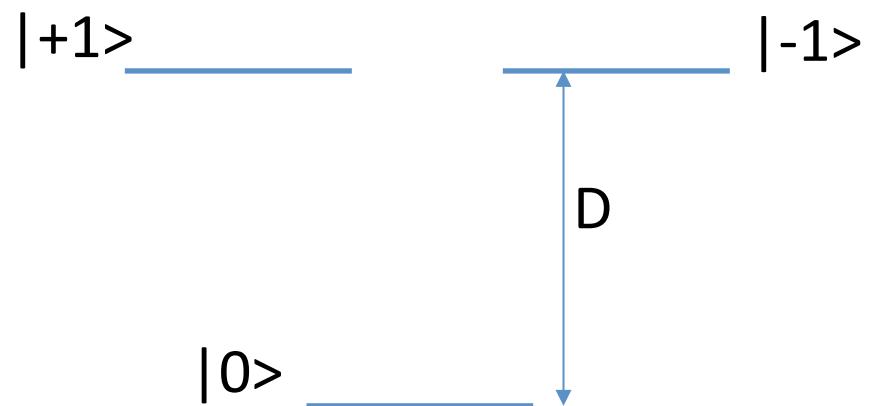


No cavity,
no cooling.



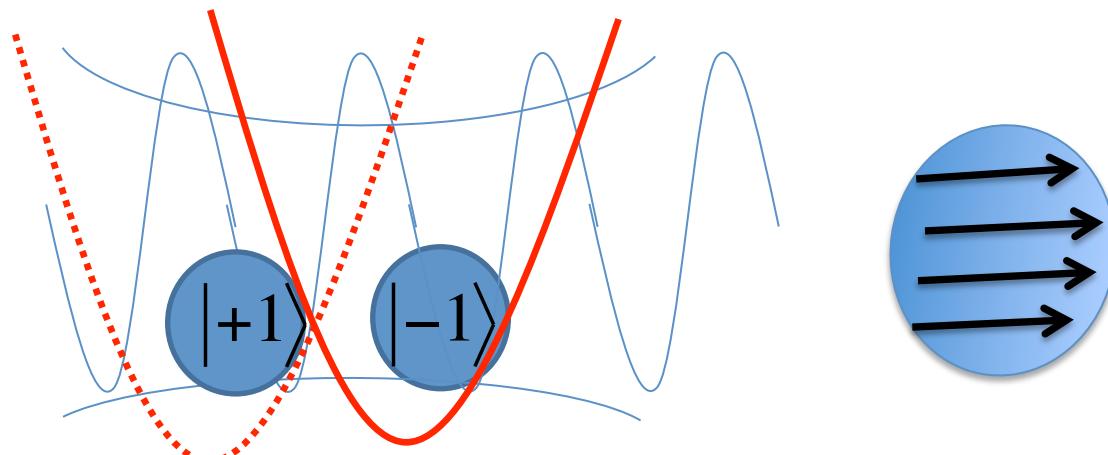
Step 1:

$$|\beta\rangle(|+1\rangle + |-1\rangle)$$



Ramsey Interferometry with a Levitated Thermal Mesoscopic Object

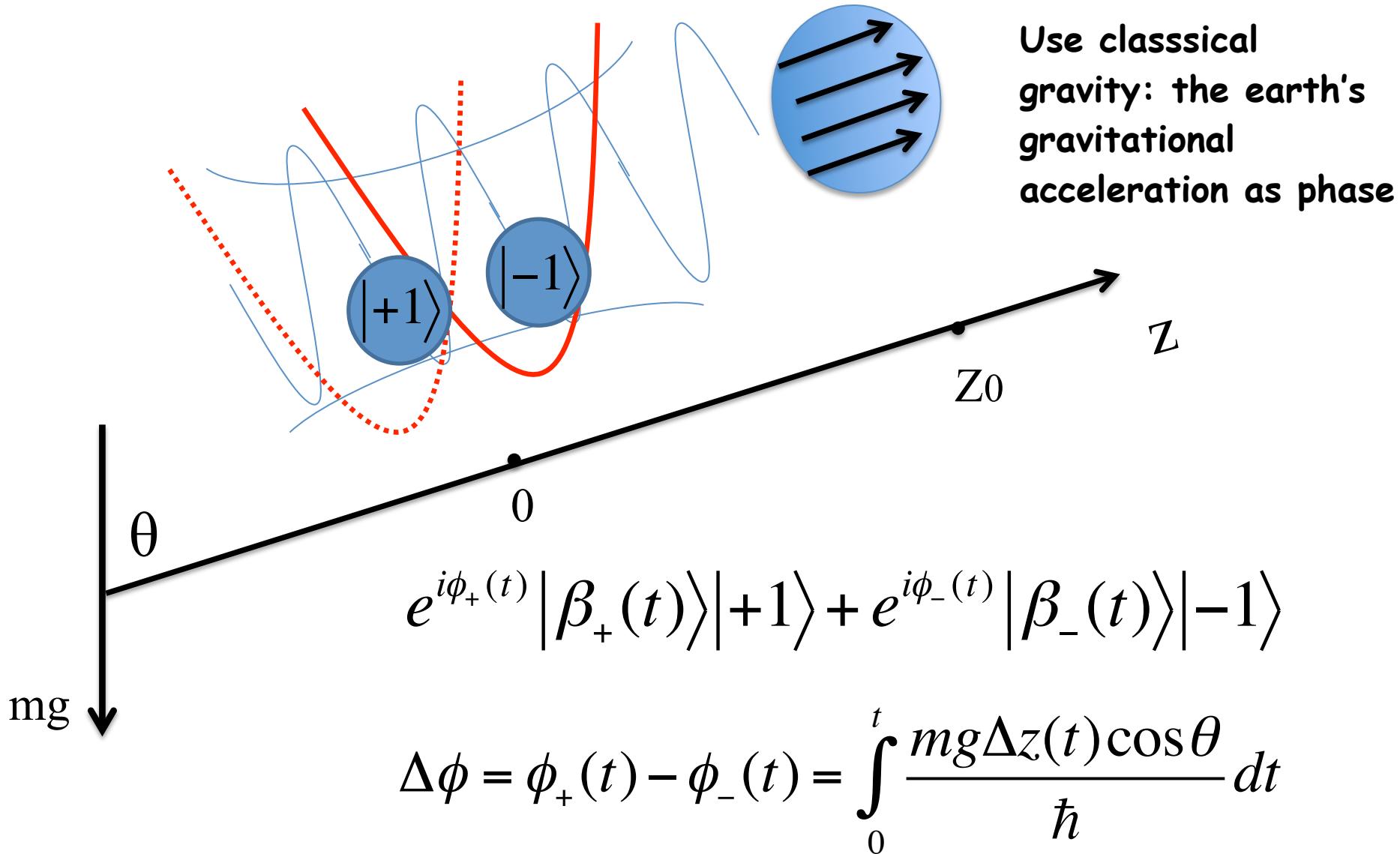
Diamond bead trapped in an optical trap. The bead contains a spin-1 NV center.



Time Evolution:

$$e^{i\phi_+(t)} |\beta_+(t)\rangle |+1\rangle + e^{i\phi_-(t)} |\beta_-(t)\rangle |-1\rangle$$

Ramsey Interferometry with a Levitated Thermal Mesoscopic Object



Measuring the relative phase shift between superposed components

Step 3: apply the same very rapid mw pulse as in step 1,

The presence of $\Delta\phi$ gives a modulation of the population of $|S_z=0\rangle$ according to:

$$|+1\rangle + e^{i\Delta\phi} |-1\rangle \rightarrow \cos \frac{\Delta\phi}{2} |0\rangle + \dots$$

For $m = 10^{10}$ amu (nano-crystal), superposition over 1 pm, the phase $\sim O(1)$

- M. Scala, M. S. Kim, G. W. Morley, P. F. Barker, S. Bose, Phys. Rev. Lett. **111**, 180403 (2013).
- Comment: F. Robicheaux, Phys. Rev. Lett. 118, 108901 (2017).
- Response: S. Bose et al, Phys. Rev. Lett. 118, 108902 (2017).

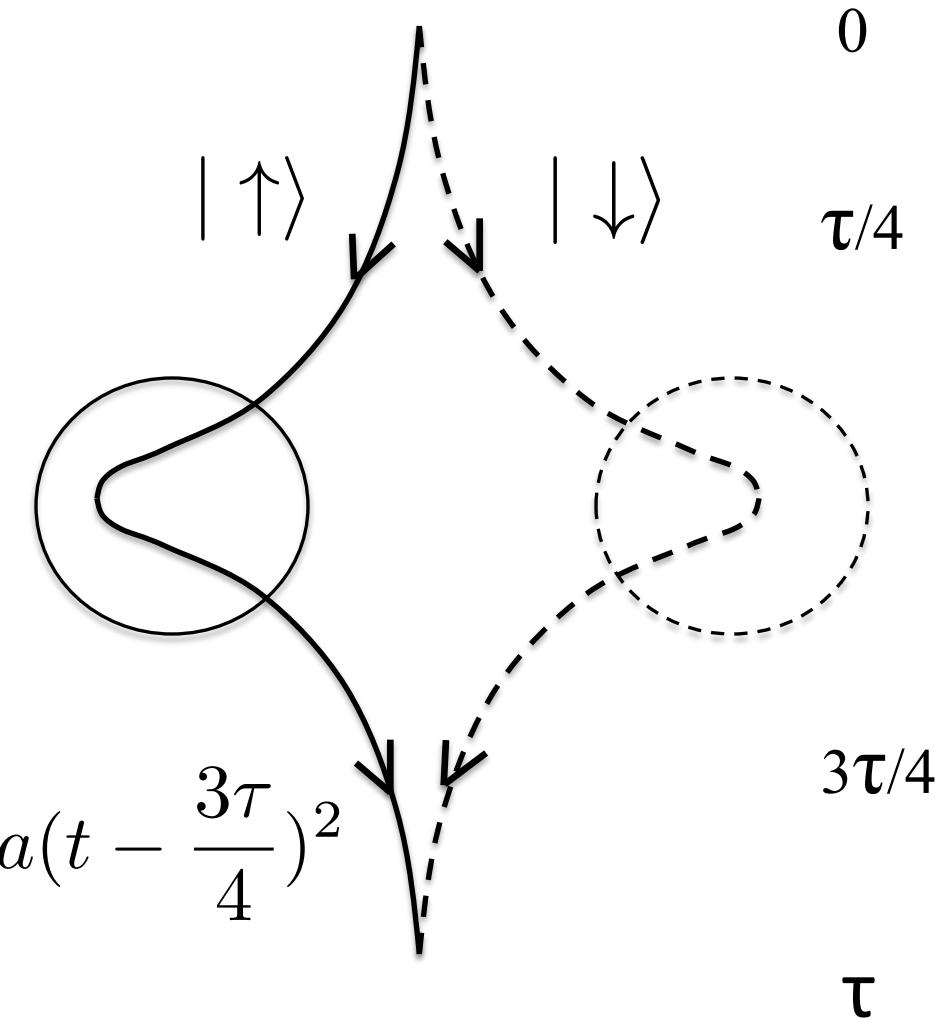
How can we increase the scale of the superposition?

Free particle in an inhomogeneous magnetic field (acceleration $+a$ or $-a$)

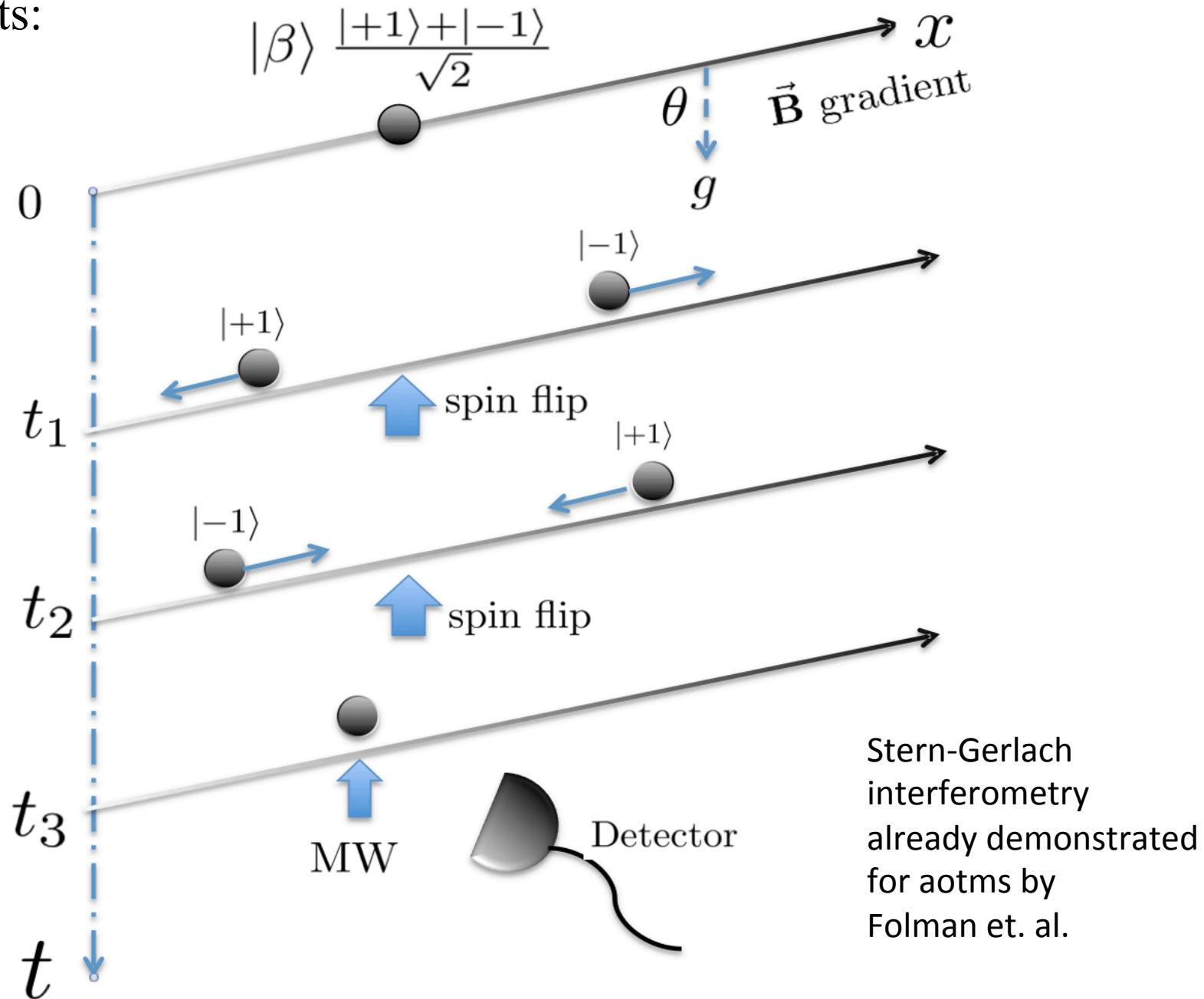
$$x_\sigma(t, j) = x_j(0) \pm \frac{1}{2}at^2$$

$$= \frac{a\tau}{4}(t - \frac{\tau}{4}) \mp \frac{1}{2}a(t - \frac{\tau}{4})^2$$

$$= \frac{1}{2}a(\frac{\tau}{4})^2 \mp \frac{a\tau}{4}(t - \frac{3\tau}{4}) \pm \frac{1}{2}a(t - \frac{3\tau}{4})^2$$



Free flight scheme able to achieve 100 nm separation among superposed components:



$$|\Psi(t_3)\rangle = \frac{1}{\sqrt{2}} |\psi(t_3)\rangle (|+1\rangle + e^{-i\phi_g} |-1\rangle)$$

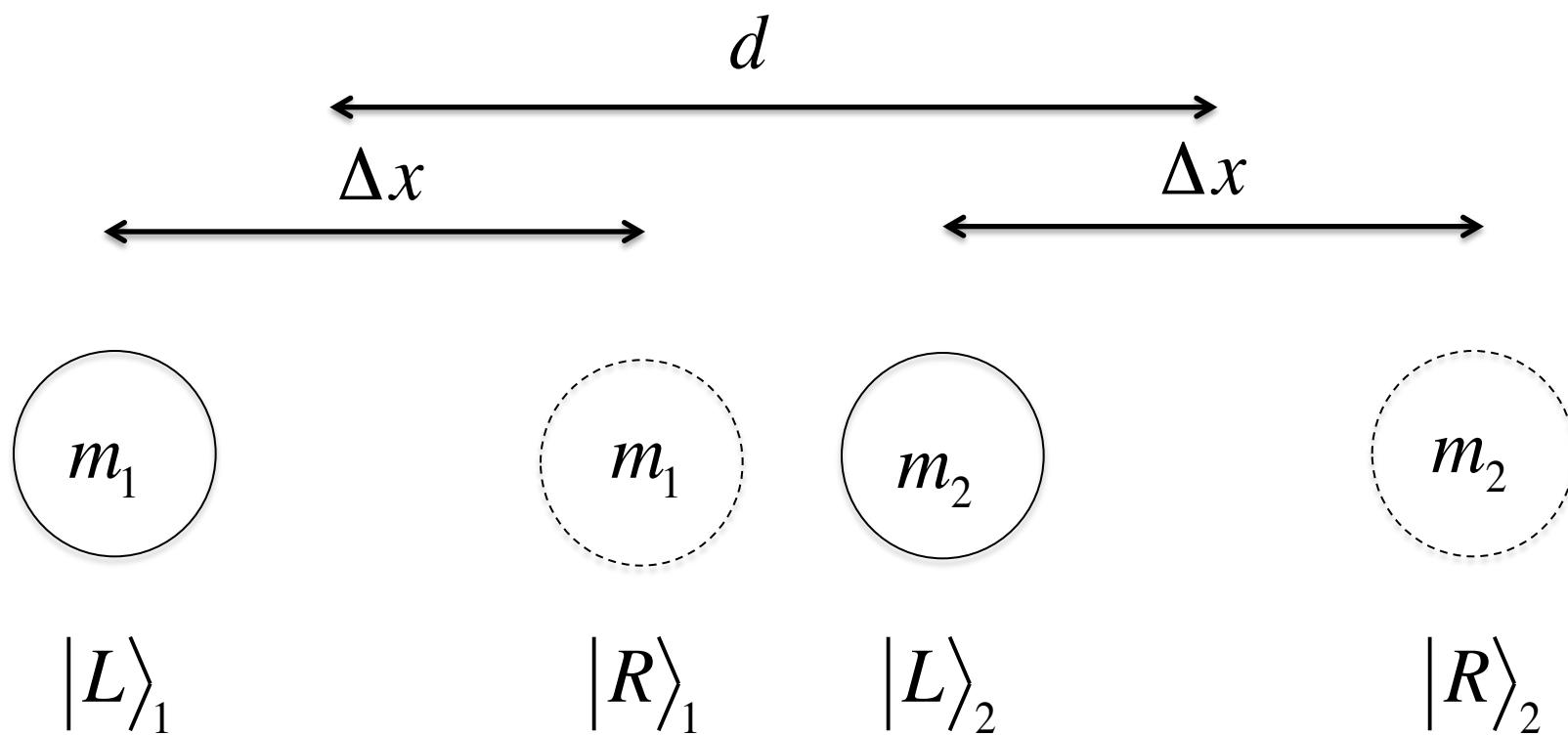
$$\langle x | \psi(t_3) \rangle = e^{-ip_0x} e^{-[(x-x_0-p_0t_3/m-g\cos\theta t_3^2/2)^2/2(\sigma')^2]}$$

$$\phi_g = (1/16\hbar) g t_3^3 g_{\text{NV}} \mu_B (\partial B / \partial x) \cos \theta$$

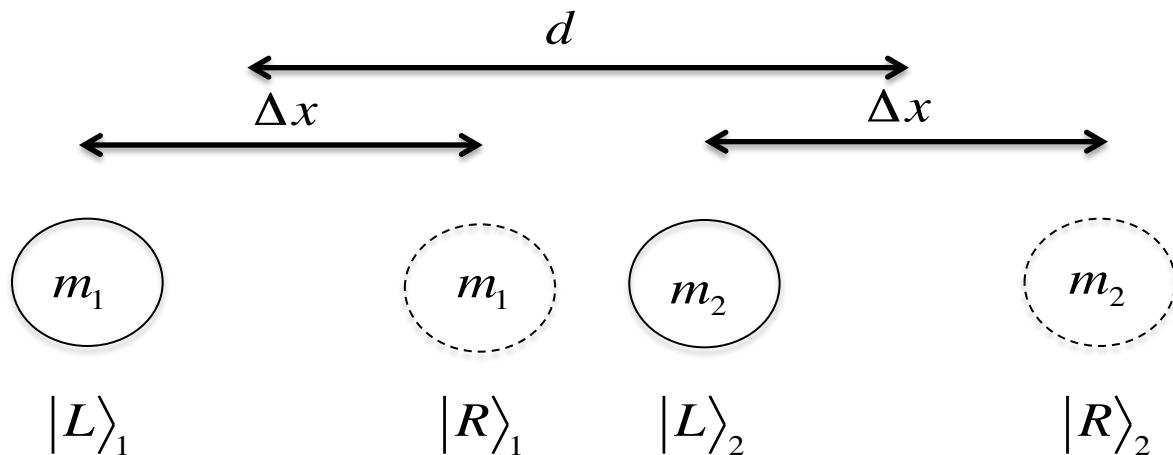
$$\Delta x_M = 2 \times \frac{1}{2m} g_{\text{NV}} \mu_B \frac{\partial B}{\partial x} (t_3/4)^2$$

10¹⁰ amu mass can be placed in a superposition of states separated by 100 nm.

A Schematic of two matter-wave interferometers near each other



Consider two neutral test masses *held* in a superposition, each exactly as a path encoded qubit (states $|L\rangle$ and $|R\rangle$), near each other.



$$|\Psi(t=0)\rangle_{12} = \frac{1}{\sqrt{2}}(|L\rangle_1 + |R\rangle_1) \frac{1}{\sqrt{2}}(|L\rangle_2 + |R\rangle_2)$$

$$= \frac{1}{2}(|L\rangle_1|L\rangle_2 + |L\rangle_1|R\rangle_2 + |R\rangle_1|L\rangle_2 + |R\rangle_1|R\rangle_2)$$

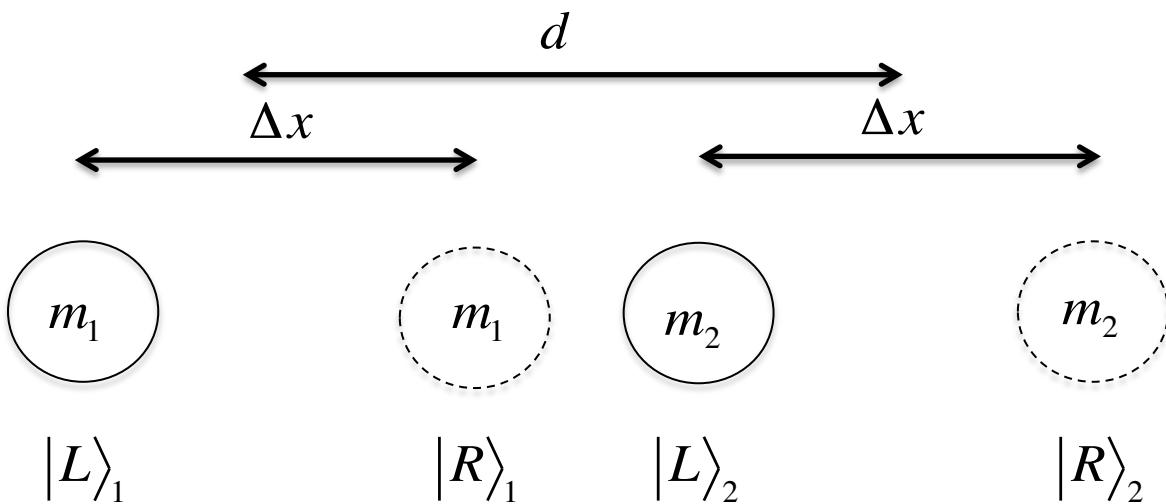
$$\rightarrow |\Psi(t=\tau)\rangle_{12} = \frac{1}{2}(e^{i\phi_{LL}}|L\rangle_1|L\rangle_2 + e^{i\phi_{LR}}|L\rangle_1|R\rangle_2 + e^{i\phi_{RL}}|R\rangle_1|L\rangle_2 + e^{i\phi_{RR}}|R\rangle_1|R\rangle_2),$$

where

$$\phi_{RL} \sim \frac{Gm_1m_2\tau}{\hbar(d - \Delta x)}, \quad \phi_{LR} \sim \frac{Gm_1m_2\tau}{\hbar(d + \Delta x)},$$

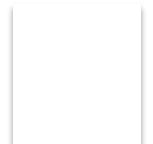
$$\phi_{LL} = \phi_{RR} \sim \frac{Gm_1m_2\tau}{\hbar d}$$

If they interact *only* through the gravitational force

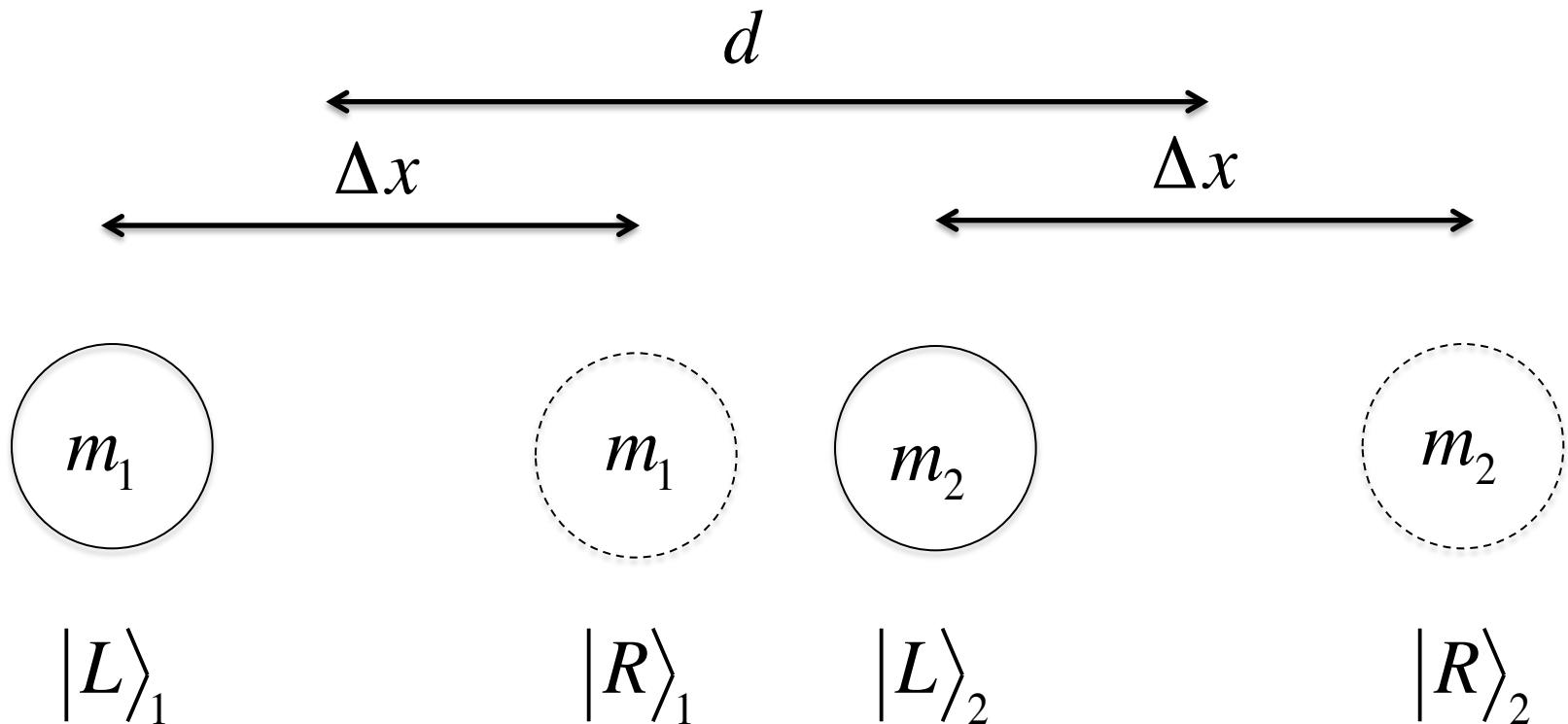


If they interact *only* through the gravitational force

$$\begin{aligned}
 |\Psi(t = \tau)\rangle_{12} &= \frac{1}{2}(e^{i\phi_{LL}}|L\rangle_1|L\rangle_2 + e^{i\phi_{LR}}|L\rangle_1|R\rangle_2 \\
 &\quad + e^{i\phi_{RL}}|R\rangle_1|L\rangle_2 + e^{i\phi_{RR}}|R\rangle_1|R\rangle_2) \\
 &= \frac{e^{i\phi_{RR}}}{\sqrt{2}}\{|L\rangle_1\frac{1}{\sqrt{2}}(|L\rangle_2 + e^{i\Delta\phi_{LR}}|R\rangle_2) \\
 &\quad + |R\rangle_1\frac{1}{\sqrt{2}}(e^{i\Delta\phi_{RL}}|L\rangle_2 + |R\rangle_2)\}
 \end{aligned}$$



The above state is maximally entangled when $\Delta\phi_{LR} + \Delta\phi_{RL} \sim \pi$.



For

$$d - \Delta x \ll d, \Delta x,$$

we have

$$\Delta\phi_{RL} \sim \frac{Gm_1m_2\tau}{\hbar(d - \Delta x)} \gg \Delta\phi_{LR}, \Delta\phi_{LL}, \Delta\phi_{RR}$$

For

$$d - \Delta x \ll d, \Delta x,$$

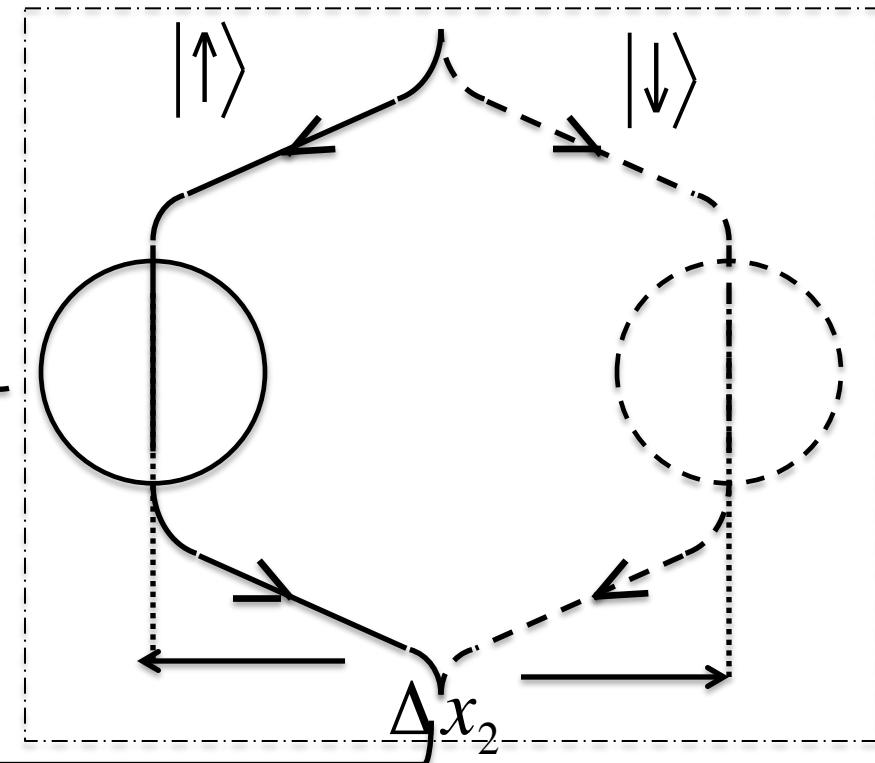
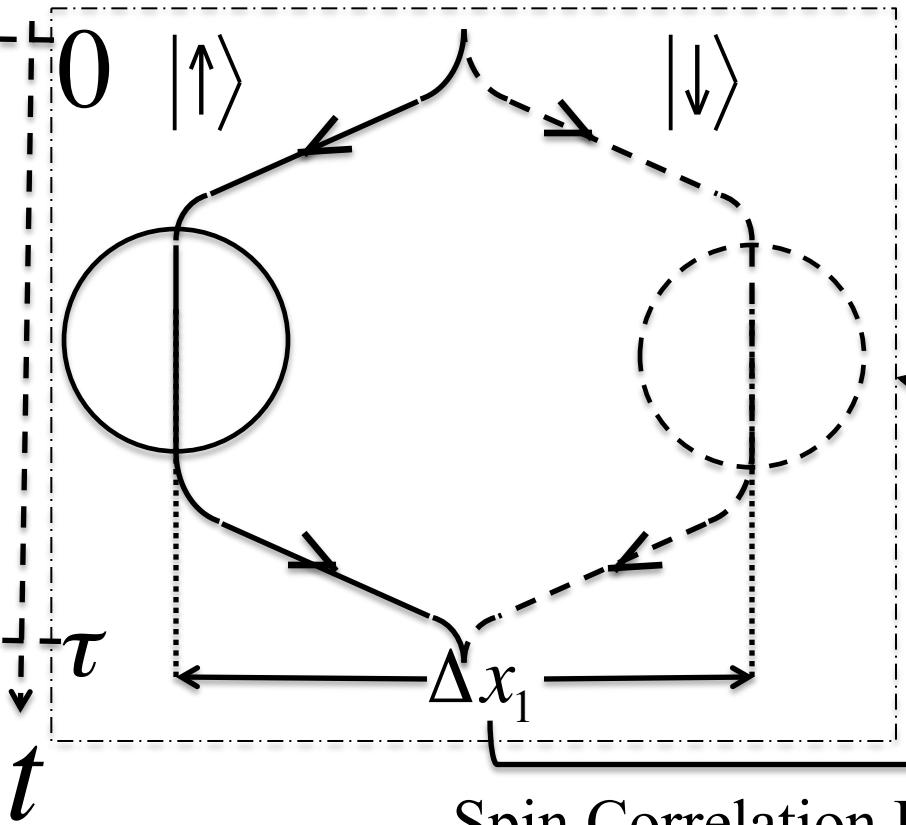
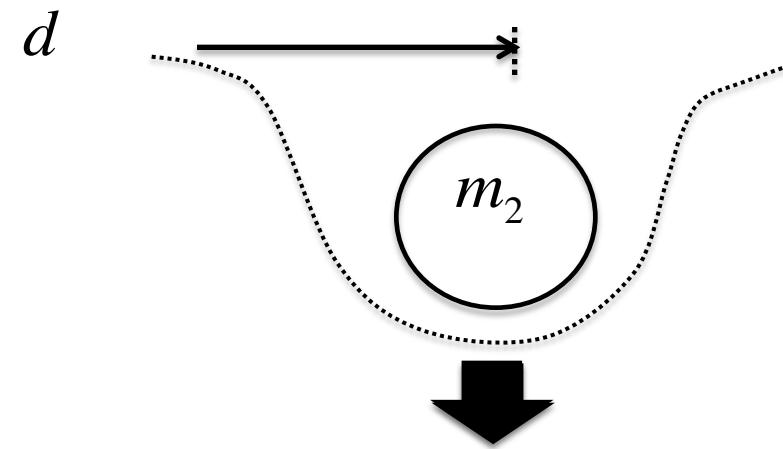
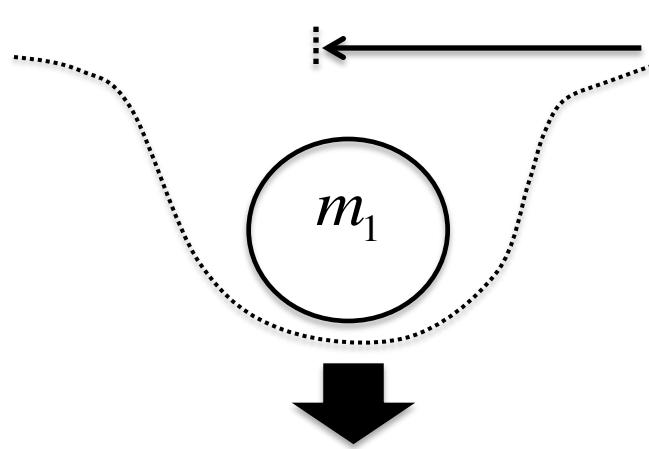
we have

$$\Delta\phi_{RL} \sim \frac{Gm_1m_2\tau}{\hbar(d - \Delta x)} \gg \Delta\phi_{LR}, \Delta\phi_{LL}, \Delta\phi_{RR}$$

For mass $\sim 10^{-14}$ kg (microspheres), separation at closest approach of the masses ~ 200 microns (to prevent Casimir interaction), **time ~ 1 seconds**, gives:

Scale of superposition ~ 100 microns, **Delta phi_{RL} ~ 1**

Planck's Constant fights Newton's Constant!



Spin Correlation Functions Certifying Entanglement

Spin Entanglement Witness:

Step 1: SG splitting:

$$|C\rangle_j \frac{1}{\sqrt{2}}(|\uparrow\rangle_j + |\downarrow\rangle_j) \rightarrow \frac{1}{\sqrt{2}}(|L,\uparrow\rangle_j + |R,\downarrow\rangle_j)$$

Step 2: Gravitational interaction induced phase accumulation on
the joint states of masses 1 & 2 (*mapped to nuclear spins*)

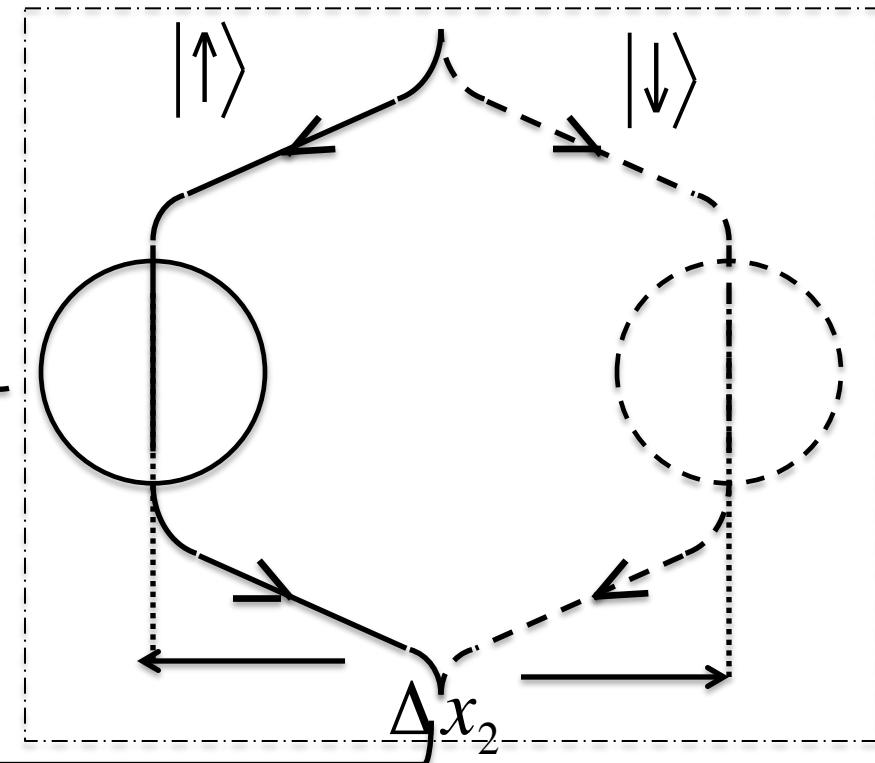
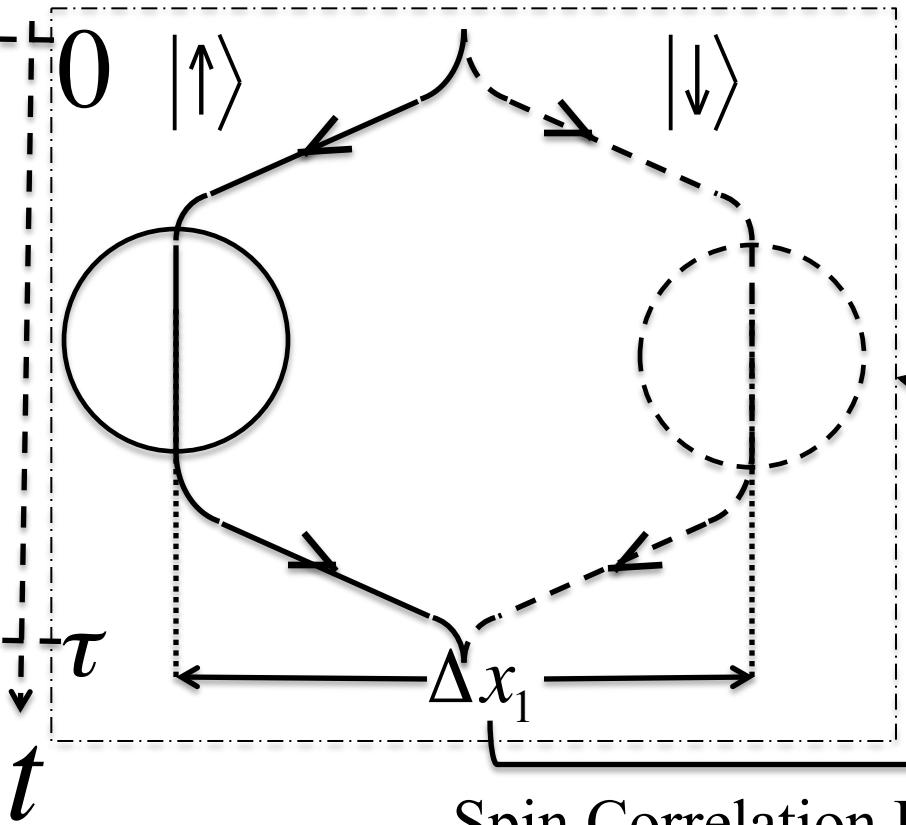
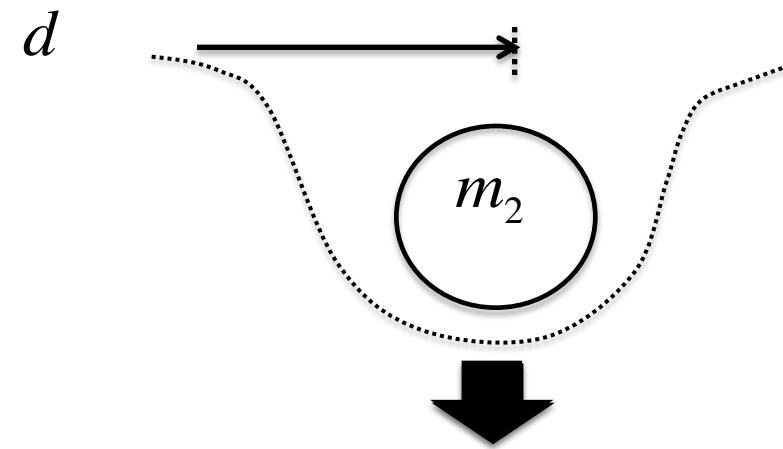
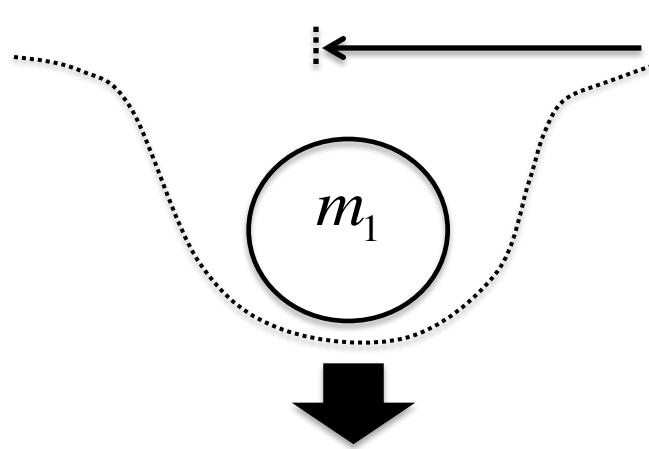
Step 3: SG recombination: $|L,\uparrow\rangle_j \rightarrow |C,\uparrow\rangle_j$, $|R,\downarrow\rangle_j \rightarrow |C,\downarrow\rangle_j$

Step 4: Witness spin entangled state:

$$\begin{aligned} |\Psi(t = t_{\text{End}})\rangle_{12} = & \frac{1}{\sqrt{2}}\{|\uparrow\rangle_1 \frac{1}{\sqrt{2}}(|\uparrow\rangle_2 + e^{i\Delta\phi_{LR}}|\downarrow\rangle_2) \\ & + |\downarrow\rangle_1 \frac{1}{\sqrt{2}}(e^{i\Delta\phi_{RL}}|\uparrow\rangle_2 + |\downarrow\rangle_2)\}|C\rangle_1|C\rangle_2 \end{aligned}$$

through the correlations:

$$\mathcal{W} = |\langle \sigma_x^{(1)} \otimes \sigma_z^{(2)} \rangle - \langle \sigma_y^{(1)} \otimes \sigma_z^{(2)} \rangle|$$



Spin Correlation Functions Certifying Entanglement

How is this related to Quantum Gravity?

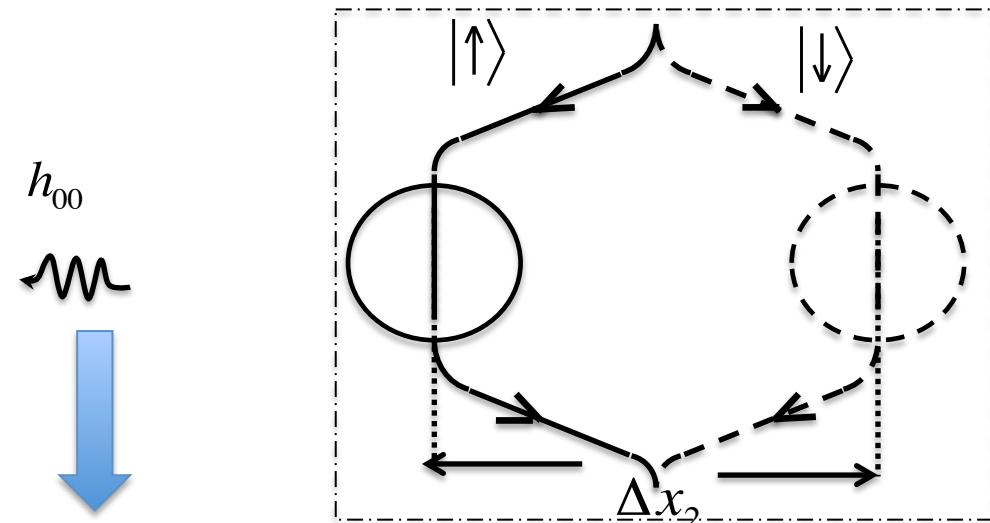
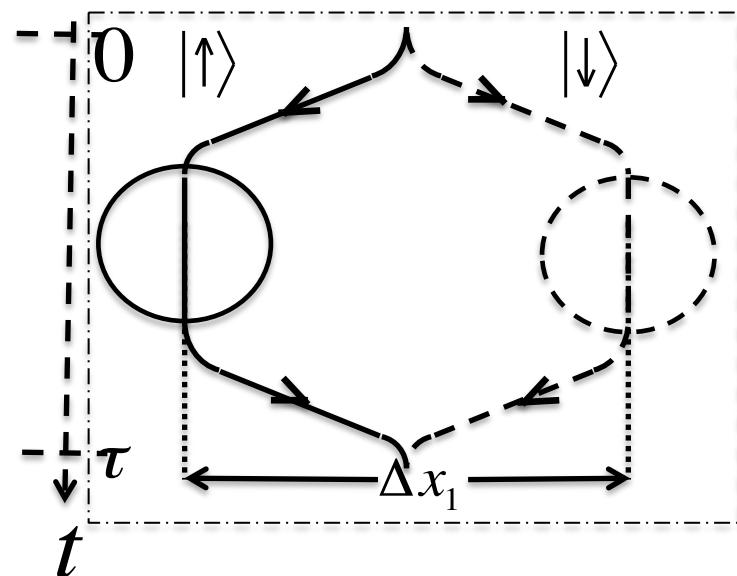
LOCC Maps keep separable states separable (*cannot* create entanglement!)



Local Operations and Classical Communication (LOCC)

- 1. Unitary evolution
- 2. Measurement

Classical mediator of information/bits

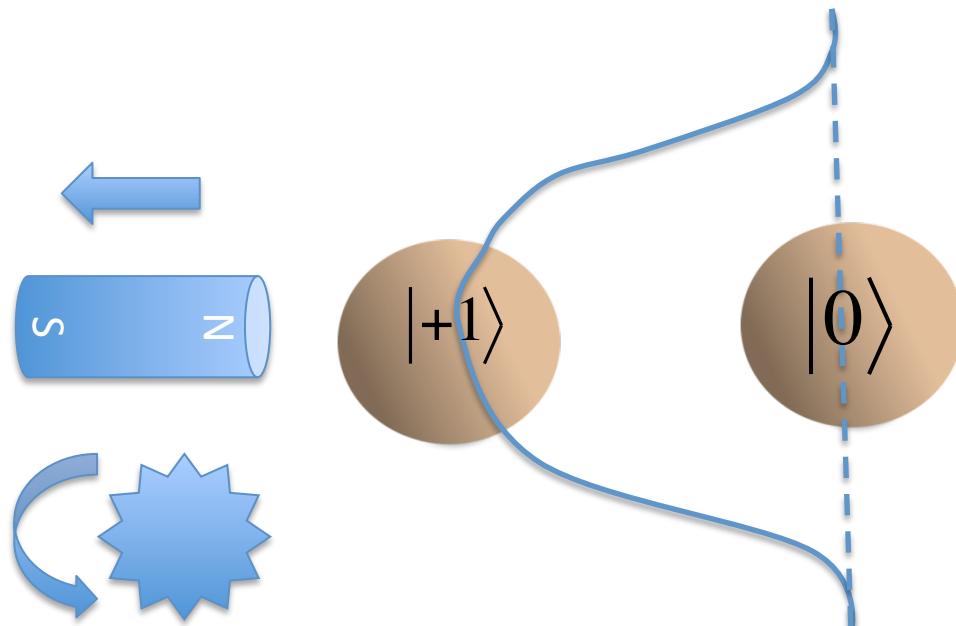


Must be quantum if the spins in the masses get entangled

Challenges (technical):

~ 200 micron superposition:

Literally pull one of the spin components?



Gavin Morley's idea

Other demands:

1. Spin the masses to average charge multipoles.
2. Internal cooling to 77 K, External pressure 10^{-15} Pascal, 0.15 K temperature

What does it imply in the context of **low energy effective field theory**?

$$\mathcal{H} = \sum_{j,\xi} m_j c^2 a_{j,\xi}^\dagger a_{j,\xi} + \sum_{\mathbf{k}} \hbar \omega_k b_{\mathbf{k}}^\dagger b_{\mathbf{k}}$$

$$- \hbar \sum_{j,\mathbf{k},\xi} g_{j,\mathbf{k}} a_{j,\xi}^\dagger a_{j,\xi} (b_{\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{r}_{j,\xi}} + b_{\mathbf{k}}^\dagger e^{-i\mathbf{k}\cdot\mathbf{r}_{j,\xi}})$$

Coherent States of the gravitational field

Superposition

$$|\Psi(t)\rangle_{\text{mat+grav}} = \frac{1}{2} \sum_{\xi, \xi' \in \{L, R\}} a_{1,\xi}^\dagger a_{2,\xi'}^\dagger |0\rangle$$

$$\otimes \prod_{\mathbf{k}} e^{i \frac{(g_{1,\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{r}_{1,\xi}} + g_{2,\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{r}_{2,\xi'}})^2}{\omega_k} t} |\alpha_{\mathbf{k},\xi,\xi'}\rangle_{\mathbf{k}},$$

$$g_{j,\mathbf{k}} = m_j c^2 \sqrt{\frac{2\pi G}{\hbar c^3 kV}}$$

$$\frac{g_{1,\mathbf{k}} g_{2,\mathbf{k}}}{\omega_k} \propto \frac{1}{k^2}$$

Superpositions of *distinct* (?) coherent states of the gravitational field

Summary

- Large mass, small scale of superpositions:

Stern-Gerlach based Ramsey interferometry in a trap:

M. Scala, M. S. Kim, G. W. Morley, P. F. Barker, S. Bose, Phys. Rev. Lett. **111, 180403 (2013).** [related work by Tongcang Li et. al.]

- Large mass, large scale superpositions:

Free flight Stern-Grlach based Ramsey interferometry:

C. Wan, M. Scala, G. W. Morley, ATM. A. Rahman, H. Ulbricht, J. Bateman, P. F. Baker, S. Bose, M. S. Kim, Phys. Rev. Lett. 117, 143003 (2016).

- *Spin Entanglement Witness for Quantum Gravity:*

S. Bose, A. Mazumdar, G. W. Morley, H. Ulbricht, M. Toros, M. Paternostro, P. F. Barker, A. Geraci, M. S. Kim, G. J. Milburn, Phys. Rev. Lett. 119, 240401 (2017). Related work: C. Marletto and V. Vedral

